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1 Measurement of Λ hyperon spin-spin correlations in p+p
2 collisions by the STAR experiment

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5 Polarization of Λ hyperons has been observed in various collision sys-
6 tems over a wide range of collision energies over the last 50 years since
7 its discovery in Fermilab in the 70's. The existing experimental and theo-
8 retical techniques were not able to provide a conclusive answer about the
9 origin of the polarization. In these proceedings, we discuss the possibility
10 to use a new experimental method which utilizes measurement of $\Lambda\bar{\Lambda}$, $\Lambda\Lambda$,
11 and $\bar{\Lambda}\bar{\Lambda}$ pair spin-spin correlations. With this new approach, it should be
12 possible to distinguish if the polarization originates from early stage ef-
13 fects, such as initial state parton spin correlation, or if it is a final state
14 effect originating from hadronization. Furthermore we, study the feasi-
15 bility to perform this measurement in $p + p$ collisions at $\sqrt{s} = 200$ GeV
16 collected by STAR in 2012 which should provide sufficient statistics of
17 $\Lambda\bar{\Lambda}$, $\Lambda\Lambda$, and $\bar{\Lambda}\bar{\Lambda}$ pairs to perform this measurement.

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1 Introduction

An interesting discovery by Fermilab was published in 1976. They observed that Λ hyperons produced in $p+\text{Be}$ collisions with a 300 GeV proton beam are polarized [1]. This observation is surprising because neither the proton beam nor the beryllium target was polarized. As a result, experimentalists and theorists all around the world started investigating this phenomenon.

The Λ hyperon polarization can be measured through reconstruction of a hadronic decay channel $\Lambda^0 \rightarrow p\pi^-$ (and charge conjugate) and subsequent measurement of the angle (θ_p , or θ^*) between the decay proton momentum in the Λ rest frame (p) and a normal vector to the production plane (\hat{n}). This decay channel is selected, because the proton is preferentially emitted in the direction of the Λ polarization in the Λ rest frame. A cartoon illustrating the production plane determination, using the variables defined above, is shown in Fig. 1.

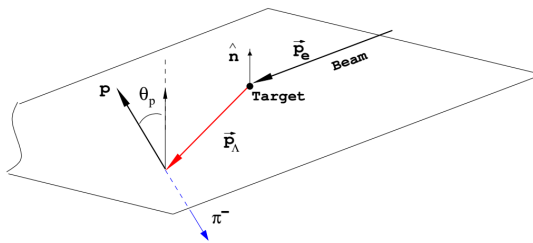


Figure 1: Cartoon illustrating production plane determination. The plane is defined by momentum of the beam (\vec{p}_e) and momentum of Λ (\vec{p}_Λ). Normal vector \hat{n} to the plane is then given by $\hat{n} = \vec{p}_e \times \vec{p}_\Lambda$. The Λ polarization is then quantified by measuring the angle θ_p between the momentum of proton in the rest frame of its mother and \hat{n} . Taken from Ref. [2].

The polarization (P_Λ) is then extracted from the angular distribution of the protons according to formula

$$\frac{dN}{d \cos(\theta^*)} = 1 + \alpha P_\Lambda \cos(\theta^*), \quad (1)$$

where α is the weak decay constant of the Λ hyperon. This method was used in the first Λ hyperon polarization measurement from Ref. [1] and other measurements that followed. It is also possible to measure the Λ hyperon polarization with respect to a different reference direction, e.g. a jet axis, or polarization of the beam for measurements with polarized beams. A brief overview of experimental results using these traditional methods is provided in Sec. 2.

Section 3 provides a description of a new method for Λ hyperon polarization measurement, which relies on the determination of Λ hyperon pair spin-spin correlation.

46 In addition, the section shows first steps of analysis utilizing this method in $p + p$
 47 collisions at $\sqrt{s} = 200$ GeV by the STAR experiment.

