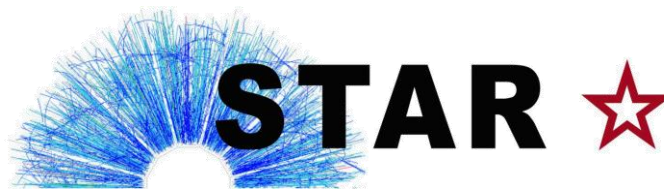


Measurement of transverse single-spin asymmetries for dijet production in polarized p+p collisions at $\sqrt{s} = 200$ GeV at STAR

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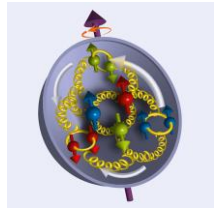


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The Siverts Effect in pp Dijet Production

Proton Spin



$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma$$

$$+ \Delta G$$

$$+ (\sum_q \mathcal{L}_q + \mathcal{L}_g)$$

Quark polarization : ~30% contribution

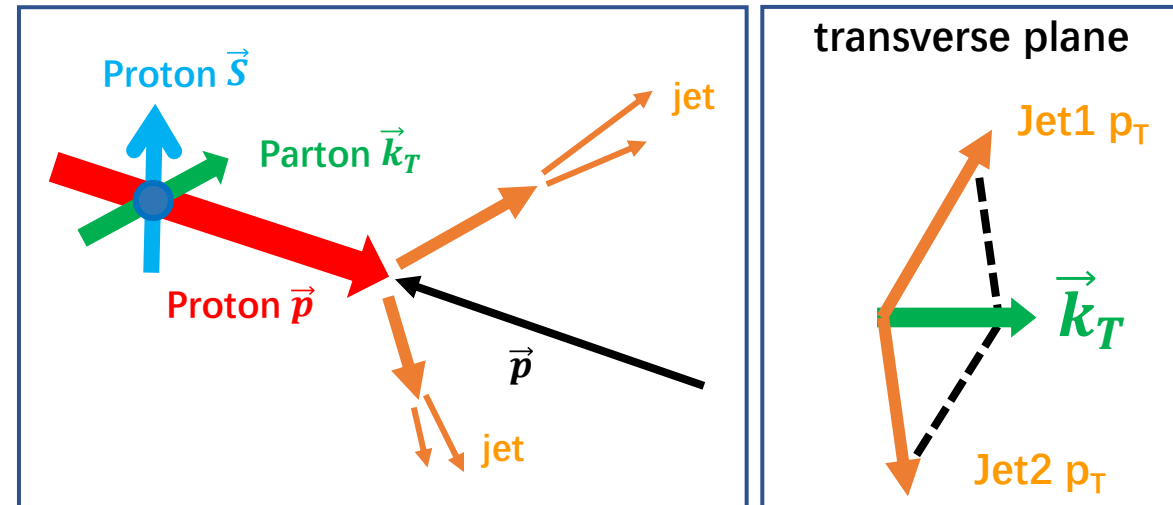
Gluon polarization : comparable to $\Delta\Sigma$, less well constrained

Orbital angular momentum (OAM) : largely unconstrained

- The parton OAM can be manifested via the Siverts Effect, a spin-dependent average transverse momentum:

$$\langle \vec{S}_{proton} \cdot (\vec{p}_{proton} \times \vec{k}_T) \rangle \neq 0$$

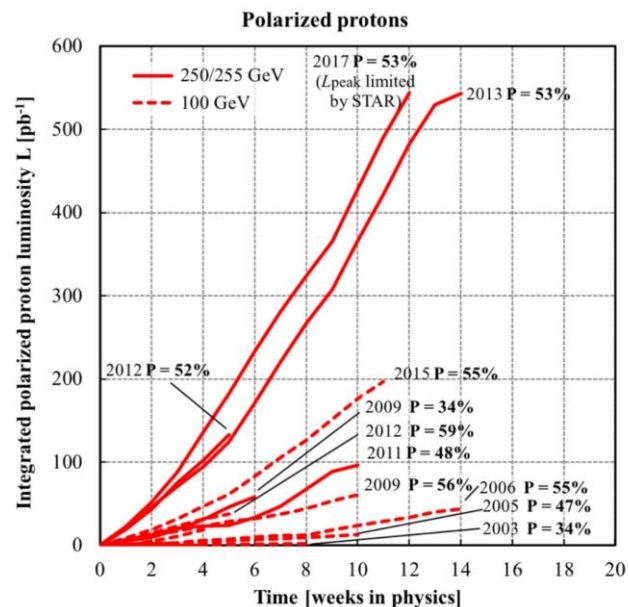
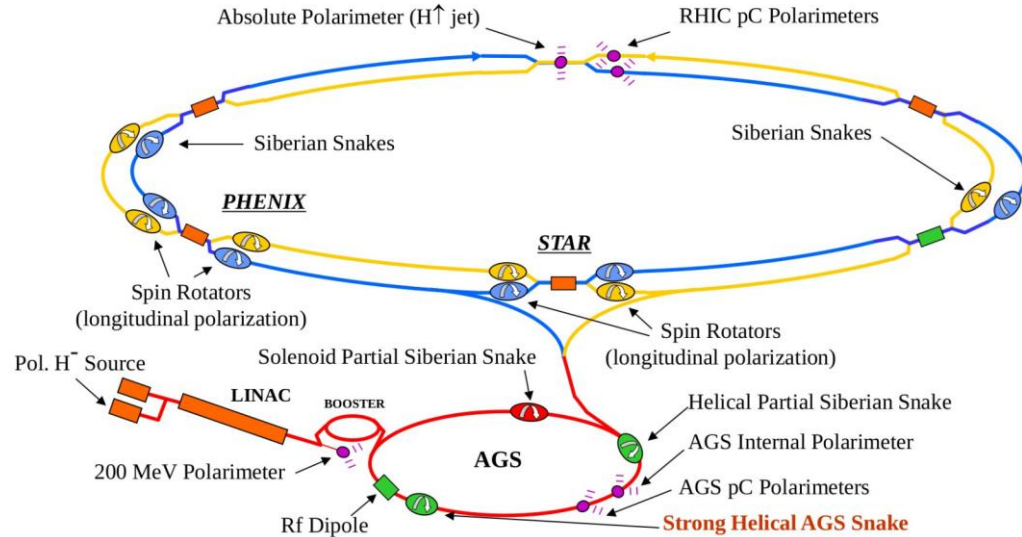
- The u -quark \vec{k}_T and the d -quark \vec{k}_T are expected to be opposite in sign and different in magnitude.
- The Siverts effect can be measured in dijet production by examining the tilt in the back-to-back dijet opening angle.



