



***Azimuthal distributions of high-pt  
direct  $\gamma$  and  $\pi^0$  w.r.t reaction plane  
at STAR***

*Ahmed Hamed*

*For the  Collaboration*

***Hot Quarks 2010***

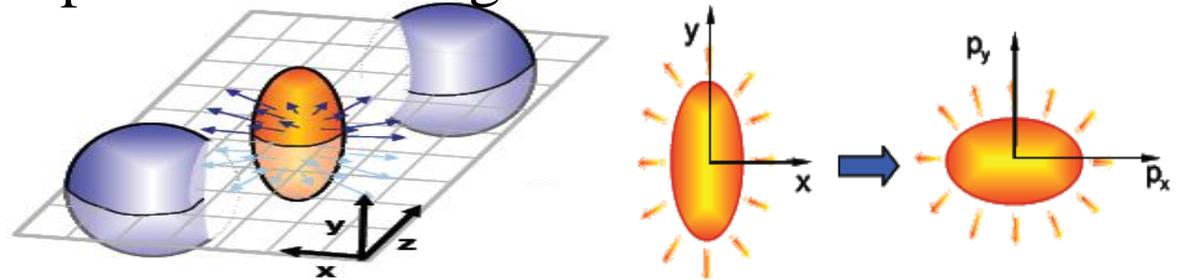
***La Londe Les Maures, 21-26th June, 2010***

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## Elliptic flow at low $p_T$

- The azimuthal distributions of the produced particles in heavy-ion collisions are considered to be sensitive to the initial geometric overlap of the colliding nuclei.

$$\varepsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$



$$dN/d\phi \propto 1 + \sum_n 2v_n \cos n(\phi - \Psi_{RP})$$

- Eccentricity in spatial coordinate is preserved and mapped into the momentum coordinates **if the produced particles are not freely streaming.**
- At low  $p_T$  the produced particles are not freely streaming and the collectivity is built even before hadronization.

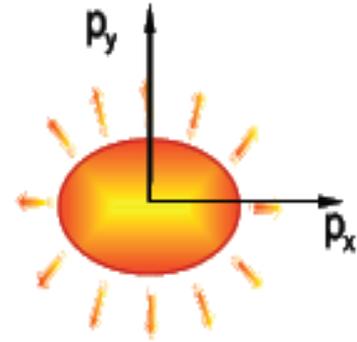
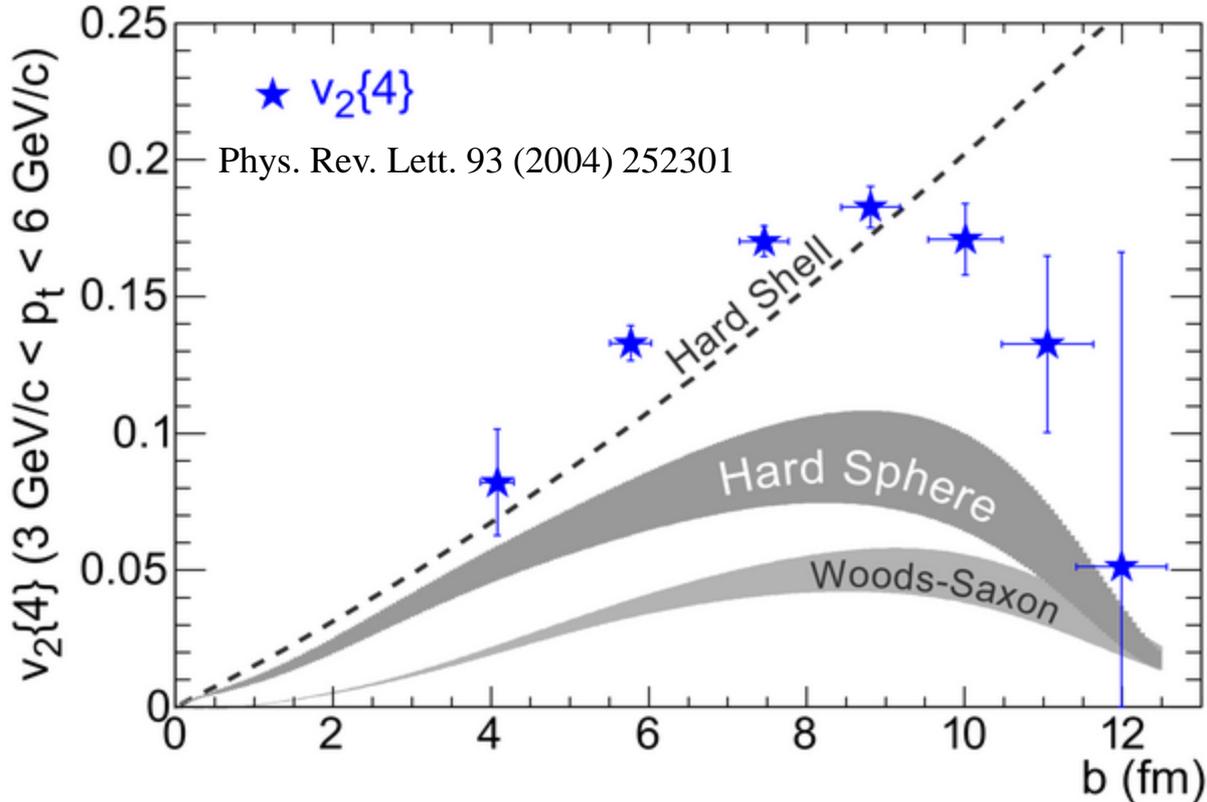
### Non-flow

$$\left\langle \sum_i \cos 2(\phi_{p_t} - \phi_i) \right\rangle = M v_2(p_t) \bar{v}_2 + \{\text{non-flow}\}$$

# Elliptic flow at high $p_T$

$V_2$  at high  $p_T$  is finite positive!

