

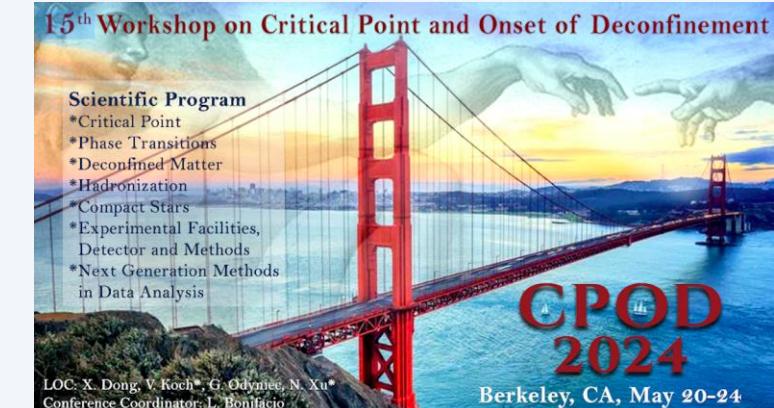


# Measurements of d- $\Lambda$ correlations from STAR

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for the STAR collaboration

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U.S. DEPARTMENT OF  
**ENERGY** 2024.05.22

# Outlines

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- ❖ Introductions & Lednicky-Lyuboshitz (L-L) approach
- ❖ Particle identification
- ❖ p- $\Lambda$  & d- $\Lambda$  correlation function
  - ❖ Source size with L-L approach
  - ❖ Correlation function & spin states
  - ❖ Scatterings length ( $f_0$ ) and effective range ( $d_0$ )
  - ❖  $\Lambda$  separation energy of  $^3\Lambda$ H
- ❖ Summary & Outlooks

# QCD Dense Matter & Nucleon-Nucleon/Hyperon Interactions

## INSIDE A NEUTRON STAR

A NASA mission will use X-ray spectroscopy to gather clues about the interior of neutron stars — the Universe's densest forms of matter.

### Outer crust

Atomic nuclei, free electrons

### Inner crust

Heavier atomic nuclei, free neutrons and electrons

### Outer core

Quantum liquid where neutrons, protons and electrons exist in a soup

### Inner core

Unknown ultra-dense matter. Neutrons and protons may remain as particles, break down into their constituent quarks, or even become 'hyperons'.

### Atmosphere

Hydrogen, helium, carbon

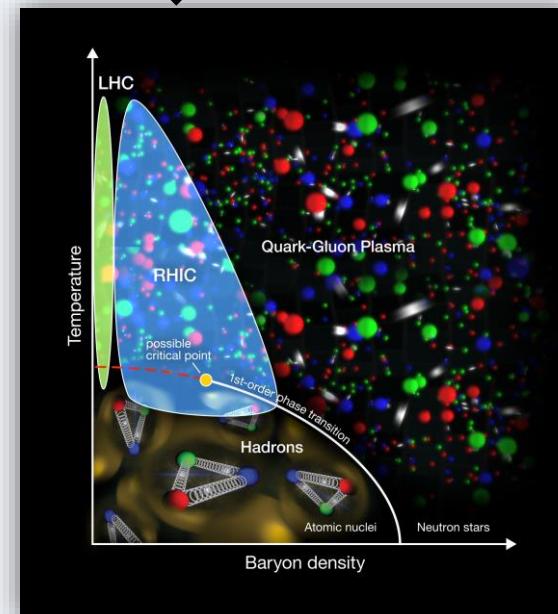
Beam of X-rays coming from the neutron star's poles, which sweeps around as the star rotates.

Credit: Source: Adapted from NASA Goddard SVS

Nature volume 546, page18 (2017)

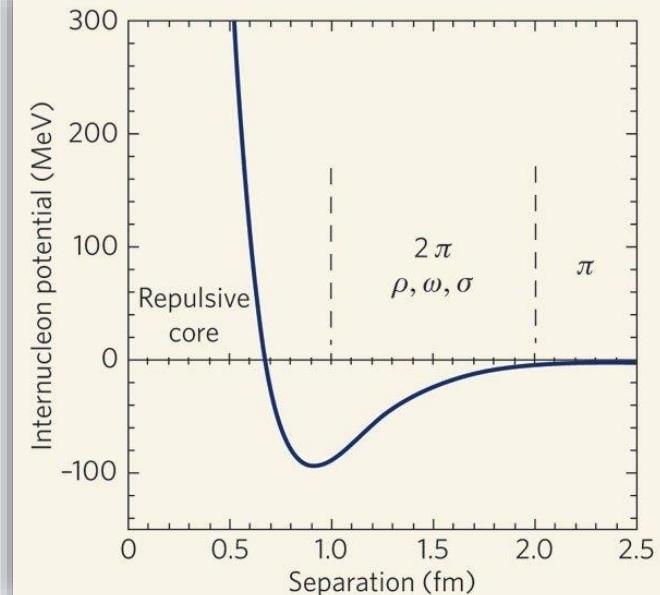
## State of Matters

Phase diagram



## Interaction of Matters

Nuclear potential



- ❖ Structure of nuclear and hyper-nuclei matter
- ❖ Role of Nucleon-Nucleon (N-N) and Hyperon-Nucleon (Y-N) interactions in the Equation-of-State

<https://www.bnl.gov/newsroom/news.php?a=219079>

<https://www.quora.com/What-does-the-potential-function-for-the-strong-nuclear-force-look-like>

# Low-E scattering experiment & Effective Range Expansion

Low energy elastic scatterings:

$$k \cot(\delta(k)) = -\frac{1}{a} + \frac{1}{2} r_0 k^2 + O(k^4)$$

$a$ : phase shift

$r_0$ : Fermi scattering length at zero energy

$O$ : effective range

$d\sigma$ : higher order contribution

Cross section:

$$\lim_{k \rightarrow 0} \sigma_e = 4\pi a^2$$

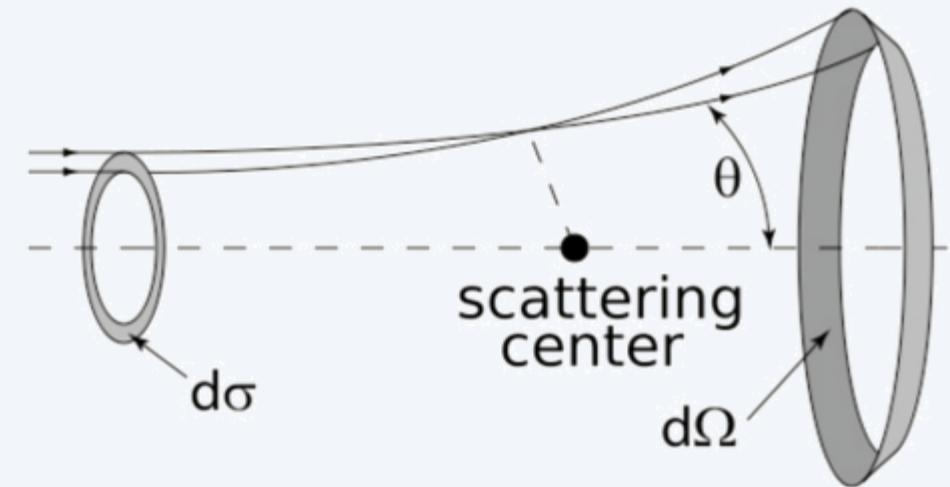
Binding energy:

$$\frac{1}{a} = \gamma - \frac{1}{2} r_0 \gamma^2$$

$$\diamond B = \frac{\gamma^2}{2\mu}$$

$\diamond \mu$ : reduced mass

$\diamond \gamma$ : binding momentum



H. A. Bethe, Phys. Rev. 76 (1949) 38

For the n-p scattering:

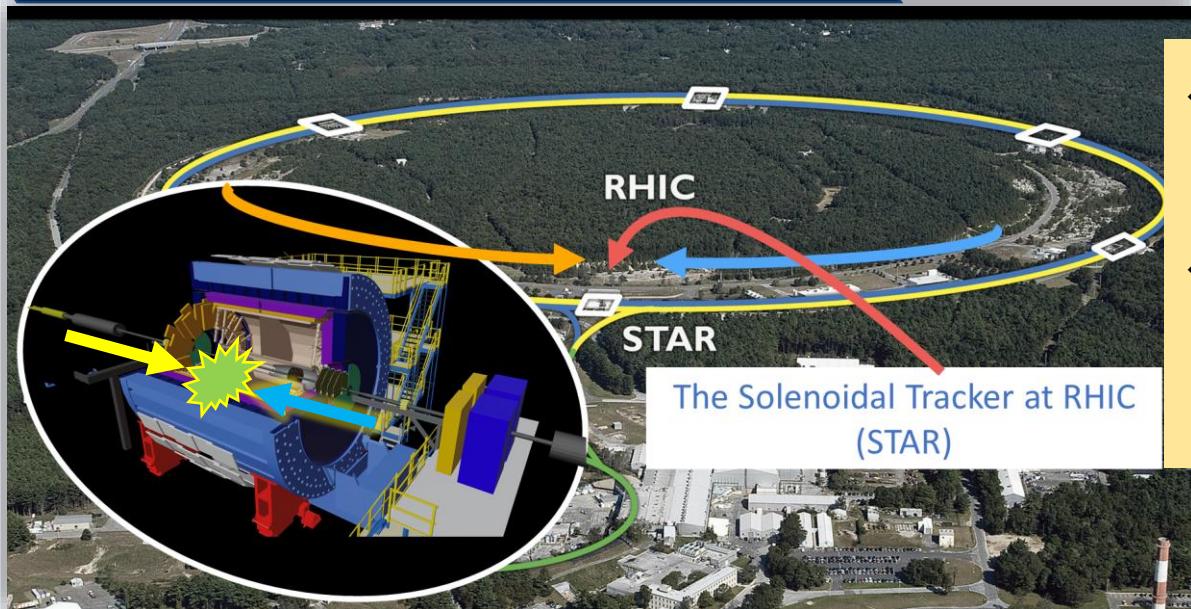
$$S_0: \quad a = -23.714 \text{ fm} \quad r_0 = 2.73 \text{ fm}$$