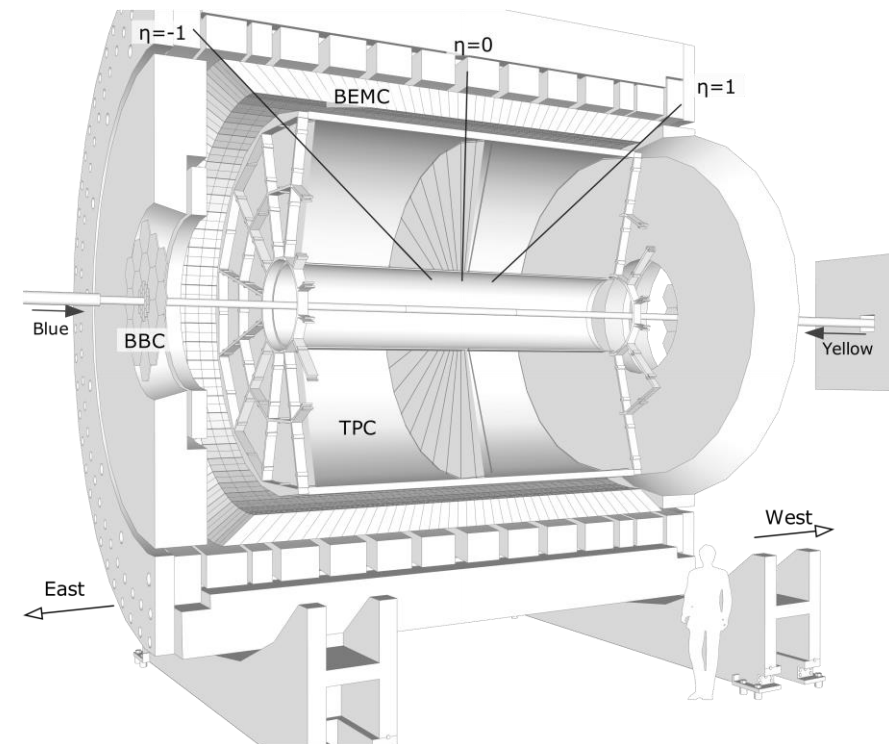
- Measuring the Siverts effect in dijet production :
 - Non-zero effects indicate possible contributions from partonic angular momentum to the proton spin.
 - Test the expected features of the Siverts effect (sign, magnitude) for flavor-separated partons
 - Help constraining the Siverts function, and explore the Siverts effect at a larger Q^2 scale than SIDIS

RHIC & STAR Detector

Relativistic Heavy Ion Collider



Solenoidal Tracker At RHIC



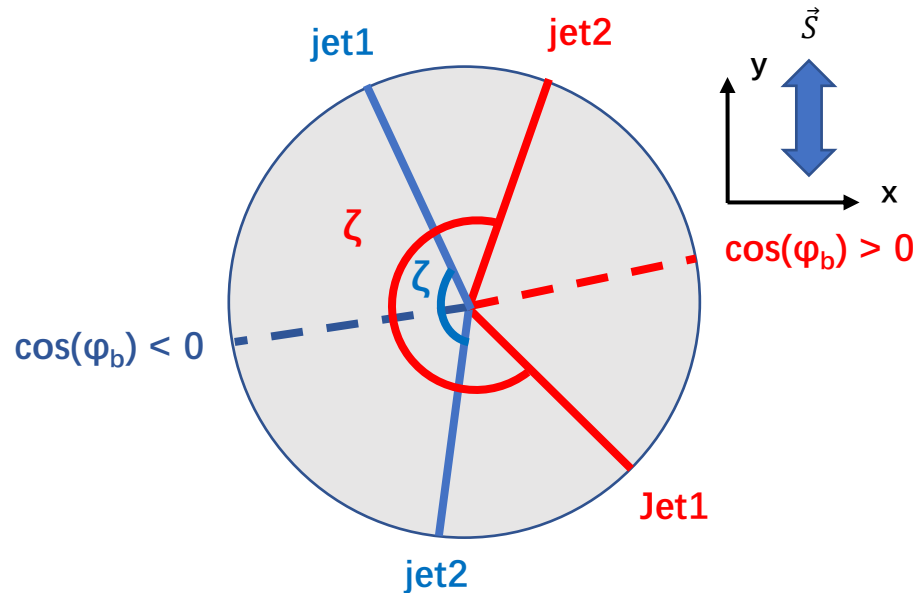
- **RHIC** – colliding transversely and longitudinally polarized pp at different energies (200, 500, 510 GeV, etc).
- **STAR detector** is capable of reconstructing tracks, identifying charged particles in $|\eta| < 1.3$, and measuring EM particle energies in $-1 < \eta < 2$.

Observable for Probing the Sivvers Effect in Dijet Event

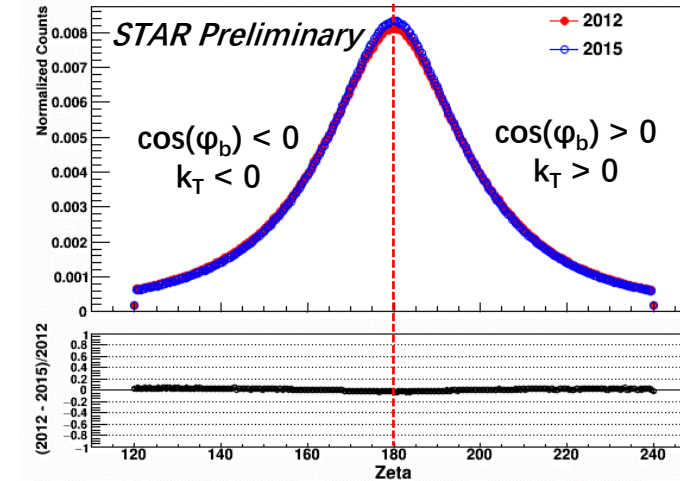
The Sivvers asymmetry can be probed via the signed opening angle ζ .

