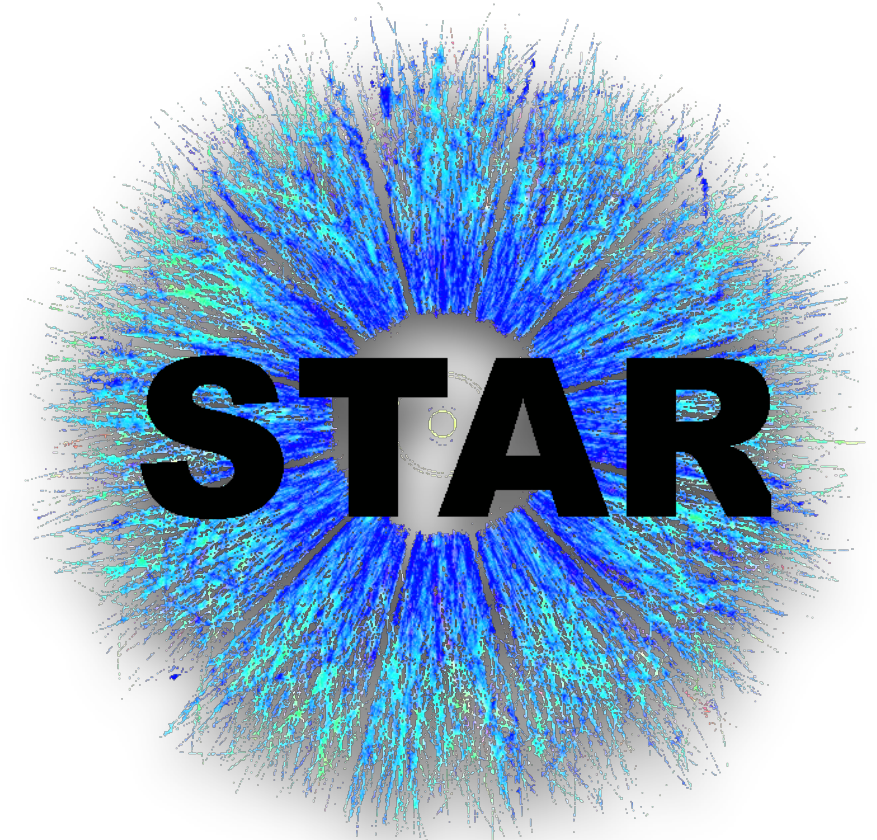


# Deuteron Number Fluctuations and Proton-deuteron Correlations in High Energy Heavy-ion Collisions in STAR Experiment at RHIC



Debasish Mallick,  
On Behalf of the STAR Collaboration,  
National Institute of Science Education and Research, HBNI,  
Jatni, India



## Outline:

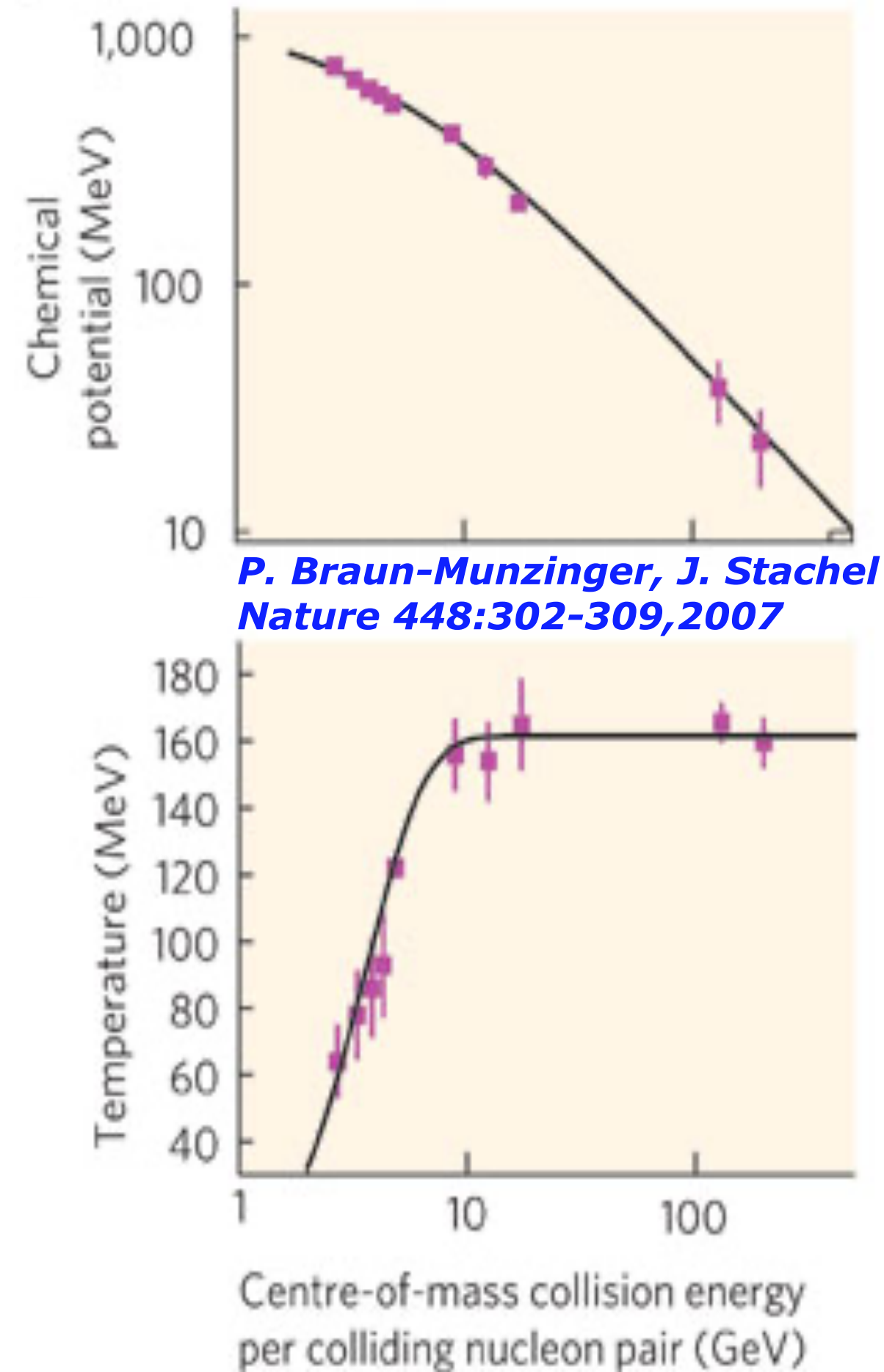
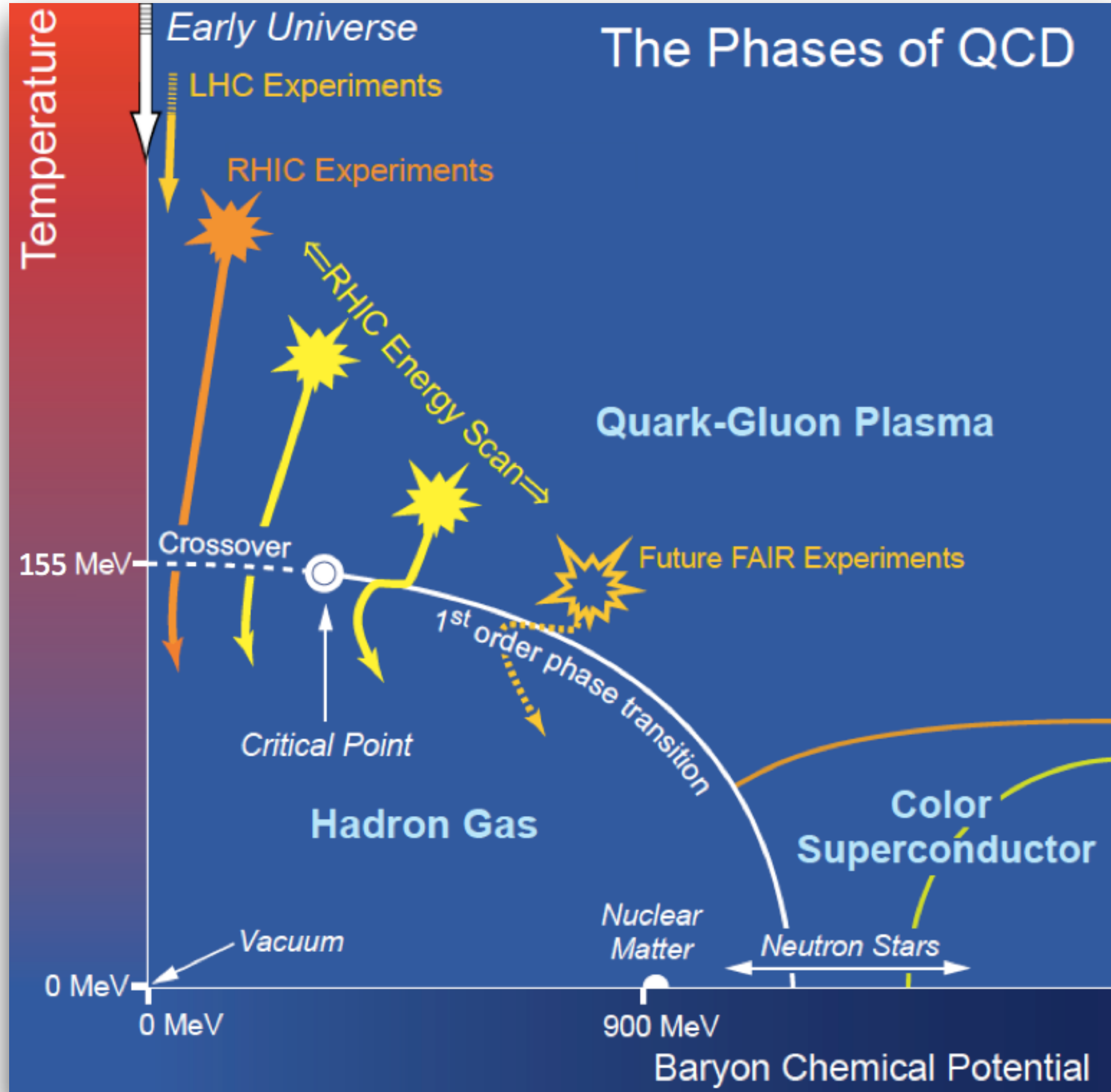
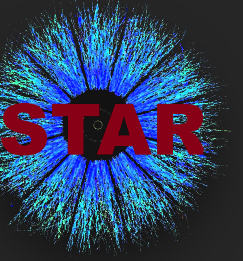
- Introduction
- Motivation and Observables
- STAR and Analysis Method
- Results
- Summary



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Science



# Introduction

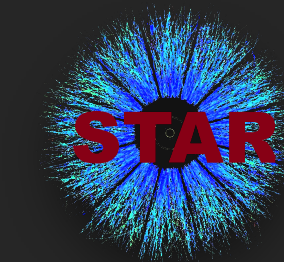


<https://drupal.star.bnl.gov/STAR/starnotes/public/sn0493>  
[https://drupal.star.bnl.gov/STAR/files/BES\\_WPII\\_ver6.9\\_Cover.pdf](https://drupal.star.bnl.gov/STAR/files/BES_WPII_ver6.9_Cover.pdf)

Goal: Study the phase diagram of QCD.  
 Beam Energy Scan (BES): Varying beam energy varies temperature ( $T$ ) and baryon chemical potential ( $\mu_B$ ).  
 Fluctuations in conserved quantities are sensitive to phase structure and critical point.



# Light Nuclei Synthesis

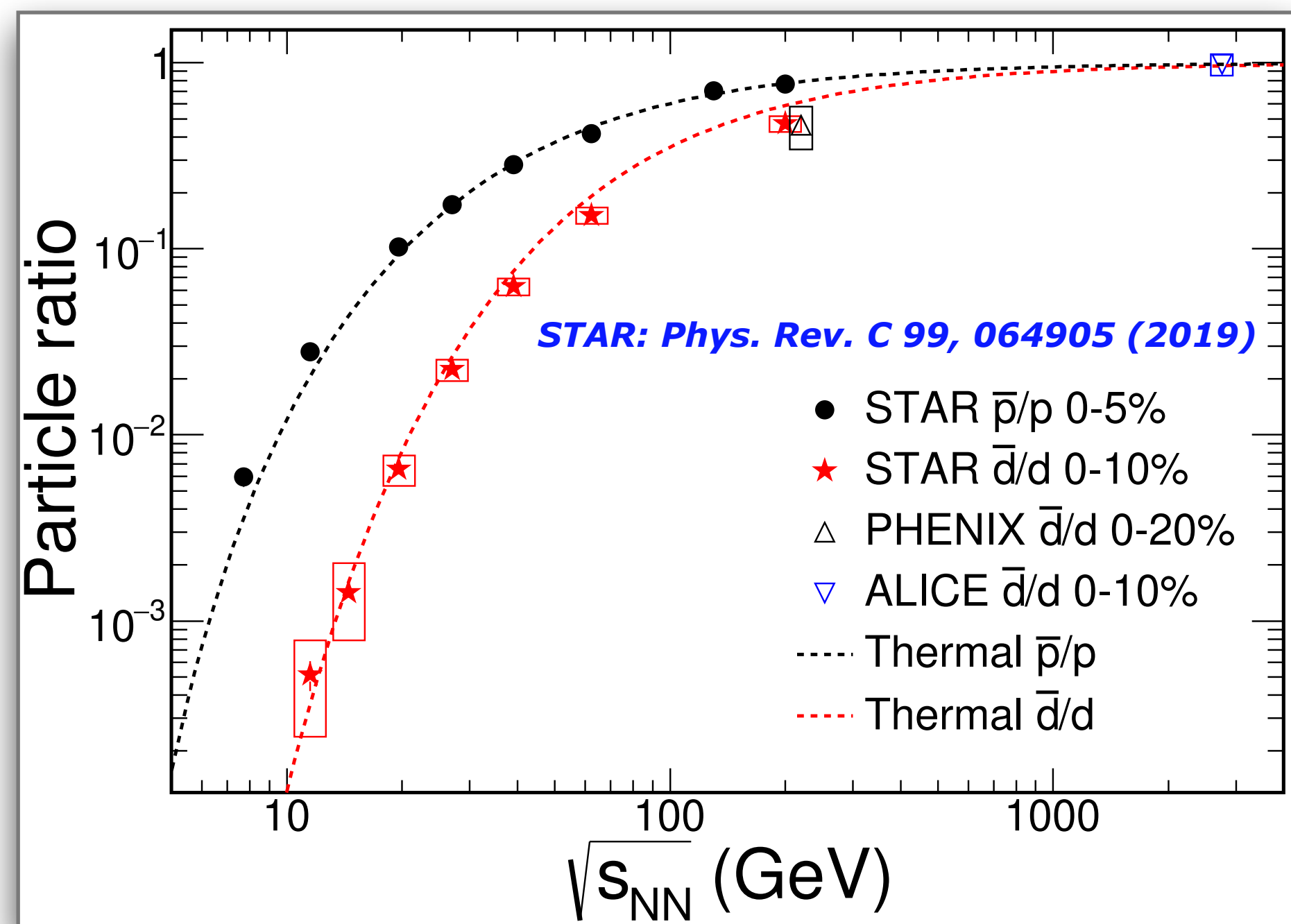


## GCE Thermal Model

**Yield of deuteron:**  $N_d = \frac{g_d V}{\pi^2} m_d^2 T K_2(m_d/T) \exp(\mu_d/T)$

where,  $g_d$ : degeneracy,  $\mu_d$ : chemical potential.

- ☑ Deuteron is treated as a free and point particle.
- ☑ Degeneracy, mass and baryon number are inputs.