48 2 Overview of Λ hyperon polarization results

49 One of the results presented in the first Λ hyperon polarization paper is shown in
 50 Fig. 2. The polarization αP_Λ of the Λ hyperons produced in $p+\text{Be}$ collisions with
 51 300 GeV proton beam on a Be target, rises with the Λ transverse momentum (p_T).

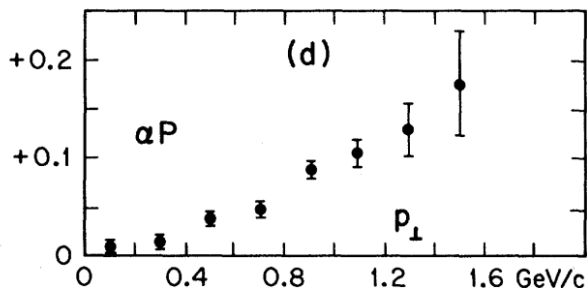


Figure 2: The first observation of Λ hyperon polarization as a function of transverse momentum, measured at Fermilab in $p+\text{Be}$ collisions with 300 GeV proton beam. Taken from Ref [1].

52 In pursuit of an explanation of the origin of Λ hyperon polarization, a number of
 53 independent investigations were performed over a wide range of collision systems and
 54 energies. A selection of such results is shown in Fig. 3. In this case, the polarization
 55 P_Λ is plotted as a function of $x_F = p_z^\Lambda/p_{\text{beam}}$. The key observation here is that the
 56 polarization appears to depend primarily on the x_F and not on the collision type or
 57 energy.

58 All presented results are from collisions of unpolarized particles. It is also im-
 59 portant to investigate, if the polarization of the produced Λ hyperons is correlated
 60 with polarization of the beams, for example, in polarized $\vec{p} + \vec{p}$ collisions at the STAR
 61 experiment. An example of such measurement is presented in Fig. 4 which shows
 62 the longitudinal spin transfer D_{LL} of Λ and $\bar{\Lambda}$ hyperons at positive pseudorapidity
 63 ($0 < \eta < 1.2$) measured in $\vec{p} + \vec{p}$ collisions at $\sqrt{s} = 200$ GeV. No significant longitu-
 64 dinal polarization of Λ hyperons is observed at STAR, which suggests that the beam
 65 polarization does not play a significant role for the polarization of the Λ hyperon at
 66 RHIC energies within the studied kinematic range¹.

¹The x_F in this η region at RHIC is going to be rather small which will likely lead to small Λ polarization, as seen in Fig. 3. It is not possible to make direct comparison to Fig. 3 due to polarization of the beams and also a different observable of the measurement.

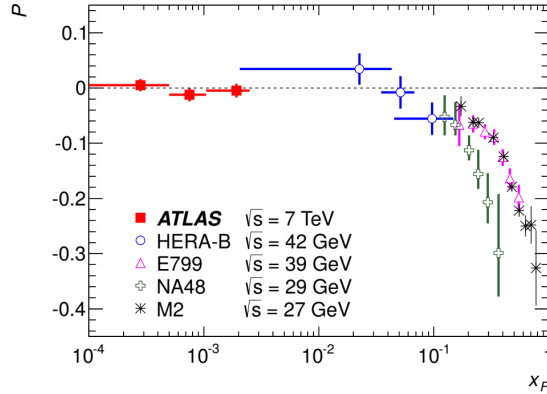


Figure 3: Comparison of the Λ hyperon polarization as a function of x_F measured in variety of collision systems and collision energies. Taken from Ref. [3].