Definition of ζ

$\zeta > \pi$ when $\cos(\varphi_b) > 0$
 $\zeta < \pi$ when $\cos(\varphi_b) < 0$
 where φ_b is dijet bisector angle



Distribution of ζ



Extraction of asymmetry

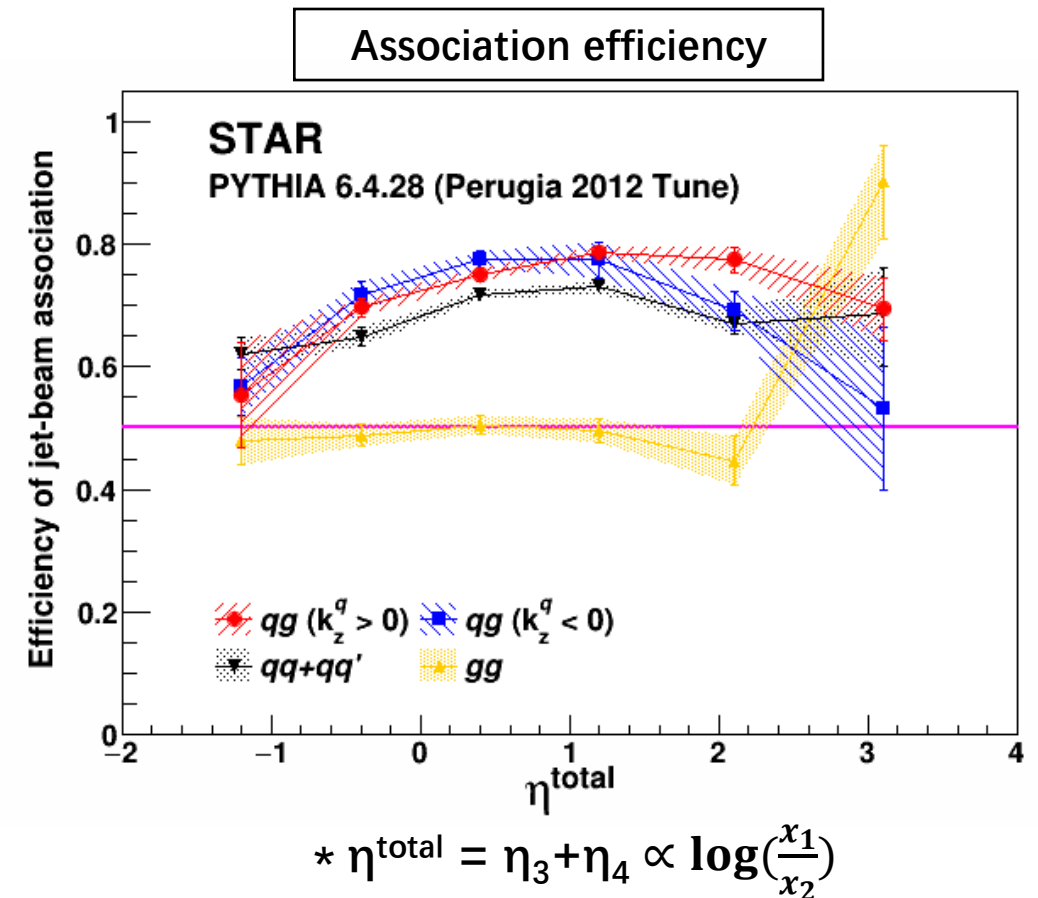
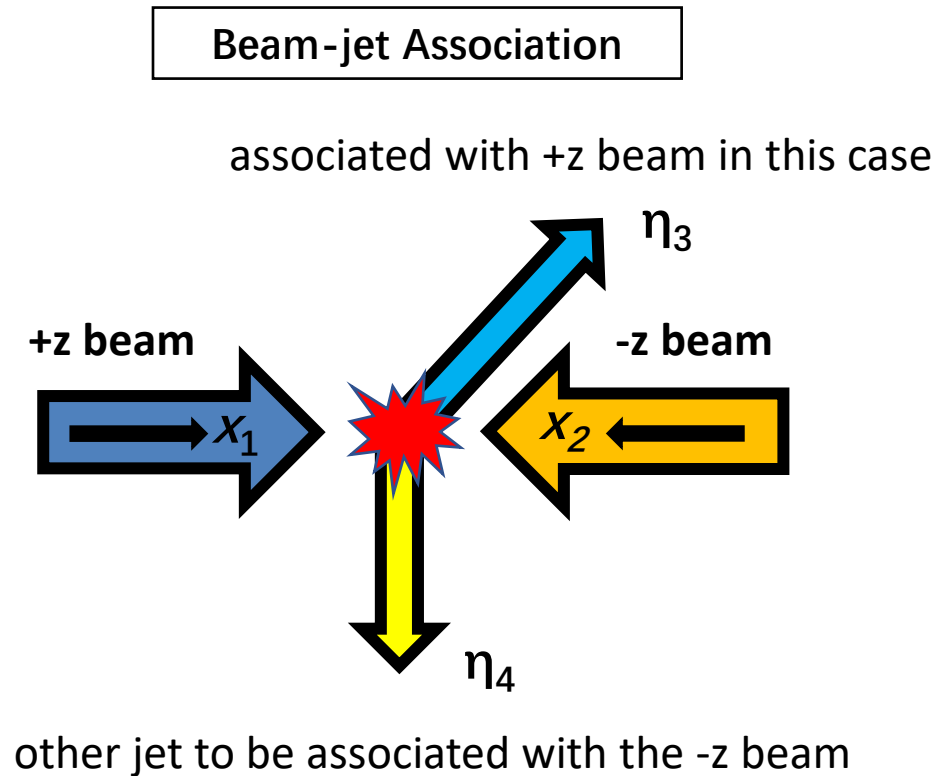
The Sivvers effect leads to a spin-dependent centroid shift of ζ , so we define the asymmetry as:

$$\Delta\zeta = \frac{\langle\zeta\rangle^+ - \langle\zeta\rangle^-}{P}$$

where $\langle\zeta\rangle^{+/-}$ is the centroid of ζ for spin-up and spin-down states, and P is the beam polarization.