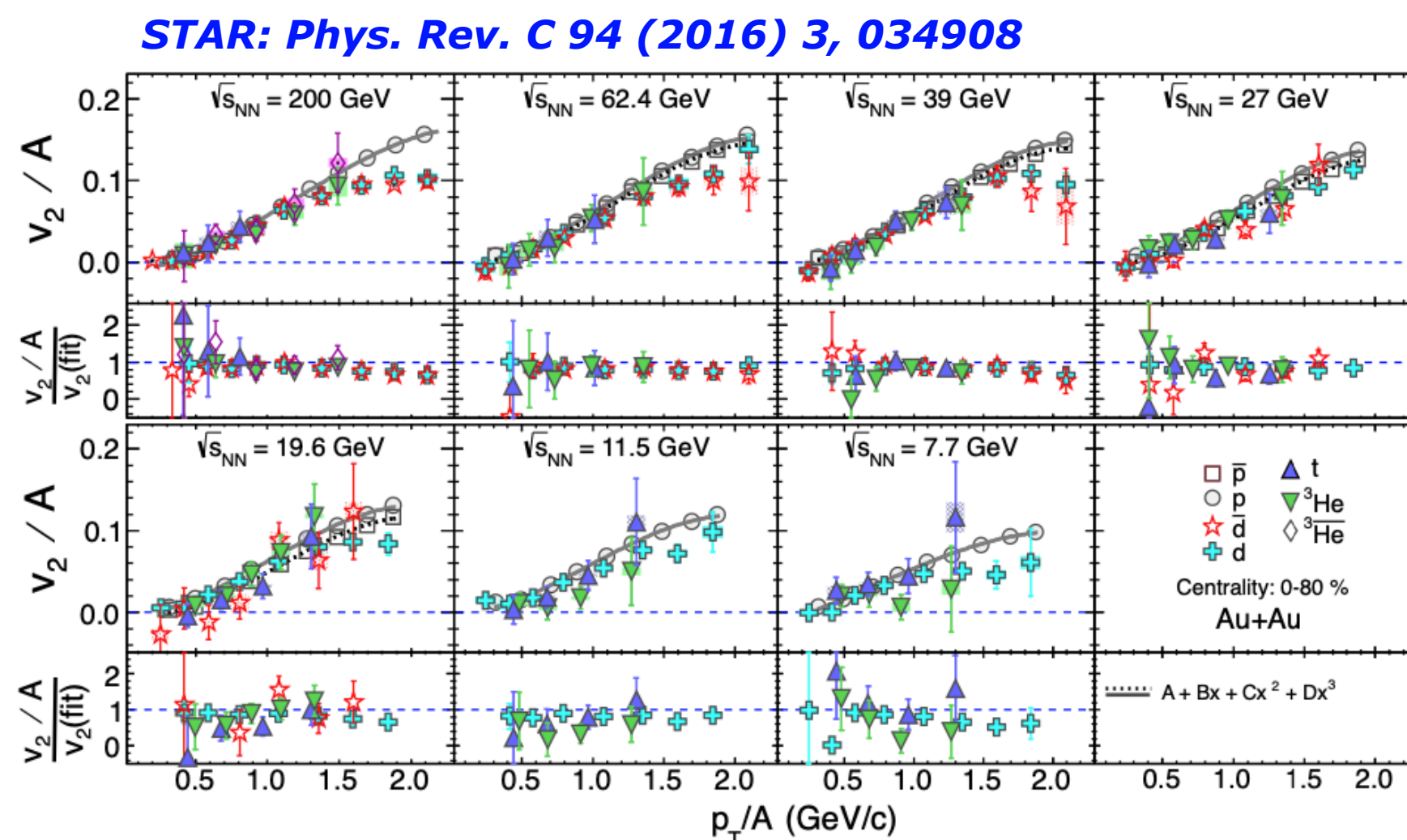
- ☑ Anti-particle to particle ratio well explained by thermal model for a wide range of  $\sqrt{s_{NN}}$ .

## Coalescence Model

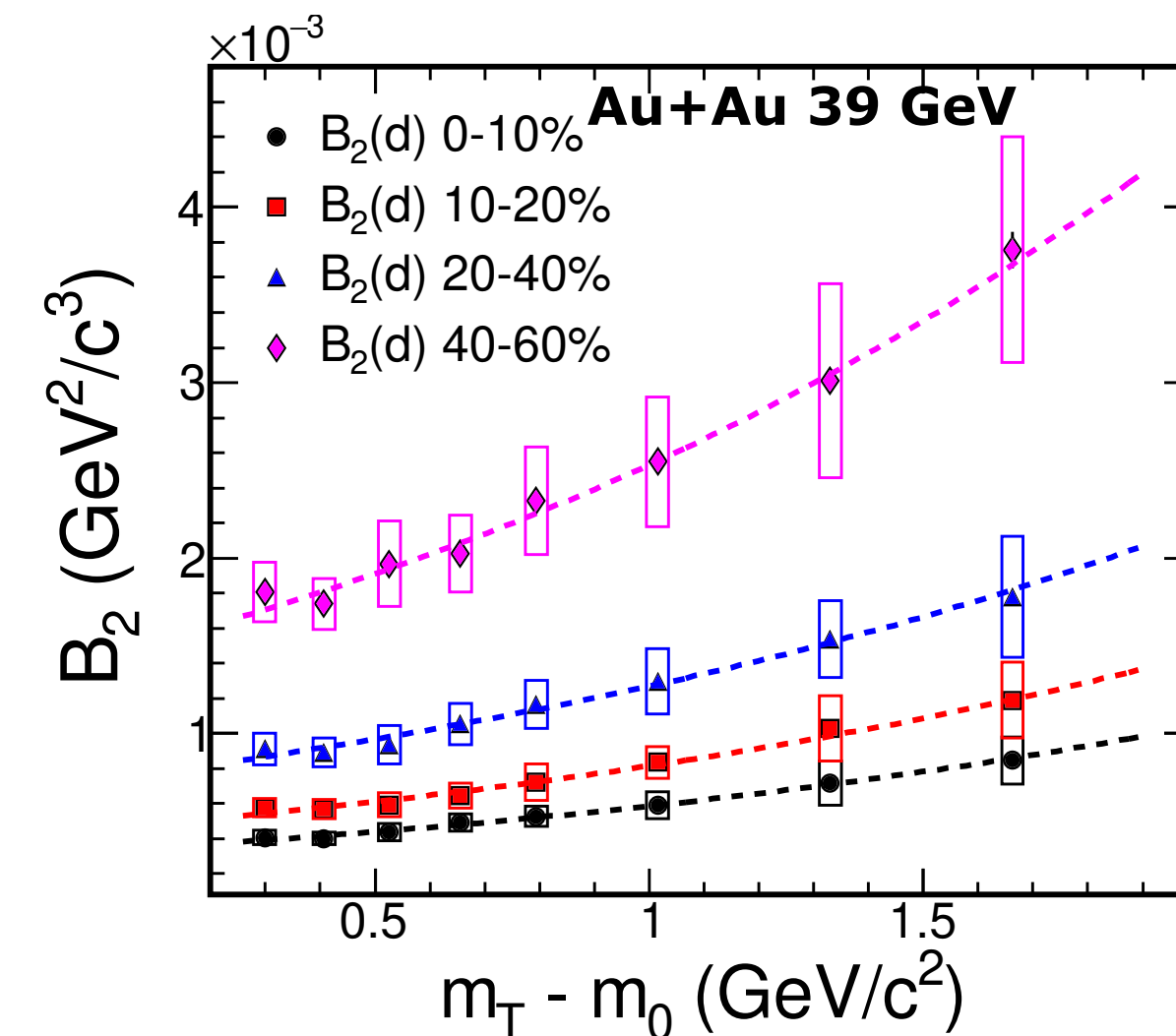
**Invariant Yield:**  $E_d \frac{d^3 N_d}{dp_d^3} = B_2 \left( E_p \frac{d^3 N_p}{dp_p^3} \right) \left( E_n \frac{d^3 N_n}{dp_n^3} \right)$

**Elliptic Flow:**  $v_2^d(p_T) \approx 2v_2^p \left( \frac{p_T}{2} \right)$

- ☑ Light nuclei created using protons and neutrons.
- ☑  $B_2$  extracted as a function of centrality,  $m_T$ , and  $\sqrt{s_{NN}}$ .



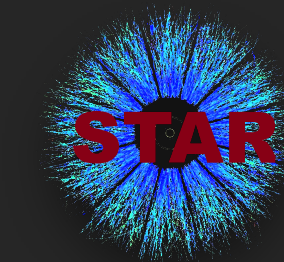
- ☑ Nucleon coalescence picture works up to  $p_T/A \leq 1.5$  GeV/c.



**STAR: Phys. Rev. C 99, 064905 (2019)**

- ☑  $B_2 \propto e^{(m_T - m)}$
- ☑  $B_2 \propto (4/3)\pi p_0^3$   
 $p_0$  is the radius in momentum space.

# Light Nuclei Synthesis



## GCE Thermal Model

**Yield of deuteron:**  $N_d = \frac{g_d V}{\pi^2} m_d^2 T K_2(m_d/T) \exp(\mu_d/T)$

where,  $g_d$ : degeneracy,  $\mu_d$ : chemical potential.

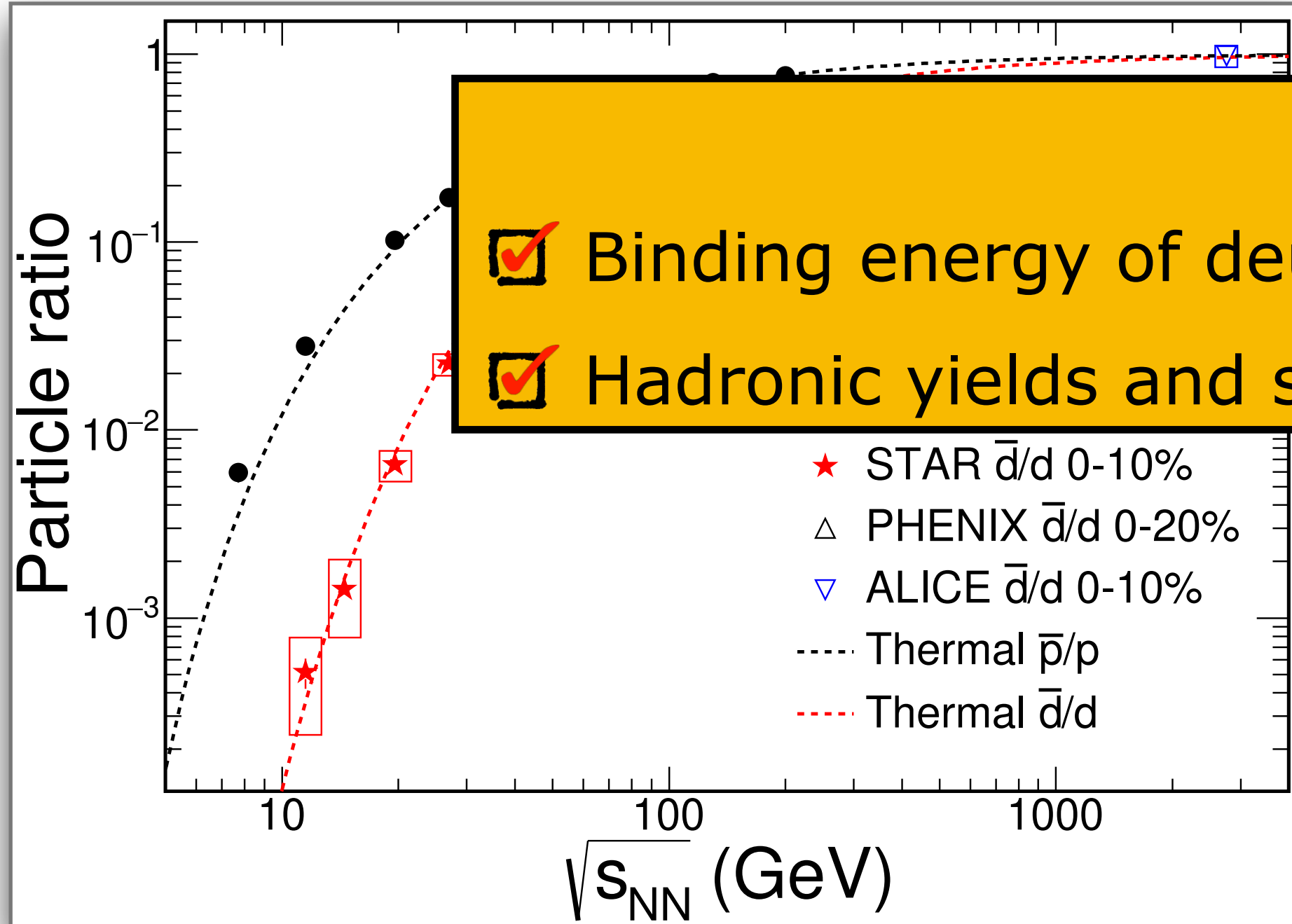
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- Degeneracy, mass and baryon number are inputs.

## Coalescence Model

**Invariant Yield:**  $E_d \frac{d^3 N_d}{dp_d^3} = B_2 \left( E_p \frac{d^3 N_p}{dp_p^3} \right) \left( E_n \frac{d^3 N_n}{dp_n^3} \right)$

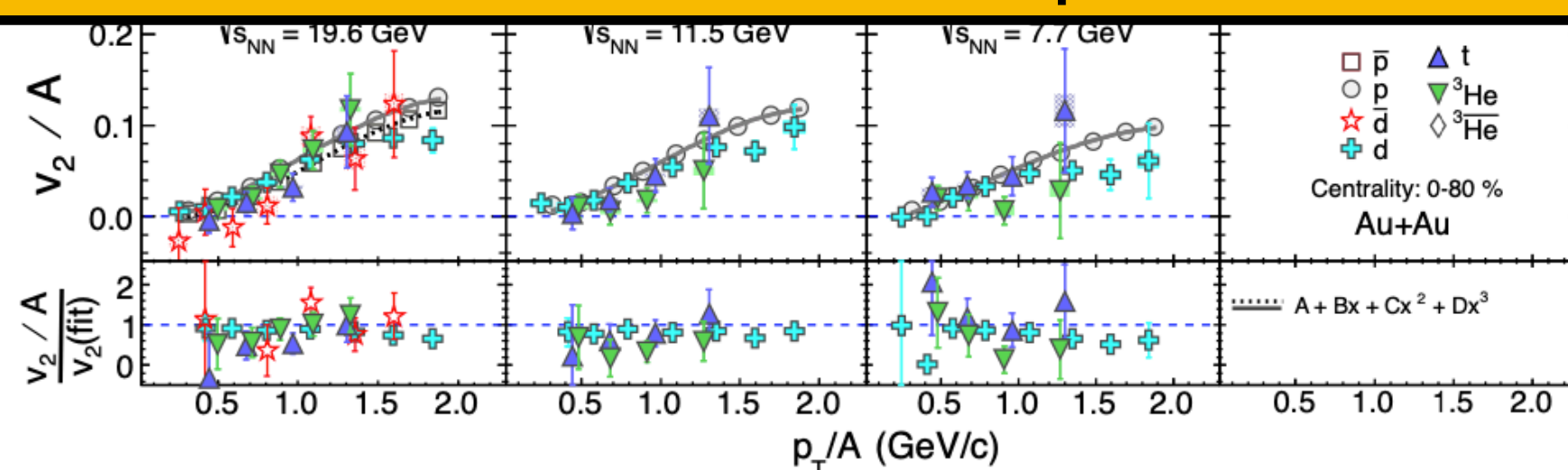
**Elliptic Flow:**  $v_2^d(p_T) \approx 2v_2^p \left( \frac{p_T}{2} \right)$

- Light nuclei created using protons and neutrons.
- $B_2$  extracted as a function of centrality,  $m_T$ , and  $\sqrt{s_{NN}}$ .