$$S_1: \quad a = 5.425 \text{ fm} \quad r_0 = 1.749 \text{ fm}$$

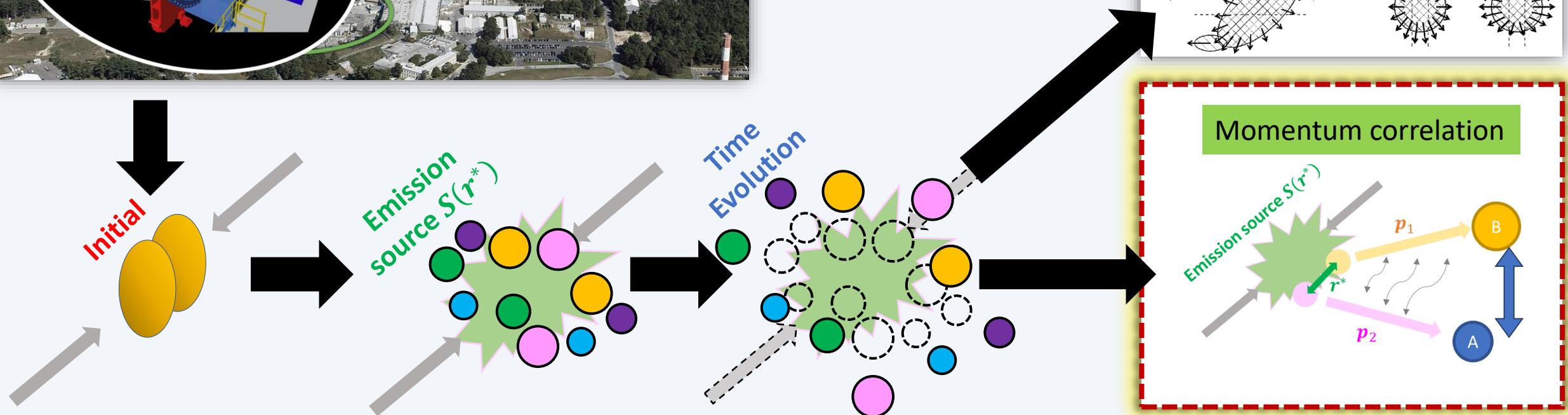
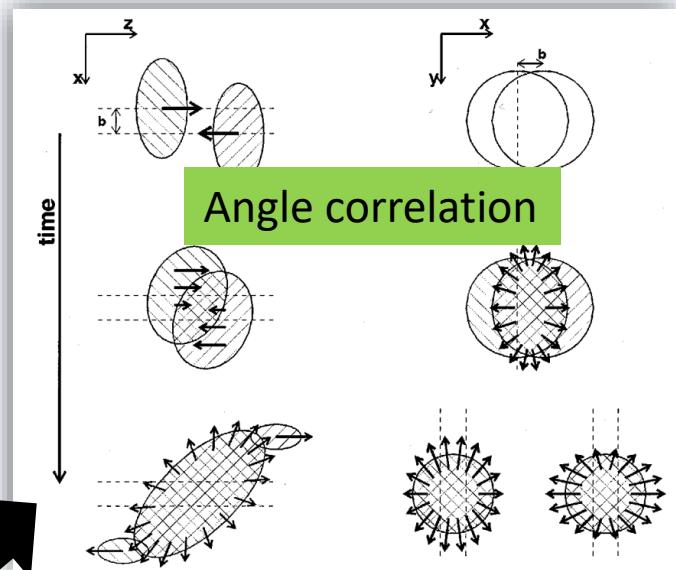
$$\rightarrow B_d = 2.2 \text{ MeV}$$

# Heavy Ion Collision Experiment

Annu. Rev. Nucl. Part. Sci. 1999.



- ❖ Space and time evolution of particle-emitting source
- ❖ Final state interactions
  - ❖ N(-N)-Y interactions
  - ❖ hypernuclei structure



# Baryon Correlation Function (CF)

Momentum correlation function:

$$C(\mathbf{p}_1, \mathbf{p}_2) \equiv \frac{P(\mathbf{p}_1, \mathbf{p}_2)}{P(\mathbf{p}_1) \cdot P(\mathbf{p}_2)}$$

Single-particle momentum

Statistical

Normalization factor

$k^*$ : particle momentum in the pair rest frame

Experimental

$$C(k^*) = \mathcal{N} \frac{A(k^*)}{B(k^*)}$$

A( $k^*$ ) Same events  
B( $k^*$ ) Mixed events

Approximating the emission process and the momenta of the particles:

$$C(k^*) = \int d^3r^* S(r^*) |\Psi(r^*, k^*)|^2$$

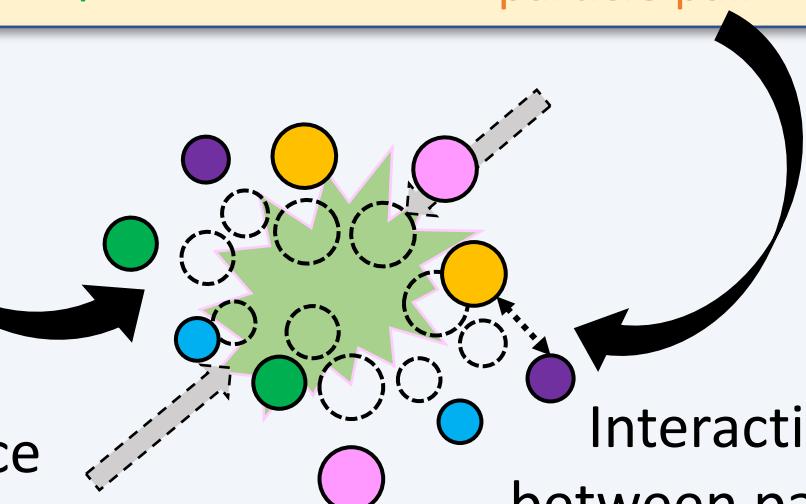
Distribution of the relative distance of particle pair

Relative wave function of the particle pair

Size of the emitting source

Interactions between particles

Modeling



# Lednicky-Lyuboshitz (L-L) Approach

R. Lednicky, et al. Sov.J.Nucl.Phys. 35 (1982) 770

J. Haidenbauer, Phys.Rev.C 102 (2020) 3, 034001

L. Fabbietti, et al., Ann.Rev.Nucl.Part.Sci. 71 (2021) 377-402

Michael Annan Lisa, et al., Ann.Rev.Nucl.Part.Sci. 55 (2005) 357-402

Approximating the emission process and the momenta of the particles:

*Modeling*

$$C(\mathbf{k}^*) = \int d^3r^* S(r^*) |\Psi(r^*, \mathbf{k}^*)|^2$$

Distribution of the relative distance of particle pair

Relative wave function of the particle pair

Major Assumptions

Source

❖ Smoothness approximation for source function\*

❖ Static and spherical Gaussian source

- Single particle source:  $S_i(x_i, p_i^*)$

- Pair source (radius  $R_G$ ):  $S(x, p^*) \propto e^{-x^2/2R_G^2} \delta(t - t_0)$

Wave function

❖ S-wave scattering wave

❖ Effective range expansion for  $\Psi(r^*, \mathbf{k}^*)$

❖ Approximate the wave function by its asymptotic form

Gaussian source approximation:

$$S(r^*) = (2\sqrt{\pi}R_G)^{-3} e^{-r^{*2}/4R_G^2}$$

Scattering amplitude:

Consider only S-wave     $\Psi(r^*) = e^{-ir^*\cdot\mathbf{k}^*} + \frac{f(\mathbf{k}^*)}{r^*} e^{ir^*\cdot\mathbf{k}^*}$

$$f(\mathbf{k}^*) \approx \left( \frac{1}{f_0} + \frac{d_0 \mathbf{k}^{*2}}{2} - i\mathbf{k}^* \right)^{-1}$$

Scattering length:

$$a \rightarrow -f_0$$

Effective range:

$$r_0 \rightarrow d_0$$

Lednicky-Lyuboshitz (L-L) approach

$R_G$  : spherical Gaussian source of pairs

$f_0$  : scattering length

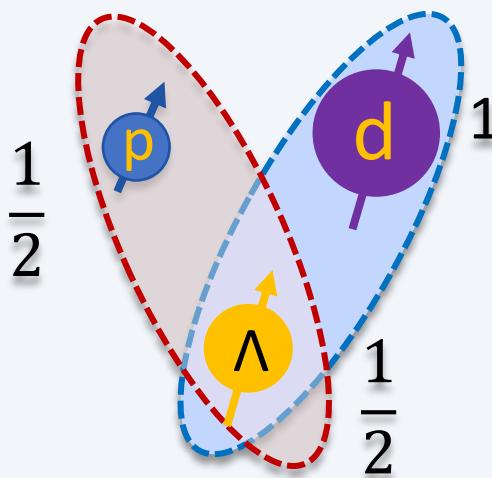
$d_0$  : effective range

\*The smoothness approximation has been checked for expanding thermal sources, found to be very reasonable for large (RHIC-like) sources, but still questionable for smaller sources

# Modeling with Separated Spin States

Modeling

Singlet State	$^1S_0$	(S)
Triplet State	$^3S_1$	(T)



Doublet State	$^2S_{1/2}$	(D)
Quartet State	$^4S_{3/2}$	(Q)

Approximating the emission process and the momenta of the particles:

$$C(\mathbf{k}^*) = \int d^3r^* S(\mathbf{r}^*) |\Psi(\mathbf{r}^*, \mathbf{k}^*)|^2$$

Source      Wave function

Spin averaged

$|\Psi(\mathbf{r}^*, \mathbf{k}^*)|^2$  expanded with averaged parameters:  $\bar{f}_0$  and  $\bar{d}_0$

$$|\Psi(\mathbf{r}^*, \mathbf{k}^*)|^2 \rightarrow \boxed{f_{S1}} |\Psi_{S1}(\mathbf{r}^*, \mathbf{k}^*)|^2 + \boxed{f_{S2}} |\Psi_{S2}(\mathbf{r}^*, \mathbf{k}^*)|^2$$