Beam-to-Jet Association

- To figure out the “parton flow” from beam to jets, a beam-jet association is performed.
- We assume the **more forward jet (largest $|\eta|$)** is associated with a fragmenting parton from the **beam moving in that direction**.



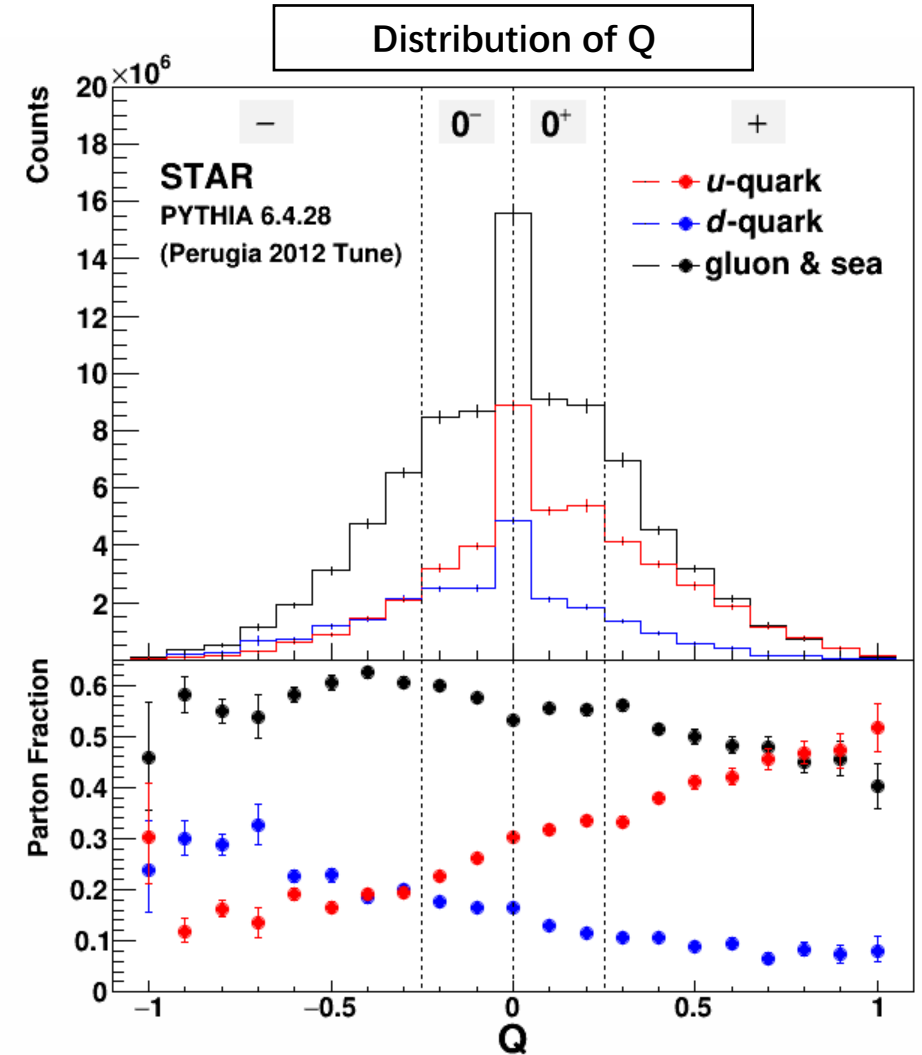
Jet Charge Tagging

We use the **Jet Charge (Q)** of the associated jets to enhance the fraction of u -quarks and d -quarks separately.

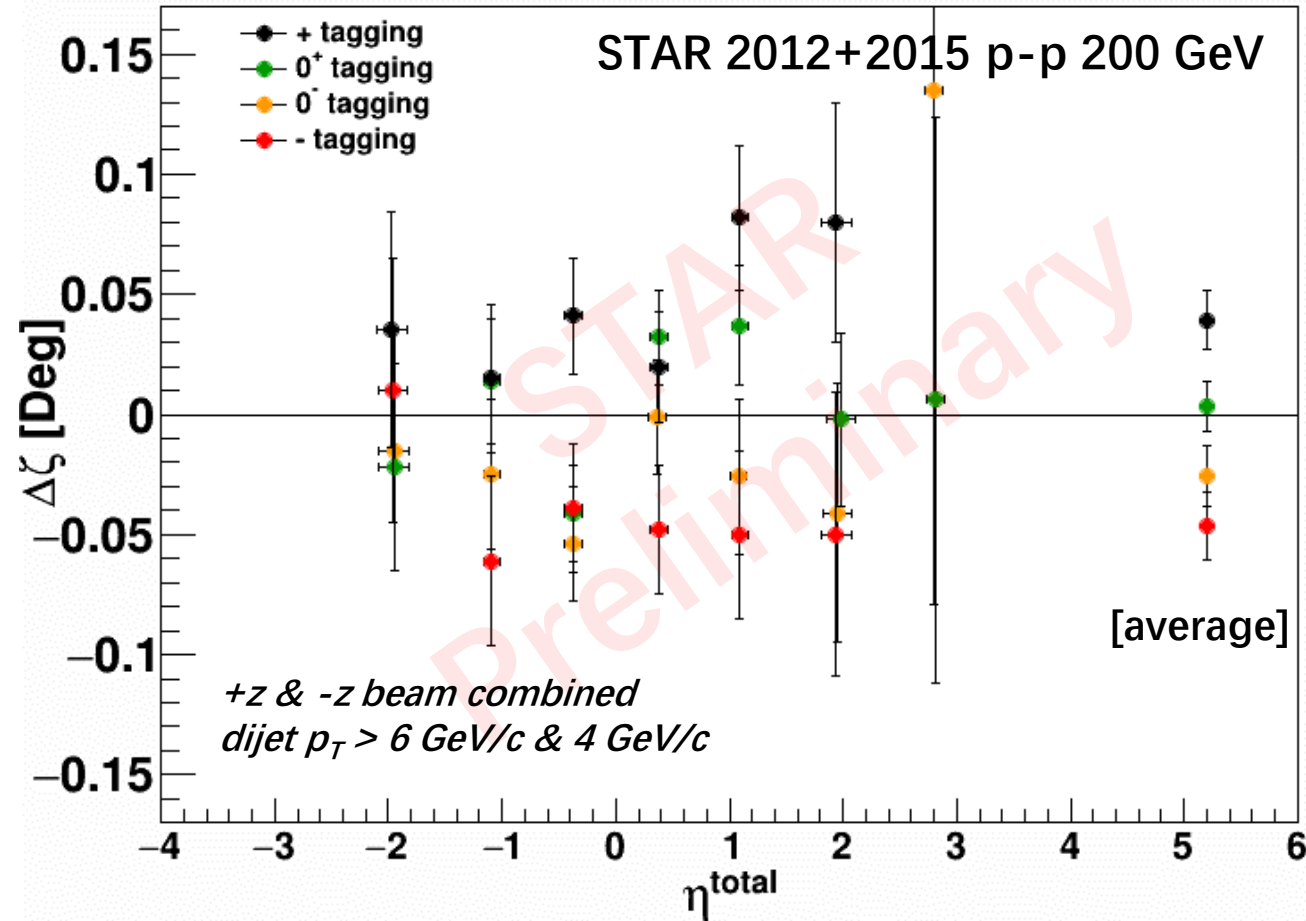
$$Q = \sum_{\substack{\text{all the tracks} \\ \text{with } p_T > 0.8 \text{ GeV}/c}} \frac{\text{track } |p|}{\text{jet } |p|} \cdot \text{track charge}$$