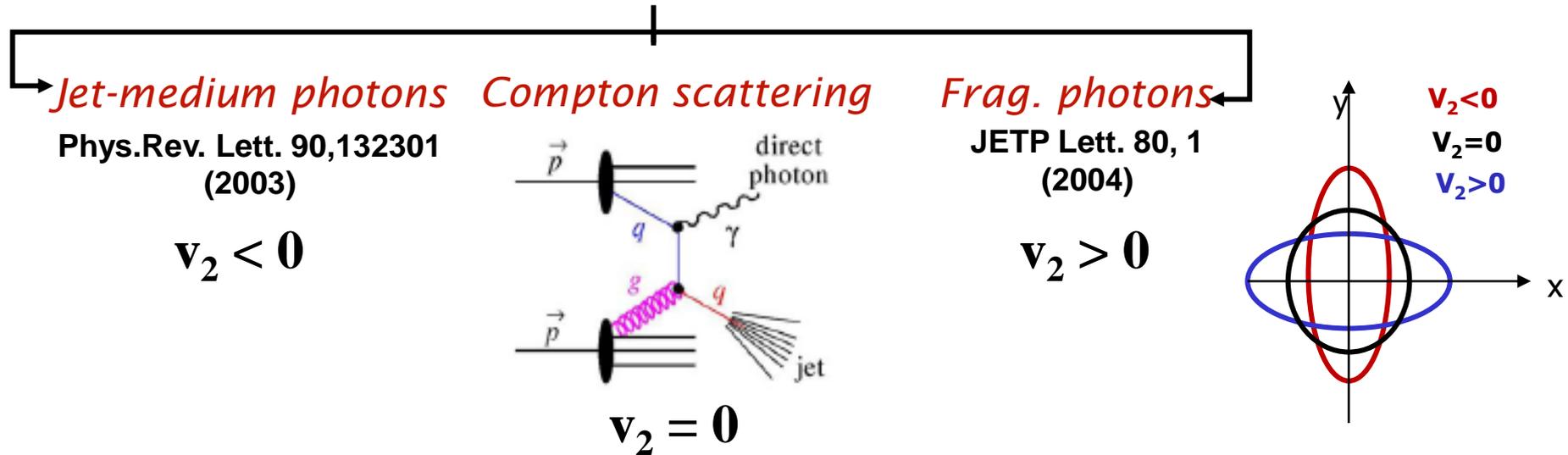
Jet quenching : energy loss dependence of path length



- Azimuthal anisotropy at large  $p_T$  seem to be too large for a pure “jet quenching” Phys. Rev. C. 66, 027902 (2002)
- Surface emission is not consistent with the measured value of elliptic flow at high  $p_T$  Phys. Rev. Lett. 93 (2004) 252301

# Why elliptic flow of direct photons?

- $v_2$  of electromagnetically interacting particles.
- $v_2$  measurements at higher  $p_T$ .
- Production mechanisms of photons



## ➤ Path length dependence of parton energy loss.

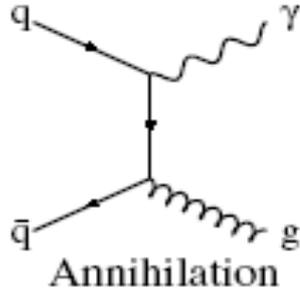
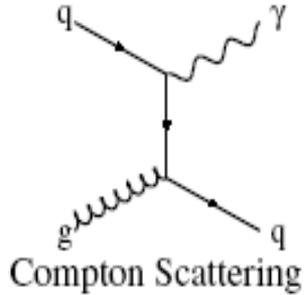
$v_2 < 0$  : Particles preferred to traverse through the longer path “out-of-plane”

$v_2 = 0$  : No preferred direction w.r.t. reaction plane

$v_2 > 0$  : Particles preferred to traverse through the shorter path “in-plane”

# LO production of direct $\gamma$ and fragmentation $\gamma$

## Direct photons



➤ Very challenging measurements due to the S/B ratio,  $\pi^0$  is the major source of bg.

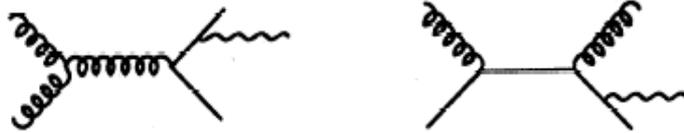
➤ The Compton-scattering process

$$\rightarrow \gamma/\pi^0 > \alpha_{em}$$

## Fragmentation photons

Fragmentation photons  $\gamma^{frag}$

$\gamma^{frag}$  seems to be accompanied by additional hadrons.



Example of Bremsstrahlung diagrams

$$zD_{\gamma/q}(z, Q^2) = e_q^2 \frac{\alpha}{\gamma_\pi} [1 + (1-z)^2] \ln(Q^2/\Lambda^2) \sim \alpha_{em}/\alpha_s$$

$$zD_{\gamma/g}(z, Q^2) = 0$$

❑ The sub-process of  $\gamma^{frag}$  is of order of  $O(\alpha_s^2)$  but its yield is comparable to  $\gamma^{dir}$  LO process  $O(\alpha_s \alpha_{em})$ .

❑ The  $\gamma^{frag}$  contribution is expected to fall off more rapidly in  $x_T$  than the other lowest order of  $\gamma^{dir}$ . (G. Sterman et al. Rev. Mod. Phys. 67, 157 (1995))

❑  $\gamma^{frag} / \gamma^{dir} \sim 30-40\%$  at  $p_T^\gamma > 8 \text{ GeV}/c$  at mid-rapidity at RHIC energy. D. De Florian and W. Vogelsang, Phys. Rev. D72, 014014 (2005)



# Methods of $\gamma^{dir}$ measurements

## ✓ Standard Method

1. Measure inclusive photons.
2. Reconstruct other sources of photons “hadrons”!
3. Subtract photons from decay of  $\pi^0$ ,  $\eta$  etc.