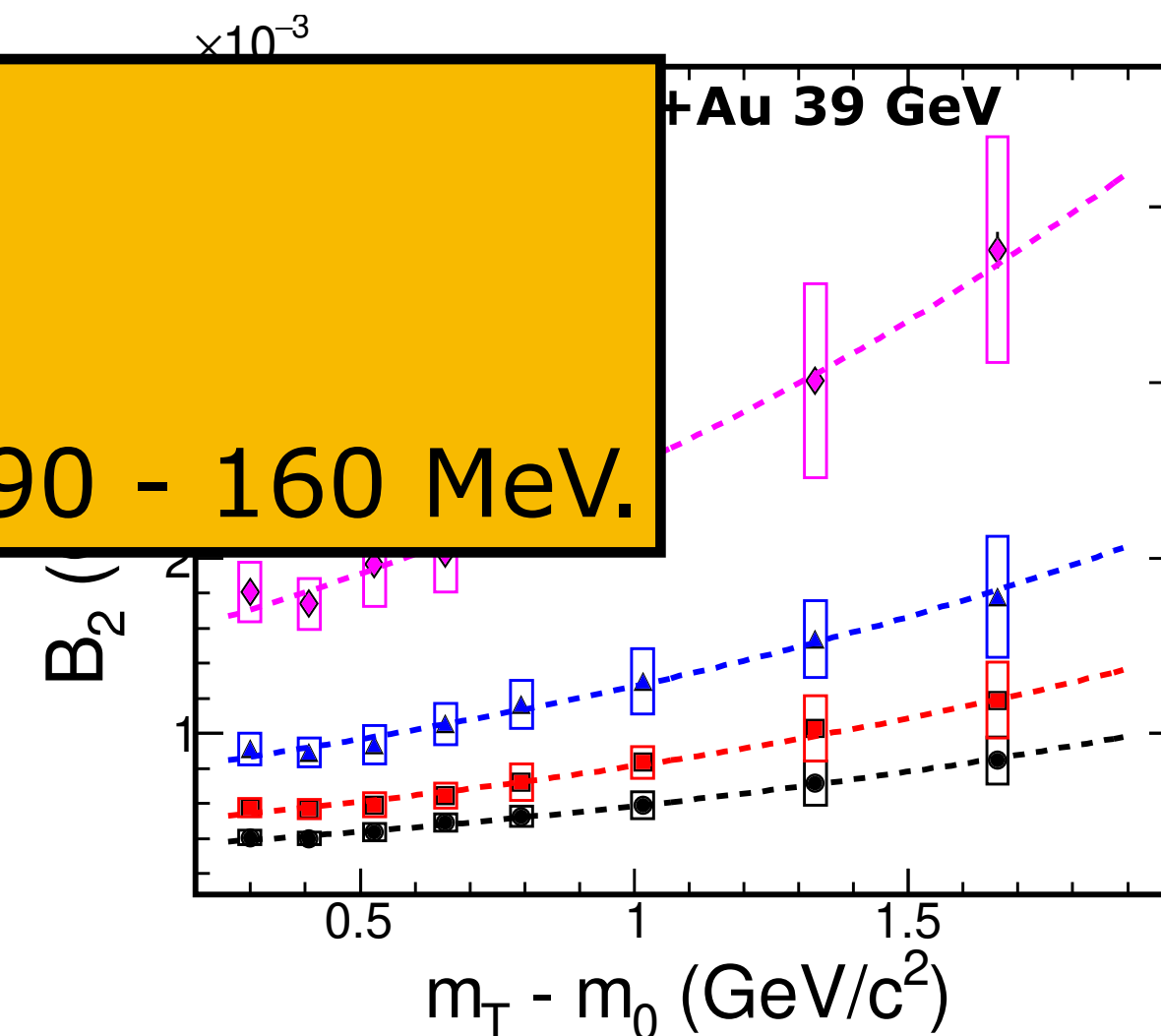


- Binding energy of deuteron  $\sim 2.2$  MeV.
- Hadronic yields and spectra are fixed around temperature  $\sim 90 - 160$  MeV.

## Typical Scales



- Nucleon coalescence picture works up to  $p_T/A \leq 1.5$  GeV/c.



STAR: Phys. Rev. C 99, 064905 (2019)

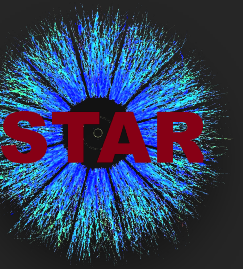
- $B_2 \propto e^{(m_T - m)}$
- $B_2 \propto (4/3)\pi p_0^3$

$p_0$  is the radius in momentum space.

- Anti-particle to particle ratio well explained by the thermal model for a range of  $\sqrt{s_{NN}}$ .



# Observables

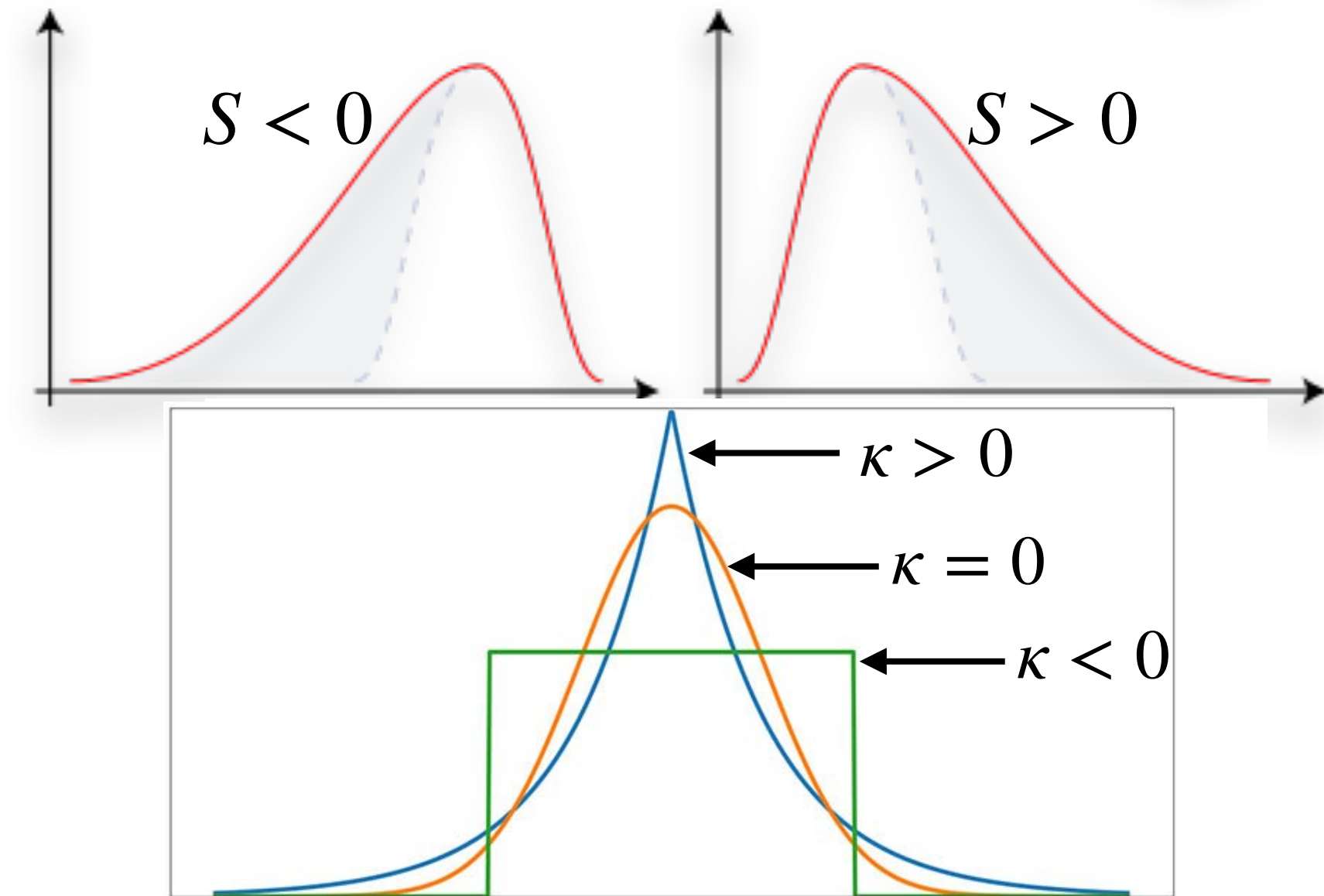


- Higher-order cumulants characterise the subtle features of a distribution.

$$\begin{aligned}
 C_1 &= \langle N \rangle \\
 C_2 &= \langle (\delta N)^2 \rangle \\
 C_3 &= \langle (\delta N)^3 \rangle \\
 C_4 &= \langle (\delta N)^4 \rangle - 3 \langle (\delta N)^2 \rangle^2
 \end{aligned}$$

$$\frac{C_2}{C_1} = \frac{\sigma^2}{M} \quad \frac{C_3}{C_2} = S\sigma \quad \frac{C_4}{C_2} = \kappa\sigma^2$$

$M$  = Mean  
 $\sigma^2$  = Variance  
 $S$  = Skewness  
 $\kappa$  = Kurtosis



- Higher order cumulants of conserved number distributions are, in general, sensitive observables.
  - Related to the correlation length and susceptibilities.
  - Deuteron cumulants add more information on baryon number fluctuation.

$$C_2 \sim \xi^2 \quad C_4 \sim \xi^7 \quad \text{*Quantitative numbers - Model dependent}$$

$$\frac{\chi_q^{(4)}}{\chi_q^{(2)}} = \kappa\sigma^2 = \frac{C_{4,q}}{C_{2,q}} \quad \frac{\chi_q^{(3)}}{\chi_q^{(2)}} = S\sigma = \frac{C_{3,q}}{C_{2,q}}$$

[S. Ejiri, F. Karsch, K. Redlich, Phys. Lett. B633 \(2006\) 275-282](#)  
[M. A. Stephanov, Phys. Rev. Lett. 102, 032301 \(2009\)](#)  
[R.V. Gavai, S. Gupta, Phys. Lett. B696:459-463,2011](#)  
[A. Bazavov et. al, Phys. Rev. Lett. 109, 192302 \(2012\)](#)  
[A. Bzdak et. al, Physics Reports 853 \(2020\) pp. 1-87](#)  
[S. Borsanyi et. al, Phys. Rev. Lett. 111, 062005 \(2013\)](#)

Pearson correlation coefficient

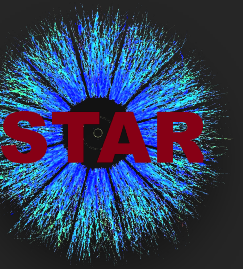
$$\rho(N_x, N_y) = \frac{\langle (\delta N_x \delta N_y) \rangle}{\sigma_x \sigma_y}$$

$\rho$  measures linear correlation between two variables.

$\rho > 0$ : Positive correlation

$\rho < 0$ : Anti-correlation





## Coalescence Toy Model

Z. Fecková, J. Steinheimer, B. Tomášik and M. Bleicher: *Phys. Rev. C* 93, 054906 (2016)

Probability of deuteron formation,  $\lambda_d = B_2 n_p n_n$

Assume, proton ( $n_p$ ) and neutron ( $n_n$ ) follow Poisson distributions,

- At low  $\sqrt{s_{NN}}$ ,  $B_2$  increases. *STAR: Phys. Rev. C* 99, 064905 (2019)
- Larger value of  $n_p$  and  $n_n$  at low  $\sqrt{s_{NN}}$ .
- Results in rise of scaled moments of deuteron number.

**Scaled Moments:**  $\sigma^2/M = C_2/C_1$ ,  $S\sigma = C_3/C_2$ ,  $\kappa\sigma^2 = C_4/C_2$

### Two assumptions in the model:

Model A: Correlated p and n ( $n_p = n_n$ ).      Model B: Independent p and n.

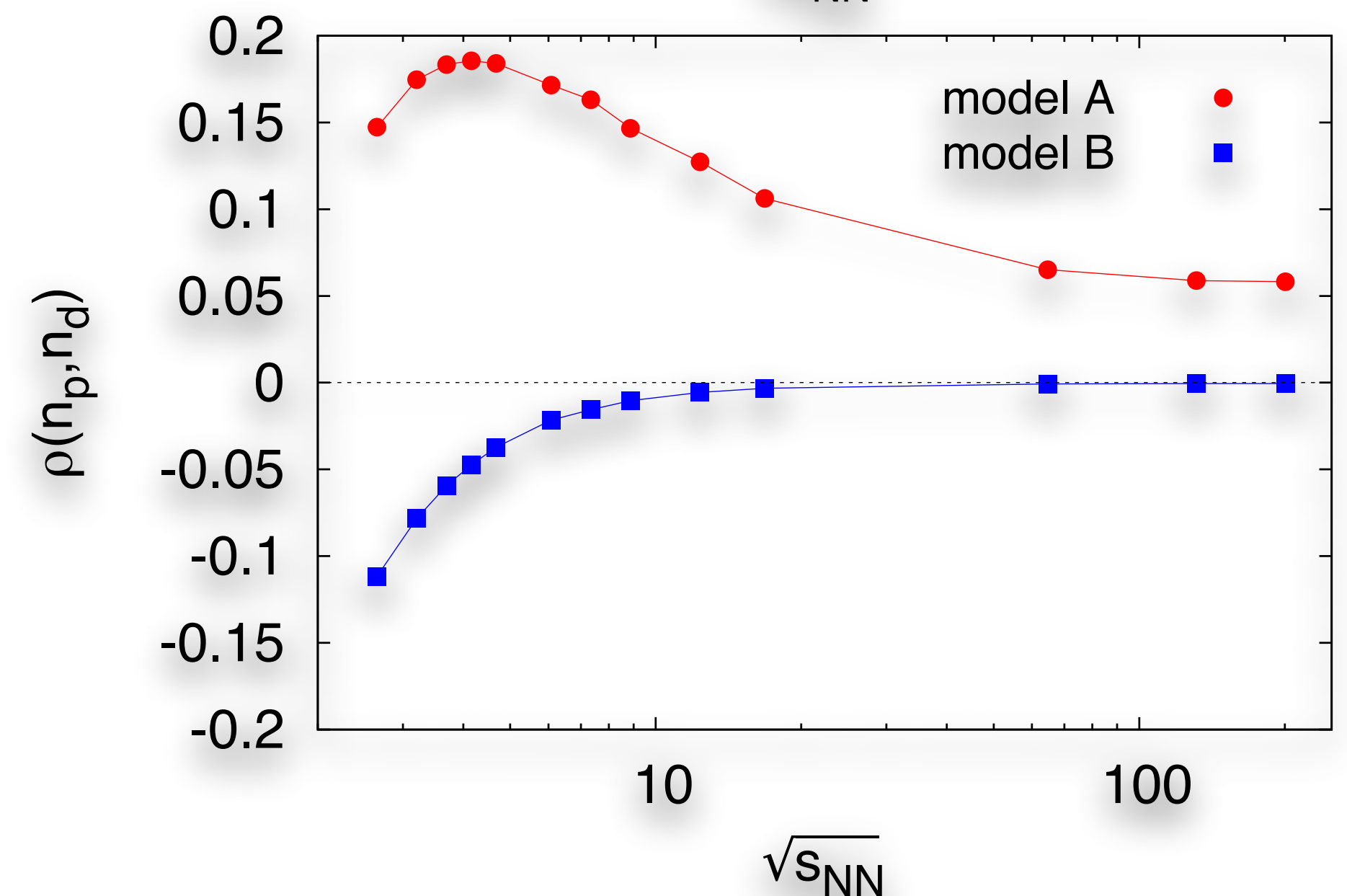
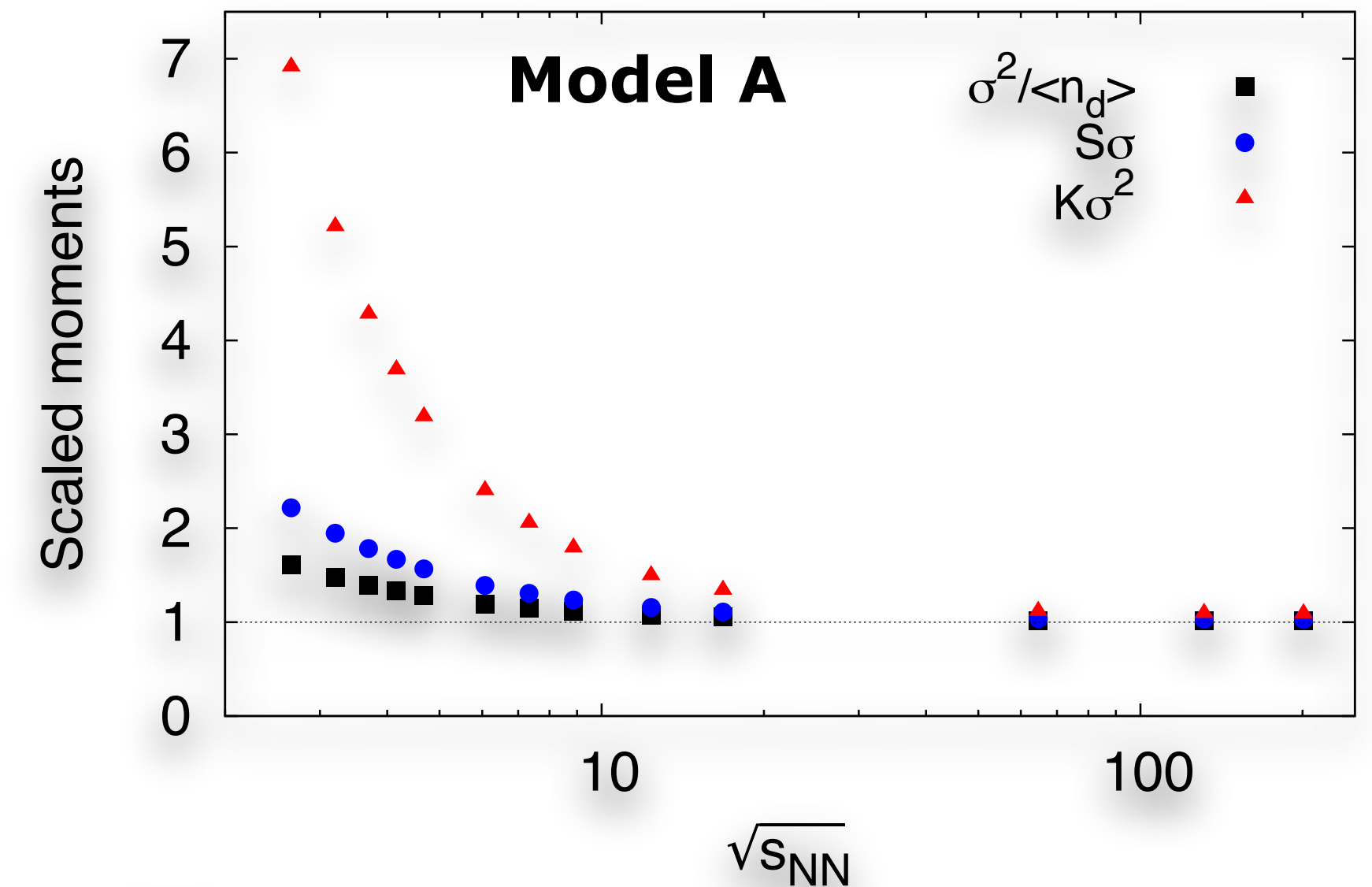
$$\lambda_d = B_2 n_p^2$$