Data is divided into four bins:

1. **Plus tagging** ($Q \geq 0.25$) : enhances u
2. **Zero+ tagging** ($0 \leq Q < 0.25$) : less enhancement to u
3. **Zero- tagging** ($-0.25 < Q < 0$) : less enhancement to d
4. **Minus tagging** ($Q \leq -0.25$) : enhances d



The $\Delta\zeta$ asymmetry

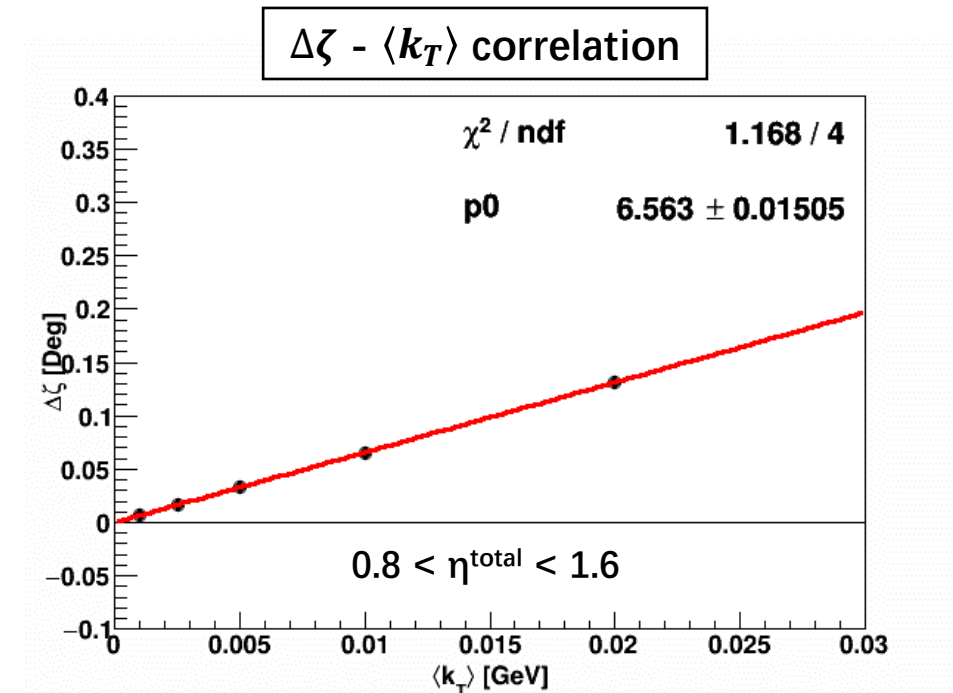
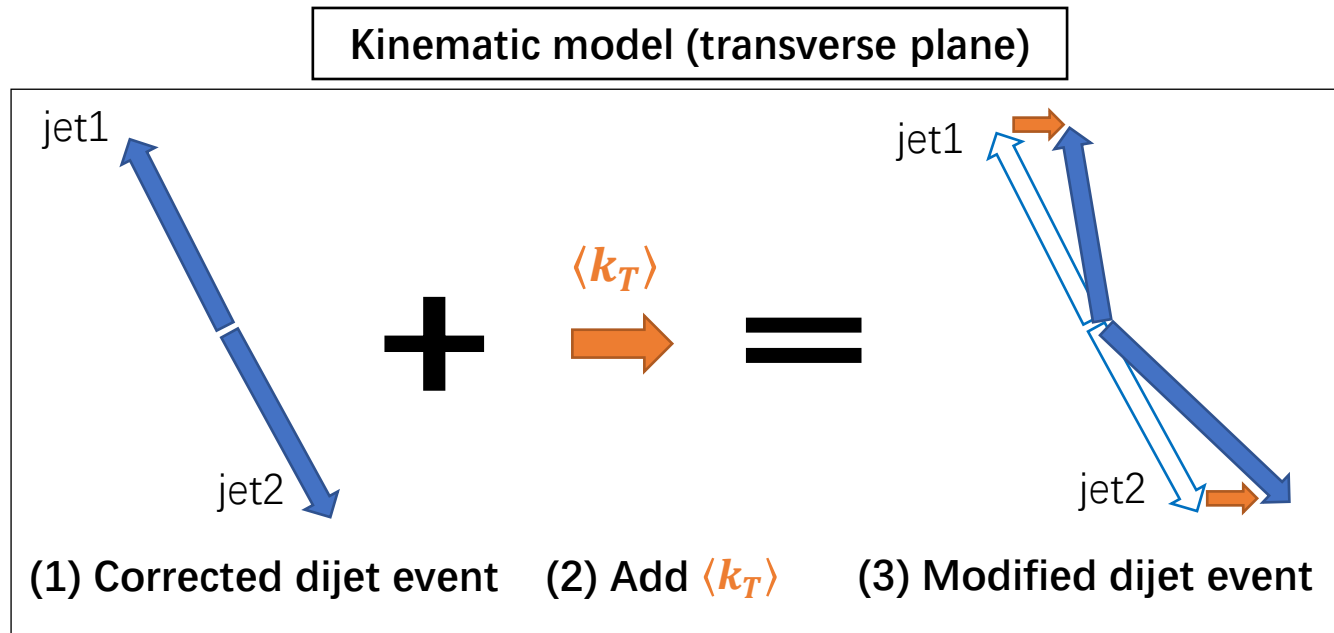