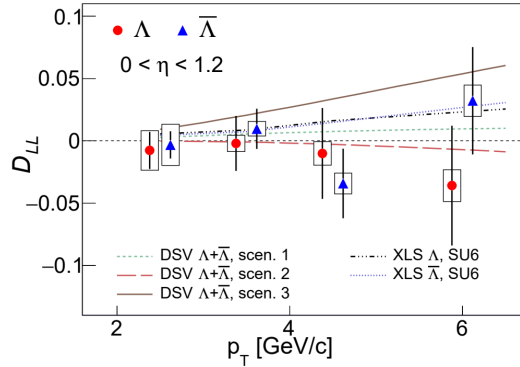


Figure 4: Longitudinal spin transfer D_{LL} of Λ (red circles) and $\bar{\Lambda}$ (blue circles) hyperons at positive pseudorapidity ($0 < \eta < 1.2$) measured in $\vec{p} + \vec{p}$ collisions at $\sqrt{s} = 200$ GeV by the STAR experiment. Taken from Ref [5].

67 The examples above are a small selection of all experimental efforts over the last 50
68 years in attempt to explain the origin of Λ hyperon polarization. Unfortunately, none
69 of the measurements, or the theoretical models, can provide a definitive answer on
70 where the Λ hyperon polarization is generated. In the following section, we investigate
71 a possibility to improve our knowledge of the phenomenon by measurement of Λ
72 hyperon pair spin-spin correlations.

73 **3 Λ hyperon spin-spin correlations**

74 In general, most experimental techniques for measurement of Λ hyperon polarization
75 are based on the same general idea. It is the measurement of an angle, usually

76 denoted θ^* , between a reference direction and momentum of the decay proton in the
 77 hyperon rest frame. The reference direction can be chosen based on specific physics
 78 considerations. As shown above, the first possibility is to use the production plane.
 79 Other common alternatives are, e.g., the polarization of the beam, or a jet axis, in
 80 case the Λ hyperon is part of a jet.

81 Another possibility is to look for events with two or more Λ or $\bar{\Lambda}$ hyperons and
 82 measure the angle θ_{12}^* between the decay (anti-)protons, both boosted into the rest
 83 frame of their mother particles. Since these (anti-) protons are preferentially emitted
 84 in the direction of their mother's polarization, such measurement gives access to $\Lambda\bar{\Lambda}$,
 85 $\Lambda\Lambda$, and $\bar{\Lambda}\bar{\Lambda}$ pair spin-spin correlations [6, 7]. For this method, it is possible to use
 86 the following formula (see also Ref. [7]):

$$\frac{dN}{d\cos(\theta_{12}^*)} = 1 + \alpha_1\alpha_2 P_{\Lambda_1\Lambda_2} \cos(\theta_{12}^*), \quad (2)$$

87 where α_1 and α_2 are weak decay constants of Λ hyperons in the pair and $P_{\Lambda_1\Lambda_2}$ is the
 88 polarization of the pair. Λ_1 and Λ_2 can both be either Λ or $\bar{\Lambda}$.

89 One of the key advantages of this approach is that it should be able to identify if
 90 the polarization comes from initial stage effects, such as spin-spin correlation of the
 91 initial stage partons, or if it is a final state effect originating from e.g. hadronization.
 92 The initial state correlation should be seen in data as spin-spin correlation of the
 93 $\Lambda\bar{\Lambda}$ pairs, as those likely originate from a single $s\bar{s}$ quark pair produced in the hard
 94 partonic scattering. At the same time, no strong correlation is expected for $\Lambda\Lambda$ and
 95 $\bar{\Lambda}\bar{\Lambda}$ pairs, as those cannot originate from a single $s\bar{s}$ quark pair.

96 In order to investigate the Λ hyperon pair spin-spin correlations in $p+p$ collisions
 97 at $\sqrt{s} = 200$ GeV at STAR, it is important to verify that there are no other known
 98 mechanisms to generate non-zero $P_{\Lambda_1\Lambda_2}$. This was done with PYTHIA 8.3 simulation
 99 of $p+p$ collisions at $\sqrt{s} = 200$ GeV. The extracted $1/N_{\text{evt}}dN/d\cos(\theta_{12}^*)$ distributions
 100 as a function of $\cos(\theta_{12}^*)$ for $\Lambda\bar{\Lambda}$ and $\Lambda\Lambda$ pairs are shown in Fig. 5. The distributions
 101 are fitted with a linear function which is used to extract the value of $P_{\Lambda_1\Lambda_2}$ using
 102 equation (2). For both combinations, the polarization is zero, meaning that pure
 103 PYTHIA does not predict any Λ hyperon spin-spin correlations at mid-rapidity in
 104 $p+p$ collisions at $\sqrt{s} = 200$ GeV.