$$\lambda_d = B_2 n_p n_n$$

$$\rho(n_p, n_d) = \frac{\langle (n_p - \langle n_p \rangle)(n_d - \langle n_d \rangle) \rangle}{\sigma_p \sigma_d}$$

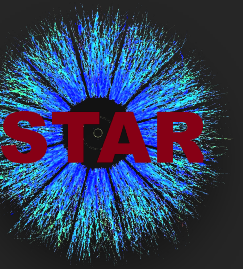
Model A:  $\rho > 0$

Model B:  $\rho < 0$

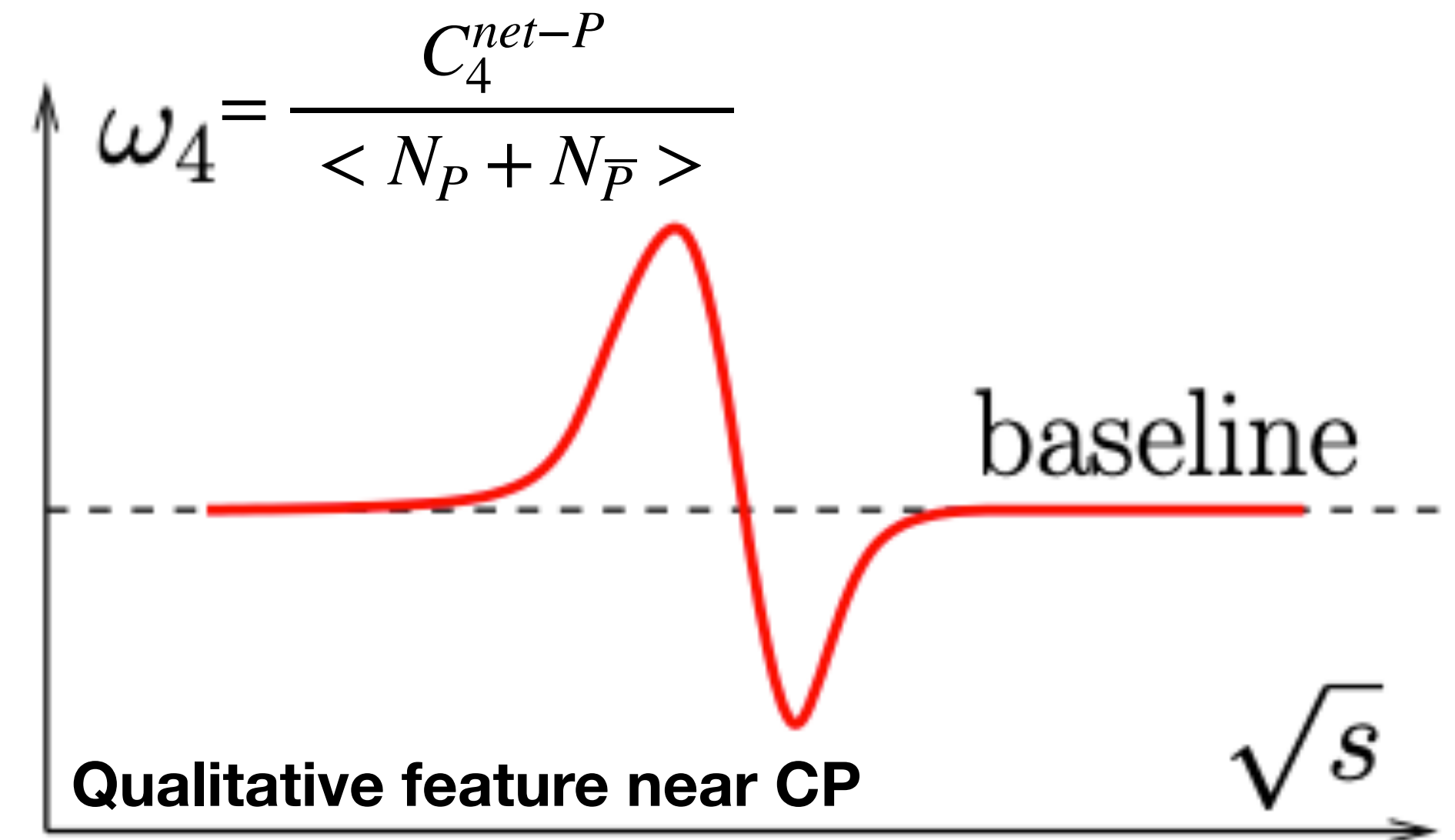
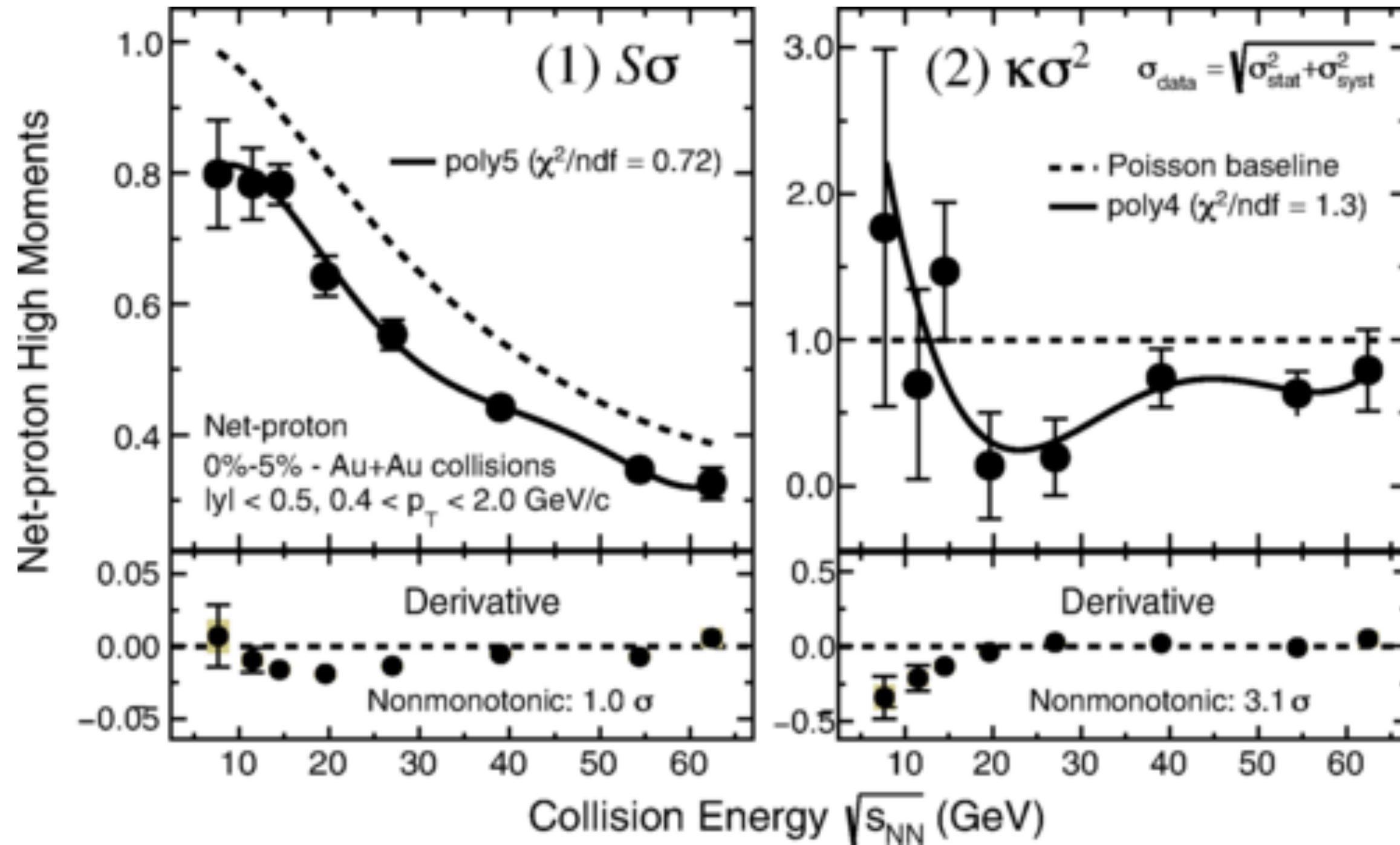




# Baryon Number Fluctuation



STAR: Phys. Rev. Lett. 126 (2021) 092301



M. A. Stephanov Phys. Rev. Lett. 107, 052301 (2011)

M. A. Stephanov 2011 J. Phys. G: Nucl. Part. Phys. 38 124147

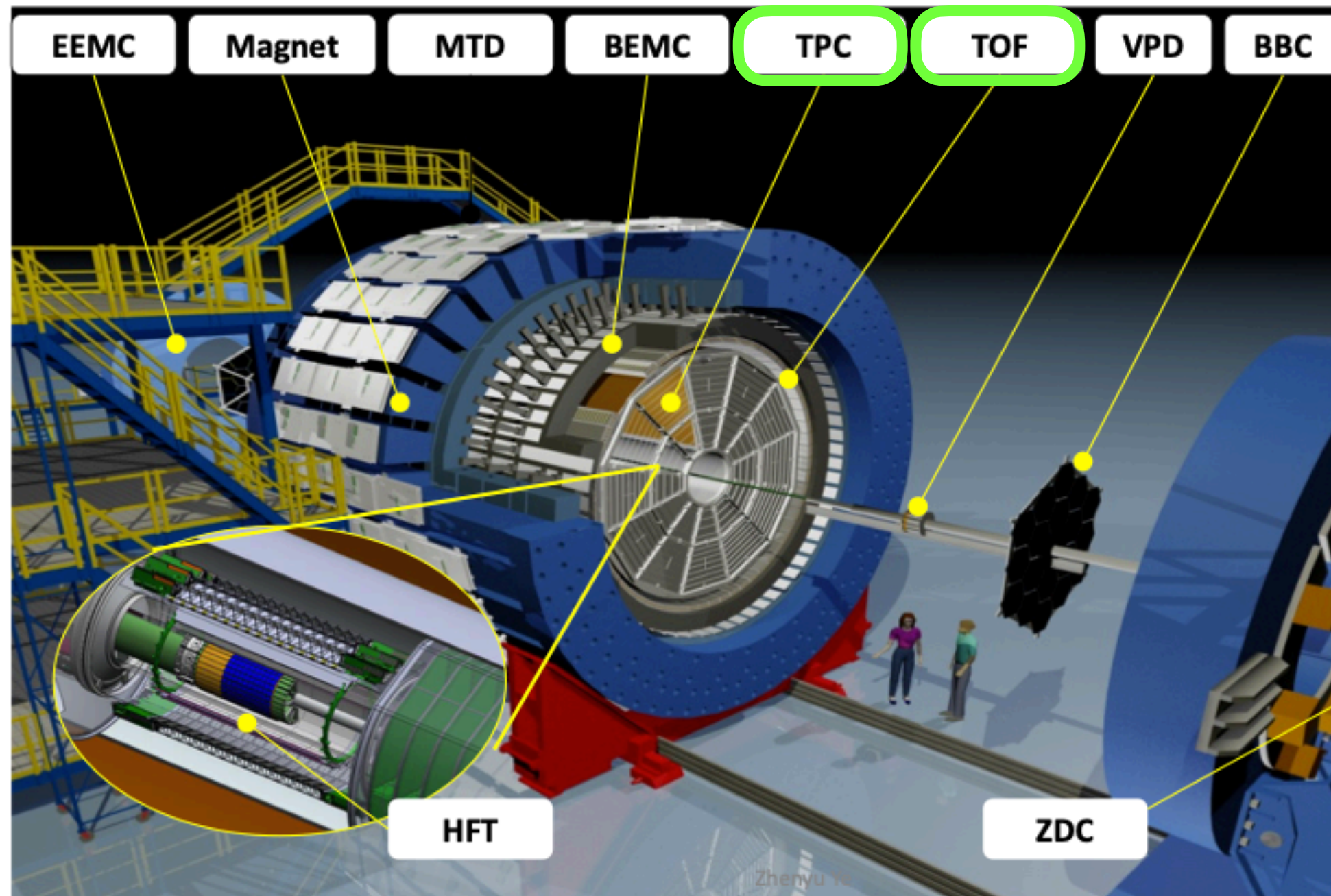
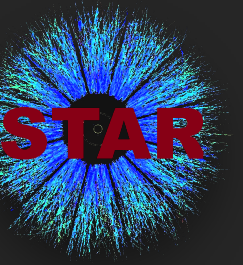
- ✓ Cumulants of deuteron number distribution and proton-deuteron correlation are sensitive to production mechanism.
  - ✓ Until now studies have been done only with baryons of  $|B|=1$ .
  - ✓ QCD critical point leads to large density fluctuation within certain correlation length.
- Deuteron production might be affected by local density fluctuations.

