Investigating the Chiral Magnetic Wave at RHIC-STAR

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Abstract

In heavy-ion collisions, the Chiral Magnetic Wave (CMW) is theorized to produce 1 an electric quadrupole moment, leading to differences in elliptic flow between 2 positively and negatively charged particles. This CMW signal can be detected 3 by examining the correlation between charge-dependent elliptic flow and event л charge asymmetry. This study focuses on the difference in covariance of elliptic 5 flow (v_2) and charge asymmetry (A_{ch}) for positively and negatively charged par-6 ticles across various collision centralities. A comparison between the two isobar systems $\binom{96}{44}$ Ru + $\frac{96}{44}$ Ru and $\frac{96}{40}$ Zr + $\frac{96}{40}$ Zr) is conducted to investigate any potential 8 enhancement of the CMW signal in Ru+Ru collisions, which may be attributed 9 to the stronger magnetic field generated in Ru. Additionally, correlations between 10 triangular flow (v_3) and A_{ch} will be presented to assess background contributions. 11

Keywords:

Heavy-ion collisions, Quark Gluon Plasma, Chiral Magnetic Wave

1. Introduction

In non-central relativistic heavy-ion collisions, interaction of produced quarks 12 with topologically non-trivial gluon fields and the strong magnetic field generated 13 by spectator protons can give rise to novel phenomena such as the Chiral Magnetic 14 Effect (CME) and the Chiral Separation Effect (CSE). The CME refers to the sep-15 aration of electric charge along the magnetic field due to a nonzero axial charge 16 density, while the CSE is described as the separation of chiral charges along the 17 magnetic field in the presence of a finite vector charge density. The interplay be-18 tween these two effects induces an oscillatory propagation of electric and chiral 19 charge densities, leading to a long-wavelength collective excitation known as the 20 Chiral Magnetic Wave (CMW) [1]. The CMW results in the formation of an elec-21 tric quadrupole moment where the "poles" ("equator") of the produced fireball 22

²³ acquire additional positive (negative) charge, which leads to charge-dependent el-²⁴ liptic flow (v_2) . The isobar (Ru+Ru and Zr+Zr) collision systems, offers a unique ²⁵ opportunity to study the CMW signal. The stronger magnetic fields in Ru+Ru is ²⁶ expected to yield a larger signal, while both isobars have similar backgrounds.

2. Data selection and analysis method

The present analysis is based on ~ 1.6 billion minimum-bias events from each 27 Ru+Ru and Zr+Zr collisions at $\sqrt{s_{NN}}$ =200 GeV, recorded by STAR detector at 28 RHIC. Events are selected with primary vertex within $-35 < V_z < 25$ cm and 29 within 2 cm in the transverse plane. Tracks must have pseudorapidity range of $|\eta|$ 30 < 1 and distance of closest approach (DCA) to the primary vertex less than 3 cm. 31 The CMW induced quadrupole leads to the difference in v_2 of negatively and 32 positively charged particles, Δv_2 , which is expected to be proportional to the 33 charge asymmetry (A_{ch}) , i.e., $\Delta v_2 = v_2^- - v_2^+ \approx r_2 A_{ch}$. Here r_2 denotes the slope 34 parameter of Δv_2 versus A_{ch} and $A_{ch} = (N_+ - N_-)/(N_+ + N_-)$ with $N_+ (N_-)$ denoting 35 the number of positive (negative) particles in a given event. In addition to measur-36 ing v_n^{\pm} as a function of A_{ch} , the covariance of v_n^{\pm} and A_{ch} , i.e., $\langle v_n^{\pm} A_{ch} \rangle - \langle v_n^{\pm} \rangle \langle A_{ch} \rangle$ 37 also known as three particle integral correlator can also be measured to quantify 38 strength of CMW signal [2]. It is more robust against detector efficiency correc-39 tion and statistically reduce fluctuations as sample need not be divided into further 40 $A_{\rm ch}$ bins to calculate slope of Δv_2 versus $A_{\rm ch}$. To study CMW signal equivalently to 41 r_2 parameter, an alternative observable, the delta integral covariance can be used, 42 defined as $\Delta IC = (\langle v_n^- A_{ch} \rangle - \langle v_n^- \rangle \langle A_{ch} \rangle) - (\langle v_n^+ A_{ch} \rangle - \langle v_n^+ \rangle \langle A_{ch} \rangle)$, where the angular 43 bracket denotes the average over the events. The Q-cumulant method [3] with η 44 gap of 0.3 between particle of interest and reference particles is used to calculate 45 anisotropic flow coefficients (v_n) . 46

3. Results

Fig. 1 (left) shows covariance of v_2 and A_{ch} for positive and negative charges 47 in Ru+Ru and Zr+Zr collisions versus centrality. Both systems exhibit similar 48 trends, covariance increases from central to peripheral collisions for both positive 49 and negative particles. ΔIC also increases from central to peripheral collisions. 50 Fig. 1 (right) presents $\Delta IC/\sigma_{A_{ch}}^2$ (for v_2), the values from both the systems agree 51 within errors, no enhanced CMW signal is seen in Ru+Ru over Zr+Zr despite Ru 52 having four more protons compared to Zr. The covariance of v_3 and A_{ch} for posi-53 tive and negative charges (Fig. 2 left) and $\Delta IC / \sigma_{A_{ch}}^2$ (Fig. 2 right) are independent 54 of collision centrality. 55



Figure 1: Comparison of covariance (Left) between v_2 and A_{ch} , and (Right) $\Delta IC/\sigma_{A_{ch}}^2$ for v_2 in Ru+Ru and Zr+Zr collisions. Points are horizontally shifted for clarity.



Figure 2: Comparison of covariance (Left) between v_3 and A_{ch} , and (Right) $\Delta IC/\sigma_{A_{ch}}^2$ for v_3 in Ru+Ru and Zr+Zr collisions. Points are horizontally shifted for clarity.

4. Summary

The charge-dependent correlation between elliptic flow and A_{ch} is investigated 56 in Ru+Ru and Zr+Zr collisions at $\sqrt{s_{NN}}$ =200 GeV. The covariance of v_2 and A_{ch} 57 increases from central to peripheral collisions for both positive and negative par-58 ticles. However, no charge-dependent effects are observed in the third-harmonic. 59 No enhancement of the CMW signal is found in Ru+Ru collisions, despite the 60 stronger magnetic field in Ru+Ru collisions. To further explore the CMW con-61 tribution, techniques such as Event Shape Engineering will be used to estimate 62 CMW signal. 63

References

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