

# 1 $J/\psi$ production in Ru+Ru and Zr+Zr collisions at $\sqrt{s_{NN}} =$ 2 200 GeV with the STAR experiment

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In this contribution, we present the measurements of the inclusive  $J/\psi$  yield and elliptic flow ( $v_2$ ) at midrapidity ( $|y| < 1.0$ ) in Ru+Ru and Zr+Zr collisions at  $\sqrt{s_{NN}} = 200$  GeV by the STAR experiment. The centrality and transverse momentum ( $p_T$ ) dependences of the nuclear modification factor ( $R_{AA}$ ) are measured with much better precision compared to the previous measurements in Au+Au collisions at the same collision energy. The  $p_T$  dependence of inclusive  $J/\psi$  elliptic flow ( $v_2$ ) in the same collision systems at  $|y| < 1.0$  are also presented. Two methods, the Scalar-Product method and the event-plane method are used individually to extract the  $J/\psi$   $v_2$  in Zr+Zr and Ru+Ru collisions. The two methods' results are consistent in the overlap  $p_T$  range. The measurements of the inclusive  $J/\psi$  yield and  $v_2$  are compared to the same results in Au+Au collisions at  $\sqrt{s_{NN}} = 200$  GeV, and then physics implications are discussed.

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## 9 1. Introduction

10 Lattice QCD, the discrete formulation of quantum chromodynamics, predicts a new state of  
11 matter created in ultra-relativistic heavy-ion collisions, where quarks and gluons are no longer  
12 confined within hadrons, called the quark-gluon plasma (QGP). The heavy quarks (charm and  
13 beauty), which are produced via initial hard partonic scatterings and then experience the full  
14 evolution of the QGP, are ideal probes to study the properties of the QGP [1]. Quarkonia are  
15 bound states of heavy quarks and their anti-quarks, for example the  $J/\psi$  meson ( $c\bar{c}$ ). It has been  
16 predicted that the  $J/\psi$  production should be suppressed due to the presence of the hot medium  
17 and the color-screening effect [2, 3]. The degree of suppression of the quarkonia depends on its  
18 binding energy and the properties of the QGP. On the other hand, uncorrelated or correlated charm  
19 quarks and anticharm quarks could bind into new  $J/\psi$  mesons in the deconfined medium, called the  
20 regeneration effect. The dissociation and regeneration effects are expected to depend on the size  
21 of collision systems. In 2018, the STAR experiment collected a large statistics sample of isobaric  
22 collisions ( $^{96}_{44}\text{Ru}+^{96}_{44}\text{Ru}$  and  $^{96}_{40}\text{Zr}+^{96}_{40}\text{Zr}$ ) at  $\sqrt{s_{NN}} = 200$  GeV. The system size of the Ru+Ru and Zr+Zr  
23 collisions is larger than Cu+Cu but smaller than Au+Au collisions, so it is a unique opportunity to  
24 study the collision system dependence of the inclusive  $J/\psi$  production.

25 Besides dissociation and regeneration effects, feed-down from the excited states and cold  
26 nuclear matter effects also contribute to the final  $J/\psi$  yield. Studying the properties of the QGP via  
27 the  $J/\psi$  production requires a good understanding of all the hot and cold nuclear matter effects. A  
28 precise measurement of  $J/\psi$  yield in a wide kinematic range is crucial to have a better understanding  
29 of different effects.