Ed. Shuryak et. al, Phys. Rev. C 101 (2020) 3, 034914

K.J. Sun et. al, Phys. Lett. B 774 (2017) 103-107



# STAR Detector

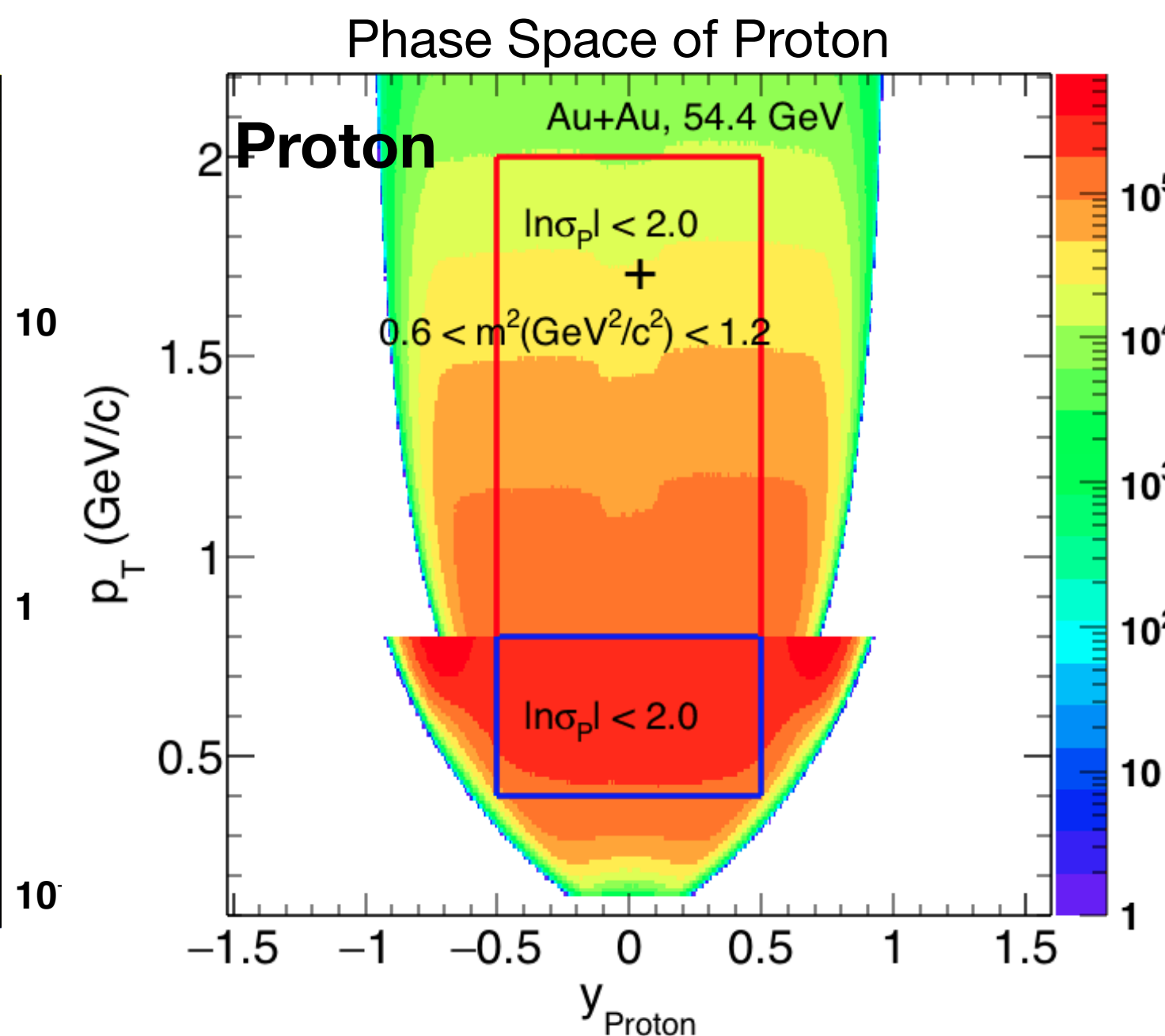
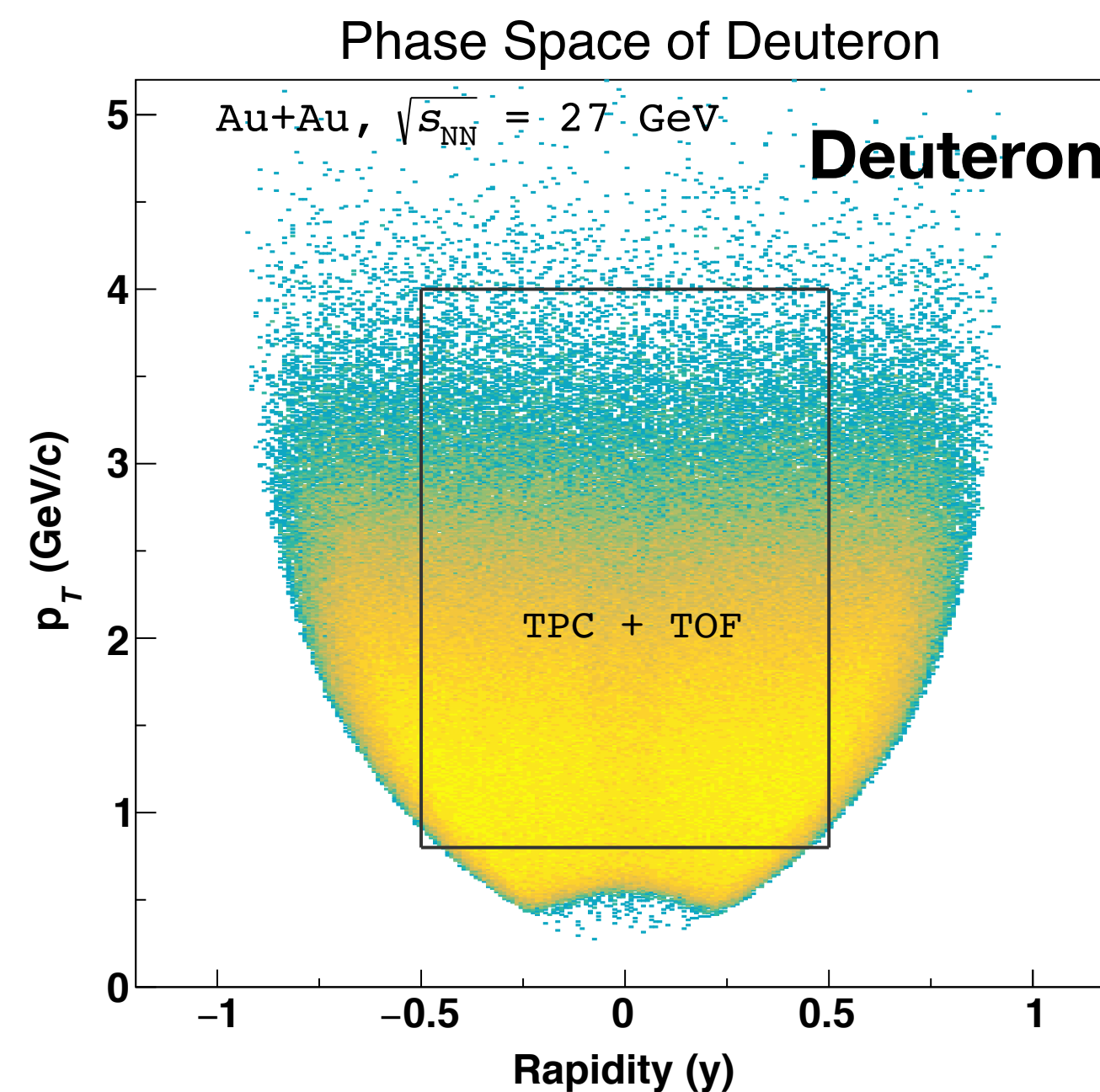
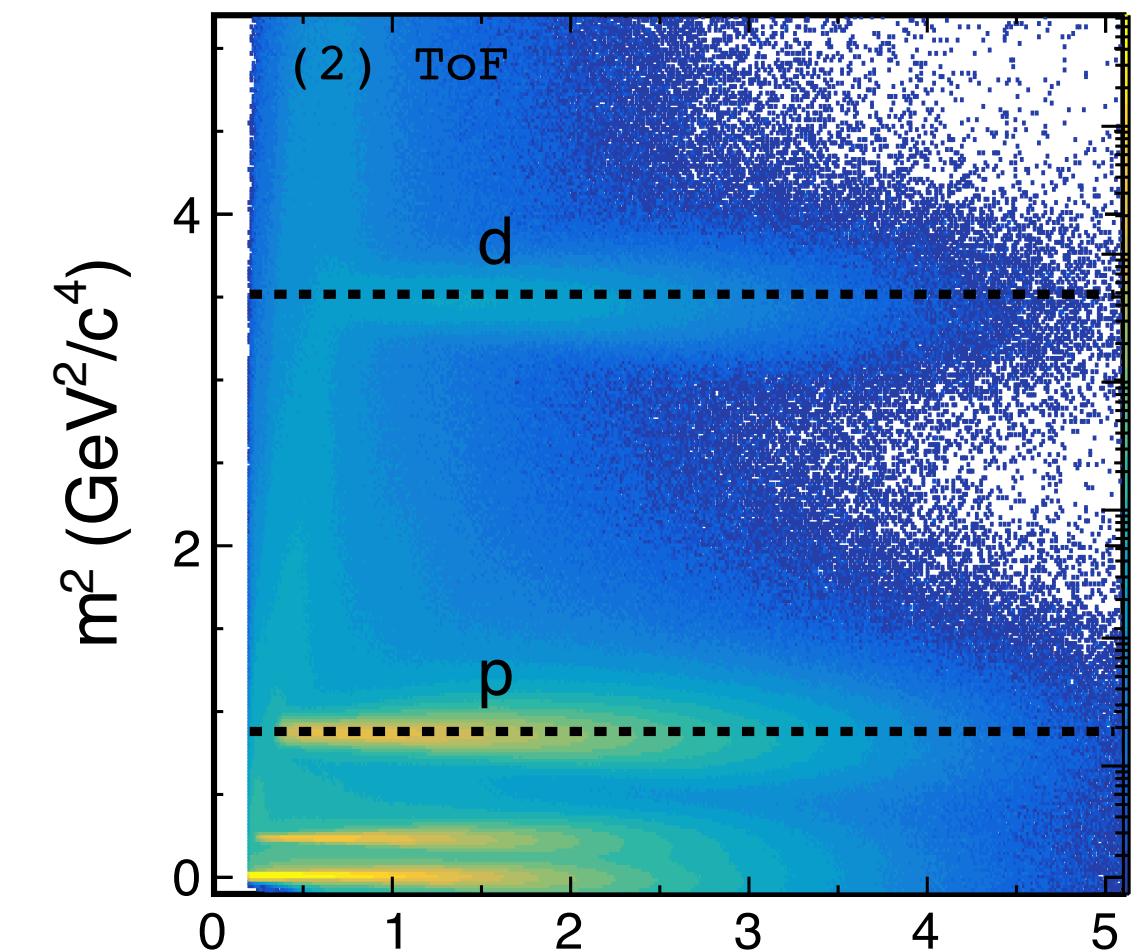
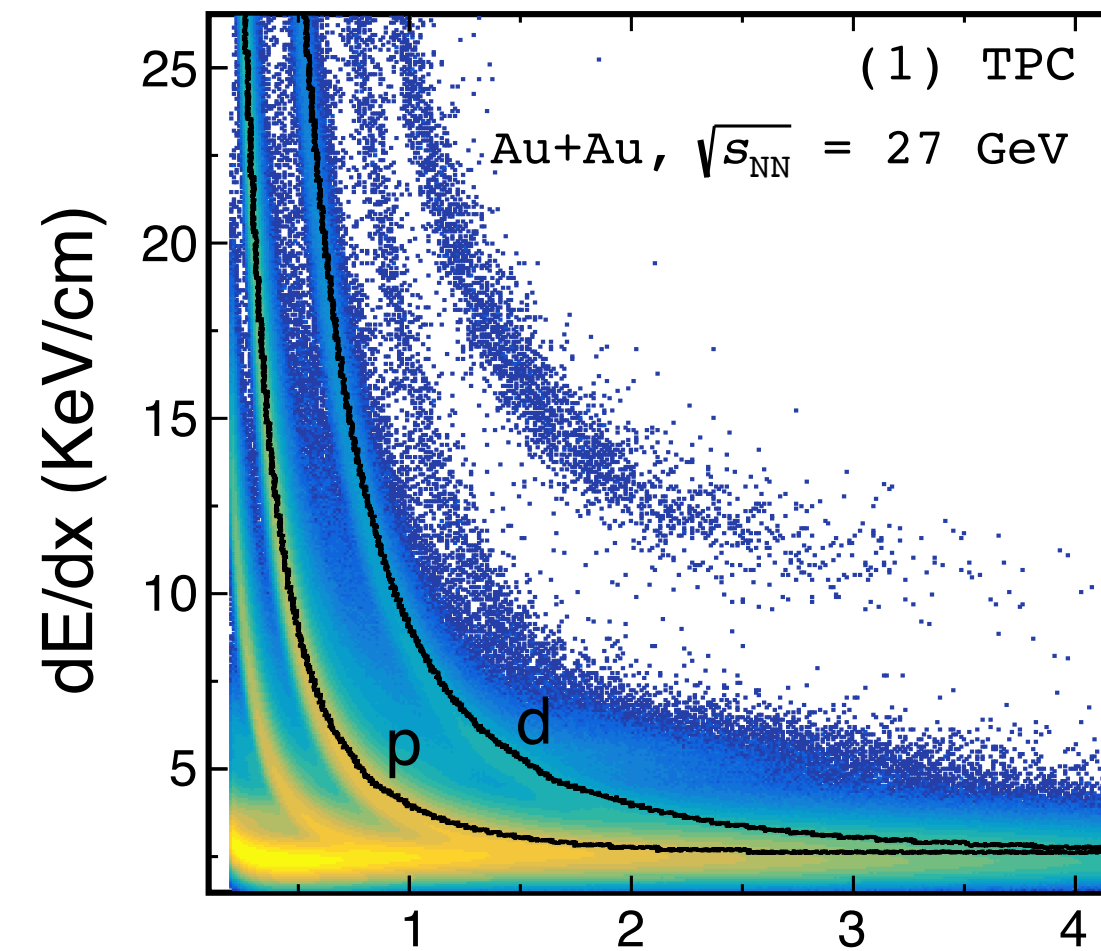


STAR: Nucl.Instrum.Meth.A 499 (2003) 624-632

**PID and Centrality:** Using both Time Projection Chamber (TPC) and Time-of-Flight (ToF) detectors.  
**Uniform coverage for full azimuth and  $|\eta| < 1$ .**  
**Excellent PID capability.**

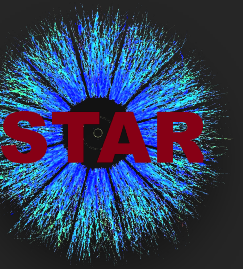
### Dataset: BES-I

Collision system: Au+Au collision (centrality: 0-5% , 70-80%)  
 CoM energy: 7.7, 11.5, 14.5, 19.6, 27, 39, 54.4, 62.4, 200 GeV  
 Year : 2010 — 2017



