# Measurement of Transverse Single Spin Asymmetry at Forward Rapidity by the STAR Experiment in p+p Collisions at $\sqrt{s}$ = 200 and 500 GeV

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Preprint available in arXiv:2012.11428 Accepted for publication in PRD

## Motivation

- Transverse single spin asymmetry( $TSSA/A_N$ )
- The large forward TSSA was first found in 1970s and can not be explained by LO QCD calculation



Aidala et al. Rev. Mod. Phys., 85,655(2013)

A lot of work was done to explore the underlying mechanisms in the past few decades



## **Motivation**

- Transverse momentum dependent PDF(TMD)
- Collinear twist-3 factorization

These two models have different energy scale requirements, but they share some similarities

- A decomposition of the contributions to TMD
  - Initial state effect: asymmetry originates from PDF

 $\hat{f}_{q/p^{\dagger}}(x, \mathbf{k}_{\perp}) = f_{q/p}(x, k_{\perp}) + \frac{1}{2} \Delta^{N} f_{q/p^{\dagger}}(x, k_{\perp}) \mathbf{S} \cdot \left(\hat{\mathbf{P}} \times \hat{\mathbf{k}}_{\perp}\right) \qquad \text{Sivers function}$ 

**Final state effect**: asymmetry originates from fragmentation

Transversity  $\otimes$  **Collins function** 

Both effects can contribute to the TSSA.

• Experimental data are very important in validating the factorization and constraining the PDFs

#### **Motivation**

Jet TSSA – sensitive to the initial state effect.Collins asymmetry – sensitive to the final state effect.



#### **Experiment Setup- RHIC & STAR**

The Relativistic Heavy Ion Collider at BNL provides unique opportunity to study spin physics because it is the world's only polarized proton-proton collider.
Polarized protons





EM-Calorimeter made of 1000+ lead glass cells
 Large pseudo-rapidity range in the forward direction 2.6-4.1
 Two cell types

### **Analysis- Dataset**

#### Dataset:

Transversely polarized proton-proton collisions

Year	Energy	Events
2011	500 GeV	165M
2015	200 GeV	569M

#### Beam polarization:

52 / 57% (500 / 200 GeV)

#### Trigger:

FMS-Board-sum and FMS-Jet-patch, both based on energy deposition in a defined region of the FMS





#### **Analysis- Asymmetry calculation**

The **luminosity** and **detector efficiency** can be difficult to determine.

$$N^{\uparrow}(\phi) = \epsilon \mathcal{L}^{\uparrow} \sigma^{\uparrow}$$
$$= \epsilon \mathcal{L}^{\uparrow} (1 + pol * A_N \cos \phi) \sigma$$

• "Cross-ratio" method help eliminate those factors



$$pol \cdot A_N^{\text{raw}} \cos \phi = \frac{\sqrt{N^{\uparrow}(\phi)N^{\downarrow}(\phi+\pi)} - \sqrt{N^{\downarrow}(\phi)N^{\uparrow}(\phi+\pi)}}{\sqrt{N^{\uparrow}(\phi)N^{\downarrow}(\phi+\pi)} + \sqrt{N^{\downarrow}(\phi)N^{\uparrow}(\phi+\pi)}}$$

#### Background subtraction

The fraction comes from the fitting of the mass spectrum Signal/background shapes are from simulation

$$A_N^{\text{raw}_{sig}} = f_{\text{sig}_{sig}} * A_N^{\pi^0} + (1 - f_{\text{sig}_{sig}}) * A_N^{bkg}$$
$$A_N^{\text{raw}_{sb}} = f_{\text{sig}_{sb}} * A_N^{\pi^0} + (1 - f_{sig_{sb}}) * A_N^{bkg}$$

## **Analysis- Collins Asymmetry**

VS.

 $\pi^{0} / \text{EM-jet TSSA}$  $N^{\uparrow}(\phi) = \epsilon \mathcal{L}^{\uparrow} \sigma^{\uparrow}$  $= \epsilon \mathcal{L}^{\uparrow} (1 + pol * A_{N} \cos \phi) \sigma$ 

- > Azimuthal angle
- > All  $\pi^0$  candidates
- > Background subtraction for  $\pi^0$

Collins asymmetry  $N^{\uparrow}(\phi_c) = \epsilon \mathcal{L}^{\uparrow} \sigma^{\uparrow}$   $= \epsilon \mathcal{L}^{\uparrow} (1 + pol * A_{UT} \sin \phi_c) \sigma$ 

Collins angle

- > Only  $\pi^0$  within a jet
- > No background subtraction

For jet reconstruction: For  $\pi^0$  in a jet :

- Anti-kt R=0.7  $\Delta R = \sqrt{\Delta \phi^2 + \Delta \eta^2}$
- $p_T > 2 \text{ GeV} > 0.04$