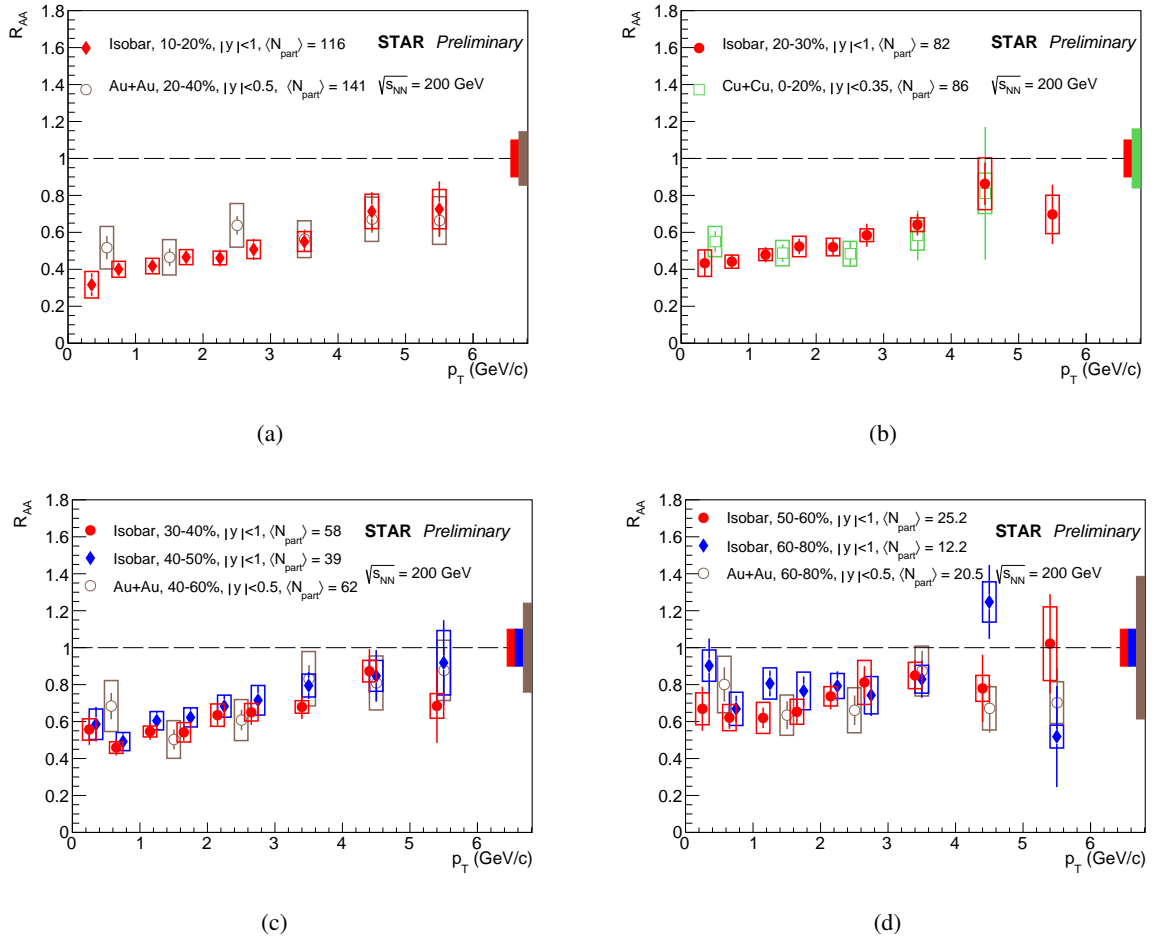
30 Another important observable for studying the properties of the QGP is the azimuthal depen-  
31 dence of  $J/\psi$  production.  $J/\psi$  mesons produced from regeneration will inherit the flow of charm  
32 quarks, while the  $J/\psi$  mesons from initial hard partonic scatterings are predicted to have very lim-  
33 ited flow. A significant elliptic flow of  $J/\psi$  mesons has been observed at LHC energy [4, 5], but in  
34 Au+Au collisions at the top RHIC energy, the elliptic flow of  $J/\psi$  mesons is consistent with zero  
35 within the large uncertainties [6].