PHENIX is well-adapted for this method due to the calorimeter granularity and the distance between the calorimeter to the interaction point  $\rightarrow \pi^0$  reconstruction in central Au+Au up to  $p_T \sim 20$  GeV/c

- Limited at very high  $p_T$ , effective method for both symmetric and asymmetric hadron decays

## ✓ Transverse Shower Profile Method

STAR is well-suited for the transverse shower shape analysis due to the Shower Maximum Detector  $\rightarrow \gamma/\pi^0$  discrimination up to  $p_T \sim 26$  GeV/c. M. Beddo et al., Nucl. Instrum. Meth. A499, 725 (2003)

- Effective at very high  $p_T$ , but limited only for the symmetric hadron decays



# Analysis technique

1. Electromagnetic neutral cluster ( $\pi^0, \eta, \rho^0, \omega, \dots, \gamma^{\text{frag}}, \gamma^{\text{dir}}$ )

2. Reaction plane measurements

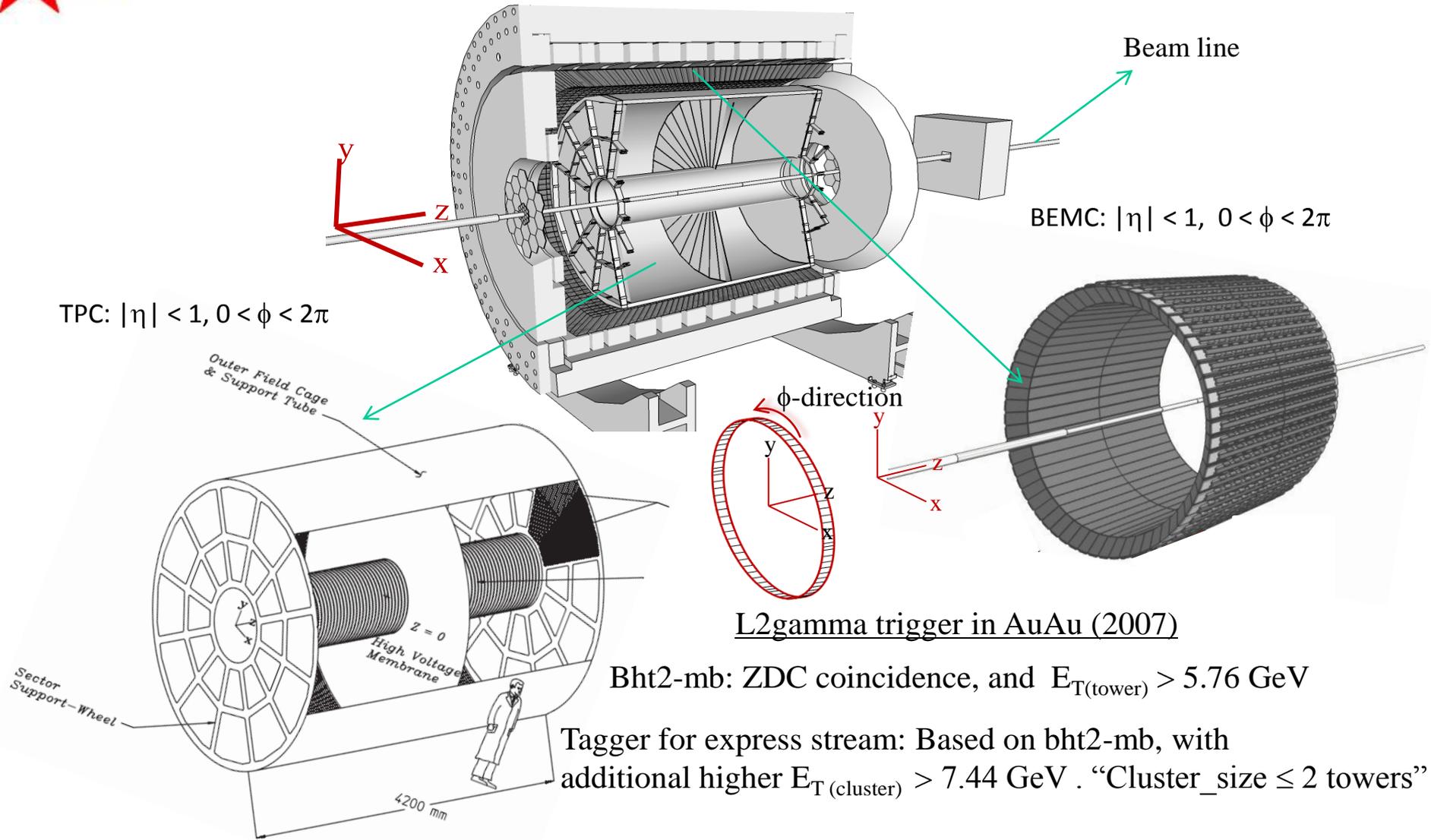
3.  $v_2$  neutral cluster vs.  $v_2$  charged particles

4. Transverse shower profile to obtain sample rich/free of  $\gamma^{\text{dir}}$

5. Obtain  $v_2$  of  $\gamma^{\text{dir}}$



# STAR detector and on-line $\gamma$ -rich event selections



L2gamma trigger in AuAu (2007)

Bht2-mb: ZDC coincidence, and  $E_{T(\text{tower})} > 5.76 \text{ GeV}$

Tagger for express stream: Based on bht2-mb, with additional higher  $E_{T(\text{cluster})} > 7.44 \text{ GeV}$ . "Cluster\_size  $\leq 2$  towers"

