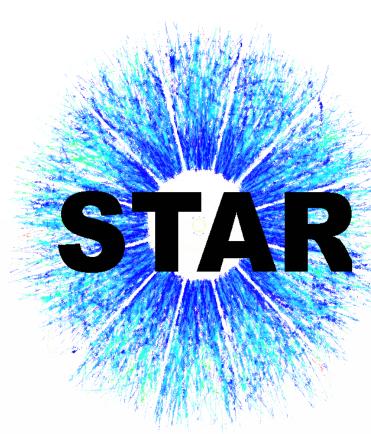
# System Size and Shape Dependence of Anisotropic Flow



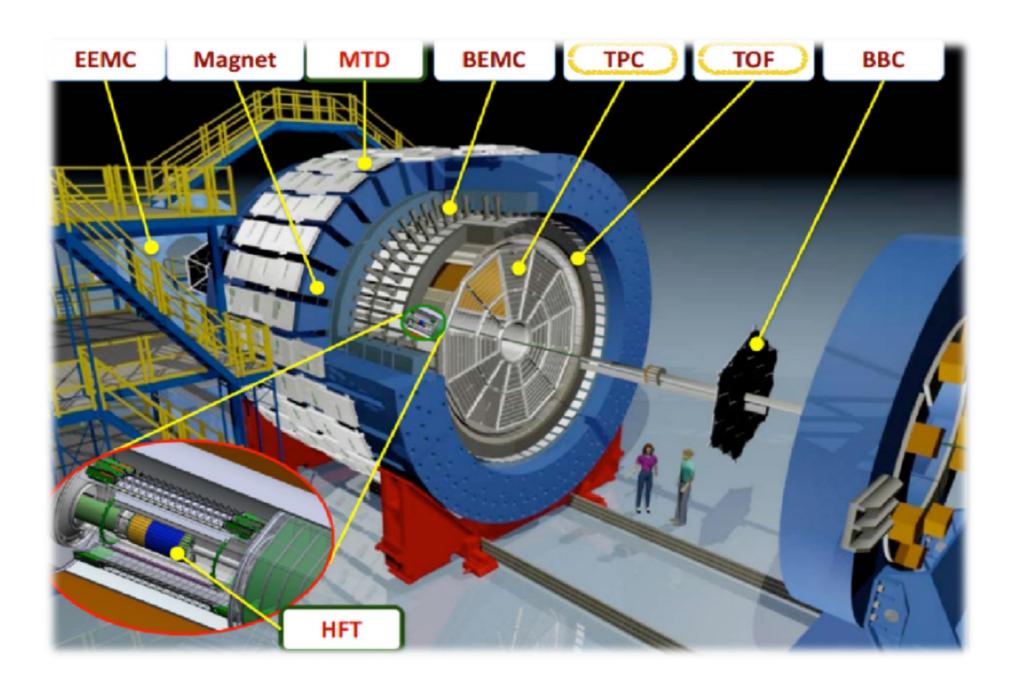
# Niseem Magdy, for the STAR Collaboration



### **Abstract**

In this work, we studied the first three flow harmonics,  $v_1^{even}$ ,  $v_2$  and  $v_3$ , as a function of mean multiplicity, (Mult), in U+U, Au+Au, Cu+Au, Cu+Cu, d+Au and p+Au collisions at  $\sqrt{s_{NN}}$  ~200 GeV. The measurements confirm the impacts of initial geometry (shape and dimensionless size) on the flow harmonics. Such an effect is consistent with the dispersion relation for sound propagation in the hot and dense medium created in these collisions. Our measurements indicate that  $v_1^{even}$  and  $v_3$ are system independent and the scaled  $v_2$  shows a common trend for all systems.

#### **STAR Detector**



Uniform acceptance in  $|\eta| < 1$ 

### Two particle correlation function $Cr(\Delta \varphi)$ used in this analysis

$$Cr(\Delta\varphi) = \frac{dN/d\Delta\varphi(same)}{dN/d\Delta\varphi(mix)}$$

➤ Non-flow signals, as well as some residual detector effects suppressed with  $|\Delta \eta| > 0.7$  cut.

