Multiplicity and Rapidity Dependent Study of (Multi-)strange hadrons in small collision system using the STAR Detector

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Abstract. Strangeness enhancement has long been considered as a signature of the quark-gluon plasma formation in heavy-ion collisions. Strangeness enhance-2 ment has also been observed in small systems at the Large Hadron Collider 3 (LHC), but the underlying physics still needs to be fully understood. This mo-4 tivates us to study the strange hadron production in small systems at RHIC. We 5 present new measurements of (multi-)strange hadrons (K_{c}^{0} , Λ , Ξ and Ω) in d+Au 6 collisions at $\sqrt{s_{NN}} = 200$ GeV, collected by STAR in 2016. We report the mul-7 tiplicity and rapidity dependence of strange hadron transverse momentum $(p_{\rm T})$ 8 spectra, $p_{\rm T}$ -integrated yields dN/dy, average transverse momentum ($\langle p_{\rm T} \rangle$), yield 9 ratios of these strange hadrons to pions, nuclear modification factors and rapidity 10 asymmetry in d+Au collisions.

Keywords: Heavy-ion collisions, Quark Gluon Plasma, RHIC, STAR

Introduction 1 13

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The production of strange hadrons in high-energy hadronic interactions provides a way 14 to investigate the properties of quantum chromodynamics (QCD), the theory of strongly 15 interacting matter. Strangeness enhancement in heavy-ion collisions with respect to 16 p+p collisions has been suggested as a signature of quark-gluon plasma (QGP) for-17 mation [1]. However, the creation of QGP in small systems is still under intense debate. 18 Small collision systems like proton-nucleus (p-A) and deuteron-nucleus (d-A) can 19 be considered as control experiments where the formation of an extended QGP phase is 20 not expected. These collision systems are used for baseline measurements to study the 21 possible effects of cold nuclear matter and to disentangle them from hot dense matter 22 effects signifying the QGP formation in heavy-ion collisions. The process of generating 23 hadrons can be affected by various factors that include modifications in parton distri-24 bution functions within nuclei, the possibility of parton saturation, multiple scatterings 25 of the partons traversing the nucleus, and the radial flow [2, 3]. It is anticipated that the 26 magnitude of these effects may vary with the rapidity of the produced particles. 27

Nuclear effects can be quantified using variables such as nuclear modification factor 28 and rapidity asymmetry. The nuclear modification factor is defined as the ratio of the 29 yield of a particle in heavy-ion collisions to its yield in p+p collisions, scaled by the 30 number of binary nucleon-nucleon inelastic collisions [4]. 31

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$$R_{AB}(p_{T}) = \frac{\text{Yield}_{AB}(p_{T})}{\langle N_{bin} \rangle \text{Yield}_{pp}(p_{T})}$$
(1)

where $\langle N_{bin} \rangle$ is the average number of binary nucleon-nucleon collisions. A comparative study of particle production in forward and backward rapidity regions uses rapidity

asymmetry (Y_{Asym}) [4]. Y_{Asym} is defined as

$$Y_{Asym}(p_T) = \frac{Y_B(p_T)}{Y_F(p_T)}$$
(2)

where Y_B and Y_F are particle yields in backward and forward rapidity regions, respectively.

37 2 Analysis Details

A successful run of *d*+Au collisions at $\sqrt{s_{NN}} = 200$ GeV was carried out in 2016 at RHIC. In this analysis utilizing the *d*+Au dataset, we focus on K_s^0 , $\Lambda(\bar{\Lambda})$, $\Xi(\bar{\Xi})$, $\Omega(\bar{\Omega})$ that are weakly decaying particles that travel a certain distance before decaying into daughter particles ($K_S^0 \rightarrow \pi^+ + \pi^-$, $\Lambda(\bar{\Lambda}) \rightarrow p(\bar{p}) + \pi^-(\pi^+)$, $\Xi^-(\bar{\Xi}^+) \rightarrow \Lambda(\bar{\Lambda}) + \pi^-(\pi^+)$, $\Omega^-(\bar{\Omega}^+) \rightarrow \Lambda(\bar{\Lambda}) + K^-(K^+)$).

The K_s^0 , $\Lambda(\bar{\Lambda})$, $\Xi(\bar{\Xi})$, $\Omega(\bar{\Omega})$ signals are extracted using the invariant mass method. 43 The decay topology can be used to enhance the signal and suppress the background. 44 We have used double Gaussian and second order polynomial functions to describe the 45 signal and background invariant mass distributions, respectively. The signal mean and 46 the width are extracted from the functional fitting. Raw yield is determined within the 47 mass window of $M_0 \pm 3\sigma$, where M_0 and σ is the fitted mean and fitted width of K_S^0 (or 48 $\Lambda/\Xi/\Omega$) respectively. The raw yield for each p_T interval is corrected for branching ratio, 49 acceptance and efficiency to obtain corrected pT spectra. Weak decay contributions from 50 Ξ are subtracted from the yields. 51

3 Results and Discussions

In Fig. 1, dN/dy scaled by the average number of nucleon participants ($\langle N_{part} \rangle$) for K⁰_s, $\Lambda(\bar{\Lambda})$, $\Xi^{-}(\bar{\Xi}^{+})$ and $\Omega^{-} + \bar{\Omega}^{+}$ as a function of $\langle N_{part} \rangle$ is *d*+Au collisions is shown. *d*+Au collisions fill the gap between *p*+*p* collisions and peripheral Au+Au and Cu+Cu collisions. The yields of strange hadrons in *d*+Au collisions lie in trend with peripheral Cu+Cu and Au+Au collisions. Strangeness enhancement is observed for K⁰_s, $\Lambda(\bar{\Lambda})$, $\Xi^{-}(\bar{\Xi}^{+})$ and $\Omega^{-} + \bar{\Omega}^{+}$ in 0-20% and 20-50% *d*+Au collisions at 200 GeV.