- Large separation ($\sim 5\sigma$) between plus-tagging and minus-tagging.
- First observation of non-zero Sivers asymmetries in dijet production of polarized proton collisions!
- Comparison of asymmetries for the two beams show that the systematic uncertainty is well under control.
- Asymmetry systematically shifting from “+” to “-” values when u quark fraction goes down and d quark fraction goes up.
- η^{total} dependency in plus-tagging, possibly due to :
 - x -dependency in PDFs
 - potential x -dependency in Sivers $\langle k_T \rangle$

Converting the $\Delta\zeta$ asymmetry to $\langle k_T \rangle$

Three steps are taken to convert the $\Delta\zeta$ asymmetry to $\langle k_T \rangle$:

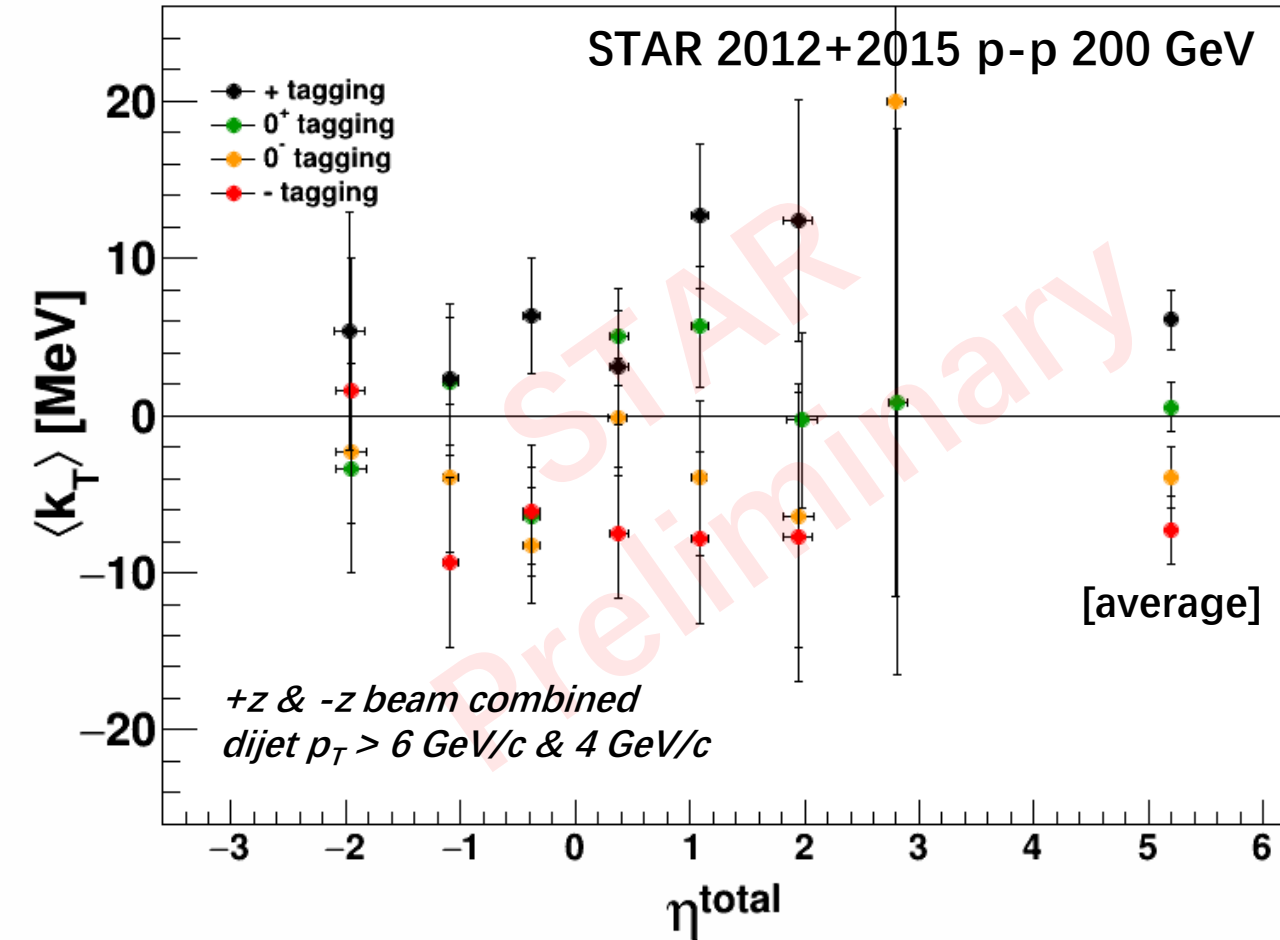
- I. Correct detector jet p_T to parton p_T with machine learning.
- II. Use simple kinematic modeling of $\langle k_T \rangle$, calculate $\Delta\zeta$ with corrected p_T , and get $\Delta\zeta$ - $\langle k_T \rangle$ correlation.