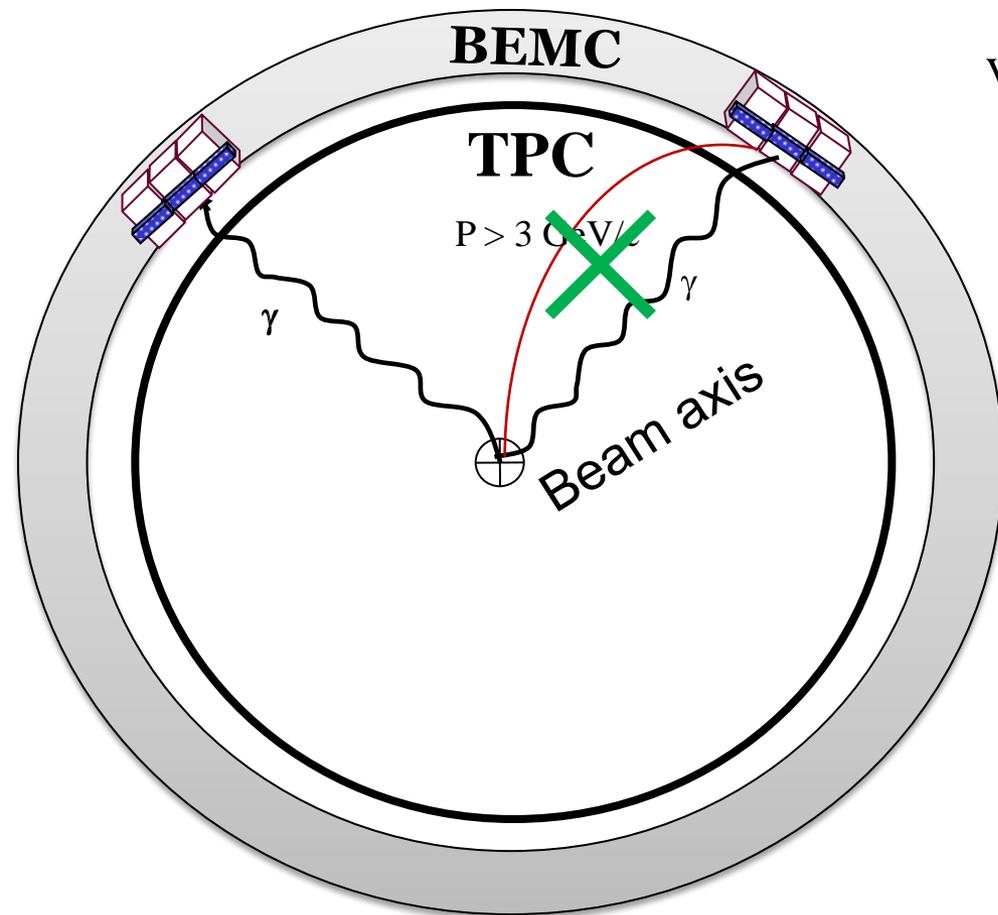
1  $\gamma$ -triggered event each 5k minbias event  $\rightarrow \sim 500 \mu\text{b}^{-1}$  of AuAu 2007 @ 200GeV

6k events of minimum bias trigger



# STAR detector and off-line neutral cluster selections

- ✓ Select neutral clusters “triggers” (BEMC-BSMD) using charged-particle veto (TPC)



vertex within  $\pm 55$  cm of the center of TPC.

At least one cluster with  $E_T > 8$  GeV,  
 $E_{\text{smd}\eta} > 0.5$  GeV,  $E_{\text{smd}\phi} > 0.5$  GeV,  
and no track with  $p > 3$  GeV/c  
pointing to that cluster.

In Au+Au: 28% of the integrated  
luminosity  
has  $E_T > 8$  GeV of which **96.5%** left at  
least 0.5 GeV on each planes of  
SMD of which **93%** has no track  
with  $p > 3$  GeV/c pointing to it.

## Event plane from TPC

$$\psi = \frac{1}{2} \tan^{-1} \left( \frac{\sum_i \sin(2\delta_i)}{\sum_i \cos(2\delta_i)} \right)$$

## Shift method for event plane flattening

$$\Psi' = \Psi + \sum_n \frac{1}{n} [-\langle \sin(2n\Psi) \rangle \cos(2n\Psi) + \langle \cos(2n\Psi) \rangle \sin(2n\Psi)]$$

## Sub-event method for reaction plane resolution

$$\sigma_{\text{RP}} = C \sqrt{\langle \cos[2(\psi^A - \psi^B)] \rangle}, C = \sqrt{2}$$

