# Determining $\pi^{0} A_{\text {LL }}$ from STAR 2012 Endcap Calorimeter Data 

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## Background

The STAR Experiment, located at the Relativistic Heavy Ion Collider (RHIC) at the Brookhaven National Laboratory, uses longitudinally polarized proton-proton collisions to produce neutral pions ( $\pi^{\circ}$ ) whose production is assumed to exhibit a spin asymmetry $\left(A_{L L}\right)$. By measuring $A_{L L}$ the gluon spin contribution to the spin of the proton can be studied. The STAR detector (Fig. 1) is located along the 2.3 mile RHIC circumference (Fig. 2). As seen in Fig. 3, a neutral pion decays into two photons immediately ( $\sim 10^{-16} \mathrm{~s}$ ). The analysis reported here is based on data taken in 2012 for 510 GeV center of mass
energy of the colliding protons


Fig. 1 STAR Laboratory


Fig. 2 Relativistic Heavy Ion Collider (RHIC) at Brookhaven National Laboratory

## Invariant Mass Spectra

The Endcap Electromagnetic Calorimeter (EEMC) of the STAR detector (Fig. 1) is able to measure the energy ( E ) and position of incident photon's electromagnetic shower. By using the energy and the opening angle ( $\theta$ ) between the two photons (Fig. 3), the two photon invariant mass plot can be calculated using the following equation:

$$
\text { Invariant mass }=\left(E_{1}+E_{2}\right) \sqrt{1-\left(\frac{E_{1}-E_{2}}{E_{1}+E_{2}}\right)^{2}} \sin \frac{\theta}{2}
$$



The two photon invariant mass histograms for each spin state were fitted with a skewed gaussian for the $\pi^{0}$ peak and a background function. Two different background functions were compared; a Chebyshev function and a modified Planck function. The number of $\pi^{\circ}$ s present is found by integrating the fitted Gaussian function above the background. This number is then used to find the number of $\pi^{\circ}$ s produced in each spin state. This information can then be used to find the $\pi^{0}$ asymmetry $\left(\mathrm{A}_{\mathrm{LL}}\right)$

Skewed Gaussian =
$a e^{-0.5\left(\frac{x-b}{c(1+d(x-b))}\right)^{2}}$
$c=$ sigma
$d=$ 'skewking' parameter

Modified Planck Function


$$
=\frac{1}{\left(x-c_{1}\right)^{5} e^{\frac{c_{2}}{x-c_{3}}}-c_{4}}
$$

Fig. 4 (left) Invariant mass plot at a $\mathrm{p}_{\text {T }}$ (transverse momentum of the $\pi^{0}$ ) range Planck background function (blue) and a skewed Gaussian for the $\pi^{0}$ peak (red).


Chebyshev Function
$c_{0}+c_{1}(x)+c_{2}\left(2 x^{2}-1\right)+c_{3}\left(4 x^{3}-3 x\right)+c_{4}\left(8 x^{4}-8 x^{2}+1\right)$

## False Asymmetry

False Asymmetries are a good measure to cross-check the relative luminosity determinations, and are expected to be zero. The analysis to determine the false asymmetries was done on a fill-by-fill basis. The polarizations and relative luminosity parameters in the calculations are weighted averages for fills - weighted using the number of triggers in a run. There are 4 different false asymmetries: Single Spin Symmetries associated with each of the two beams, Like-spin Double Spin Symmetry, and Unlike-spin Double Spin Symmetry. For Figs. 6 and 7 the number of $\pi^{\circ}$ s, N, used to calculate the false asymmetries was the integral of the skewed Gaussian function. The error was taken as $\sqrt{ } \mathrm{N}$.


Fig. 6 The single spin asymmetry ) plotted as a function of two beams $\mathbf{A}_{\mathrm{LY}}$ ) plotted as a function of $\mathrm{p}_{\mathrm{T}}$. The red
ine shows a fit to the false asymmetries o have the value of po $=-0.0006$ $\pm 0.0015$ which is consistent with zero. However, the reduced chi-square $\left(X^{2}\right)$
value is (11.8/6) which is higher than expected indicating a possible expected indicating ${ }^{\text {a }}$ a


Fig. 7 The Like-sign double spin asymmetry $\mathrm{p}_{\mathrm{T}}$. The red line shows a fit to the false $\mathrm{p}_{\mathrm{T}}$ asymmetries to have the value of po $=-0.0007 \pm 0.0045$ which is consistent with zero. However, the which is higher than expected indicating a possible underestimation of the errors.

## Quality Assurance

The analysis of the data is in the final stages so quality assurance of the data was necessary to assure good runs were used. QA was executed for collider fills (groups of runs) as well as individual runs. A C++ script was used to look at key parameters such as number of towers, $\pi$ mass, and signal fraction. Runs were flagged if they exceeded $2 \sigma$ to be reviewed individually. A fill-by-fill QA was done in a similar fashion.

ig. 8 (left) The $\pi^{0}$ mean mas Tlotted as a function of run number nominal mass of a $\pi^{0}$ because of nominal mas

Fig. 9 (right). The fill average $\pi^{0}$ mass plotted as a function of fill number.The declining trend of $\pi^{0}$ mass

## Determining $\mathrm{A}_{\mathrm{LL}}$

The $\pi^{0} \mathrm{~A}_{\mathrm{LL}}$ is an imbalance in the $\pi^{0}$ production in the different spin states of the colliding protons. Each proton can have positive or negative helicity (spin aligned or anti-aligned with its momentum). There are four different spin combinations for collisions, but they can be grouped into like-helicity or unlike-helicity cases. Using the following formulas, the $\mathrm{A}_{\mathrm{LL}}$ can be calculated.
$\mathrm{N}^{++}=$number of $\pi^{\circ} \mathrm{s}$ from like-helicity collisions
$\mathrm{N}^{+-}=$number of $\pi^{0}$ s from unlike-helicity collisions
$\mathrm{L}=$ luminosities at different spin states

$$
R_{3}=\frac{L^{++}+L^{--}}{L^{+-}+L^{-+}}
$$

$\mathrm{P}_{1} \& \mathrm{P}_{2}=$ polarization of the two beams
$\mathrm{R}_{3}=$ relative luminosity

$$
A_{L L}=\frac{\sum_{\text {runs }} P_{1} P_{2}\left(N^{++}-R_{3} N^{+-}\right)}{\sum_{\text {runs }} P_{1}^{2} P_{2}^{2}\left(N^{++}+R_{3} N^{+-}\right)}
$$

## Future Work

The analysis of the 2012 dataset is in its final stages so work has begun with preparing the larger, 2013 data set for analysis. Work is underway to retrieve the 2013 data from tape, reconstruct the $\pi^{\circ}$ s and store it in a ROOT tree format for further analysis. As for the 2012 dataset, a quality assurance (QA) proceeds in parallel to ensure the processing is successful and the data quality is high.