## 36 2. Analysis and Results

37 The data used in this analysis are 4 billion events from isobaric collisions at  $\sqrt{s_{NN}} = 200$  GeV,  
38 collected in 2018 by the STAR experiment.  $J/\psi$  candidates are reconstructed in the dielectron decay  
39 channel. After implementing event-level and track quality cuts, we identified the electron candidate  
40 using TPC (Time Projection Chamber), TOF (Time of Flight), and BEMC (Barrel Electromagnetic  
41 Calorimeter) detectors. The Mixed-Event Technique was utilized to address the combinatorial  
42 background [7]. In the  $v_2$  analysis, the TPC second-order event plane is used to estimate the reaction  
43 plane for the BEMC-triggered events (triggered on a high- $p_T$  electron). Then  $v_2$  is calculated using  
44 the definition from [8]. For MB events,  $v_2$  is extracted using the Scalar-Product method [9]. The  
45 Q-vector of  $J/\psi$  is defined as  $Q_{2,J/\psi} = e^{in\varphi}$ , where  $\varphi$  represents the azimuthal angle.  $Q_{2,EPD}$   
46 represents the event flow vector in the EPD (Event Plane Detector).  $v_2^{SP}$  is obtained using Eq. 1,

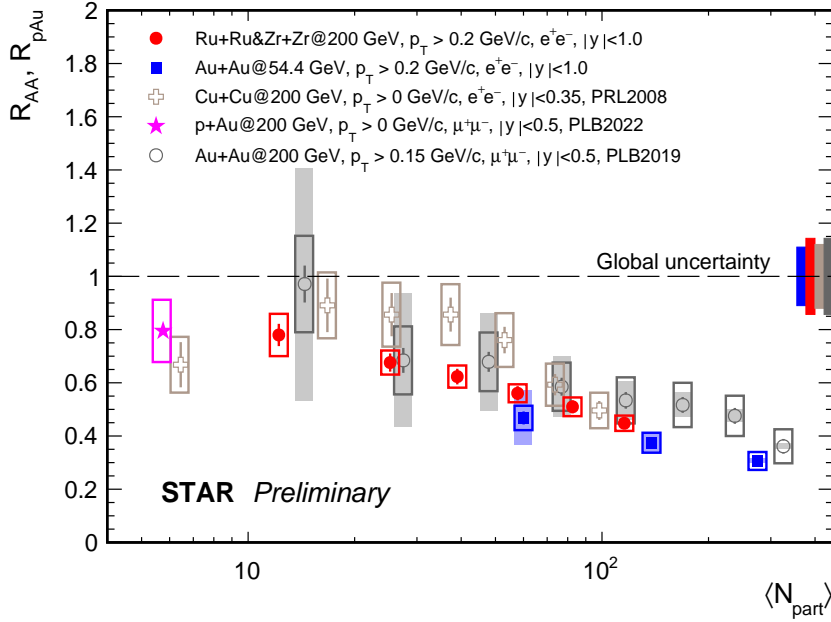
$$v_2^{SP} = \frac{\langle Q_{2,J/\psi} Q_{2,EPD}^* \rangle}{\sqrt{\frac{\langle Q_{2,EPD} Q_{2,TPCW}^* \rangle \langle Q_{2,EPD} Q_{2,TPCE}^* \rangle}{\langle Q_{2,TPCW} Q_{2,TPCE}^* \rangle}}} \quad (1)$$

47 where  $Q_{2,TPCW}$  and  $Q_{2,TPCE}$  are the 2nd harmonic event flow vectors measured in the TPC.  
48 The *J/ψ* candidate is constructed using TPC, while  $Q_{2,EPD}$  is constructed using the EPD. The large  
49 pseudorapidity ( $\eta$ ) gap between the TPC and EPD significantly reduces non-flow contributions to  
50 our final  $v_2$  results.



**Figure 1:**  $R_{AA}$  is measured as a function of  $p_T$  in 6 centralities of isobaric collisions, with the most central range in (a) and the most peripheral in (d). The statistical uncertainties are represented by the error bars, while the systematic uncertainties are denoted by the boxes. The bands around unity indicate the uncertainties originating from the  $T_{AA}$  and the  $p+p$  baselines. The results are compared to similar measurements in Au+Au and Cu+Cu collisions at  $\sqrt{s_{NN}} = 200$  GeV.