## $v_2$ of charged and neutral particles

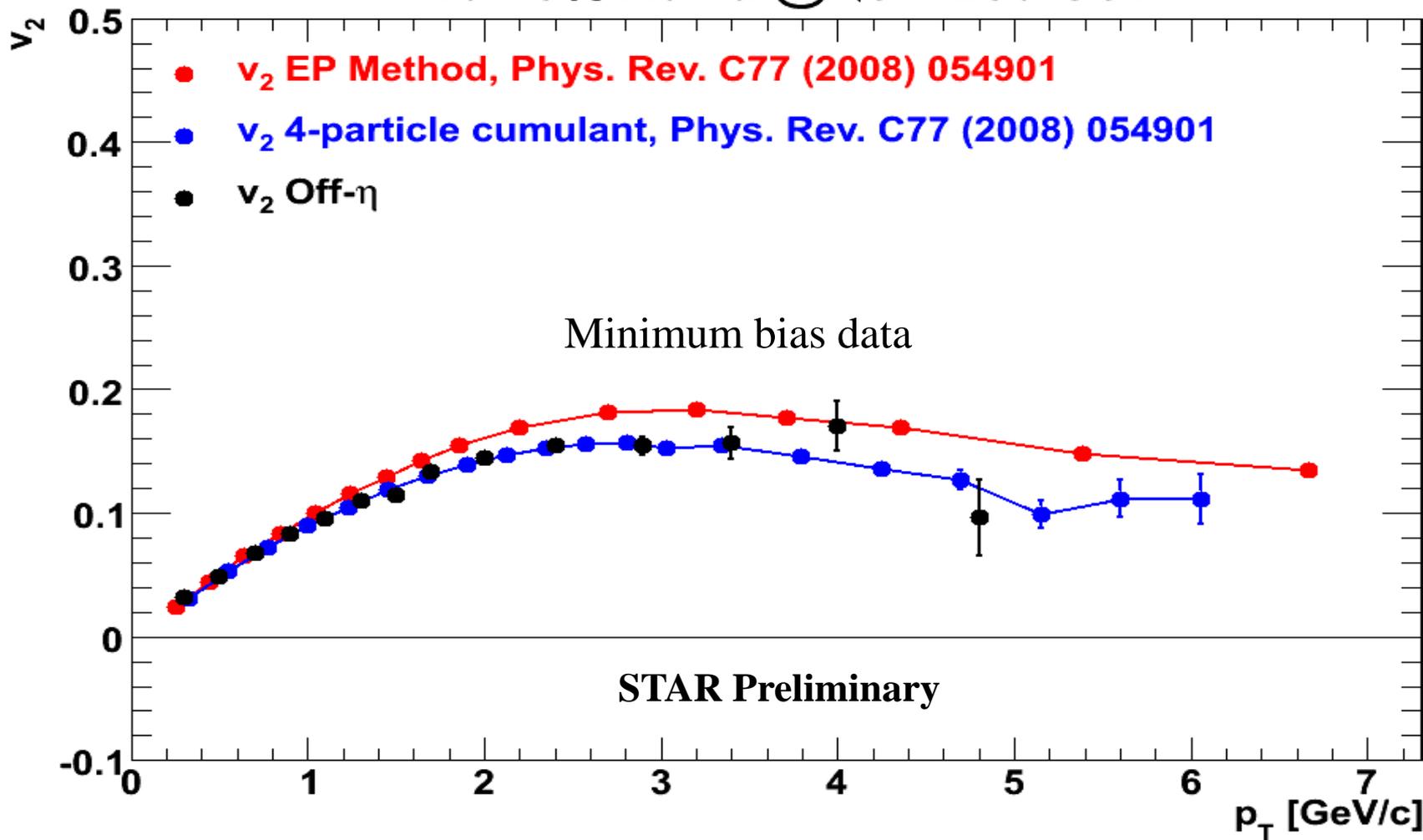
$$v_2^{\text{track,obs}} = \langle \cos(2\phi_{\text{track}} - 2\psi) \rangle$$

$$v_2^{\text{neutral,obs}} = \langle \cos(2\phi_{\text{tower}} - 2\psi) \rangle$$