- III. Convert the $\Delta\zeta$ vs. η^{total} results to $\langle k_T \rangle$ vs. η^{total} results :

$$\langle k_T \rangle = \Delta\zeta / \text{slope}$$

The Converted $\langle k_T \rangle$ Results



- Based on the simple kinematic model, we have :

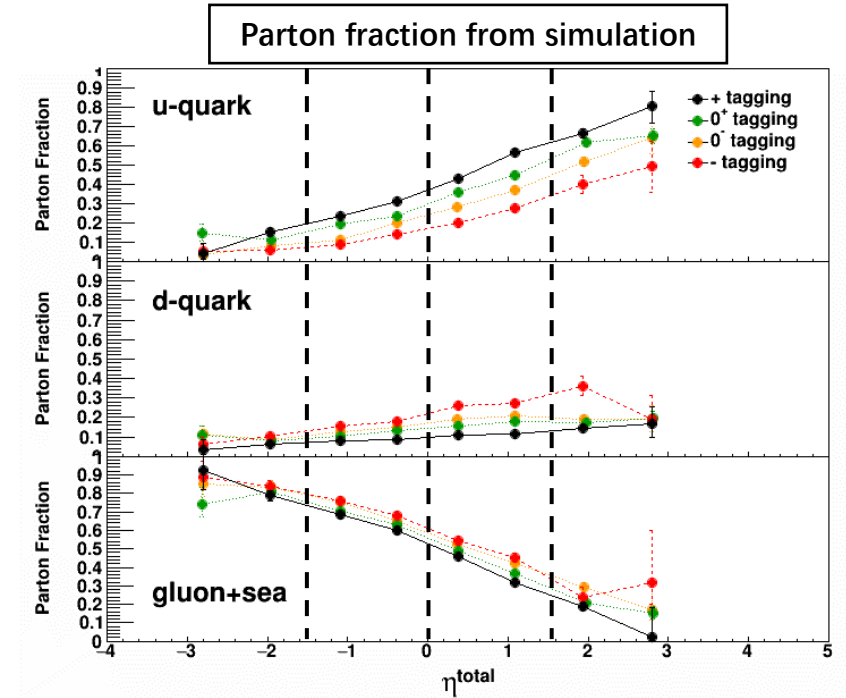
$$\langle k_T^{+tagging} \rangle = +6.1 \pm 1.9 \text{ MeV}/c$$

$$\langle k_T^{-tagging} \rangle = -7.3 \pm 2.2 \text{ MeV}/c$$

- Very small signals successfully accessed with STAR detector!
- In fact, each tagged measurement can be considered as a mixture of 3 different partonic contributions (u , d , g +sea). The 4 tagged measurements provide enough constraints to solve for the $\langle k_T \rangle$ for each parton.

Inverting the Tagged $\langle k_T \rangle$ to Individual Parton $\langle k_T \rangle$

- The parton fraction is estimated in simulation for each tagged measurement.
- Constructing the system of equations (**8X3 matrix**):
 - 4 charge-taggings** : differentiation between u and d quarks
 - Each inversion involves the data from **a pair of adjacent η^{total} bins** : parton fraction is dependent on η^{total}
- The **over-constrained system** is solved through **Moore-Penrose inverse**.



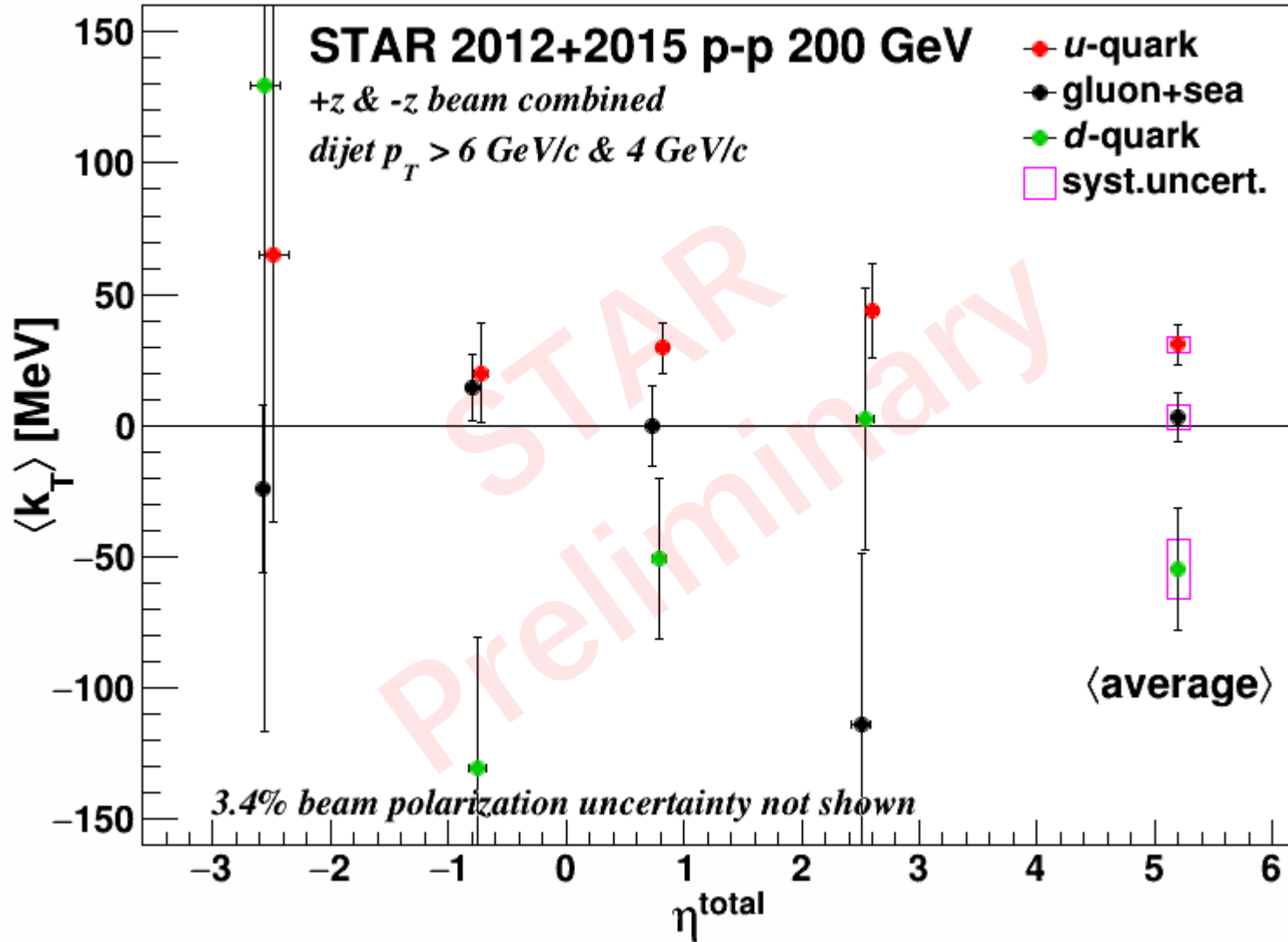
$$\begin{bmatrix} f_1^u & f_1^d & f_1^{gsea} \\ f_2^u & f_2^d & f_2^{gsea} \\ \vdots & \vdots & \vdots \\ f_8^u & f_8^d & f_8^{gsea} \end{bmatrix} \begin{bmatrix} u \\ d \\ gsea \end{bmatrix} = \begin{bmatrix} \Delta_1 \\ \Delta_2 \\ \vdots \\ \Delta_8 \end{bmatrix} \xrightarrow{\text{Moore-Penrose inverse}} \begin{bmatrix} c_1^u & c_2^u & \dots & c_8^u \\ c_1^d & c_2^d & \dots & c_8^d \\ c_1^{gsea} & c_2^{gsea} & \dots & c_8^{gsea} \end{bmatrix} \begin{bmatrix} \Delta_1 \\ \Delta_2 \\ \vdots \\ \Delta_8 \end{bmatrix} = \begin{bmatrix} u \\ d \\ gsea \end{bmatrix}$$

