

# 1 Measurement of $\Lambda$ hyperon spin-spin correlations in 2 $p+p$ collisions by the STAR experiment

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7 The  $\Lambda^0$  hyperons have an interesting property that they can be produced polarized in collisions of unpolarized particles, as was discovered in Femilab in the 70's. Since then, the  $\Lambda^0$  polarization has been measured in a variety of collision systems, including  $e^+e^-$ . Most of the measurements indicate that the polarization originates from final state effects, such as fragmentation and hadronization. It is therefore important to investigate if there is any contribution of initial state effects to the polarization, such as spin-spin correlation of initial state partons. This is possible utilizing a newly proposed experimental method which measures spin-spin correlations of  $\Lambda^0$  hyperon pairs.

In these proceedings, we present preliminary results from the first experimental measurement of spin-spin correlations of  $\Lambda^0\bar{\Lambda}^0$ ,  $\Lambda^0\Lambda^0$ , and  $\bar{\Lambda}^0\bar{\Lambda}^0$  pairs in  $p+p$  collisions at  $\sqrt{s} = 200$  GeV measured by the STAR experiment. The spin-spin correlation is measured to be consistent with zero for all three  $\Lambda^0$  hyperon pair combinations, within the uncertainties. This measurement provides the first limit on the  $\Lambda^0$  hyperon pair spin-spin correlations in  $p+p$  collisions at  $\sqrt{s} = 200$  GeV.

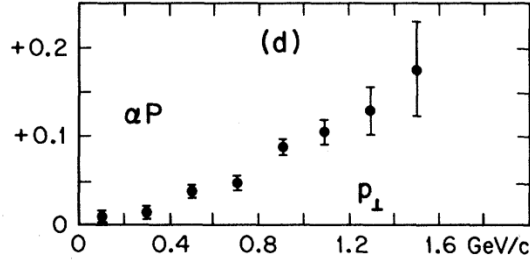
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## 8 1. Introduction

9 In the 70's, an interesting discovery has been made at Fermilab. It was observed that  $\Lambda^0$   
 10 hyperons are produced polarized in collisions of proton beam on beryllium target [1], as shown in  
 11 Fig. 1. This result was surprising, as neither the 300 GeV proton beam, nor the beryllium target  
 12 was polarized in this experiment.

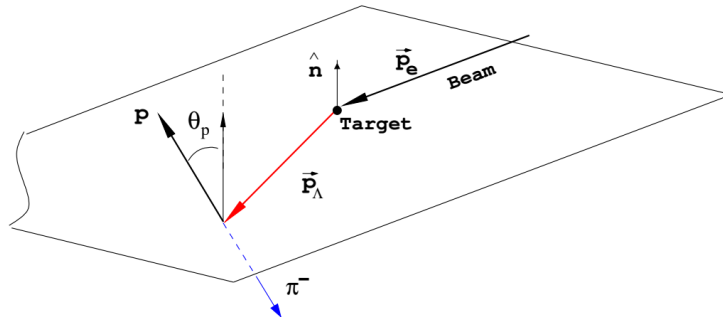


**Figure 1:** Polarization of  $\Lambda^0$  hyperons as a function of transverse momentum  $p_T$  measured by Fermilab in  $p+\text{Be}$  collisions. Taken from Ref. [1].

13 Since its discovery, the  $\Lambda^0$  hyperon polarization has been measured in a variety of collision  
 14 systems such as  $p+p$ ,  $e+p$ , or  $e^+e^-$  collisions. These measurements revealed more surprises. One  
 15 interesting observation is that the measured polarization does not heavily depend on the collision  
 16 energy or the collision system, but rather on  $x_F$  which is defined as:

$$x_F = \frac{p_{z,\Lambda}}{p_{\text{beam}}}, \quad (1)$$

17 where  $p_{z,\Lambda}$  is the component of the  $\Lambda^0$  hyperon momentum along the beam axis and  $p_{\text{beam}}$  is the  
 18 beam momentum. The polarization increases with larger  $x_F$  which is well summarized in Ref.  
 19 [2]. Another important observation is that the polarization is present also in  $e^+e^-$  collisions, as  
 20 measured by the BELLE Collaboration [3]. This result suggests that the polarization is generated  
 21 by final state effects, such as fragmentation or hadronization, as there are no partons in the initial  
 22 stage that could affect the  $\Lambda^0$  hyperon polarization in the final state.