Spin separated

$$C(\mathbf{k}^*) = \int d^3r^* S(\mathbf{r}^*) \left( \frac{1}{3} |\Psi_{1/2}(\mathbf{r}^*, \mathbf{k}^*)|^2 + \frac{2}{3} |\Psi_{3/2}(\mathbf{r}^*, \mathbf{k}^*)|^2 \right)$$

For separated spin states in d-Λ

$$\boxed{f_0(D)}$$

$$\boxed{d_0(D)}$$

$$\boxed{f_0(Q)}$$

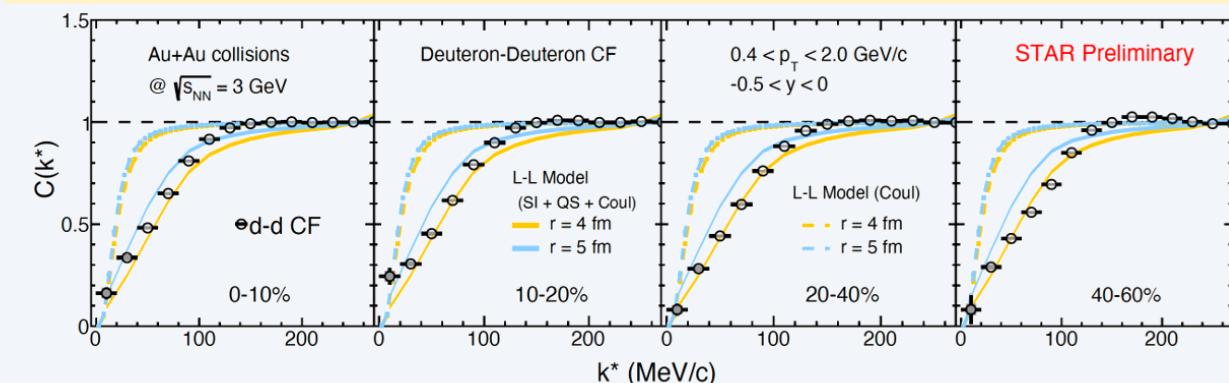
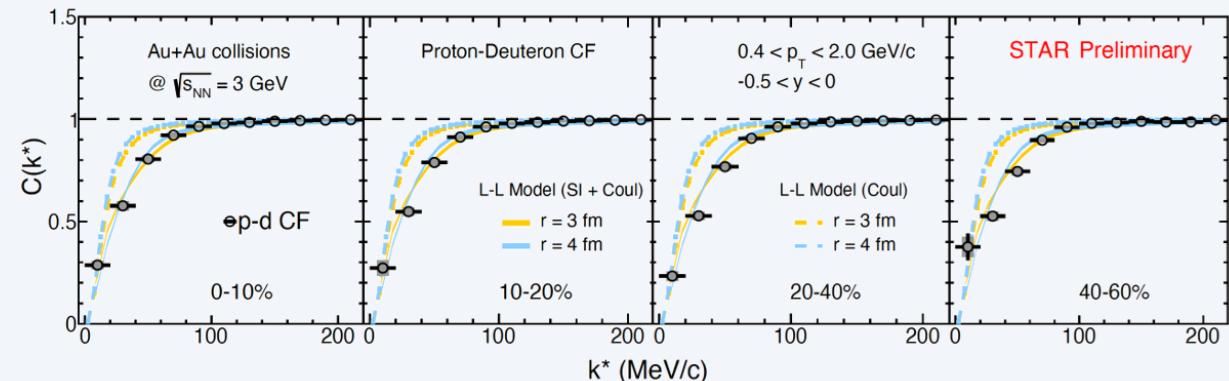
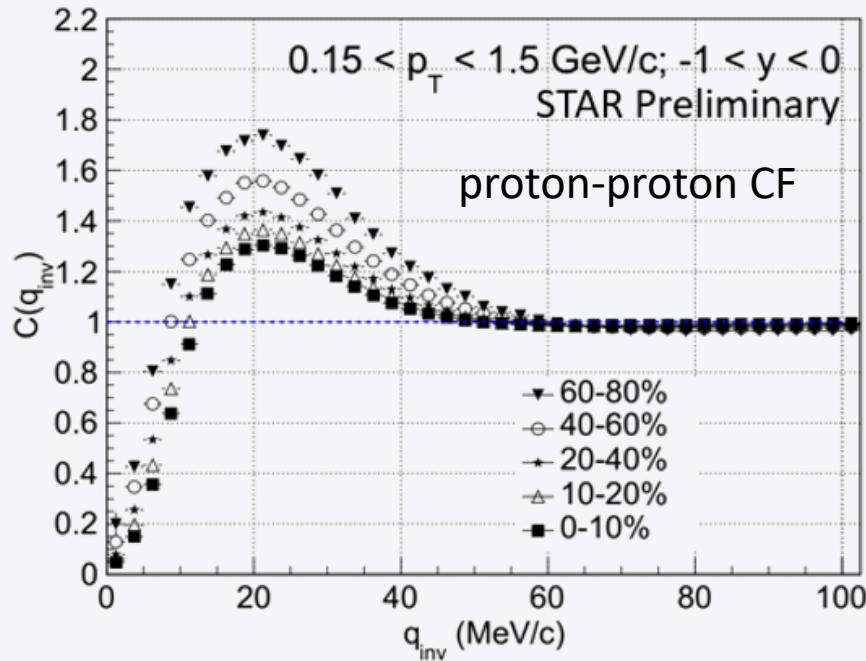
$$\boxed{d_0(Q)}$$

R. Lednický, et al. Sov.J.Nucl.Phys. 35 (1982) 770

L. Michael, et al. Ann.Rev.Nucl.Part.Sci. 55 (2005) 357-402

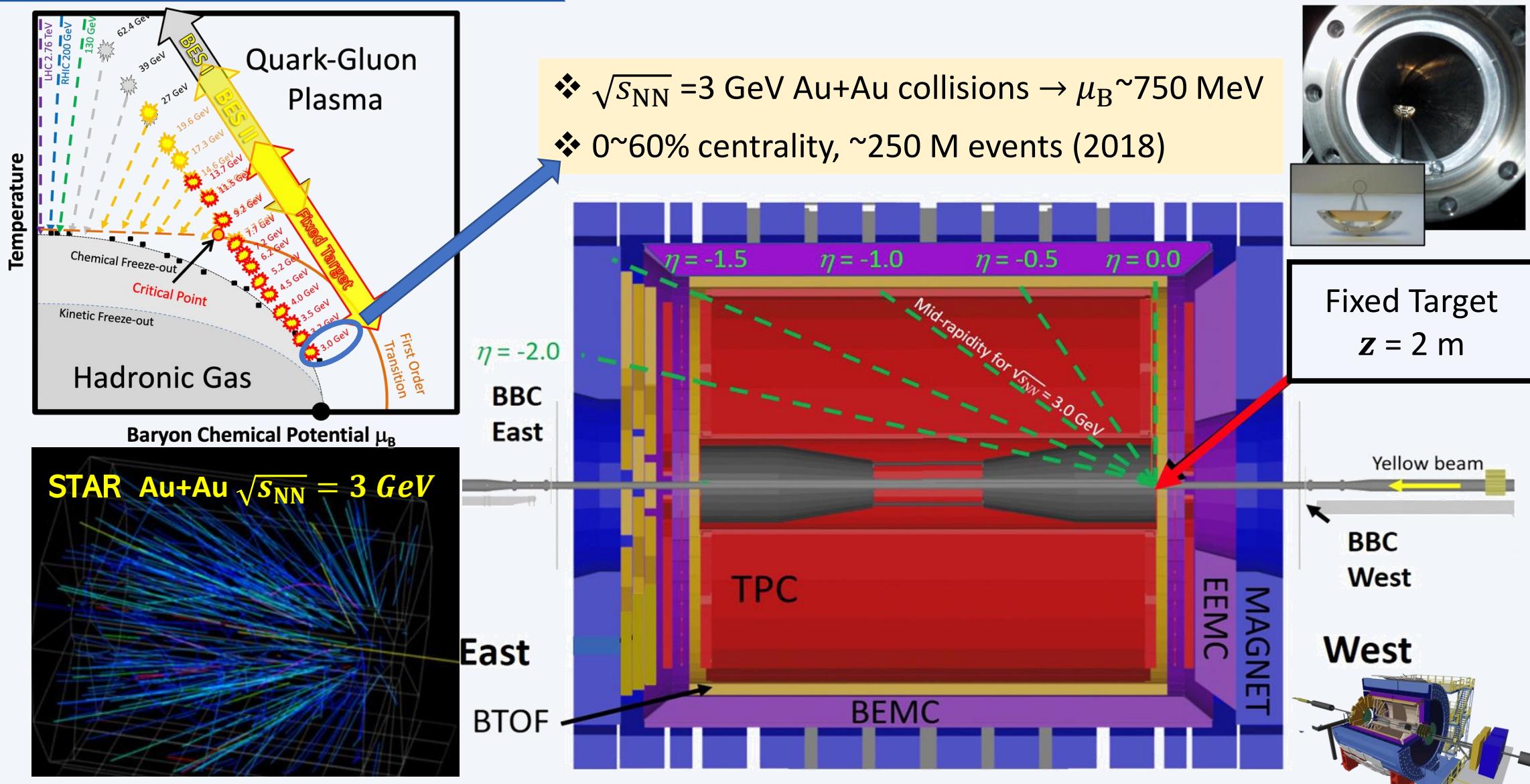
J. Haidenbauer, Phys.Rev.C 102 (2020) 3, 034001