8 x 3 matrix

3 x 8 matrix

f = parton fraction
u, d, gsea = parton $\langle k_T \rangle$
 Δ = tagged $\langle k_T \rangle$

The Unfolded Parton $\langle k_T \rangle$



- $\langle k_T^u \rangle > 0$, $\langle k_T^d \rangle < 0$, $\langle k_T^{g+sea} \rangle \sim 0$
- $\left| \frac{\langle k_T^d \rangle}{\langle k_T^u \rangle} \right| \sim 2$
- The systematic uncertainty is dominated by the uncertainty of the estimated parton fraction.
- No clear η^{total} -dependency for given statistics, suggesting a relatively weak x -dependency.

Conclusions

- The Sivers effect has been studied in dijets measured with the STAR detector, using data taken in 2012 and 2015.
- First observation of non-zero Sivers asymmetries in polarized proton collisions!
- A conversion of the asymmetry to the $\langle k_T \rangle$ results is provided based on purely kinematic model. The results are further unfolded for the $\langle k_T \rangle$ of individual partons.
- The features of the unfolded parton $\langle k_T \rangle$ are consistent with expectation:
 - $\langle k_T^u \rangle$ and $\langle k_T^d \rangle$ have different signs
 - $|\langle k_T^d \rangle / \langle k_T^u \rangle| \sim 2$
- Results provide constraints for the Sivers function at a high Q^2 scale ($Q^2 > 160 \text{ GeV}^2$).
- Several theoretical efforts are underway to make comparisons to these data.

Thank you !

BACKUP

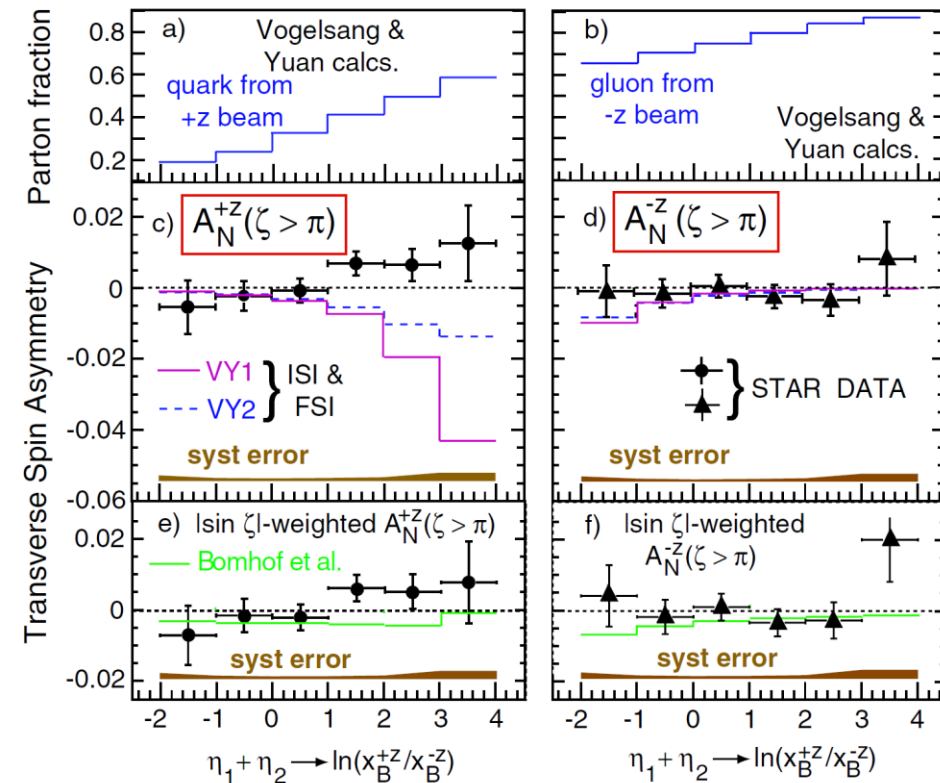
STAR 2006 Analysis Result & New 2012+2015 Improvements

In the 2006 analysis, the result was found to be consistent with zero within dominant statistical uncertainties.

With data taken in 2012 and 2015, the current analysis sees:

- 33 times larger data set
- Fully reconstructed jets (no tracking for 2006 data)
- Employ a tagging method to enhance u -quark and d -quark signals

STAR Collab. PhysRevLett 99 142003



Asymmetry is plotted as a function of the sum of dijet pseudo-rapidities ($\eta_1 + \eta_2 \propto \ln(\frac{x_1}{x_2})$) since Sivers effect is expected to be dependent on parton x .

Parton x

- $Q^2 > 160 \text{ GeV}^2$
- Parton x increases along with η^{total} , a possible x -dependence of $\langle k_T \rangle$ should manifest in the inverted results if strong enough.