**Figure 2:** Cartoon illustrating definition of the production plane in measurement of  $\Lambda^0$  hyperon polarization. Taken from Ref. [4].

23 What the aforementioned results have in common is the experimental method. The  $\Lambda^0$  hyperons  
 24 are reconstructed via hadronic decay channel  $\Lambda^0 \rightarrow p\pi^-$ . The polarization is then measured with  
 25 respect to a vector  $\hat{n}$  normal to the production plane, defined as  $\hat{n} = \vec{p}_{\text{beam}} \times \vec{p}_{\Lambda}$ , as shown in Fig.  
 26 2. What is then experimentally measured is an angle  $\theta^*$  (sometimes  $\theta_p$ ) between vector  $\hat{n}$  and the  
 27 decay proton momentum boosted into the  $\Lambda^0$  hyperon rest frame. The distribution of such boosted  
 28 protons then follows the formula:

$$\frac{dN}{d \cos \theta^*} = 1 + \alpha P_{\Lambda} \cos \theta^* \quad (2)$$

29 where  $\alpha$  is the weak decay constant of the  $\Lambda^0$  hyperon and  $P_{\Lambda}$  is the polarization.

## 30 2. $\Lambda^0$ hyperon spin-spin correlations

31 The measurements mentioned in the previous section generally aim to measure and explain  
 32 the origin of a single  $\Lambda^0$  hyperon polarization. Recently, it was proposed to investigate the  $\Lambda^0$   
 33 hyperon polarization puzzle from a slightly different perspective, that aims to determine if there is  
 34 any spin-spin correlation of  $s\bar{s}$  quark pair produced in the hard partonic scattering. This is possible  
 35 through measurement of  $\Lambda^0\bar{\Lambda}^0$ ,  $\Lambda^0\Lambda^0$ , and  $\bar{\Lambda}^0\bar{\Lambda}^0$  hyperon pair spin-spin correlations [5, 6].

36 The experimental method for this new approach is similar to the standard one described  
 37 earlier. The two key differences are that it is necessary to identify events containing at least one  
 38  $\Lambda^0$  hyperon pair (any combination), rather than single hyperons, and the reference direction used  
 39 for measurement of angle  $\theta^*$  is different. In this case, the angle  $\theta^*$  is measured between the  
 40 momenta of the decay (anti-)protons boosted into rest frame of their mother. The distribution of  
 41 such (anti-)proton pairs then follows:

$$\frac{dN}{d \cos \theta^*} = 1 + \alpha_1 \alpha_2 P_{\Lambda_1 \Lambda_2} \cos \theta^* \quad (3)$$

42 where  $\alpha_1$  and  $\alpha_2$  are the weak decay constants of  $\Lambda^0$  or  $\bar{\Lambda}^0$  and  $P_{\Lambda_1 \Lambda_2}$  is the spin-spin correlation of  
 43 the hyperon pair. In case  $P_{\Lambda_1 \Lambda_2} \neq 0$ , the spins of the  $\Lambda^0$  ( $\bar{\Lambda}^0$ ) in the pair are correlated.

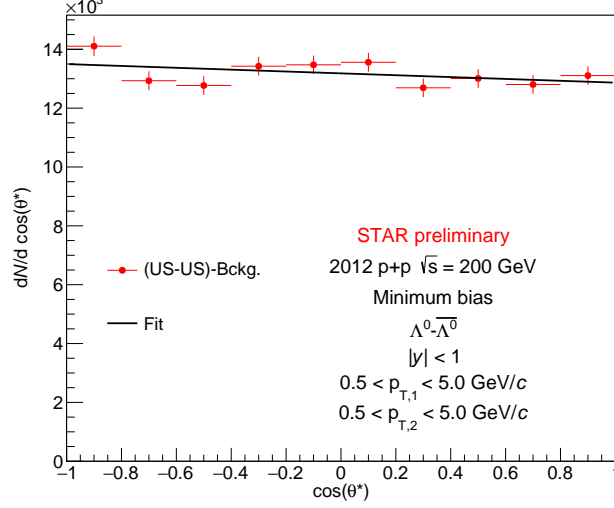
## 44 3. Results

45 The dataset used in this analysis are  $p+p$  collisions at  $\sqrt{s} = 200$  GeV collected by the STAR  
 46 experiment in 2012, which contains a total of 400 million minimum bias events. The  $\Lambda^0$  and  
 47  $\bar{\Lambda}^0$  hyperons are reconstructed at mid-rapidity ( $|y| < 1$ ) and in transverse momentum range of  
 48  $0.5 < p_T < 5.0$  GeV.