# Correlation Function & Low-E Scattering Experiment

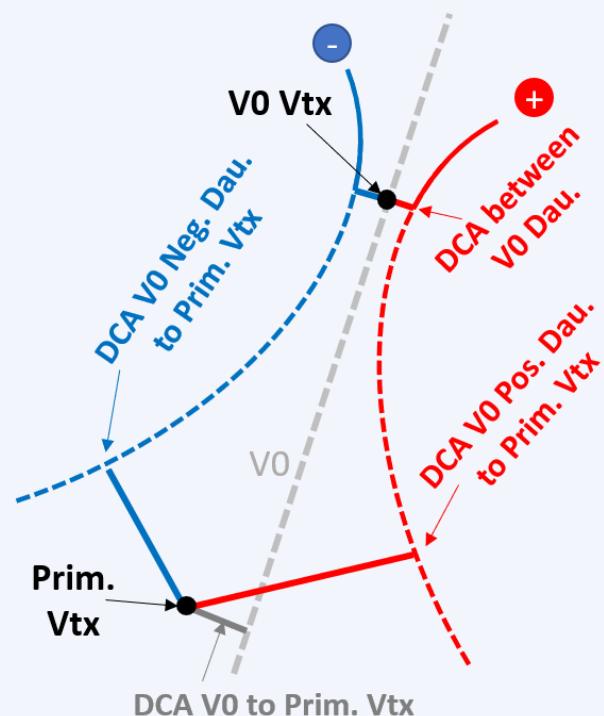
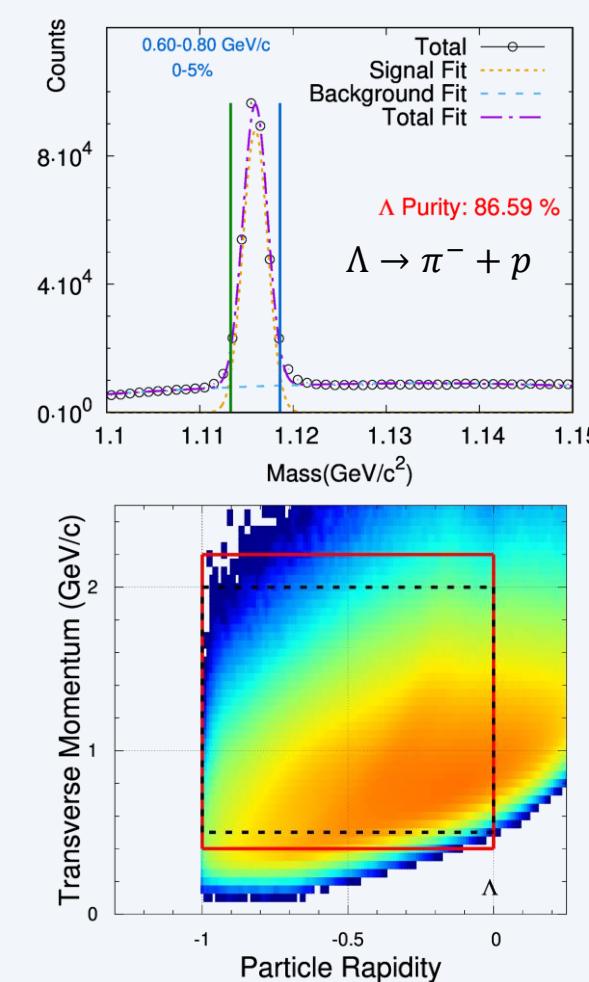
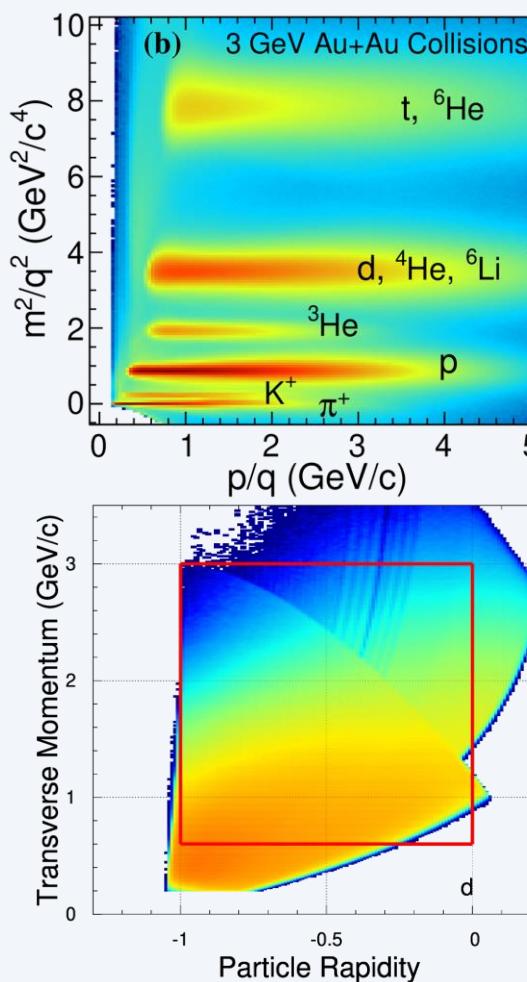
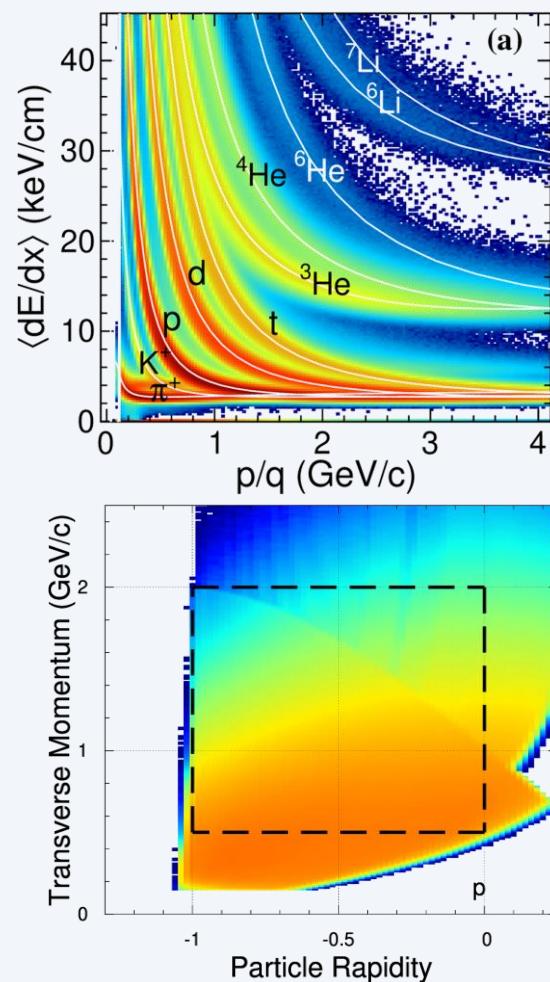


**Ongoing studies @ STAR**

# Beam Energy Scan – II & Fixed Target Setup

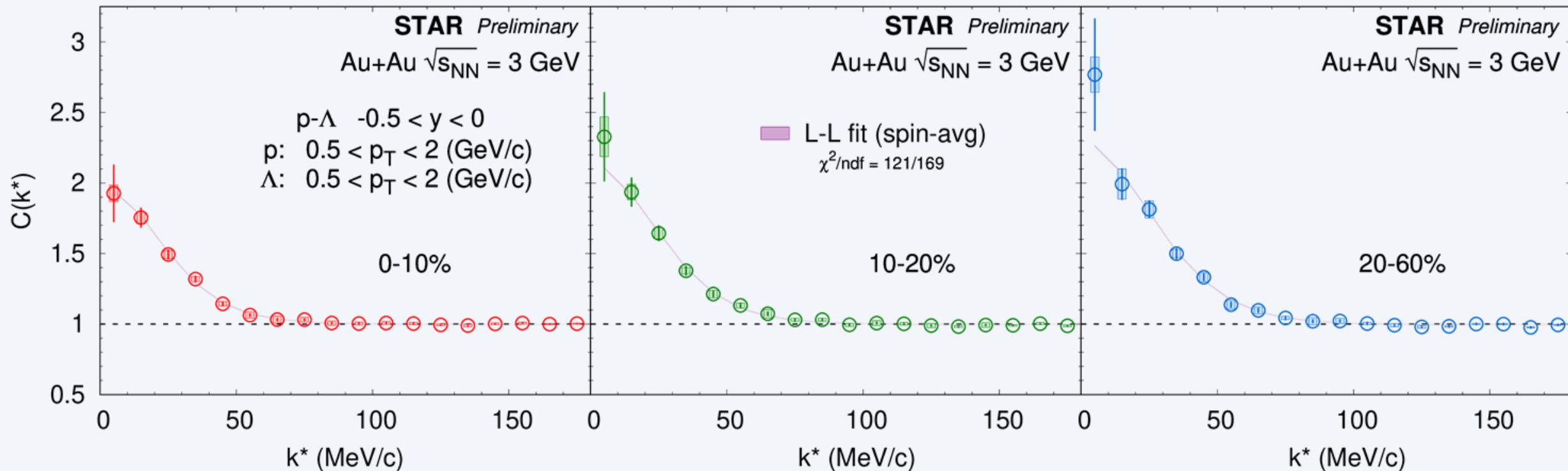


# Particle Identification & Reconstruction @ 3 GeV



- ❖  $\pi^-$ , p, and d particles are identified by TPC and TOF
- ❖ A larger  $p_T$  range is used in d- $\Lambda$  correlation measurement (red) due to statistics