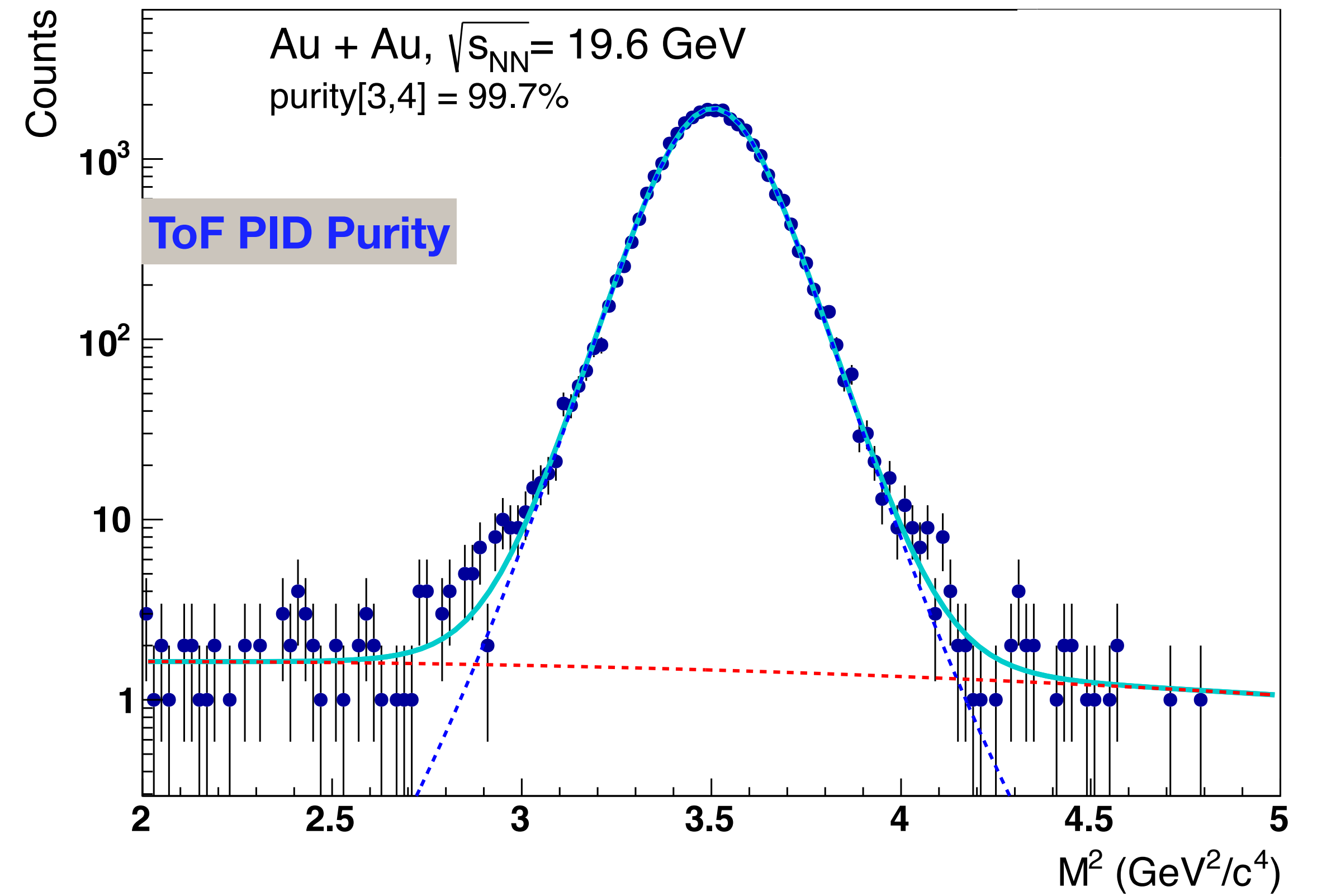
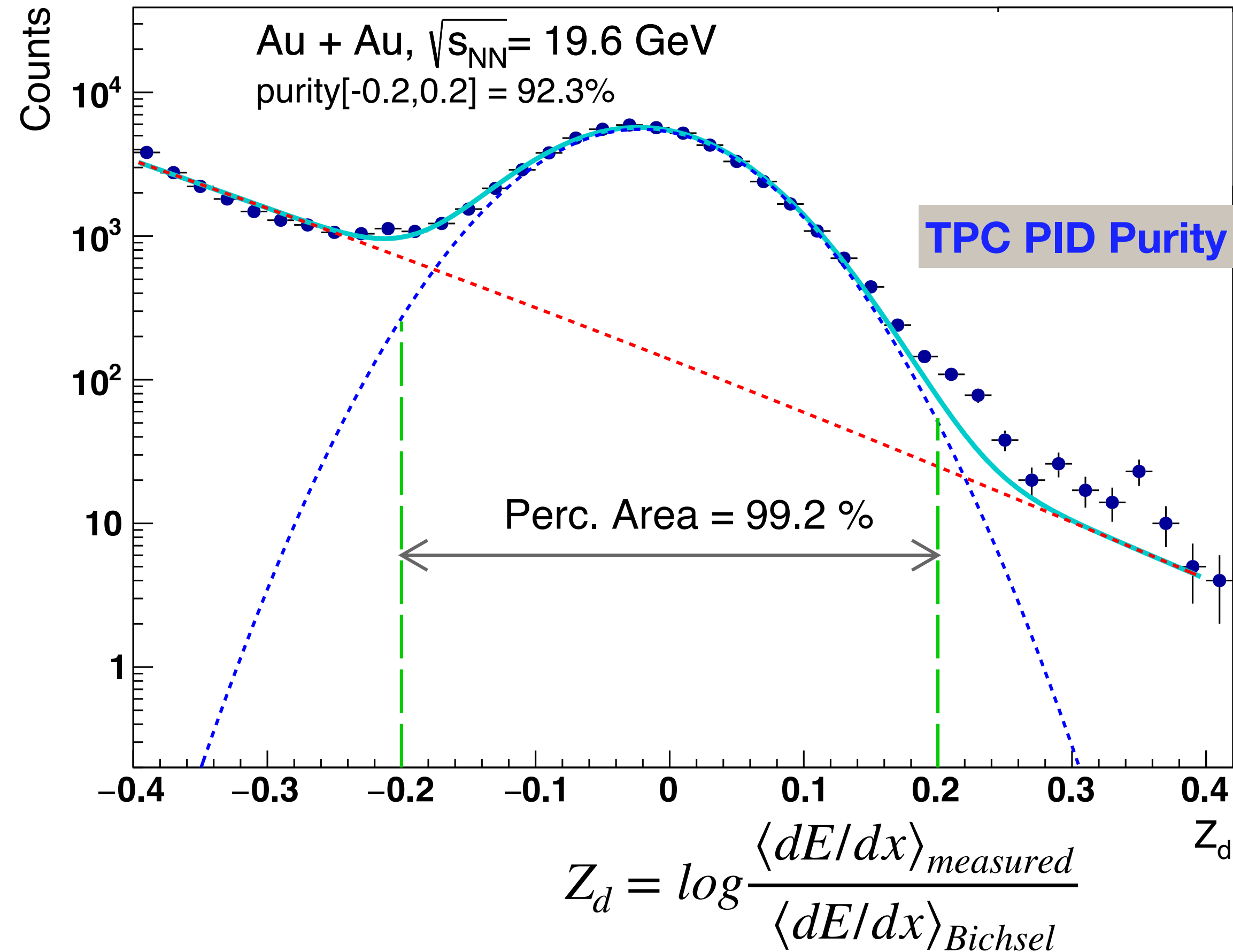


# Purity



$Z_d$  Distribution, 0-10% ,  $0.8 < p_T < 1.0$  GeV/c,  $|y| < 0.5$

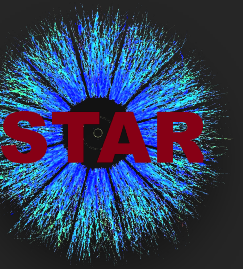
$m^2$  Distribution, 0-10% ,  $0.8 < p_T < 1.0$  GeV/c,  $|y| < 0.5$



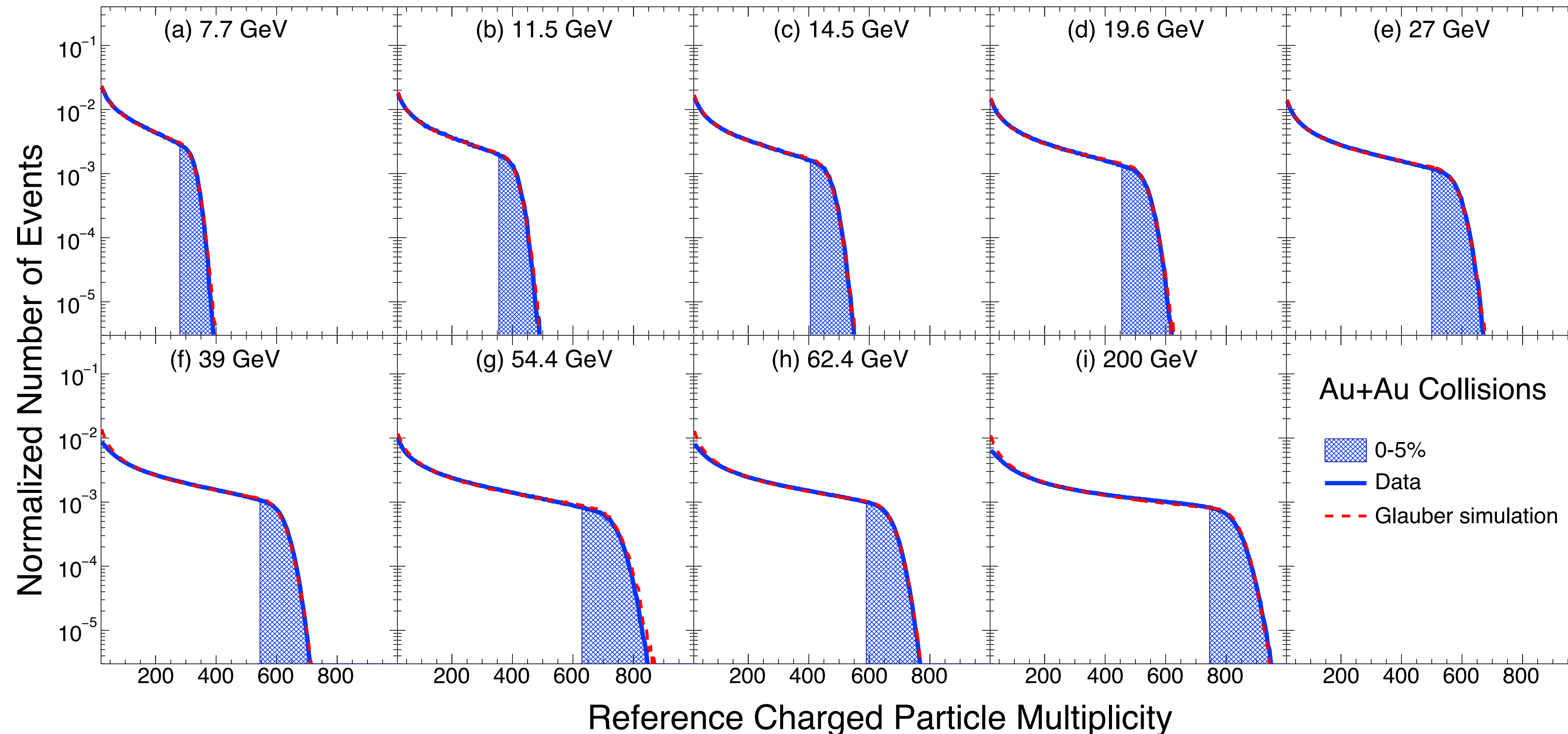
To achieve better PID purity, deuterons are identified using always both TPC and ToF detector.  
Distance of Closest Approach (DCA) is kept as  $DCA < 1$ cm to reduce the background contribution.



# Centrality Definition



Centrality using charged particles within  $|\eta| < 1.0$ , **excluding protons and deuterons**



Charged particle multiplicity is corrected for dependencies on (a) Collision vertex and (b) Beam luminosity

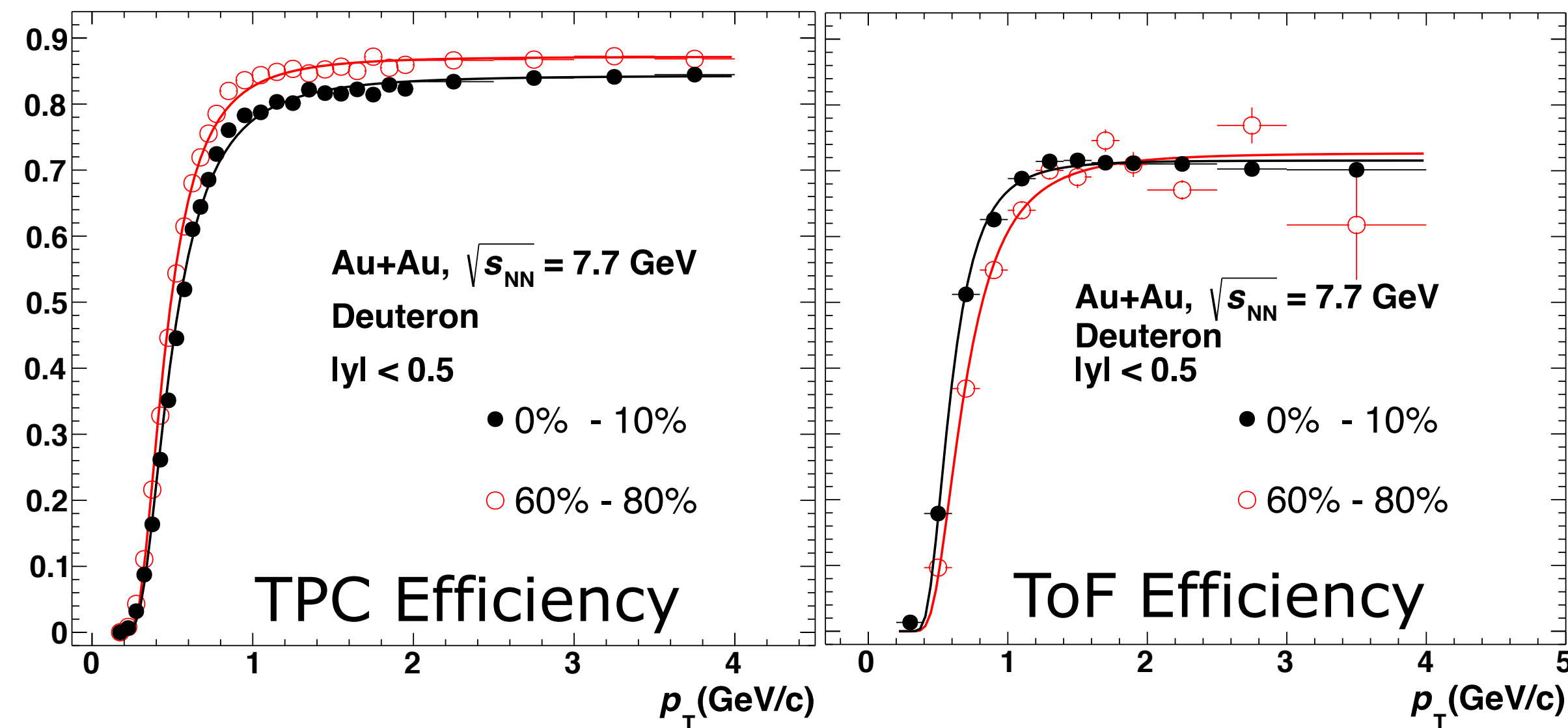
**Not** corrected for detector **efficiency**.

*STAR: Phys. Rev. C 104, 024902 (2021)*

This definition excludes self/auto-correlations between centrality and particle of interest.



## 1) Detection efficiency correction - binomial model



## 2) Centrality bin-width (CBW) correction:

❖ Effect arises from the dependence of  $C_n$  on multiplicity.

$$C_n = \sum_r \omega_r C_{n,r}, \quad \omega_r = \frac{n_r}{\sum_r n_r}$$

$n_r$  is number of events in  $r$ -th multiplicity bin.

## 3) Statistical uncertainty:

Using re-sampling technique called Bootstrap method.

$$\text{For a statistic } X, \text{Var}(X) = \frac{1}{S-1} \sum_{s=1}^S (X_s^* - \bar{X})^2.$$

$S$  is the number of samples.

$X_s^*$  is "X" measured from  $s$ -th sample.