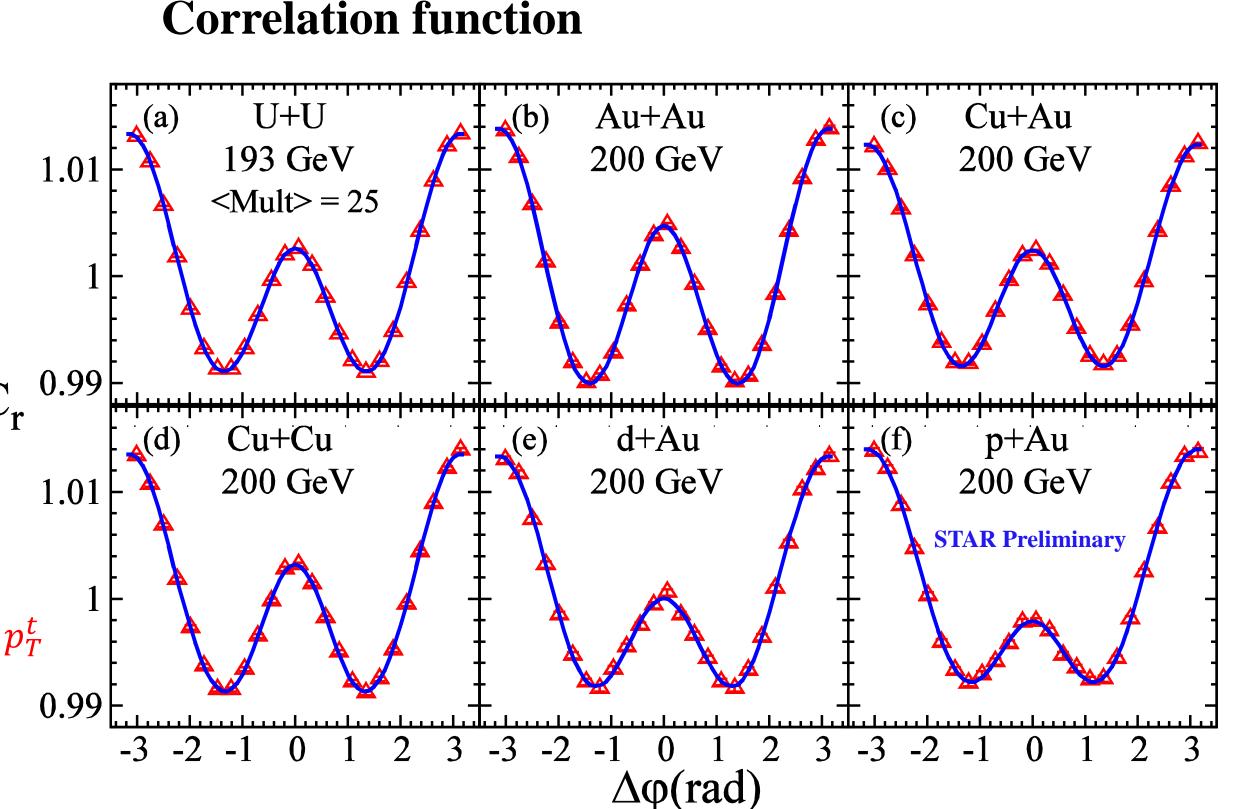
For n > 1,

$$v_{\mathrm{nn}}(p_T^a, p_T^t) = v_{\mathrm{n}}(p_T^a) \ v_{\mathrm{n}}(p_T^t)$$

For n = 1,

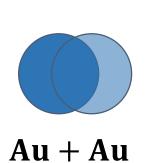
$$v_{11}(p_T^a, p_T^t) = v_1^{even}(p_T^a)v_1^{even}(p_T^t) - C p_T^a p_T^t$$

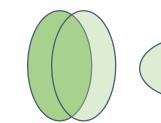
C is the momentum conservation parameter  $C \propto \frac{1}{\langle \text{Mult} \rangle \langle p_T^2 \rangle} [4].$ 



#### **Motivation**

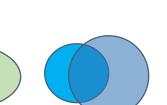
- ➤ Is the observed anisotropy in ion—ion collision a final- or initial- state effect?
- > STAR collectd data for different systems;