# p- $\Lambda$ Correlation Measurement @ STAR



## Corrections

1. Purity correction
2.  $\Lambda$  feed-down correction
3. Track splitting & merging
4. Momentum smearing effect

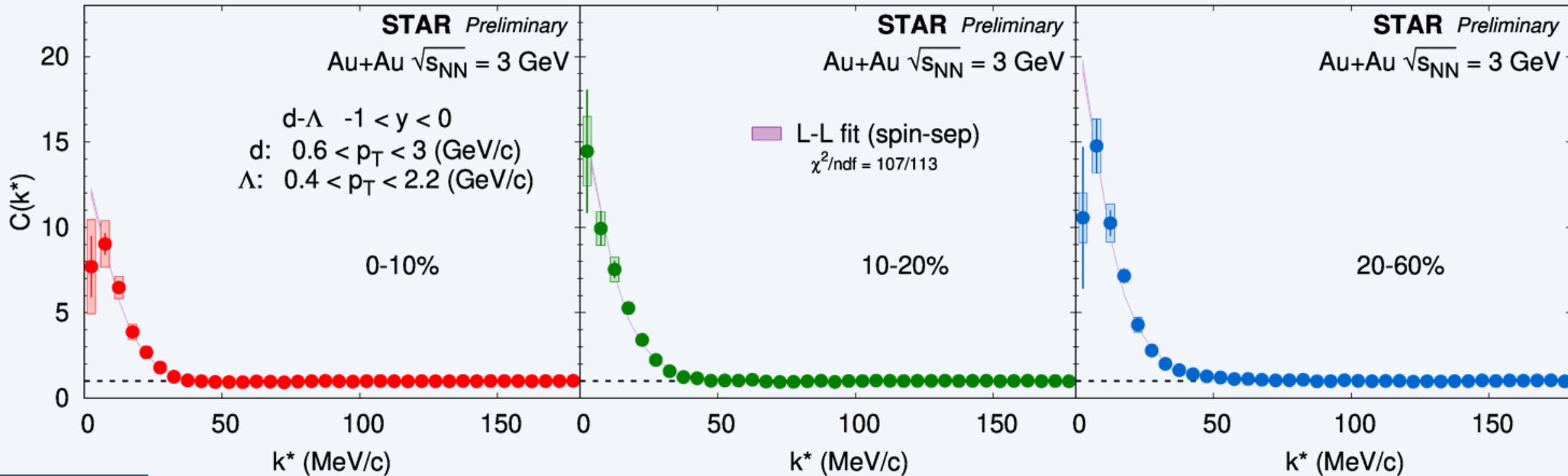
- ❖ Simultaneous fit to data in different centralities/rapidity
  - ❖  $R_G^i$ , **spin-avg**  $f_0$  and  $d_0$  with Lednicky-Lyuboshitz approach
- ❖ Spin-avg scattering length ( $f_0$ ) and effective range ( $d_0$ ):

$$f_0 = 2.32^{+0.12}_{-0.11} \text{ fm}$$

$$d_0 = 3.5^{+2.7}_{-1.3} \text{ fm}$$

$R_G^i$ : source size in different centrality

# d- $\Lambda$ Correlation Measurement @ STAR



## Corrections

1. Purity correction
2. Track splitting & merging
3. Contamination from  ${}^3\text{H} \rightarrow \pi^- + p + d$  decay

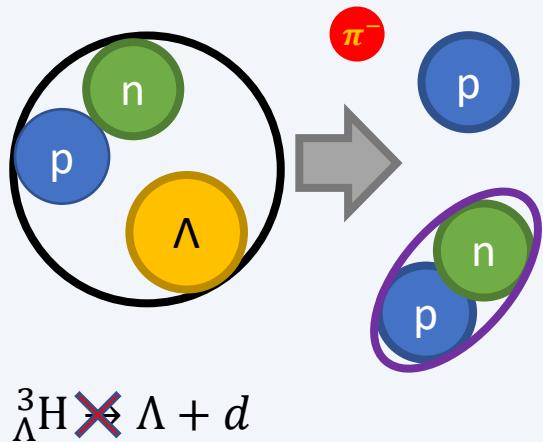
- ❖ First d- $\Lambda$  correlation measurements in the heavy-ion collision experiment
- ❖ Simultaneous fit to data in different centralities
  - ❖  $R_G^i, f_0(D), d_0(D), f_0(Q)$ , and  $d_0(Q)$  with Lednicky-Lyuboshitz approach

$f_0(D) = -20^{+3}_{-3}$ fm	$d_0(D) = 3^{+2}_{-1}$ fm
$f_0(Q) = 16^{+2}_{-1}$ fm	$d_0(Q) = 2^{+1}_{-1}$ fm

- ❖  $\Lambda$  feed-down correction not applied due to unknown d- $\Sigma/\Xi$  correlation
- ❖ Momentum smearing effect negligible

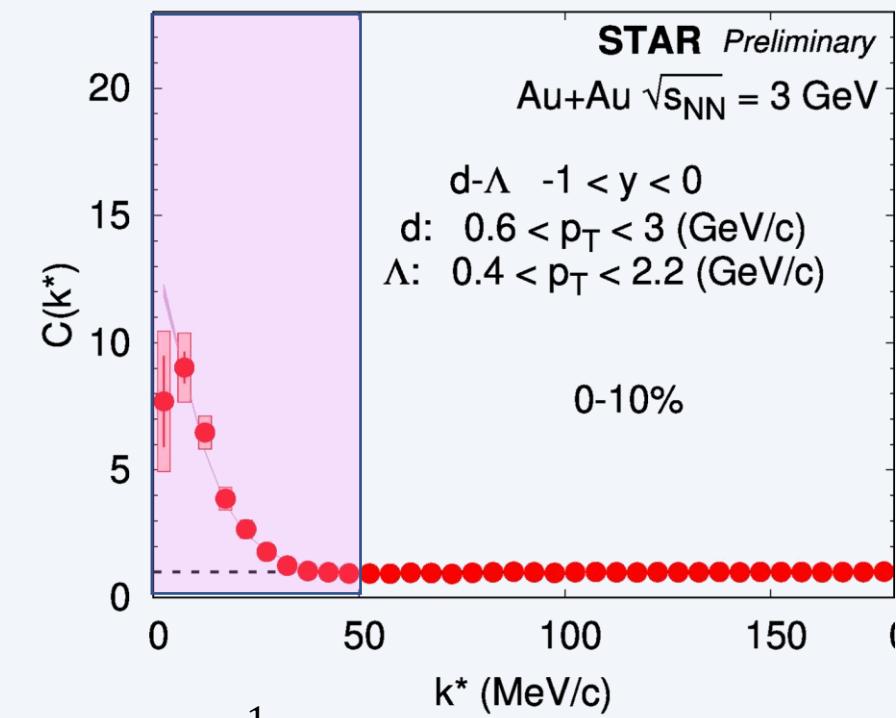
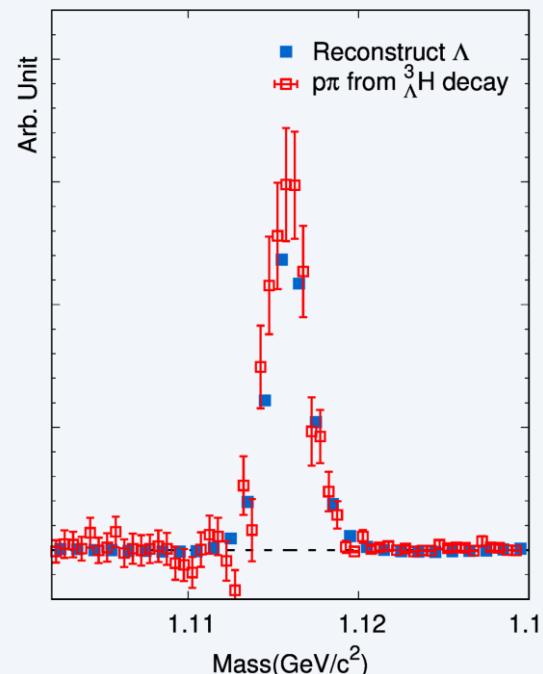
# Contamination Correction from ${}^3_{\Lambda}\text{H} \rightarrow p\pi^- + d$ Decay

${}^3_{\Lambda}\text{H} \rightarrow p + \pi^- + d$ ;  
B.R.  $\approx 40\sim 50\%$



Violation of energy conservation

- The  ${}^3_{\Lambda}\text{H}$  decayed  $p + \pi^-$  are **not experimentally distinguishable** with the reconstructed  $\Lambda$
- $(p\pi^-) - d$  from  ${}^3_{\Lambda}\text{H}$  will affect small  $k^*$  region