$$v_2 = \frac{v_2^{\text{obs}}}{\sigma_{\text{RP}}}$$

# Elliptic flow of charged particles at low $p_T$

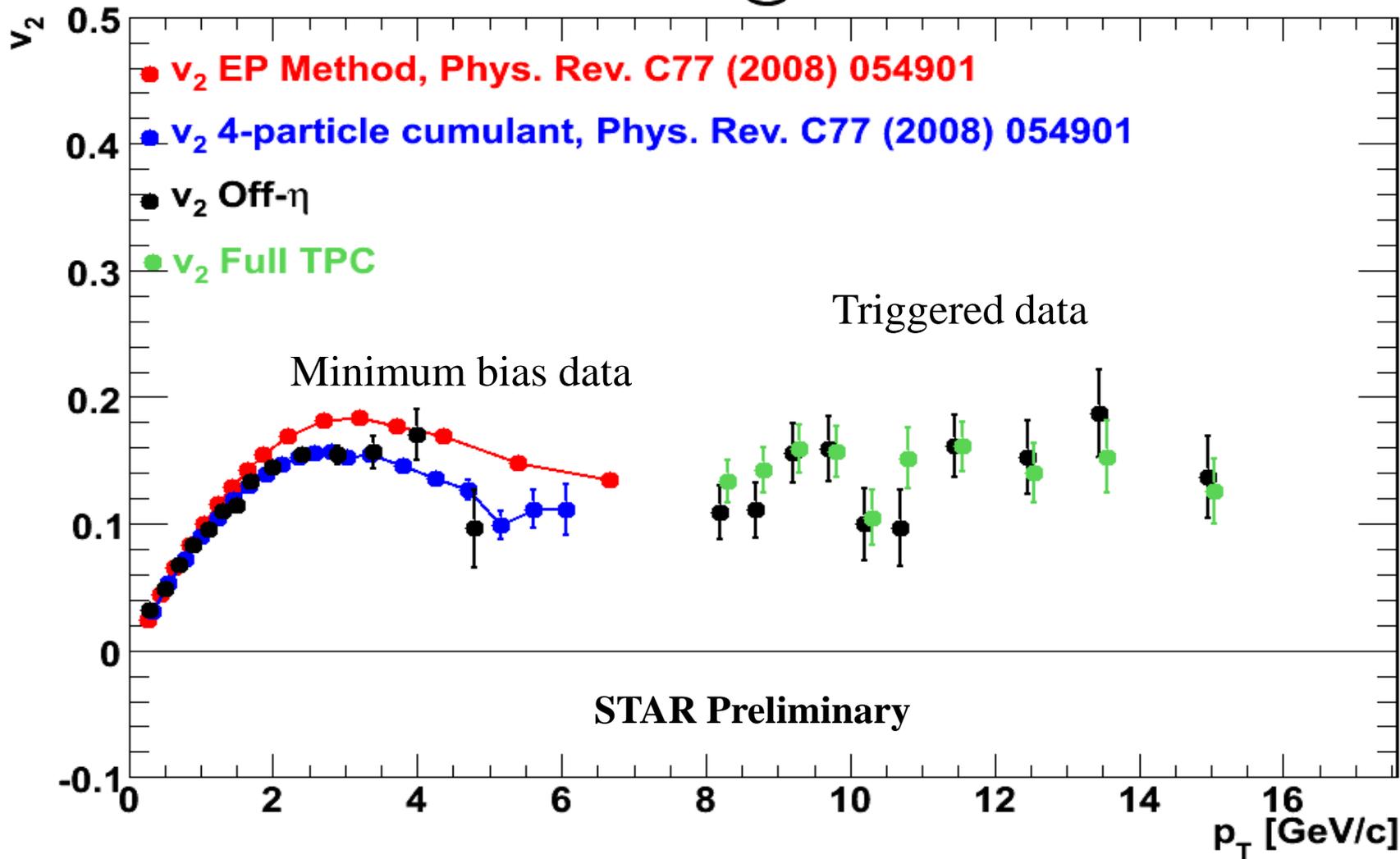
10-40% AuAu @  $\sqrt{s} = 200$  GeV



$v_2(\text{EP off-}\eta)$  reproduces the  $v_2\{4\}$  quite well

# Elliptic flow of charged particles at high $p_T$

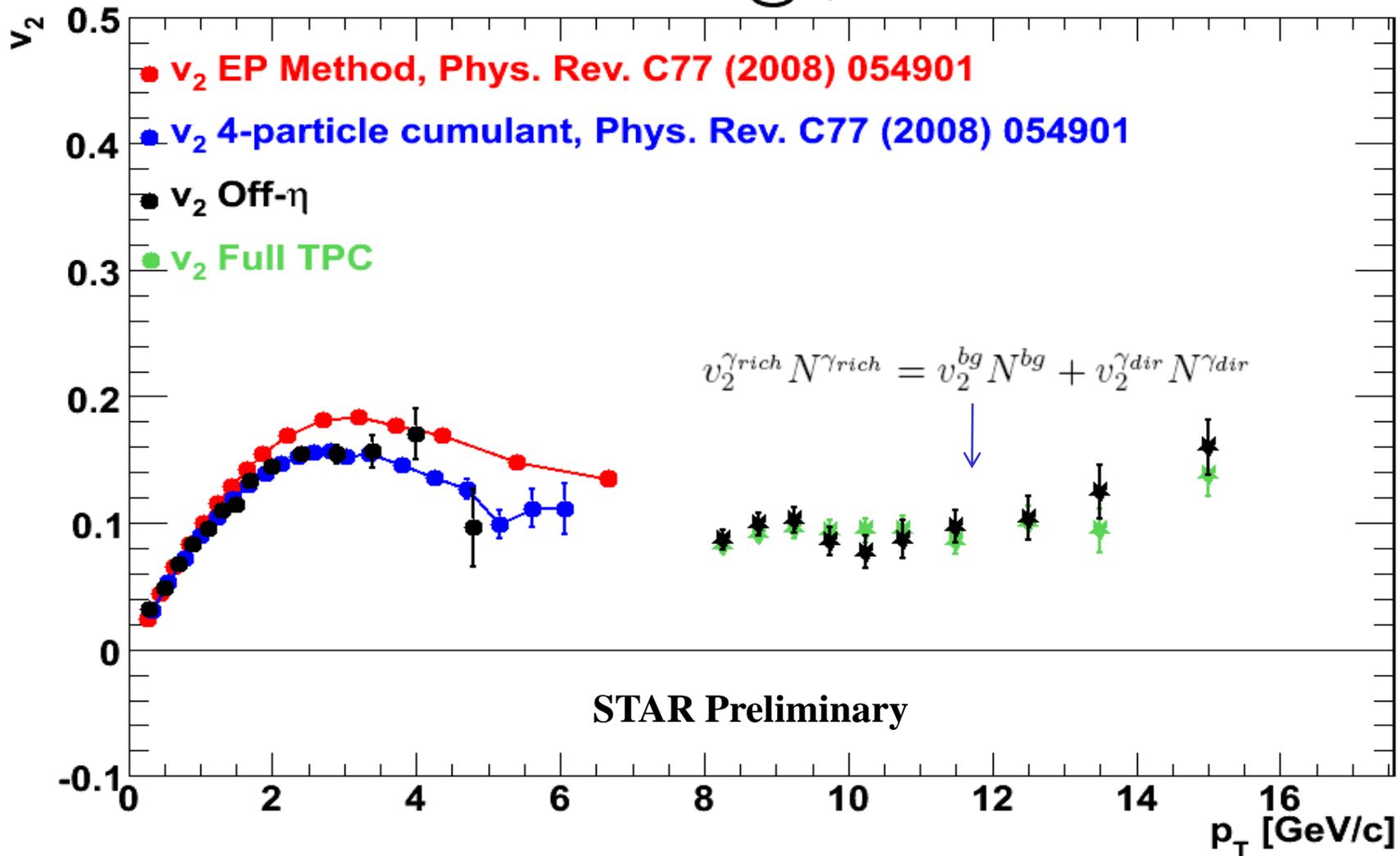
10-40% AuAu @  $\sqrt{s} = 200$  GeV



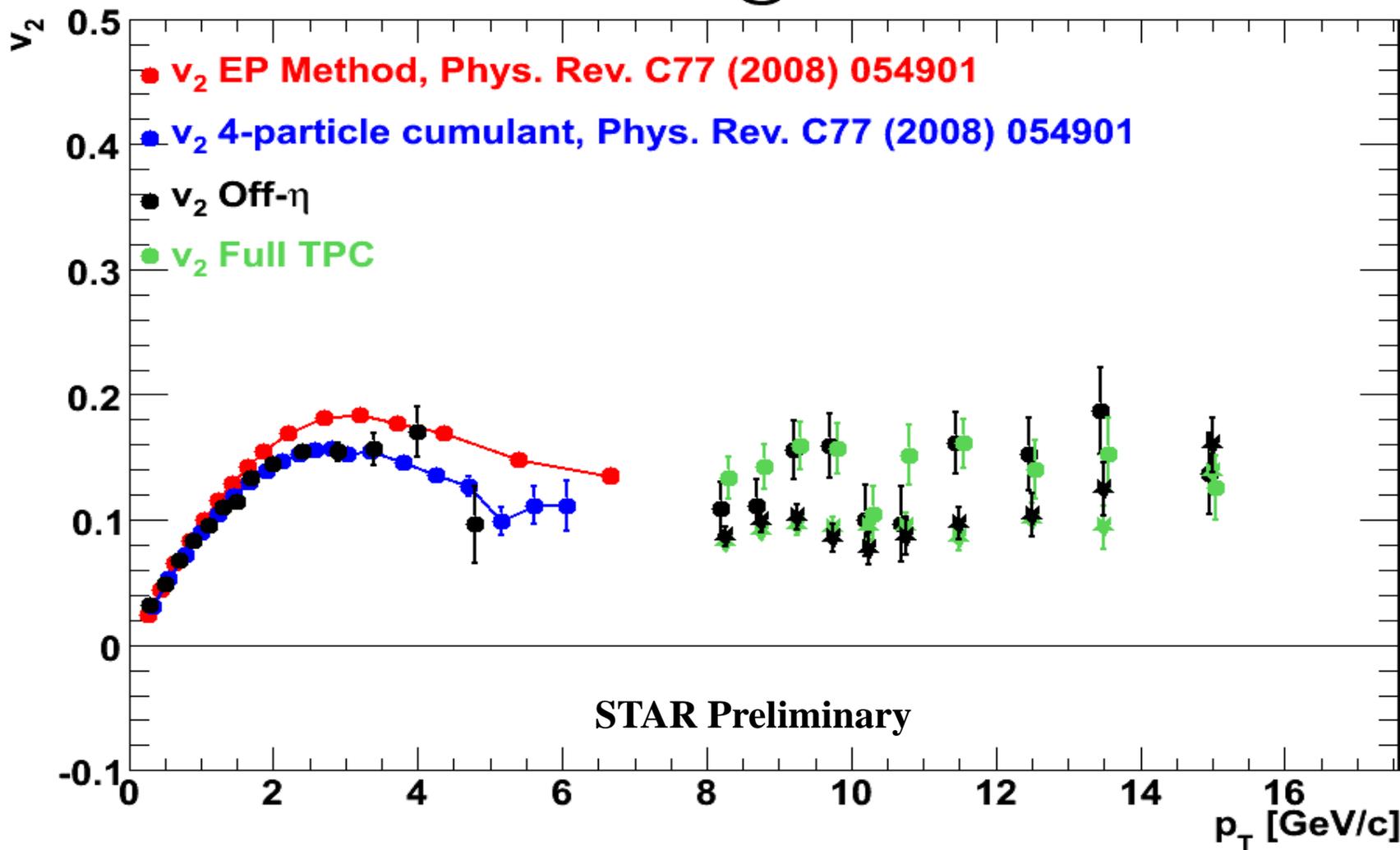