The jet is only "electromagnetic jet" 10

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# **Analysis- Systematic uncertainty**

Uncertainties:

- $\pi^0$ /jet energy scale uncertainty ( $x_F$  and  $z_{em}$ ): calibration, non-linear response, radiation damage
- $\pi^0$  TSSA: background subtraction
- Beam polarization

Corrections:

- Jet TSSA: background correction, underlying event correction, correction to particle level
- Collins asymmetry: Collins angle resolution correction

Analysis	Uncertainties types (Run-11/Run15)		
$\pi^0$ TSSA	$x_F$	Asymmetry	Beam polarization
	4.4%/3.0%	5.8%	3.4%/3.0%
Jet TSSA	$x_F$	Asymmetry	Beam polarization
	7.8%/8.5%	_	3.4%/3.0%
Collins Asymmetry	$z_{em}$	Asymmetry	Beam polarization
	8.9%/9.0%	_	3.4%/3.0%

#### **Analysis- Observables**

All measurements are done in 200 GeV (2015) and 500 GeV (2011) p+p collision

- 1)  $\pi^0$  TSSA: initial+final state effect TSSA as a function of Feynman-x  $(x_F)$ ;  $x_F = \frac{E_L^{\pi^0}}{E_{beam}}$ TSSA as a function of  $p_T$ ; Isolated/non-isolated  $\pi^0 A_N$  as a function of Feynman-x
- 2) Jet TSSA : initial state effect
- 3) Collins Asymmetry : **final** state effect

The jets used in 2) 3) are electromagnetic jet (EM-jet)

**Result-**  $\pi^0$  **TSSA vs.**  $x_F$ 



 $\Box$  The  $\pi^0$  TSSA increases with  $x_F$ .

□ Consistent between 200 GeV and 500 GeV. Energy dependence is weak.13

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# **Comparison to previous measurements**



STAR, arXiv:2012.11428

- □ Weak collision energy dependence of the  $\pi^0$  TSSA from 19.4 to 500 GeV
- □ Comparison to the previous Forward Pion Detector results at STAR shows larger TSSA in current measurement, which can be explained by the higher average  $p_T$

Result-  $\pi^0$  TSSA vs.  $p_T$ 





#### STAR, arXiv:2012.11428

 $\Box$  Overlapping  $x_F$  region between 200 GeV and 500 GeV results.

- □ The 200 GeV data shows significant increase of TSSA below 3 GeV.
- $\Box$  The 500 GeV data flattens over the  $p_T$  range.

# Result- isolated $\pi^0$ TSSA

- □ Motivation: investigate the  $\pi^0$  event topology ( $\pi^0$  with no other particle around)
- Dethod: in a surrounding area (in η- $\phi$  space, R=0.7), if the  $\pi^0$  takes most of the total energy, it is defined as isolated. The cut is placed at an energy fraction z=0.9 and 0.98

Fractions of different types of  $\pi^0$  event in the overall sample





# **Result- isolated** $\pi^0$ **TSSA**



□ The TSSAs of the two types of  $\pi^0$  are significantly different. Isolated  $\pi^0$  TSSA dominates.

□ The physical origin and mechanism accounting for higher TSSA of isolated  $\pi^0$  is not known yet – implication of a third origin?

Theory curves: J. Cammarota, et al. Phys.Rev.D.102.054002

# **Result- jet TSSA**

#### STAR, arXiv:2012.11428



Theory curves: L. Gamberg, Z. Kang, A. Prokudin, Phys.Rev.Lett.110,232301 □ The jet TSSA is a few times smaller than the  $\pi^0$  TSSA in the same  $x_F$  bin.

- Jets with minimum photon multiplicity requirement have significantly smaller TSSA.
- □ The  $A_N DY$  result shows the TSSA of full jets, and is consistent with the result of the EM-jet having at least 3 photons.

# **Result-** Collins Asymmetry for $\pi^0$ in a jet



 $j_T = E_{\pi}$  projection perpendicular to jet

The Collins asymmetries are very small at both energies
 This reflects the cancellation of the Collins effect of the u/d quark
 Weak jT dependence is observed

## Summary

- □ We measured the  $\pi^0$ /jet TSSA and Collins asymmetry using the FMS in STAR 200 and 500 GeV p-p data
- $\Box$  The  $\pi^0$  TSSA results show weak energy dependence through 20 to 500 GeV
- □ We investigated the  $\pi^0$  event topology. The isolated  $\pi^0$  TSSAs are significantly larger than the non-isolated  $\pi^0$ , the mechanism of which remains unclear. It offers new perspectives to the origin of TSSA
- □ We measured the jet TSSAs and Collins asymmetry to separate contributions from initial and final state effects, both of which are small
- These measurements provide important inputs for further investigation for TSSA