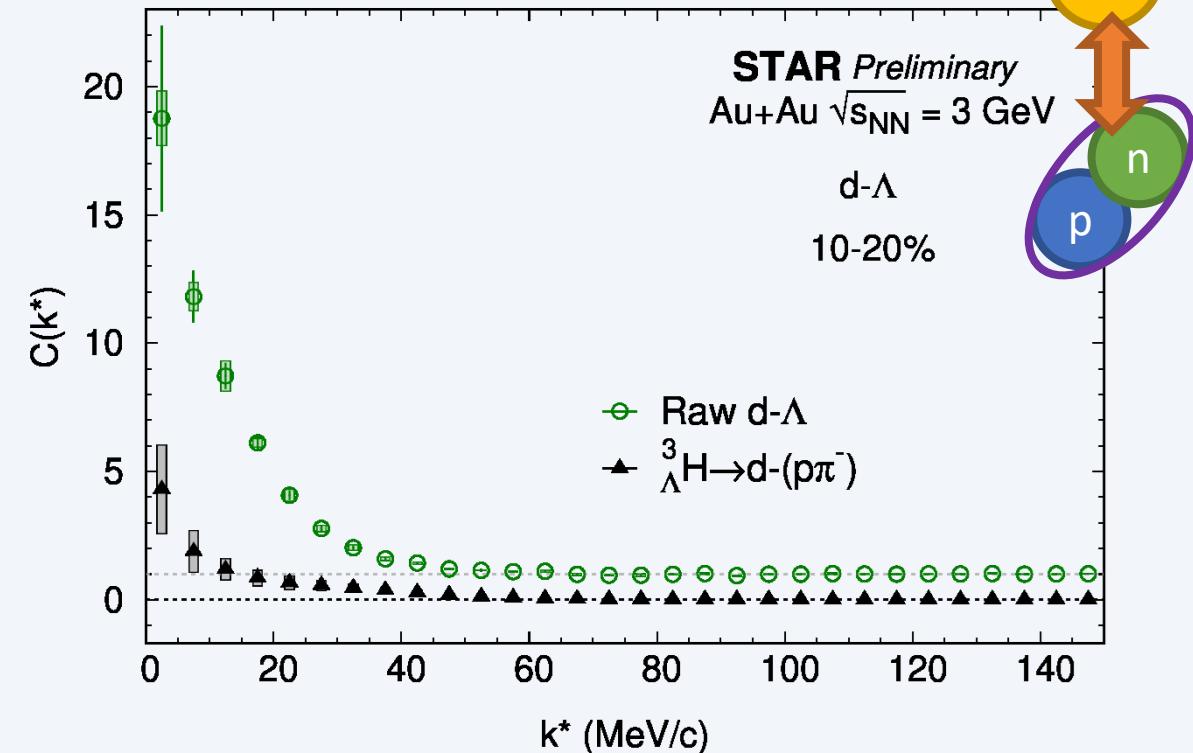
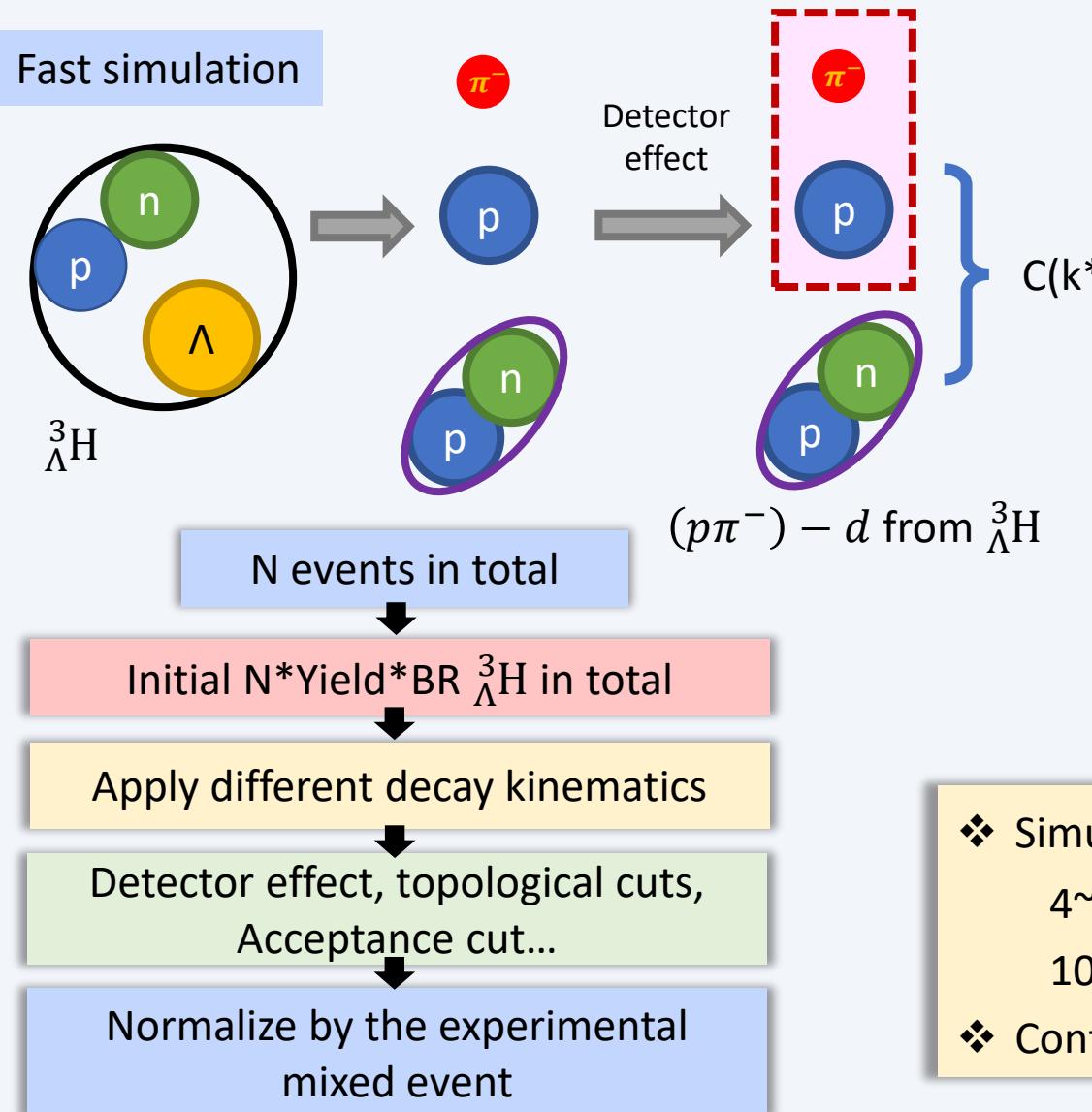
$$k^* = \frac{1}{2} |p_1^* - p_2^*|$$

$$p_1 = \frac{\sqrt{(M^2 + m_1^2 - m_2^2)^2 - 4M^2m_1^2}}{2M},$$

$$p_2 = \frac{\sqrt{(M^2 + m_2^2 - m_1^2)^2 - 4M^2m_2^2}}{2M}.$$

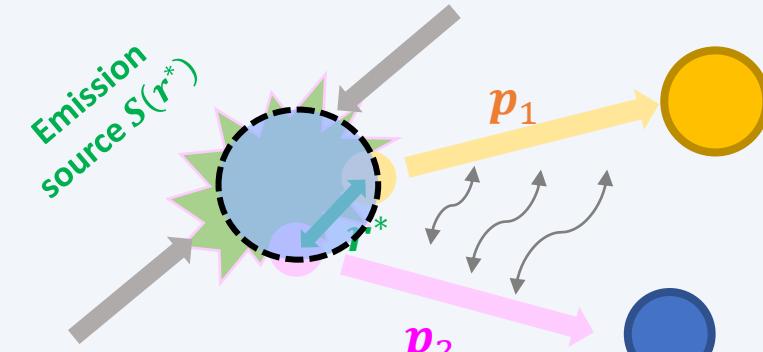
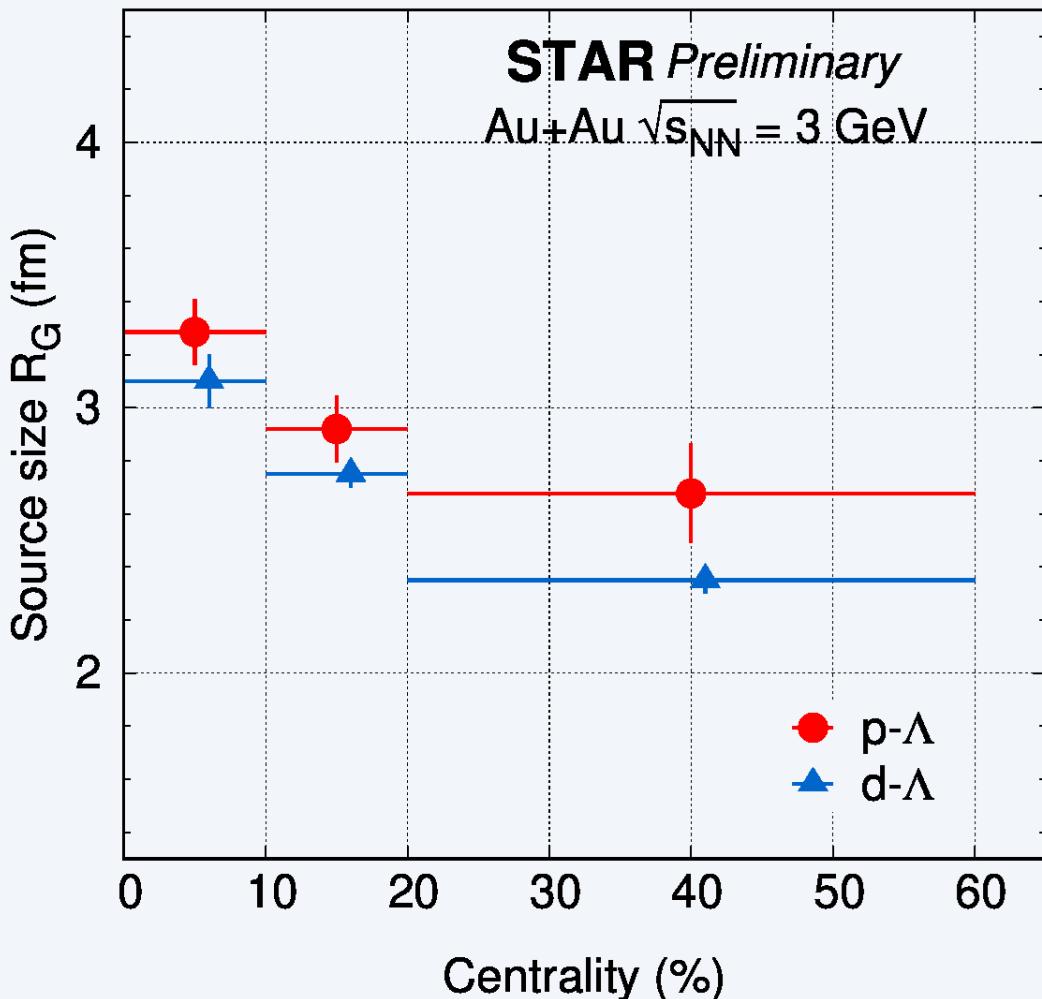
$p \approx 20 \text{ MeV/c} \Rightarrow k^* \approx 20 \text{ MeV/c}$