51 Figure 1 displays the nuclear modification factor  $R_{AA}$  as a function of  $p_T$  in different centralities  
52 compared with Au+Au and Cu+Cu [10, 11] results at  $\sqrt{s_{NN}} = 200$  GeV with a comparable  $\langle N_{part} \rangle$   
53 range. To establish the  $p+p$  baseline, measurements from both the STAR and PHENIX experiments  
54 are combined [12, 13]. The systematic uncertainties are represented by the boxes, dominated by

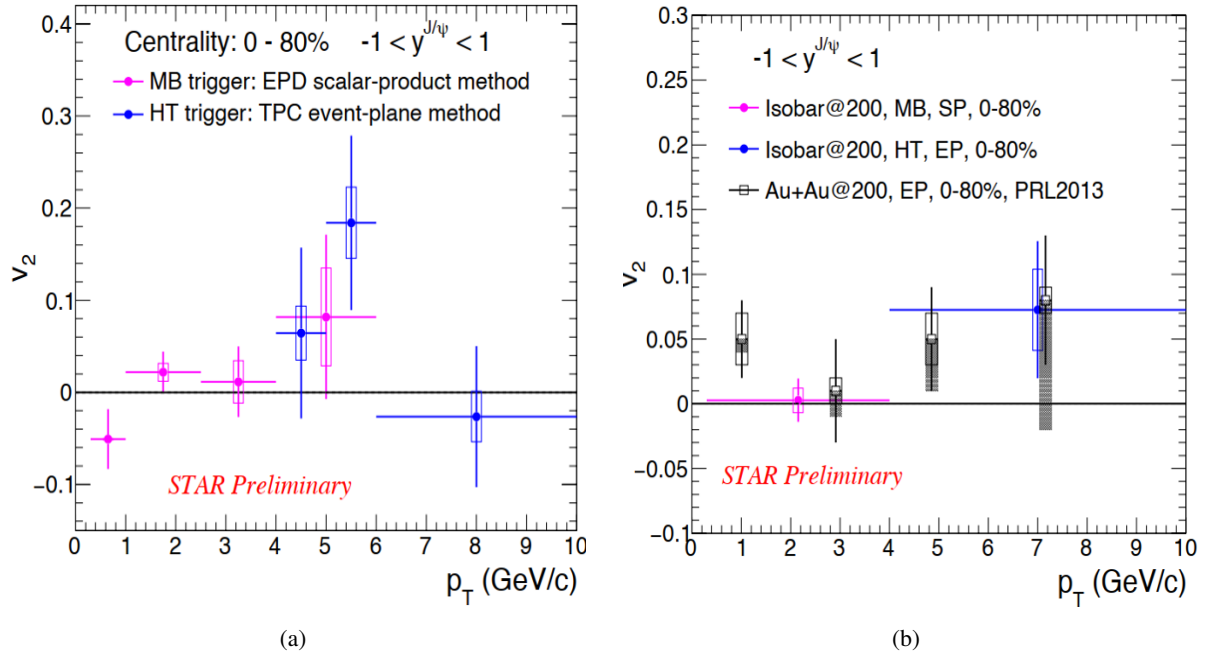


**Figure 2:**  $J/\psi$   $R_{AA}$  is measured as a function of  $\langle N_{part} \rangle$ . The error bars represent the statistical uncertainties, while the boxes represent the systematic uncertainties. The shaded bands on the data points indicate the uncertainties from the nuclear overlap function  $\langle T_{AA} \rangle$ . The bands around unity indicate the uncertainties from the  $p+p$  baselines.