Figure 2 (left) shows the ratio of integrated K_S^0 , Λ , Ξ and Ω yields to pion yields as a function of mid-rapidity multiplicity for *d*+Au collisions at $\sqrt{s_{NN}} = 200$ GeV. A smooth transition of these ratios from *p*+*p* to heavy-ion collisions is observed. Data from different collision systems are consistent with each other at similar multiplicities and yields of particles with more strangeness content decrease faster as we move from high to low multiplicities. Similar trends have been observed at the LHC.

In Fig. 2 (right), nuclear modification factors for K_S^0 , Λ and Ξ at mid-rapidity (|y| < 0.5) for *d*+Au collision is presented. Cronin-like enhancement [8] is observed at intermediate p_T , which is stronger for baryons (Ξ , Λ , p) as compared to mesons (K_S^0 , π).



Fig. 1: $(dN/dy / \langle N_{part} \rangle)$ relative to p+p as a function of $\langle N_{part} \rangle$ of $(K_s^0, \Lambda(\bar{\Lambda}), \Xi^-(\bar{\Xi}^+)$ and $\Omega^- + \bar{\Omega}^+)$ in d+Au, Cu+Cu and Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV at |y| < 0.5.



Fig. 2: Left: Integrated yield particle-to-pion ratio (left) for strange particles (K_S^0 , Λ , Ξ and Ω) as a function of multiplicity for p+p, d+Au, Cu+Cu and Au+Au collisions. Right: Nuclear modification factors (R_{dAu}) for strange particles (K_S^0 , Λ and Ξ) in d+Au collisions at $\sqrt{s_{NN}} = 200$ GeV.

In Fig. 3 (Left), integrated yield and $\langle p_T \rangle$ for $\pi^+(\pi^-)$, $K^+(K^-)$, $p(\bar{p})$, K_s^0 , $\Lambda(\bar{\Lambda})$, $\Xi^-(\bar{\Xi}^+)$ and $\Omega^- + \bar{\Omega}^+$ is shown as function of multiplicity [5-7]. We observed increase of dN/dy and a hint of increase of $\langle p_T \rangle$ as a function of multiplicity. $\langle p_T \rangle$ is larger for heavier mass particles which strongly supports the picture of collective evolution, specially radial flow. We observed smooth transition of particle production from p+p to Λ +A collisions.

⁷⁴ We have also measured the transverse momentum dependence of Y_{Asym} for K_s^0 , Λ ⁷⁵ and Ξ for different rapidity intervals (backward rapidity intervals: -0.8 < y < -0.4, -0.4 < y < 0 and forward rapidity intervals: 0 < y < 0.4, 0.4 < y < 0.8) in 0-20% *d*+Au collision ⁷⁷ at $\sqrt{s_{NN}} = 200$ GeV. Y_{Asym} values deviate from unity at low p_T (< 3 GeV/c), as shown



Fig. 3: Left: Integrated yield (dN/dy) and mean transverse momentum $\langle p_T \rangle$ for $\pi^+(\pi^-)$, $K^+(K^-)$, $p(\bar{p})$ strange particles (K_S^0 , Λ , Ξ and Ω) as a function of multiplicity for p+p, d+Au, Cu+Cu and Au+Au collisions. Right: Rapidity asymmetry (Y_{Asym}) as a function of p_T (right) for strange particles (K_S^0 , Λ and Ξ) in d+Au collisions at $\sqrt{s_{NN}} = 200$ GeV.

⁷⁸ in Fig. 3 (right), suggesting the presence of a rapidity dependence in the nuclear effects. ⁷⁹ Y_{Asym} is consistent with unity at high p_T , hinting that nuclear effects become weaker ⁸⁰ at high p_T . The asymmetry is more prominent for larger rapidity interval and heavier ⁸¹ particles.

82 4 Summary

In these proceedings, we present measurements of multiplicity and rapidity dependence 83 of (multi-)strange hadron (K_s^0 , $\Lambda(\bar{\Lambda})$, $\Xi(\bar{\Xi})$ and $\Omega(\bar{\Omega})$) production in d+Au collisions 84 at $\sqrt{s_{NN}} = 200$ GeV with the STAR experiment. Results suggest that strange particle 85 production at 200 GeV is mainly driven by multiplicity. The observed enhancement of 86 integrated yield ratio to pion increases with increasing strange content. Hint of Cronin-87 like enhancement is observed for 0-20% centrality for K_S^0 , Λ and Ξ . Y_{Asym} larger than 88 one at low $p_{\rm T}$ indicates the presence of the nuclear effects. The Y_{Asym} is more prominent 89 for larger-mass particles and higher rapidity interval. 90

91 References

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