# Contamination Correction from ${}^3\Lambda \rightarrow p\pi^- + d$ Decay



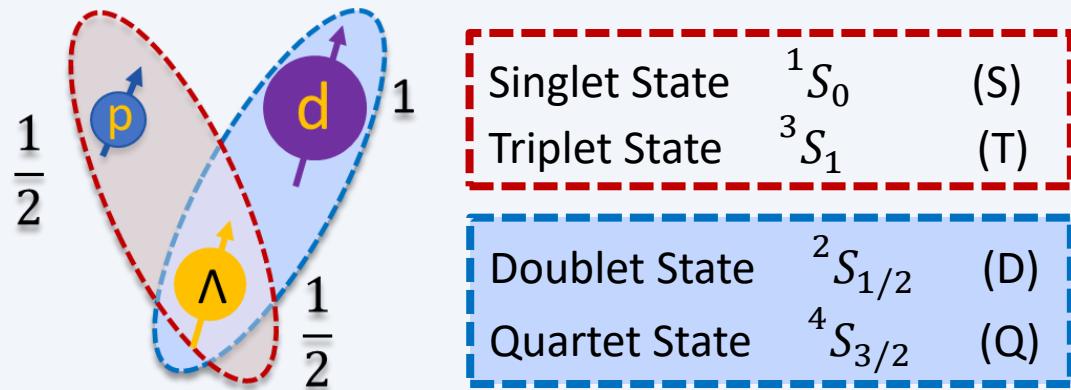
- ❖ Simulation based on STAR  ${}^3\Lambda$  yield measurement:  
4~8% of d- $\Lambda$  entries from  ${}^3\Lambda$  decay at  $k^* < 100$  MeV/c in  
10~20% centrality
- ❖ Contamination subtracted from inclusive d- $\Lambda$  correlation

# Source Size with L-L approach



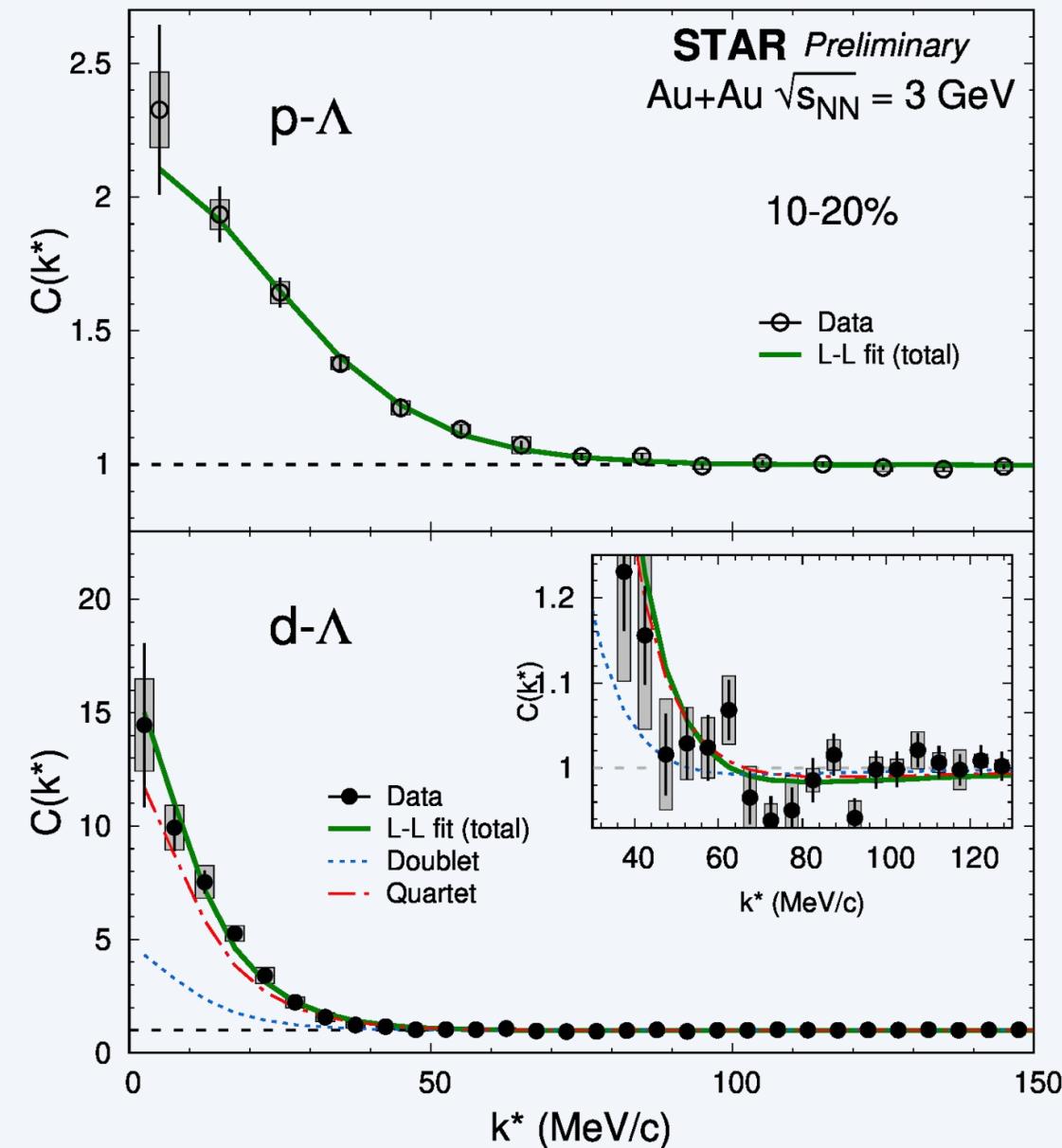
- ❖  $R_G$ : **spherical Gaussian source of pairs** by Lednicky-Lyuboshits approach
- ❖ Separation of emission source from final state interaction
- ❖ Collision dynamics as expected:
  - ❖  $R_G^{\text{central}} > R_G^{\text{peripheral}}$
  - ❖  $R_G(p - \Lambda) > R_G(d - \Lambda)$

# Correlation Function & Spin States

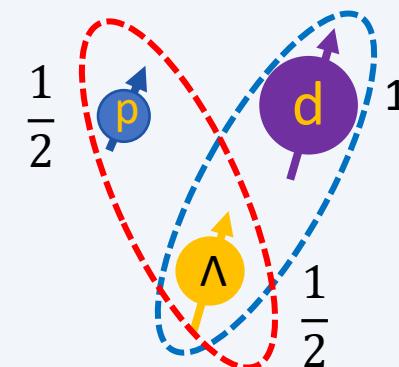
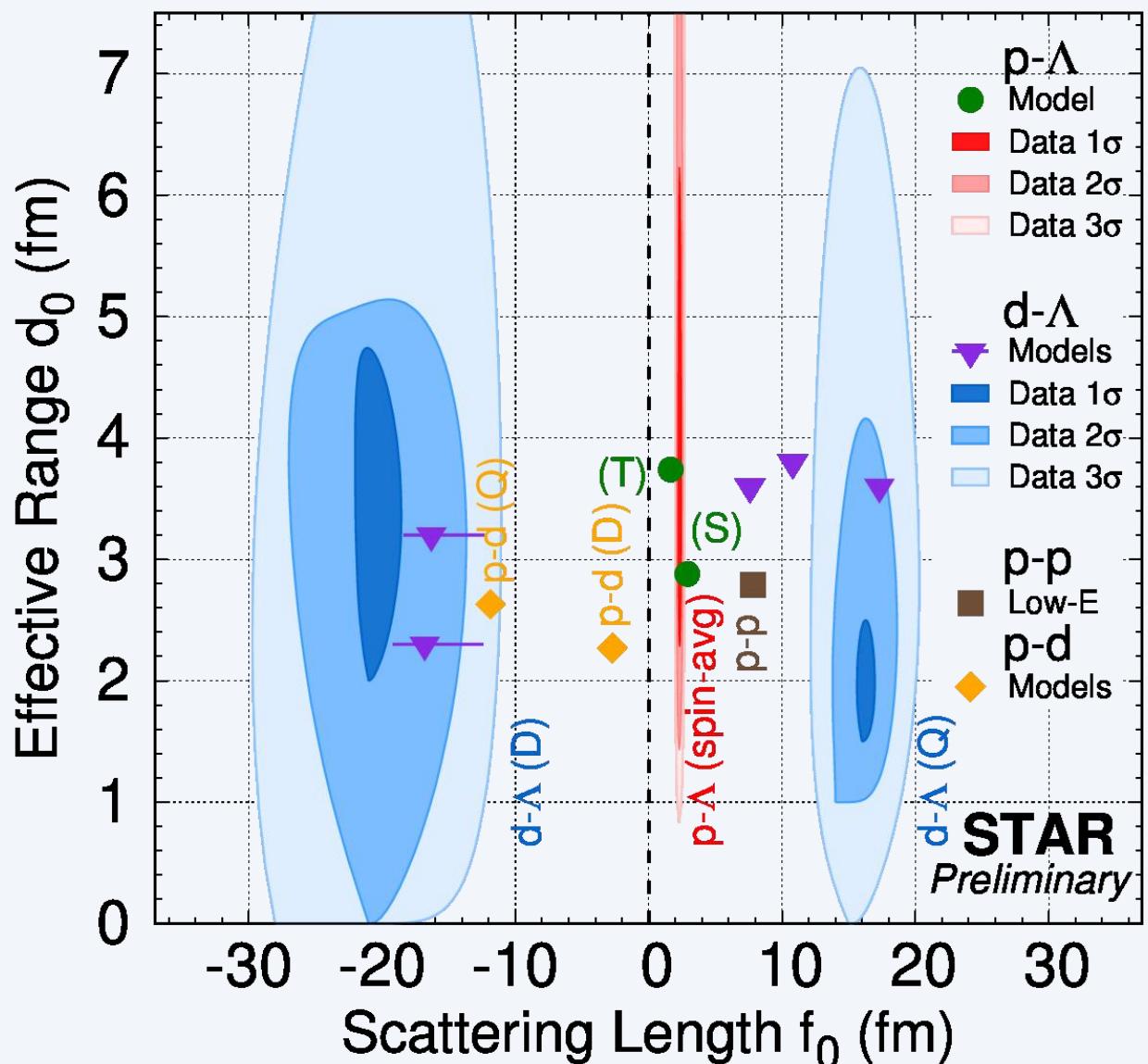


$$d-\Lambda: |\psi(r, k)|^2 \rightarrow \frac{1}{3}|\psi_{1/2}(r, k)|^2 + \frac{2}{3}|\psi_{3/2}(r, k)|^2$$

- ❖ Different spin states with different  $f_0$  and  $d_0$  parameters
- ❖ **p- $\Lambda$  correlation:** current statistics is not enough to separate two spin states → **spin-averaged fit**
- ❖ **d- $\Lambda$  correlation:** very different  $f_0$  for (D) and (Q) are predicted → **Spin-separated fit**