$v_2$  of charged particles is  $\sim 15\%$  in 10-40% (AuAu@200GeV) and constant in  $p_T$  (8-16GeV/c)

# Elliptic flow of neutral particles at high pt

10-40% AuAu @  $\sqrt{s} = 200$  GeV



$v_2$  of neutral particles is  $\sim 10\%$  in 10-40% (AuAu@200GeV) and constant in pt (8-16GeV/c)

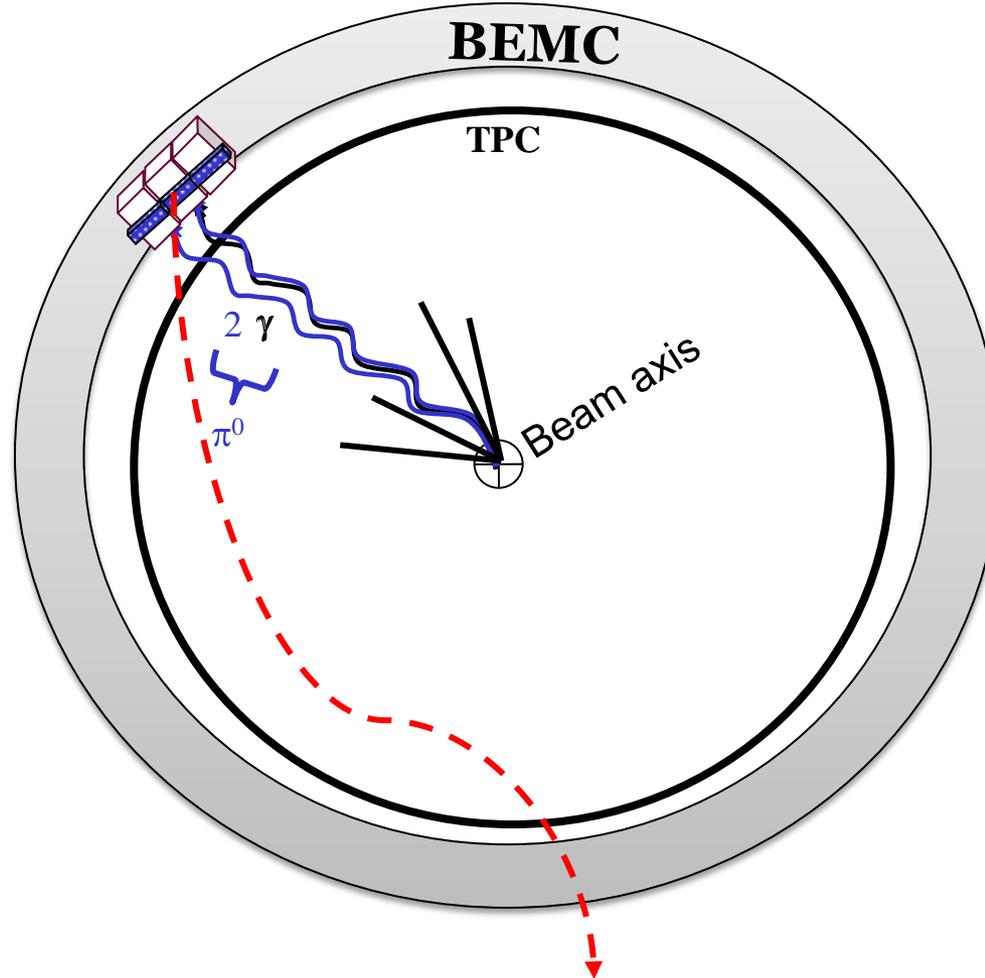


$v_2$  of neutral particles is less than  $v_2$  of charged particles due to direct photons contributions