105 The extraction of the $dN/d\cos(\theta_{12}^*)$ distributions from the data starts with se-
 106 lection of Λ and $\bar{\Lambda}$ hyperon candidates. This is done by pairing protons and pions
 107 reconstructed and identified with the STAR Time Projection Chamber. Each Λ can-
 108 didate then corresponds to one $p\pi$ pair. Events with two or more Λ candidates were
 109 considered for further analysis. For events which contain at least two such pairs, a
 110 2D distribution is filled where one axis is the invariant mass of one of the $p\pi$ pairs
 111 and the second axis is the invariant mass of the second $p\pi$ pair.

112 This is done for two combinations of the $p\pi$ pairs: an unlike-sign (US) $p\pi$ pair
 113 matched with a different US pair from the same event. This distribution contains both

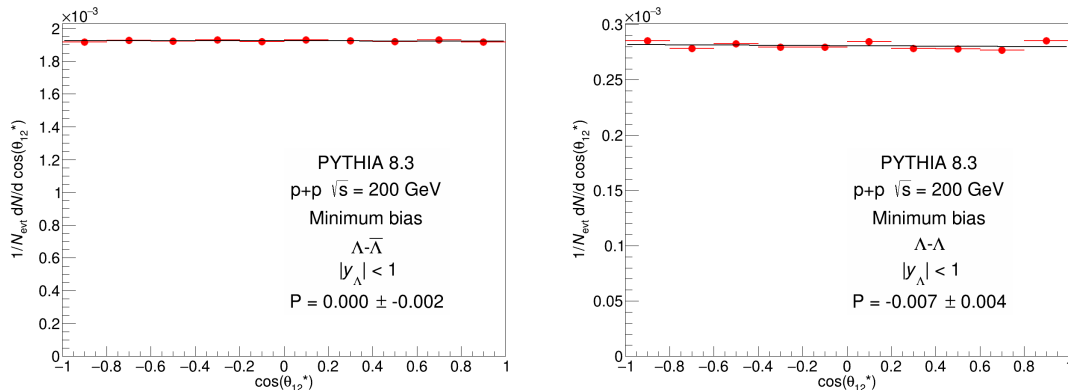


Figure 5: $1/N_{\text{evt}}dN/d \cos(\theta_{12}^*)$ as a function of $\cos(\theta_{12}^*)$ for $\Lambda\bar{\Lambda}$ (left) and $\Lambda\Lambda$ (right) pairs from PYTHIA 8.3 $p+p$ collisions at $\sqrt{s} = 200$ GeV. The simulation is fitted with a linear function which is used to extract the polarization P according to equation (2).

114 the signal and the combinatorial background. The background can be estimated using
 115 US pairs matched to the like-sign (LS) $p\pi$ pairs. The US-LS distribution is then
 116 subtracted from US-US distribution and subsequently fitted with a 2D Gaussian
 117 function to determine the Λ candidate invariant mass peak mean and width. The
 118 signal region is defined as the mean $\pm 3\sigma$, where both the mean and σ are taken from
 119 the fit. Example of two of the 2D invariant mass distributions is shown in Fig. 6.