# Scatterings Length ( $f_0$ ) and Effective Range ( $d_0$ )



$$\frac{1}{f(k)} \approx \frac{1}{f_0} + \frac{d_0 k^2}{2} - ik$$

❖ The constraint on the effective range ( $d_0$ ) is weaker

❖ The measurement is done at freeze-out

❖ Spin-avg for  $f_0$  &  $d_0$  p- $\Lambda$  system

$$f_0 = 2.32^{+0.12}_{-0.11} \text{ fm}$$

$$d_0 = 3.5^{+2.7}_{-1.3} \text{ fm}$$

❖ Successfully separate two spin states in d- $\Lambda$

$$f_0(D) = -20^{+3}_{-3} \text{ fm}$$

$$d_0(D) = 3^{+2}_{-1} \text{ fm}$$

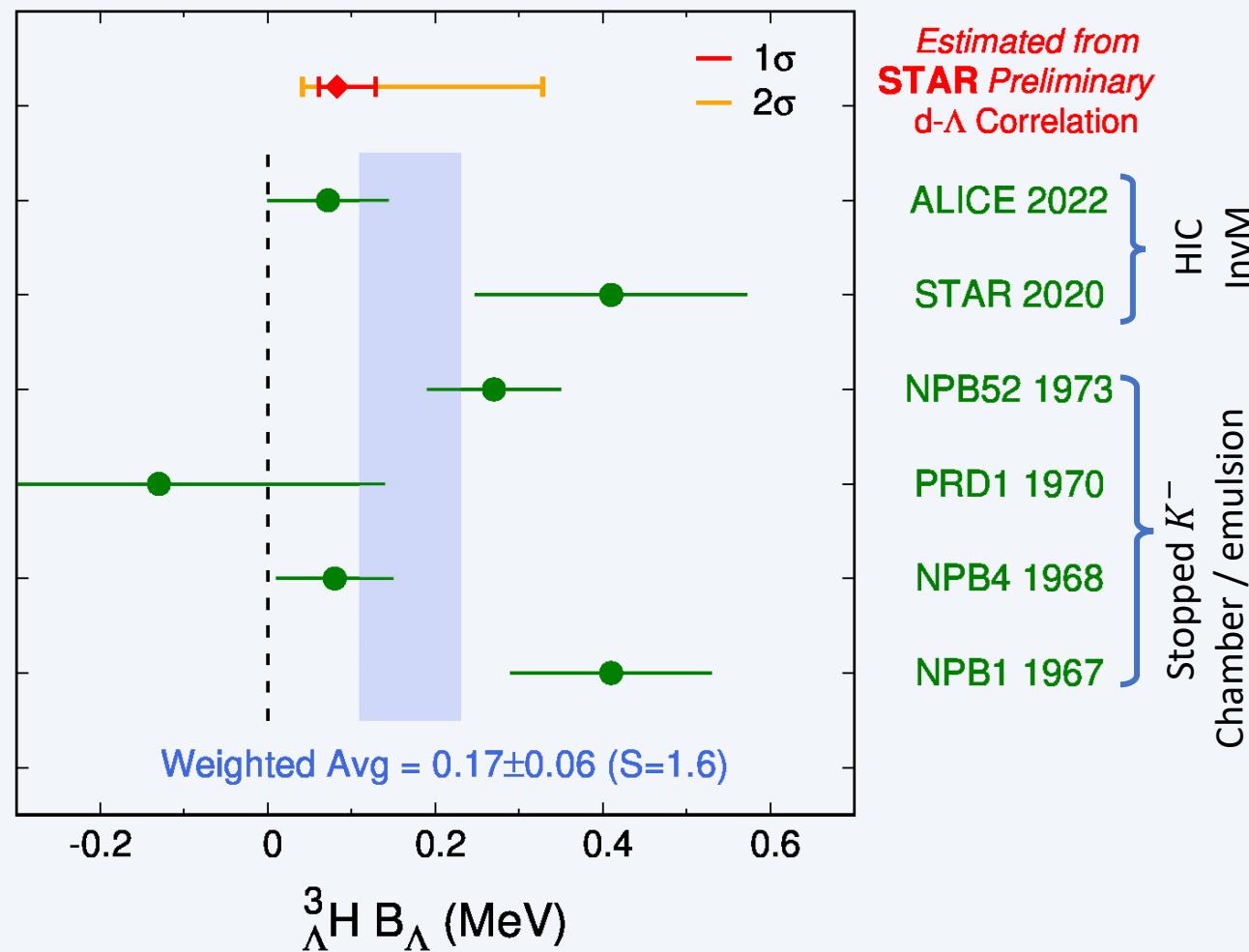
$$f_0(Q) = 16^{+2}_{-1} \text{ fm}$$

$$d_0(Q) = 2^{+1}_{-1} \text{ fm}$$

❖ Constraint fit for d- $\Lambda$ , require  $f_0(D) < 0$

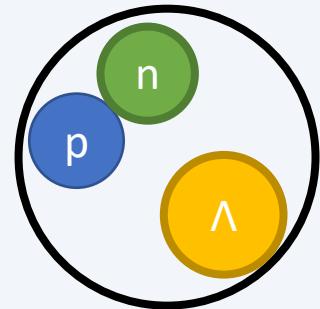
❖ Edge of d- $\Lambda$  contours are shown with Bezier smooth to improve the visibility

# ${}^3_{\Lambda}\text{H}$ Binding Energy



${}^3_{\Lambda}\text{H}$  binding energy ( $B_{\Lambda}$ ):

- ❖ Bethe formula from Effective Range Expansion (ERE) parameters  $f_0(D)$  &  $d_0(D)$



$$\frac{1}{-f_0} = \gamma - \frac{1}{2} d_0 \gamma^2$$

- ❖  $B_{\Lambda} = \frac{\gamma^2}{2\mu_{d\Lambda}}$
- ❖  $\mu_{d\Lambda}$ : reduced mass
- ❖  $\gamma$ : binding momentum

- ❖  ${}^3_{\Lambda}\text{H} B_{\Lambda} = [0.04, 0.33]$  (MeV) @ 95% CL
  - Consistent with the world average
  - ❖ A new way to constrain the  ${}^3_{\Lambda}\text{H}$  structure

# Summary and outlook

- ❖ The first d- $\Lambda$  correlation function measurements in heavy-ion collisions
- ❖ Successfully separated emission source size from final state interactions in d- $\Lambda$  correlation functions

1.  $R_G^{\text{central}} > R_G^{\text{peripheral}}$  and  $R_G(p - \Lambda) > R_G(d - \Lambda)$

2. d- $\Lambda$  correlation spin-sep:

$$f_0(D) = -20^{+3}_{-3} \text{ fm}$$

$$d_0(D) = 3^{+2}_{-1} \text{ fm}$$

$$f_0(Q) = 16^{+2}_{-1} \text{ fm}$$

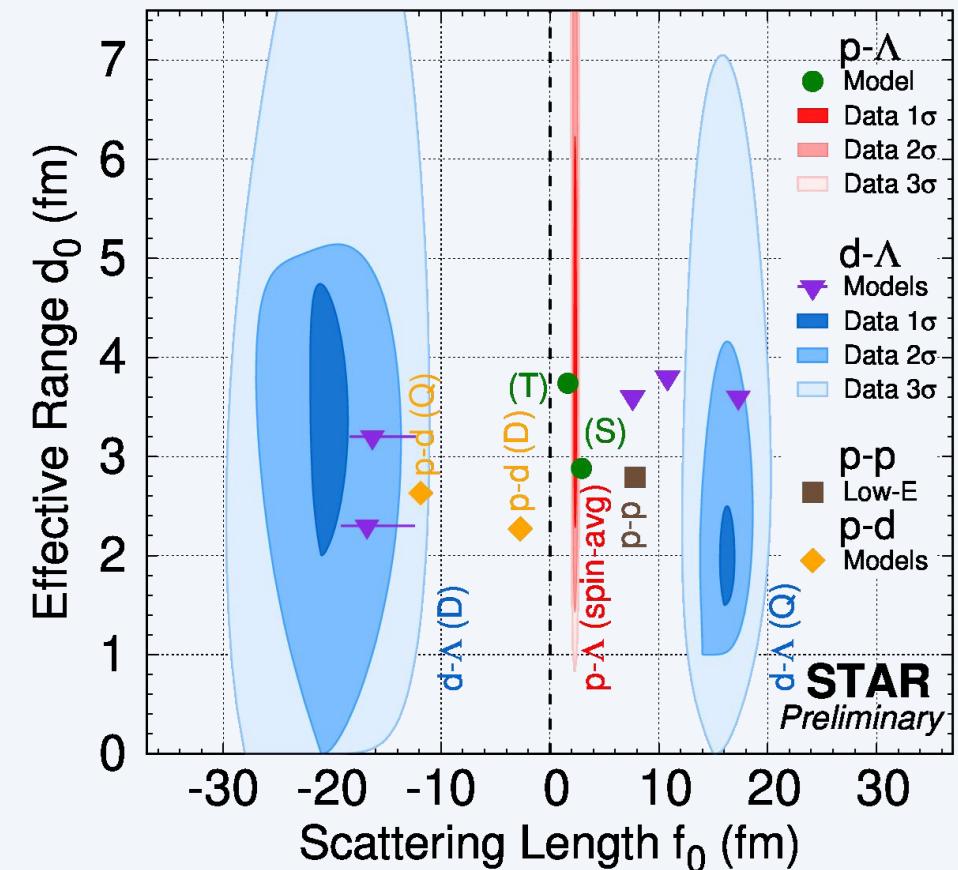
$$d_0(Q) = 2^{+1}_{-1} \text{ fm}$$

3.  ${}^3\text{H} B_\Lambda = [0.04, 0.33] \text{ (MeV)}$  @ 95% CL from d- $\Lambda$  correlation (D)

## Outlook:

More than 10 times statistics from BES-II

- ❖ Emission source size vs. energy, rapidity...
- ❖ Baryon correlations with different species





# Thank you!



2024.05.22