55 electron matching and identification uncertainties. The main feature of these measurements is that  
 56 strong suppression of the  $J/\psi$  yield is observed at central and in semi-central collisions, and no  
 57 significant  $p_T$  dependence is measured in 6 centralities. In Fig. 1(a), a comparison is made between  
 58 the 10 to 20% centrality in isobaric collisions and the 20 to 40% centrality in Au+Au collisions,  
 59 considering the similar  $\langle N_{part} \rangle$  values. The  $R_{AA}$  values in isobaric collisions are consistent with  
 60 those in Au+Au collisions, but with noticeably higher precision. Similarly, in Fig. 1(b), the  $R_{AA}$   
 61 in the 20 to 30% centrality in isobaric collisions is compared with the 0 to 20% centrality in  
 62 Cu+Cu collisions. The isobar results, characterized by their high precision, are also consistent  
 63 with the results obtained from Cu+Cu collisions. Finally, Fig. 1(c) and Fig. 1(d) demonstrate the  
 64 comparison between more peripheral isobar collisions and Au+Au collisions. Overall, the  $R_{AA}$  in  
 65 isobaric collisions is consistent with those in Au+Au and Cu+Cu collisions at comparable  $\langle N_{part} \rangle$   
 66 ranges.

67 The  $R_{AA}$  as a function of  $\langle N_{part} \rangle$  is shown in Fig. 2. The red filled circles represent the isobaric  
 68 collisions at  $\sqrt{s_{NN}} = 200$  GeV, the blue squares correspond to Au+Au collisions at  $\sqrt{s_{NN}} = 54.4$  GeV,  
 69 the crosses denote Cu+Cu collisions at  $\sqrt{s_{NN}} = 200$  GeV, and the star represents  $p$ +Au collisions at  
 70  $\sqrt{s_{NN}} = 200$  GeV. Significant suppression is observed in the larger  $\langle N_{part} \rangle$  range. All of these results  
 71 consistently demonstrate a similar trend, with the degree of suppression remaining approximately  
 72 constant for a given  $\langle N_{part} \rangle$ . This observation indicates that no significant dependence on collision  
 73 system and energy is observed at RHIC energies at similar  $\langle N_{part} \rangle$ .

74 Figure 3(a) shows  $v_2$  as a function of  $p_T$  at 0 to 80% centrality. The red markers represent



**Figure 3:**  $J/\psi$   $v_2$  as a function of  $p_T$  in the 0-80% centrality at isobaric collisions. The error bars represent the statistical uncertainties, while the boxes represent the systematic uncertainties. The results are compared to Au+Au collisions at  $\sqrt{s_{NN}} = 200$  GeV, where the shaded bands on the open squares indicate the maximum non-flow.

75  $v_2$  obtained from Minimum Bias (MB) triggers using the Scalar-Product method, while the blue  
76 markers correspond to BEMC high-tower trigger (HT) results utilizing the event-plane method.  
77 These results are consistent with each other within the overlap range. The error bars represent the  
78 statistical uncertainties, while the boxes indicate the systematic uncertainties, which are dominated  
79 by electron identification and signal extraction uncertainties. No significant  $v_2$  for the  $J/\psi$  is observed  
80 at the current level of precision. Furthermore, the result of integrating  $p_T$  over the ranges 0 to 4  
81 GeV and 4 to 10 GeV is shown in Fig. 3(b). In the low  $p_T$  region, the most precise measurement of  
82  $v_2$  at RHIC to date is obtained and indicated zero  $J/\psi$   $v_2$ . This hints at a limited regeneration effect  
83 or small charm quark flow in isobaric collisions.

### 84 3. Conclusions

85 In this contribution, high-precision  $J/\psi$   $R_{AA}$  and  $v_2$  measurements are presented in isobaric  
86 collisions. We observe significant suppression of  $J/\psi$  in isobaric collisions at  $\sqrt{s_{NN}} = 200$  GeV.  
87 No significant energy and colliding system size dependence of  $J/\psi$   $R_{AA}$  at RHIC at similar  $\langle N_{part} \rangle$   
88 is observed. The  $J/\psi$   $v_2$  is consistent with zero at the current precision. This hints at a limited  
89 regeneration effect or small charm quark flow in isobaric collisions, and strong dissociation effect  
90 in central isobaric collisions.

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