## 4) Systematic uncertainty:

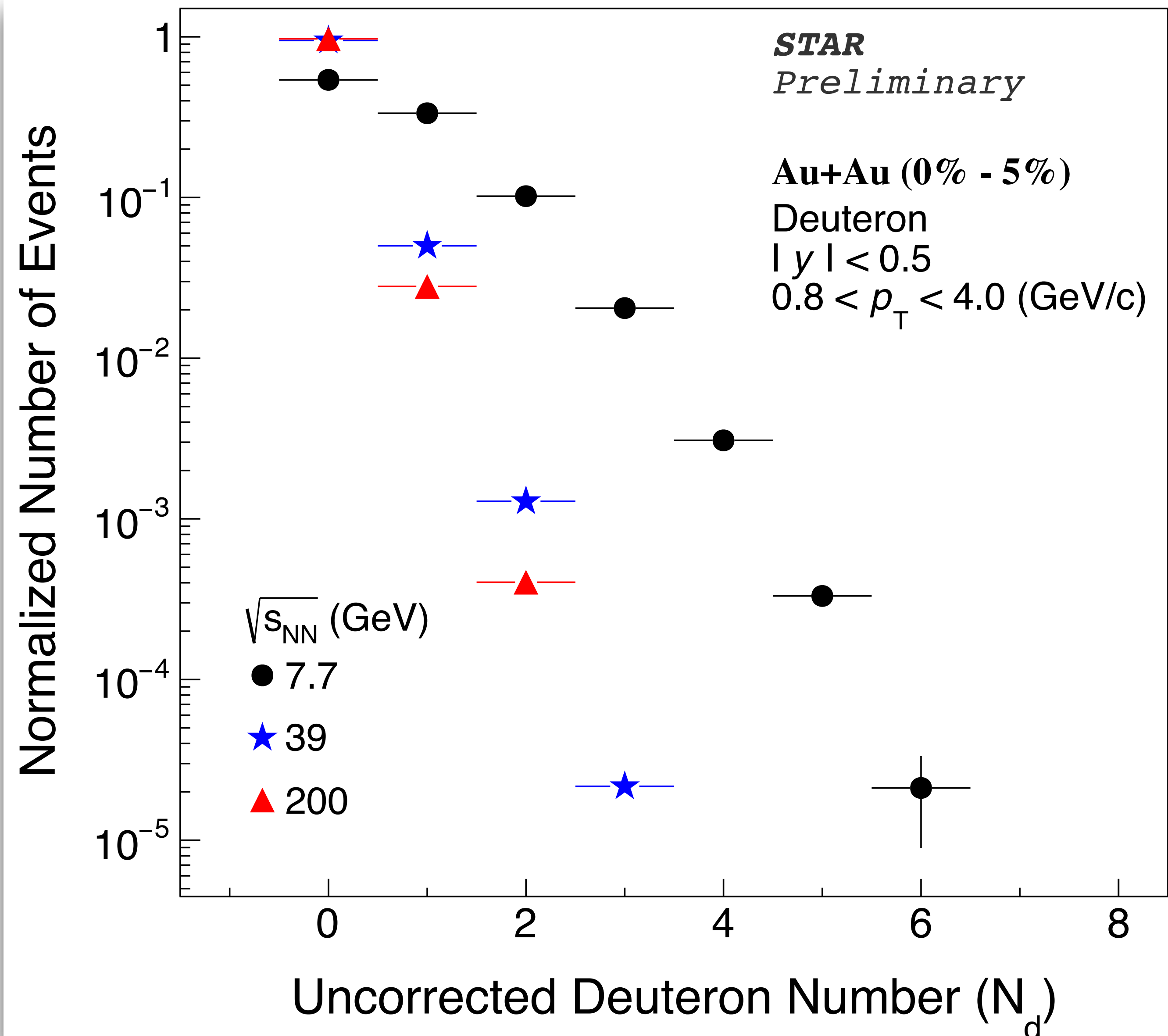
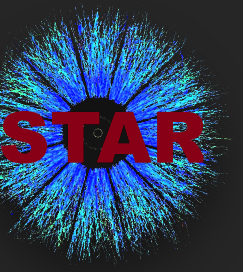
Sources:

- Particle identification from TPC and ToF
- Background/decay estimates (DCA)
- Quality cuts for track reconstruction
- Uncertainty in detection efficiency estimation

STAR: *Phys. Rev. C* 104, 024902 (2021)  
 X. Luo, *Phys. Rev. C* 91, (2015) 034907  
 T. Nonaka et al, *Phys. Rev. C* 95, (2017) 064912  
 X. Luo et al, *J. Phys. G* 40, 105104 (2013)  
 X. Luo, *J. Phys. G* 39, 025008 (2012)  
 X. Luo et al, *Phys. Rev. C* 99 (2019) no.4, 044917  
 A. Pandav et al, *Nucl. Phys. A* 991, (2019) 121608



# Raw Deuteron Number Distribution

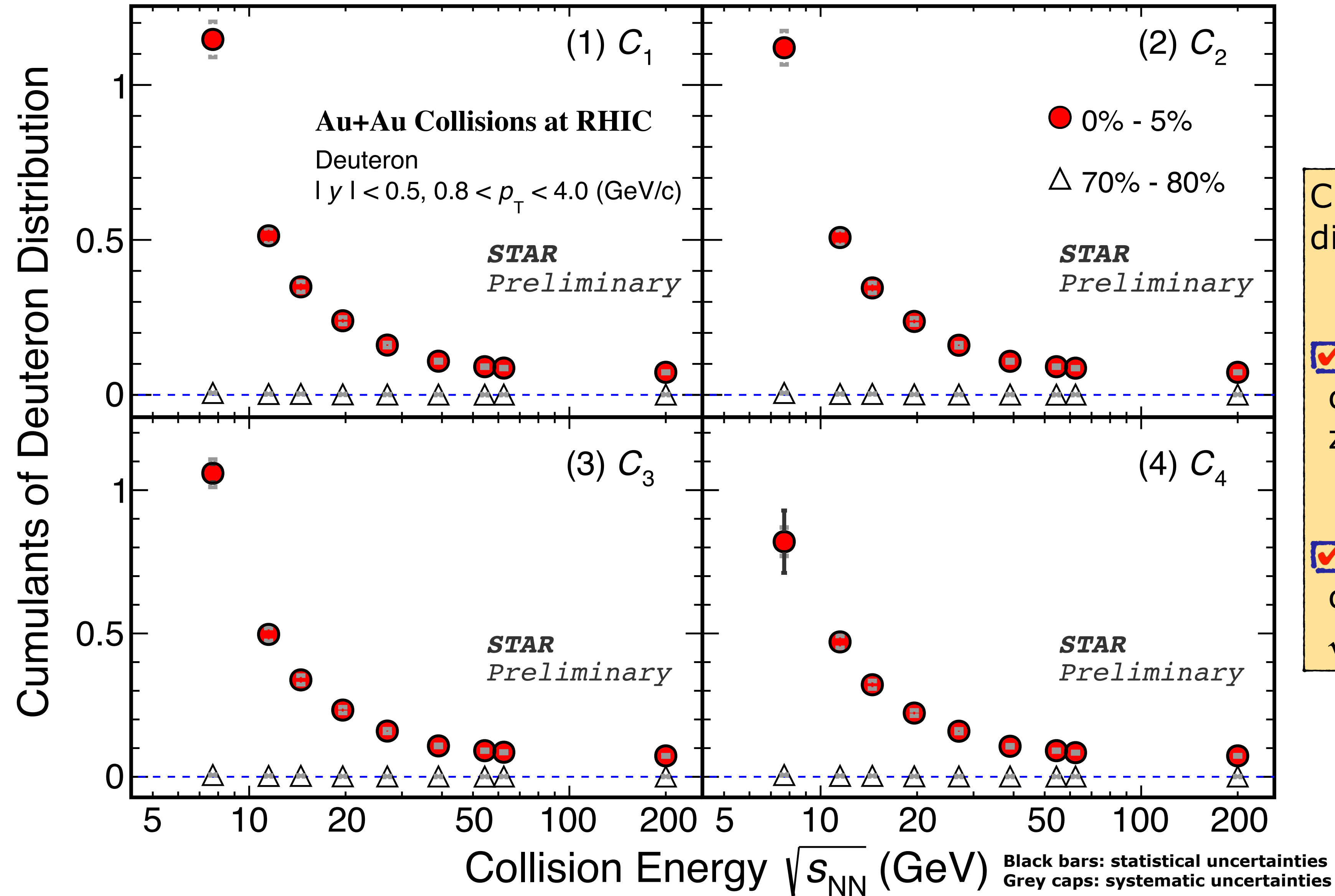
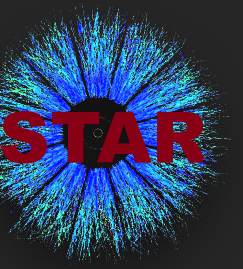


Uncorrected for efficiency and CBW effect.

Deuteron production increases towards low  $\sqrt{s_{NN}}$ .

Mean and width of distribution increase for low  $\sqrt{s_{NN}}$ .

# Cumulants of Deuteron Distribution

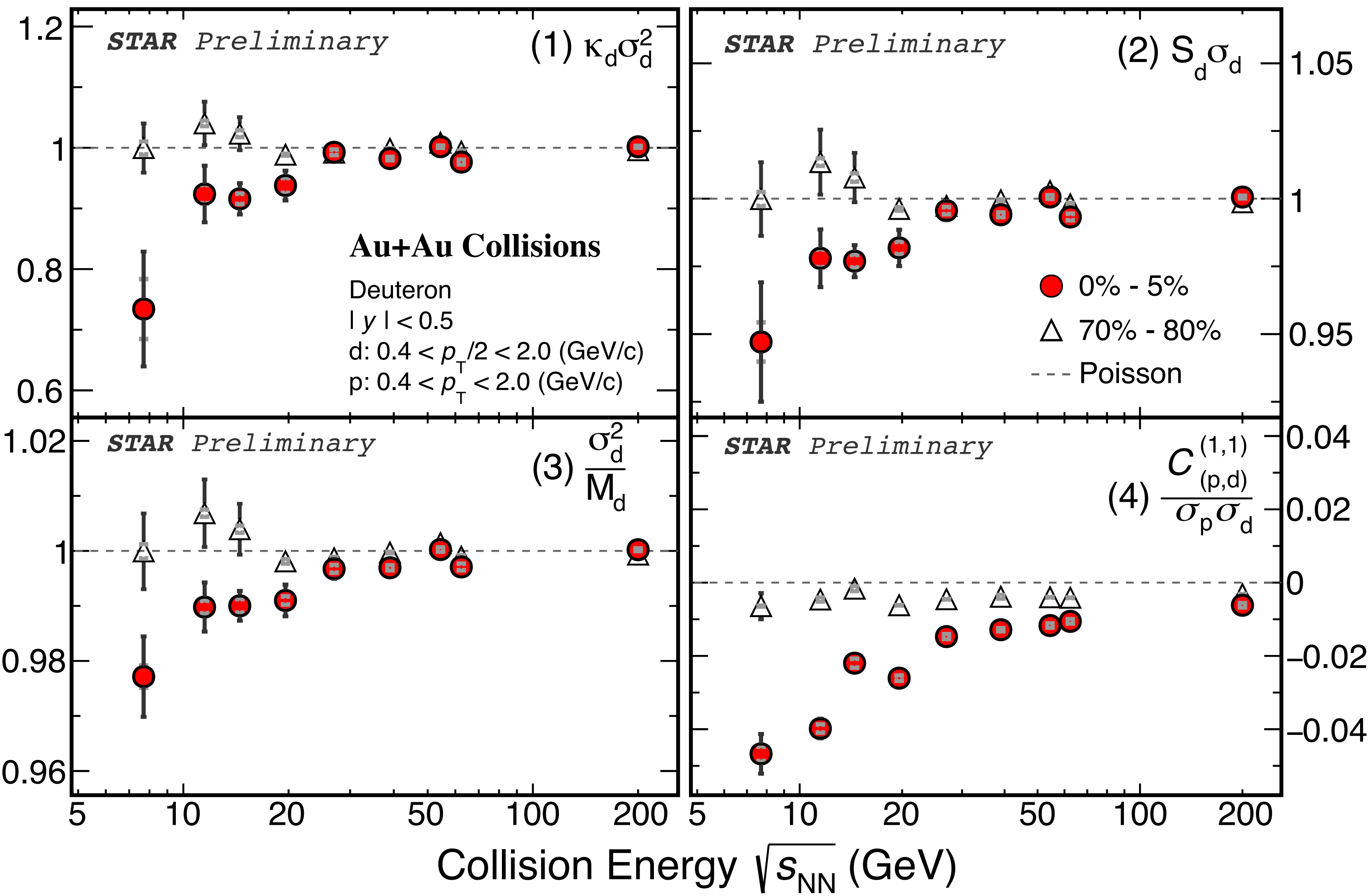
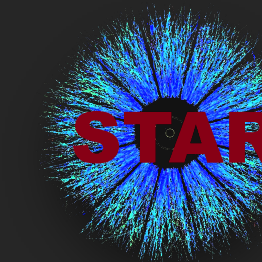


Cumulants ( $C_n$ ) of the deuteron distributions.

- For peripheral (70%-80%) Au+Au collisions, cumulants are close to zero.
- In most central (0-5%) collisions, cumulants increase as the collision  $\sqrt{s_{NN}}$  decreases.



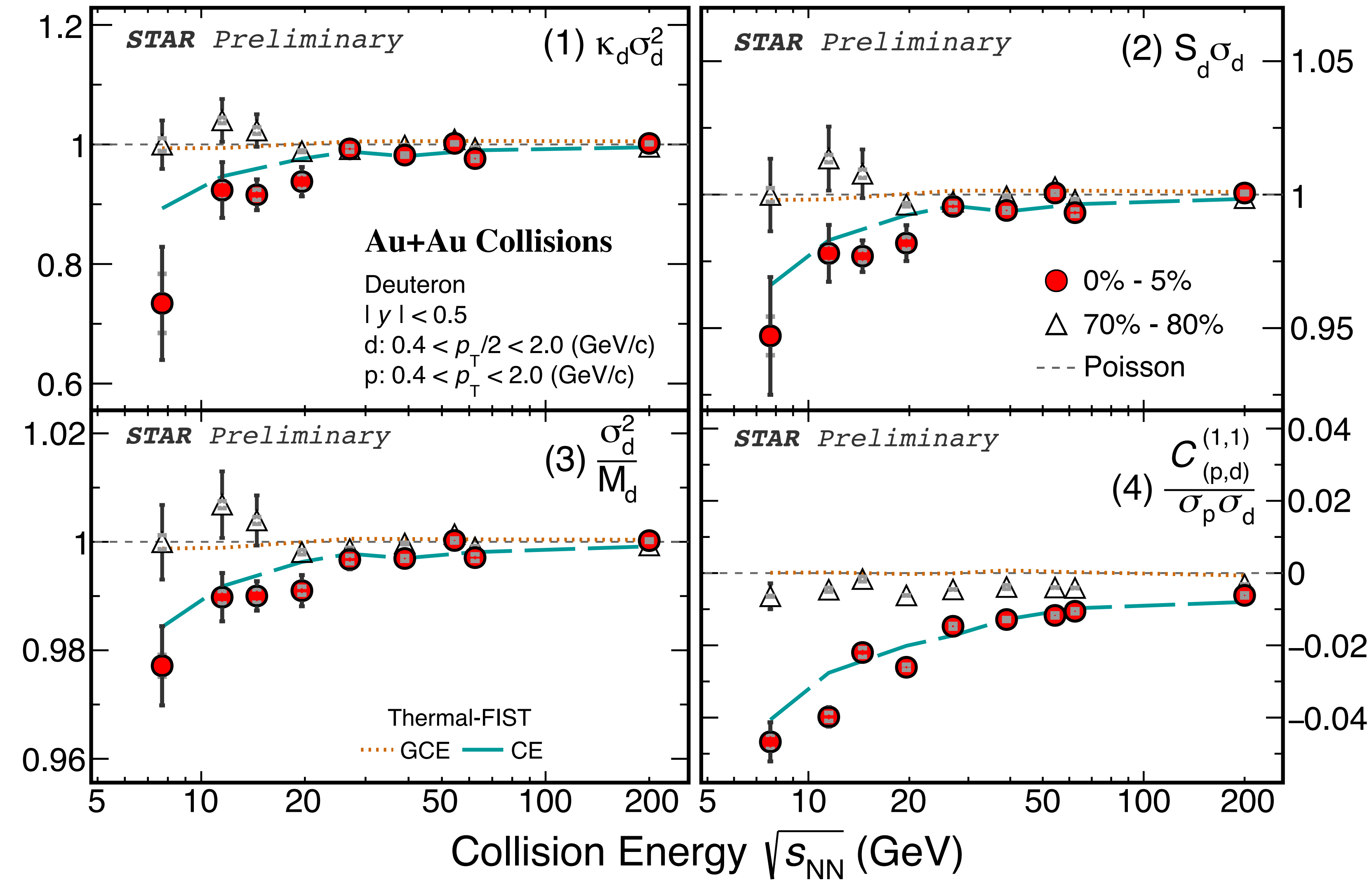
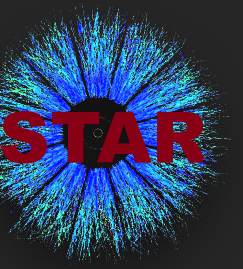
# Cumulant Ratios and p-d Correlation



- ✓ Cumulant ratios in 0-5% centrality, show monotonic dependence on  $\sqrt{s_{NN}}$ .
- ✓ Ratios in 70-80% centrality show weak  $\sqrt{s_{NN}}$  dependence and are close to 1.
- ✓ In panel(4), negative value of correlation suggests, proton and deuteron number are anti-correlated across all collision energy and centrality.
- ✓ With lowering the  $\sqrt{s_{NN}}$ , anti-correlation becomes stronger.