49 As described earlier, the  $\Lambda^0$  and  $\bar{\Lambda}^0$  are reconstructed via their hadronic decay  $\Lambda^0 \rightarrow p\pi^-$  (and  
 50 charge conjugate for  $\bar{\Lambda}^0$ ). The (anti-)protons and charged pions, identified utilizing STAR Time  
 51 Projection Chamber, are paired based on their charge into unlike-sign (US) pairs ( $p\pi^-$ ,  $\bar{p}\pi^+$ ) and  
 52 like-sign (LS) pairs ( $p\pi^+$ ,  $\bar{p}\pi^-$ ). These pairs are subsequently combined into US-US  $\Lambda^0$  hyperon  
 53 candidate pairs, containing signal and combinatorial background. The background is estimated  
 54 using US-LS pair combination.

55 The signal  $\Lambda^0$  hyperon pairs are selected from a 2D invariant mass ( $M_{\text{inv}}$ ) distribution of the  
 56  $\Lambda^0$  hyperon pairs. First the combinatorial background (US-LS)  $M_{\text{inv}}$  distribution is subtracted from

57 the US-US  $M_{\text{inv}}$  distribution. The  $M_{\text{inv}}$  distribution after background subtraction is then fitted with  
 58 a 2D Gaussian distribution. The signal is extracted based on parameters of the fit within  $\mu \pm 3\sigma$ .



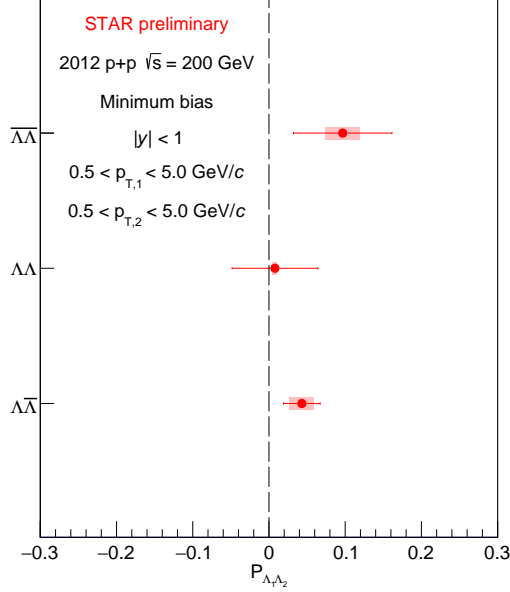
**Figure 3:**  $dN/d \cos \theta^*$  distribution of  $\Lambda^0 \bar{\Lambda}^0$  after acceptance effect correction and background subtraction. The data are fitted according to Eq. 3.

59 In the next step, the  $dN/d \cos \theta^*$  distributions of the selected  $\Lambda^0$  hyperon pairs are filled for  
 60 both US-US and US-LS charge combinations and all three possible  $\Lambda^0 (\bar{\Lambda}^0)$  pair combinations.  
 61 These measured distributions are corrected for acceptance using mixed-event (ME)  $dN/d \cos \theta^*$   
 62 distributions, separately for US-US and US-LS. After the acceptance correction, the background  
 63 (US-LS) distribution is subtracted from US-US distribution. The background subtracted distribution  
 64 is then fitted according to Eq. 3 and the spin-spin correlation  $P_{\Lambda_1 \Lambda_2}$  is extracted. The  $dN/d \cos \theta^*$   
 65 distribution of  $\Lambda^0 \bar{\Lambda}^0$  after acceptance effect correction and background subtraction is shown in Fig.  
 66 3.