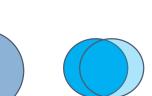


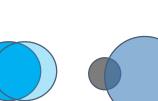


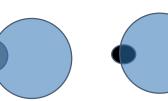
U + U

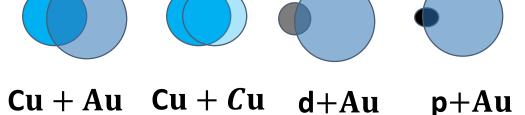












## Final-state ansatz

- $\triangleright$  The  $v_n$  measurements are sensitive to  $\varepsilon_n$ , RT and  $\left(\frac{\eta}{s}, \frac{\zeta}{s}, \dots\right)$  [1-2].
- > Acoustic ansatz
  - ✓ Sound attenuation in the viscous matter reduces the magnitude of  $v_n[1]$ .
- > Anisotropic flow attenuation;

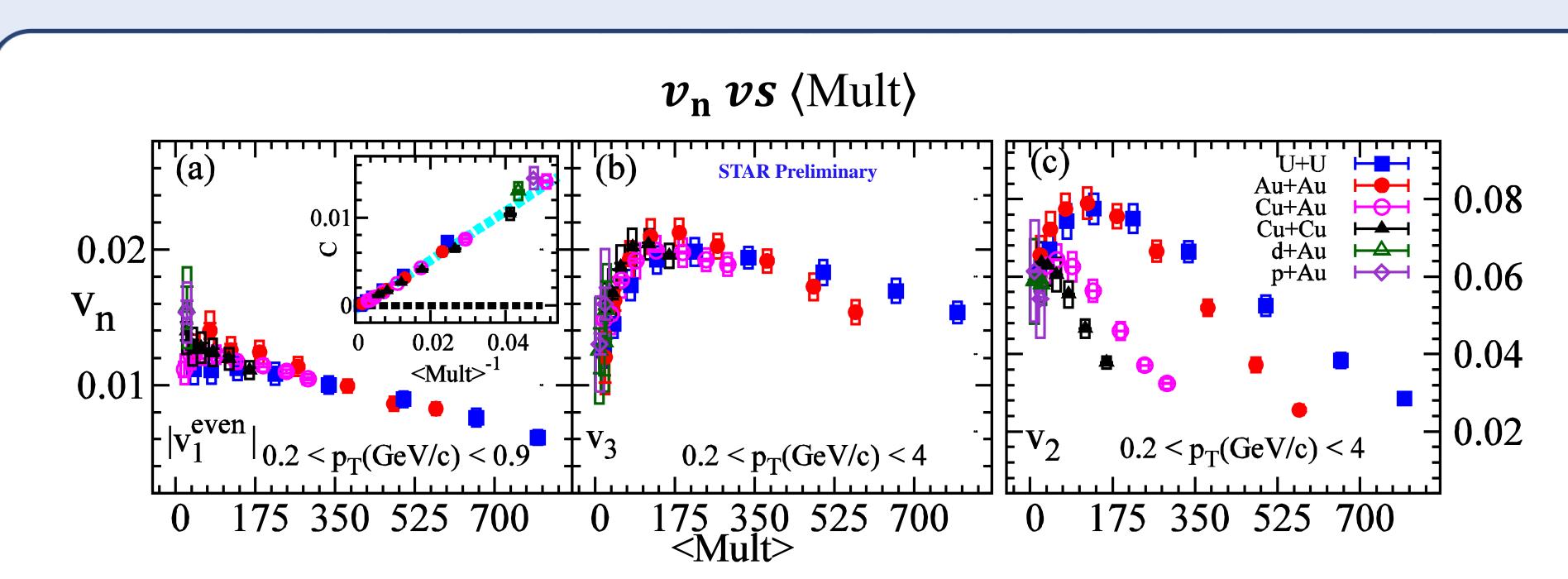
$$\frac{v_n}{\varepsilon_n} \propto e^{-\beta n^2}, \ \beta \propto \frac{\eta}{s} \frac{1}{RT} + \cdots$$

From macroscopic entropy considerations  $(RT)^3 \propto$  $\frac{dN}{d\eta}$  [3].