**Black bars: statistical uncertainties**  
**Grey caps: systematic uncertainties**

# Cumulant Ratios and p-d Correlation

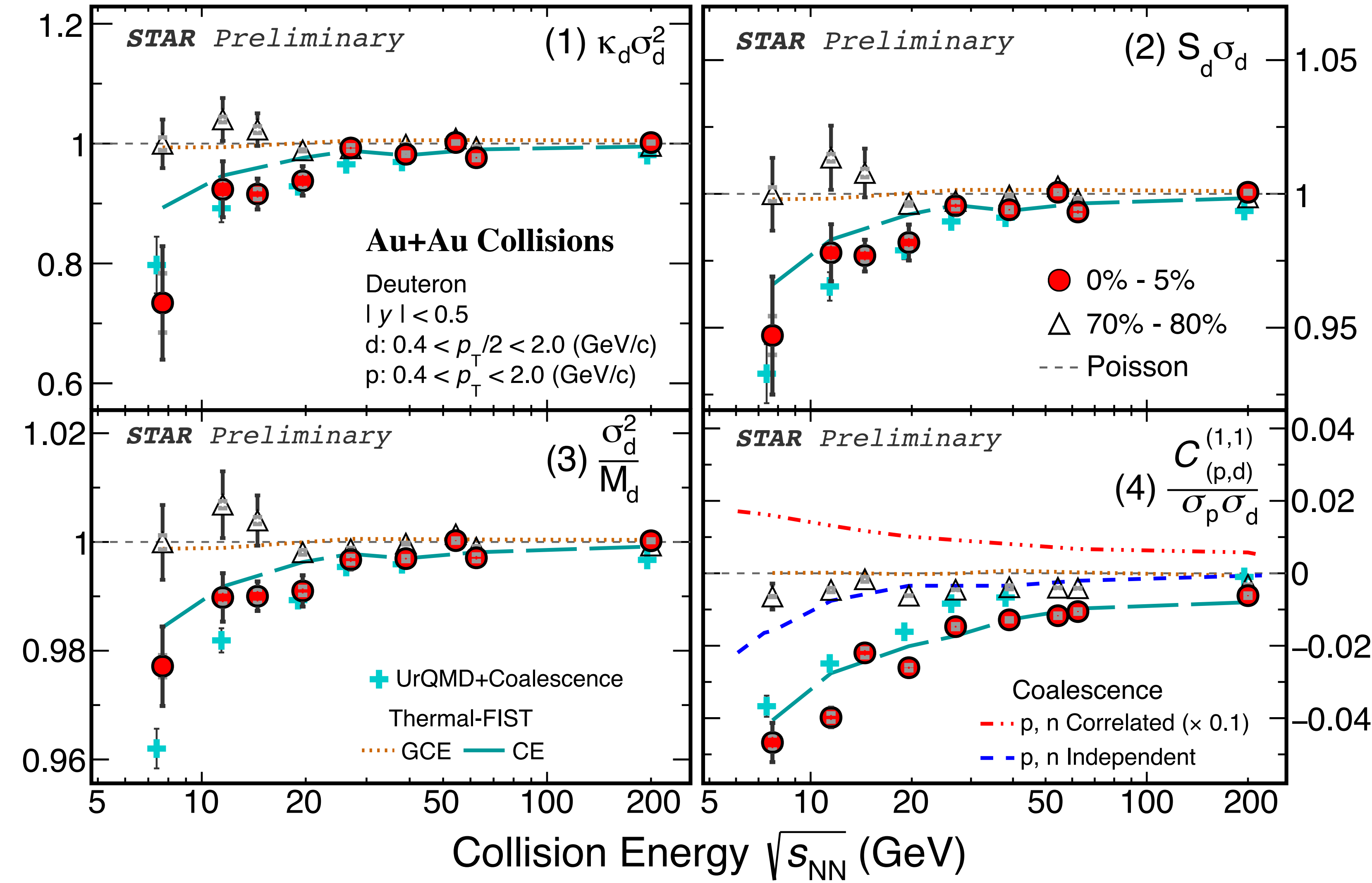
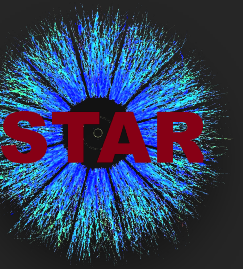


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- ✓ With lowering the  $\sqrt{s_{NN}}$ , anti-correlation becomes stronger.
- ✓ GCE thermal model seems to fail to describe the cumulant ratios for lower  $\sqrt{s_{NN}}$ .

**Black bars: statistical uncertainties**  
**Grey caps: systematic uncertainties**



# Cumulant Ratios and p-d Correlation

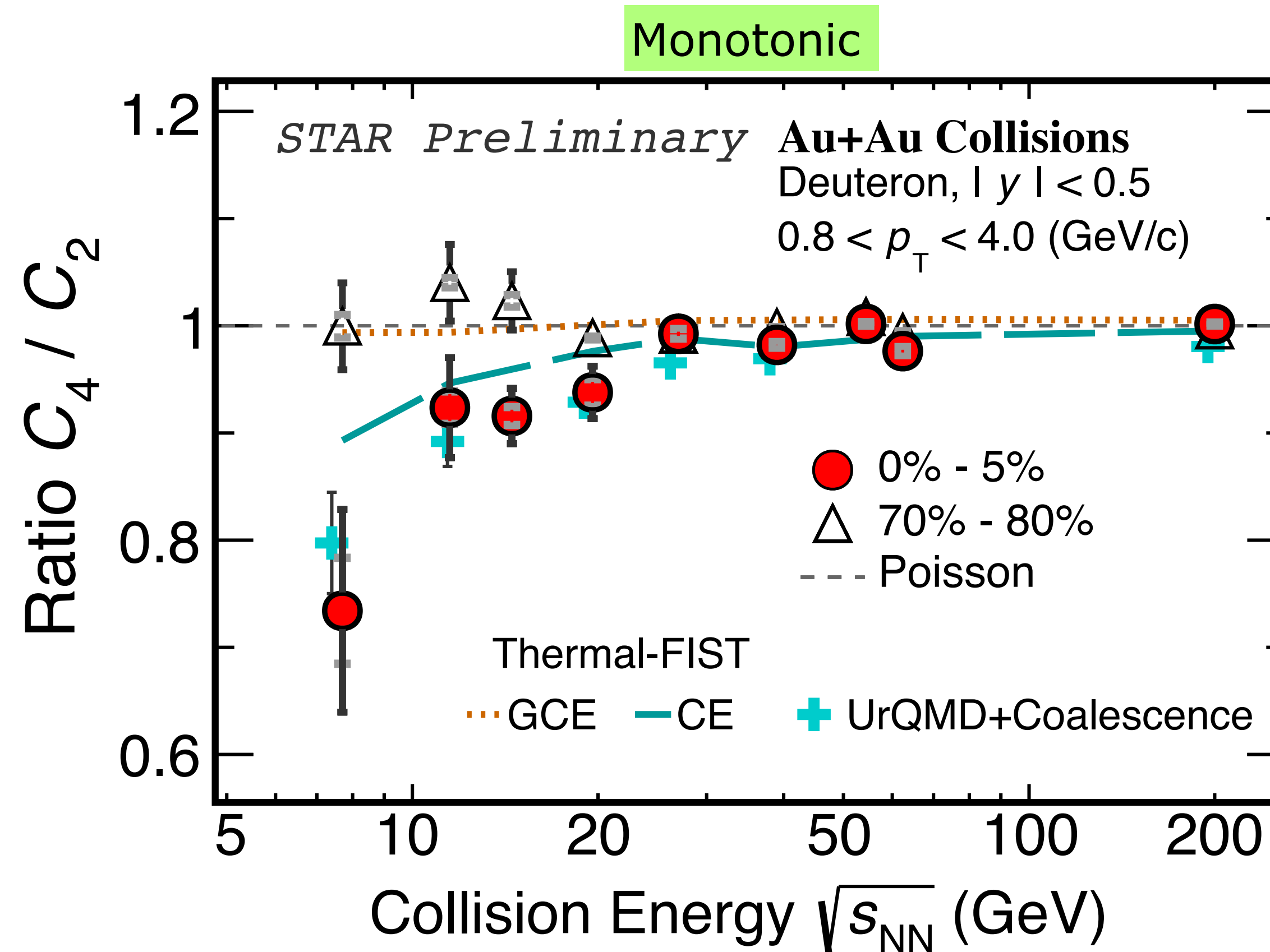
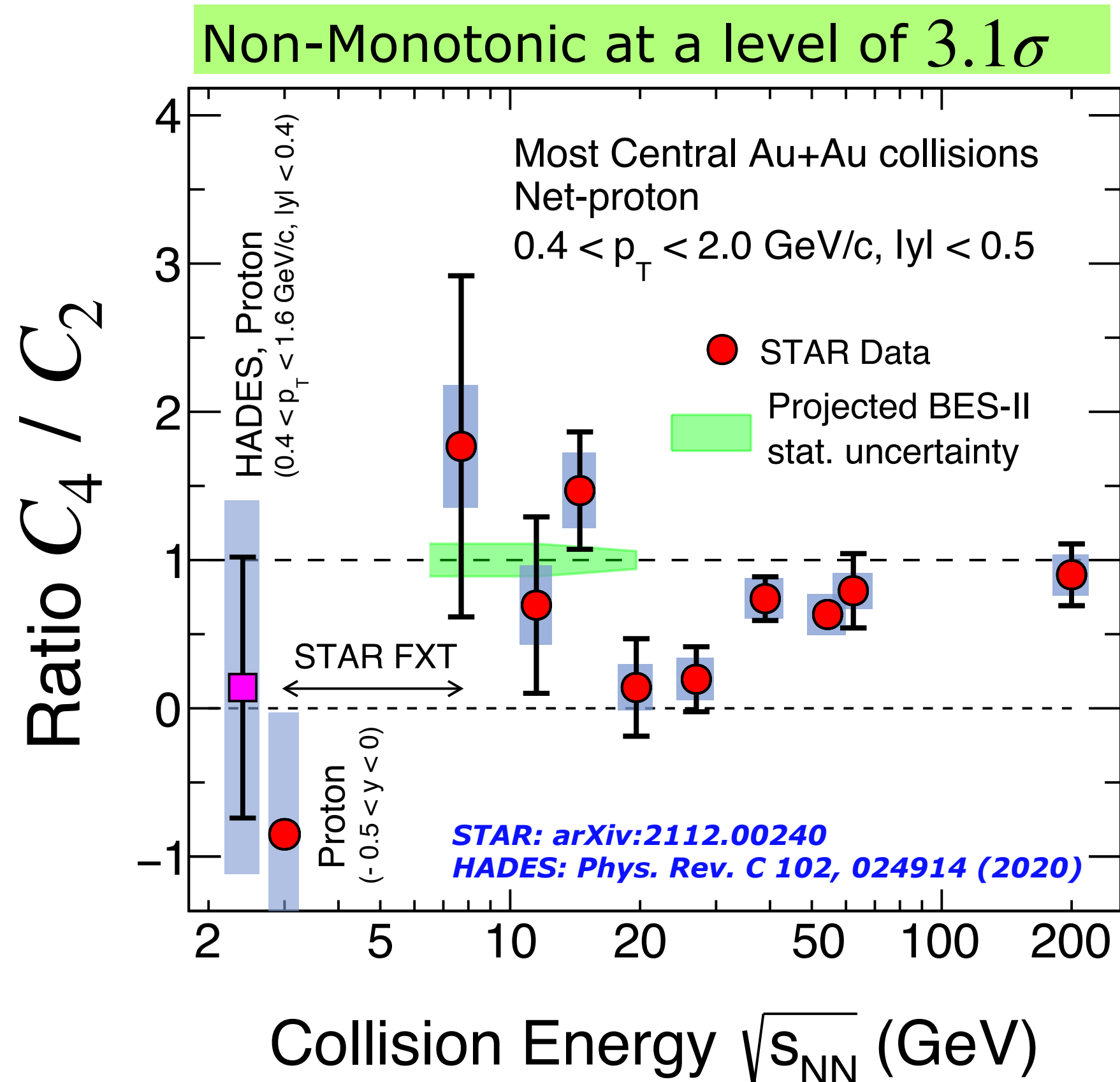
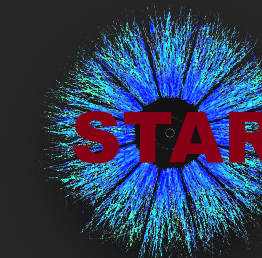


**Black bars: statistical uncertainties**  
**Grey caps: systematic uncertainties**

- ✓ Cumulant ratios in 0-5% centrality, show monotonic dependence on  $\sqrt{s_{NN}}$ .
- ✓ Ratios in 70-80% centrality show weak  $\sqrt{s_{NN}}$  dependence and are close to 1.
- ✓ In panel(4), negative value of correlation suggests, proton and deuteron number are anti-correlated across all collision energy and centrality.
- ✓ With lowering the  $\sqrt{s_{NN}}$ , anti-correlation becomes stronger.
- ✓ GCE thermal model seems to fail to describe the cumulant ratios for lower  $\sqrt{s_{NN}}$ .
- ✓ UrQMD+Coalescence and CE thermal model qualitatively reproduce collision energy dependence.
- ✓ Neither correlated nor independent assumption for proton and neutron in the toy model from [Z. Fecková et. al.; PRC 93, 054906 \(2016\)](#) reproduce the data.



# Comparison with Net-proton



Black bars: statistical uncertainties  
Grey caps: systematic uncertainties

Deuteron number  $\kappa\sigma^2$  in 0-5% centrality shows monotonic energy dependence in contrast to protons.

## Possibilities:

- **Low yield** of deuteron affecting sensitivity to critical point physics ?
- Probing **different freeze-out** surfaces ? More investigation ongoing. Theoretical inputs are also needed.



## Summary:

- ✓ We reported the first measurements of cumulants of deuteron number distribution, their ratios and proton-deuteron correlation in 0-5% and 70-80% central Au+Au collisions for  $\sqrt{s_{NN}} = 7.7 - 200$  GeV.
- ✓ UrQMD + phase-space coalescence model fairly describes the cumulant ratios and correlation for 0-5% centrality.
- ✓ For all  $\sqrt{s_{NN}}$ , proton and deuteron numbers are anti-correlated. With lowering  $\sqrt{s_{NN}}$ , anti-correlation in 0-5% centrality becomes stronger.
- ✓ HRG GCE thermal model fails to describe cumulant ratios at collision energies  $\sqrt{s_{NN}} \leq 19.6$  GeV. HRG CE and UrQMD show suppression below unity for lower  $\sqrt{s_{NN}}$ , as seen in the data, could arise from the effect of global baryon number conservation. Sensitive to the choice of ensembles.
- ✓  $\kappa\sigma^2$  of deuteron number in 0-5% centrality shows monotonic energy dependence in contrast to proton fluctuations. [STAR: Phys. Rev.Lett. 126 \(2021\) 092301](#)

## Outlook:

### Using BES-II data,

- Study contribution of  $p, d, t, He^3$  etc. together to understand net-baryon fluctuation in low  $\sqrt{s_{NN}}$ .
- Understand the production mechanism and freeze-out properties of light nuclei.