120 This procedure is done separately for $\Lambda\bar{\Lambda}$, $\Lambda\Lambda$, and $\bar{\Lambda}\bar{\Lambda}$ candidate pairs in $p+p$
 121 collisions at $\sqrt{s} = 200$ GeV measured by STAR in 2012. The extracted numbers of
 122 signal and background pairs for each of the three possible combinations is shown in
 123 Tab. 1. The number of candidate pairs is going to provide sufficient statistics to
 124 perform this measurement using the 2012 $p+p$ collisions data-set.

p_T [GeV/ c]	$\Lambda\bar{\Lambda}$		$\Lambda\Lambda$		$\bar{\Lambda}\bar{\Lambda}$	
	N_{sig}	N_{bkg}	N_{sig}	N_{bkg}	N_{sig}	N_{bkg}
Λ_1 0.5 – 5.0	26767	6929	7090	3593	4670	2495
Λ_2 0.5 – 5.0						

Table 1: Number of $\Lambda\bar{\Lambda}$, $\Lambda\Lambda$, and $\bar{\Lambda}\bar{\Lambda}$ signal candidate and background pairs in $p+p$ collisions at $\sqrt{s} = 200$ GeV measured by STAR in 2012.

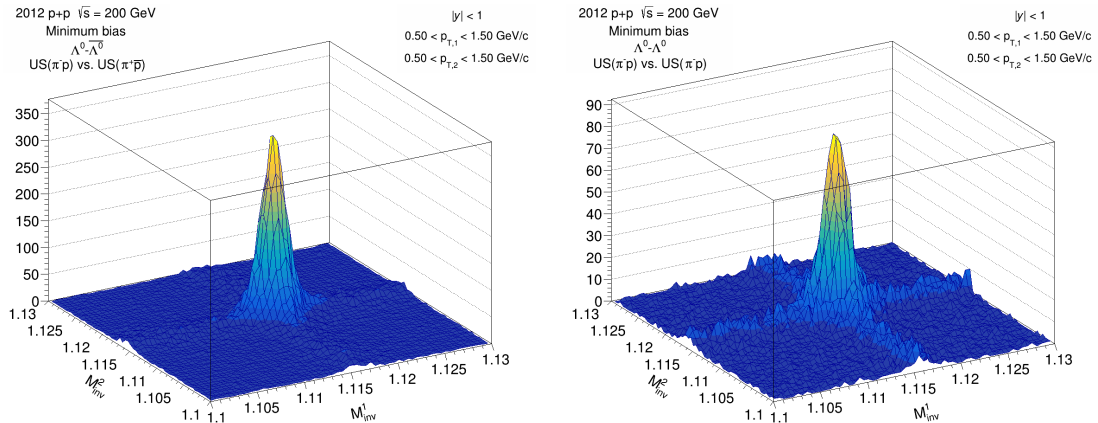


Figure 6: Two-dimensional invariant mass distributions of unlike-sign $p\pi$ pairs matched with unlike-sign $p\pi$ pairs. Left panel shows $p\pi$ pair charge combinations corresponding to $\Lambda\bar{\Lambda}$ pair candidates, the right panel shows the $\Lambda\Lambda$ pair candidates.

125 4 Summary

126 The Λ hyperon polarization puzzle is one of the main unresolved mysteries of the ex-
 127 perimental high energy particle physics. The polarization has been observed in several
 128 different collision systems at various energies. The magnitude of the polarization ap-
 129 pears to depend primarily on x_F and not much on the specific collision energy. Despite
 130 enormous experimental and theoretical efforts to explain the Λ hyperon polarization,
 131 no conclusive answer was found. In order to improve the knowledge in this field, a
 132 new method was developed which relies on measurement of Λ hyperon pair spin-spin
 133 correlations. Any non-zero signal measured in $p + p$ collisions at $\sqrt{s} = 200$ GeV at
 134 STAR would provide more insight into Λ hyperon polarization, as a simulation using
 135 PYTHIA 8.3 predicts no spin-spin correlation signal. The number of extracted $\Lambda\bar{\Lambda}$,
 136 $\Lambda\Lambda$, and $\bar{\Lambda}\bar{\Lambda}$ candidate pairs extracted from the aforementioned STAR data-set is
 137 sufficient to perform this type of measurement and thus is going to provide important
 138 additional insight into Λ hyperon polarization physics in $p + p$ collisions at RHIC
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