67 The spin-spin correlations of  $\Lambda^0 \bar{\Lambda}^0$ ,  $\Lambda^0 \Lambda^0$ , and  $\bar{\Lambda}^0 \bar{\Lambda}^0$  pairs at mid-rapidity ( $|y| < 1$ ), with  
 68 transverse momentum  $0.5 < p_T < 5.0 \text{ GeV}$ , measured in  $p+p$  collisions at  $\sqrt{s} = 200 \text{ GeV}$  are  
 69 shown in Fig. 4. The spin-spin correlations are consistent with zero for all three  $\Lambda^0$  hyperon pair  
 70 combinations, within the uncertainties. This measurement provides the first limit on the  $\Lambda^0$  hyperon  
 71 pair spin-spin correlations in  $p+p$  collisions at  $\sqrt{s} = 200 \text{ GeV}$ .

## 72 4. Summary

73 The  $\Lambda^0$  hyperons can be produced polarized in collisions of unpolarized particles, as was  
 74 discovered in the 70's at Fermilab. Various experimental techniques and theoretical models have  
 75 been developed over the last 50 years in the aim to explain the  $\Lambda^0$  hyperon polarization puzzle.  
 76 One such method, that was recently proposed, is the measurement of  $\Lambda^0$  hyperon pair spin-spin  
 77 correlations that is able to determine if there is any spin-spin correlation of the initial state  $s\bar{s}$  quark  
 78 pairs [5, 6]. In these proceedings, we have presented the first experimental measurement of the  $\Lambda^0$   
 79 hyperon pair spin-spin correlations by the STAR experiment, in  $p+p$  collisions at  $\sqrt{s} = 200 \text{ GeV}$ .



**Figure 4:** Spin-spin correlations of  $\Lambda^0\bar{\Lambda}^0$ ,  $\Lambda^0\Lambda^0$ , and  $\bar{\Lambda}^0\bar{\Lambda}^0$  pairs measured in  $p+p$  collisions at  $\sqrt{s} = 200$  GeV.

80 The spin-spin correlation of the  $\Lambda^0\bar{\Lambda}^0$ ,  $\Lambda^0\Lambda^0$ , and  $\bar{\Lambda}^0\bar{\Lambda}^0$ , measured at mid-rapidity ( $|y| < 1$ ), with  
 81 transverse momentum  $0.5 < p_T < 5.0$  GeV is found to be consistent with zero for all three  $\Lambda^0$   
 82 hyperon pair combinations, within the corresponding uncertainties. This measurement provides the  
 83 first limit on the  $\Lambda^0$  hyperon pair spin-spin correlations in  $p+p$  collisions at  $\sqrt{s} = 200$  GeV.

## 84 References

- 85 [1] Bunce G., *et al.*,  $\Lambda^0$  Hyperon Polarization in Inclusive Production by 300-GeV Protons on  
 86 Beryllium, Phys. Rev. Lett. 36, 1113 (1976)
- 87 [2] Aad G., *et al.* [ATLAS Collaboration], Measurement of the transverse polarization of  $\Lambda^0$  and  
 88  $\bar{\Lambda}^0$  hyperons produced in proton-proton collisions at  $\sqrt{s} = 7$  TeV using the ATLAS detector,  
 89 Phys. Rev. D 91, 032004 (2015)
- 90 [3] Guan Y., *et al.* [BELLE Collaboration], Observation of Transverse  $\Lambda^0/\bar{\Lambda}^0$  Hyperon Polariza-  
 91 tion in  $e^+e^-$  Annihilation at Belle, Phys. Rev. Lett. 122, 042001 (2019)
- 92 [4] Airapetian A., *et al.* [HERMES Collaboration], Transverse Polarization of  $\Lambda^0$  and  $\bar{\Lambda}^0$  Hyperons  
 93 in Quasireal Photoproduction, Phys.Rev.D 76, 092008 (2007)
- 94 [5] Gong, W., *et al.*, Measurement of Bell-type inequalities and quantum entanglement from  
 95  $\Lambda^0$ -hyperon spin correlations at high energy colliders, Phys. Rev. D 106, L031501 (2022)
- 96 [6] Törnqvist, N.A. Suggestion for Einstein-Podolsky-Rosen experiments using reactions like  
 97  $e^+e^- \rightarrow \Lambda^0\bar{\Lambda}^0 \rightarrow \pi^-p\pi^+\bar{p}$ ., Found Phys 11, 171–177 (1981)