$$ln\left(\frac{v_n}{\varepsilon_n}\right) \propto a \frac{\eta}{s} \left(\frac{dN}{d\eta}\right)^{\frac{-1}{3}}$$

$$ln(v_n) \propto a \left(\frac{\eta}{s}\right) \left(\frac{dN}{d\eta}\right)^{\frac{-1}{3}} + \ln(\varepsilon_n)$$

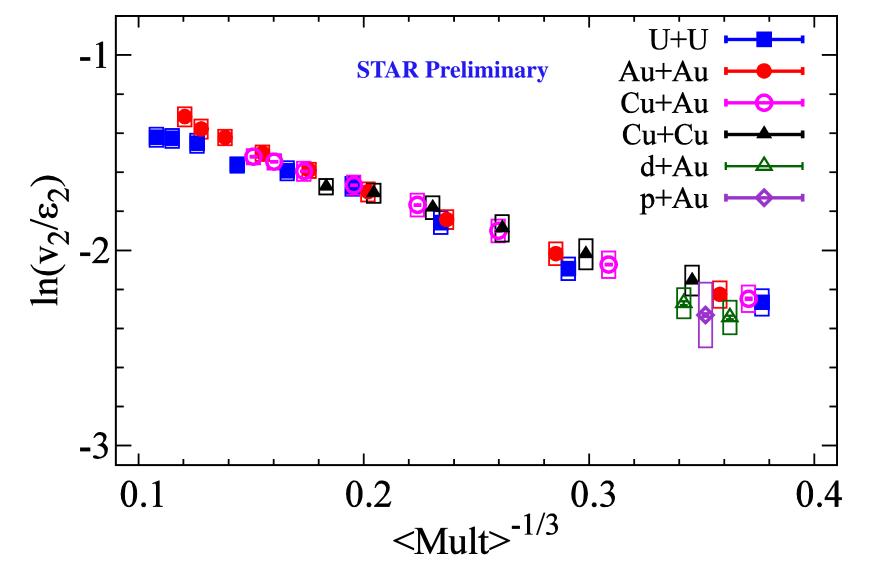
Scaling out the system size  $\left(\frac{dN}{d\eta}\right)$  and shape  $(\varepsilon_n)$ should give similar transport coefficient  $(\frac{\eta}{\epsilon})$  (i.e. similar  $v_n$ ) for different systems (final-state effect).



The measurements of  $v_1^{even}$ ,  $v_2$  and  $v_3$  as a function of (Mult) for U+U, Au+Au, Cu+Au, Cu+Cu, d+Au and p+Au collisions at  $\sqrt{s_{NN}} \sim 200$  GeV. For the same (Mult) or size,  $v_1^{even}$  and  $v_3$  are system independent, while  $v_2$  is system dependent.

## $v_2/\epsilon_2 \ vs < Mult >$

- $\triangleright$  The eccentricity-scaled  $v_2$  as a function of  $\langle Mult \rangle^{-1/3}$  for U+U, Au+Au, Cu+Au, Cu+Cu, d+Au and p+Au collisions at  $\sqrt{s_{NN}} \sim 200$  GeV.
- $\triangleright$  The scaled  $v_2$  shows a common trend for all systems.



### References

- [1] arXiv:1305.3341, Roy A. Lacey, A. Taranenko, J. Jia, et al.
- [2] PRC84 034908 (2011) P.Staig and E.Shuryak
- [3] arXiv:1601.06001, Roy A. Lacey,, et al.
- [4] PRC 86, 014907 (2012), ATLAS Collaboration

### Acknowledgment

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### Conclusion

- $\succ$  The two-particle correlation technique has been used to study  $v_1^{even}$ ,  $v_2$  and  $v_3$  as a function of (Mult) for U+U, Au+Au, Cu+Au, Cu+Cu, d+Au and p+Au collisions at  $\sqrt{s_{NN}} \sim 200$  GeV.
- $\triangleright$  At the same size ((Mult)),  $v_1^{even}$  and  $v_3$  are system independent, while  $v_2$  is system dependent.
- $\triangleright$  The scaled  $v_2$  shows a common trend for all systems.