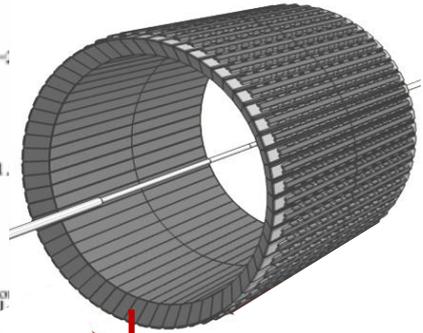
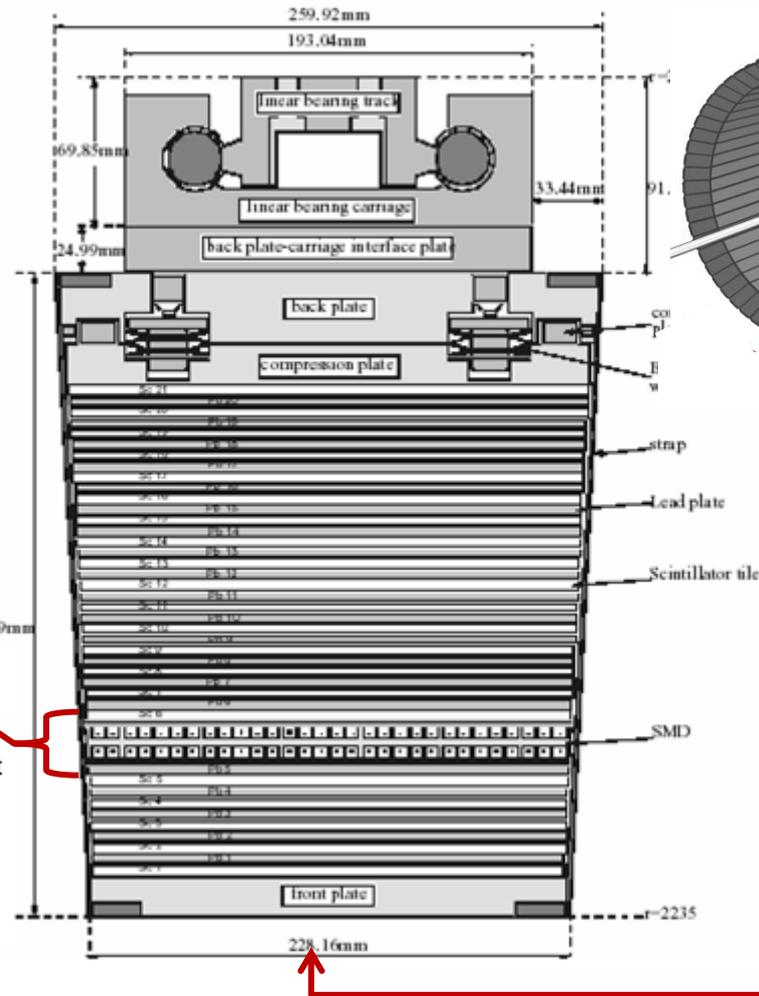
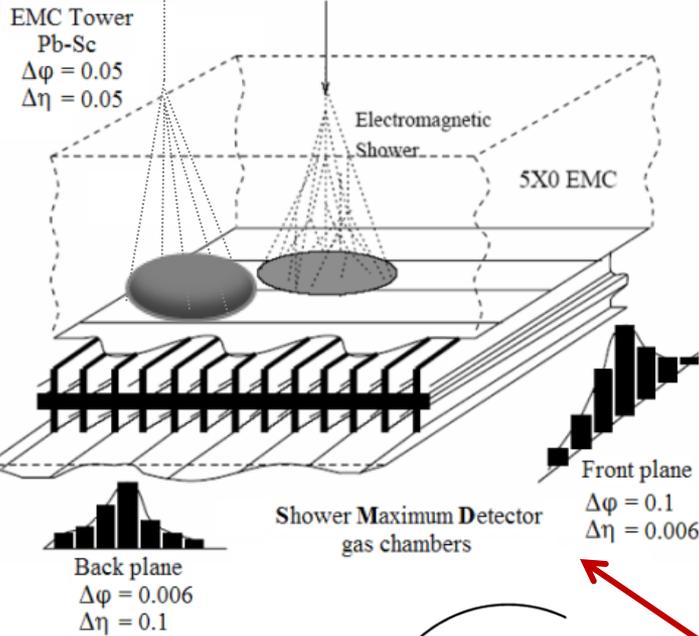
# How to separate $\gamma^{\text{dir}}$ from neutral $0\gamma$ .

$\sim 10\%$  of all  $\pi^0$  (8-16 GeV/c) decay asymmetrically with one gamma has  $p_T > 8$  GeV/c within STAR-BEMC acceptance.

$\eta$  causes similar level of background as asymmetric  $\pi^0$ .

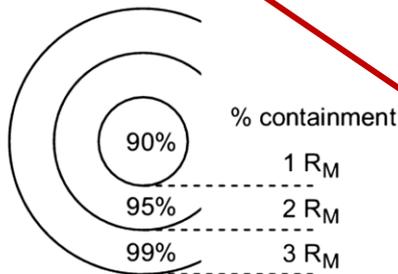


Either to reconstruct  $\pi^0$  or to use the transverse shower shape analysis to distinguish between  $\pi^0$  and  $\gamma^{\text{dir}}$



Cross section in  $\phi$

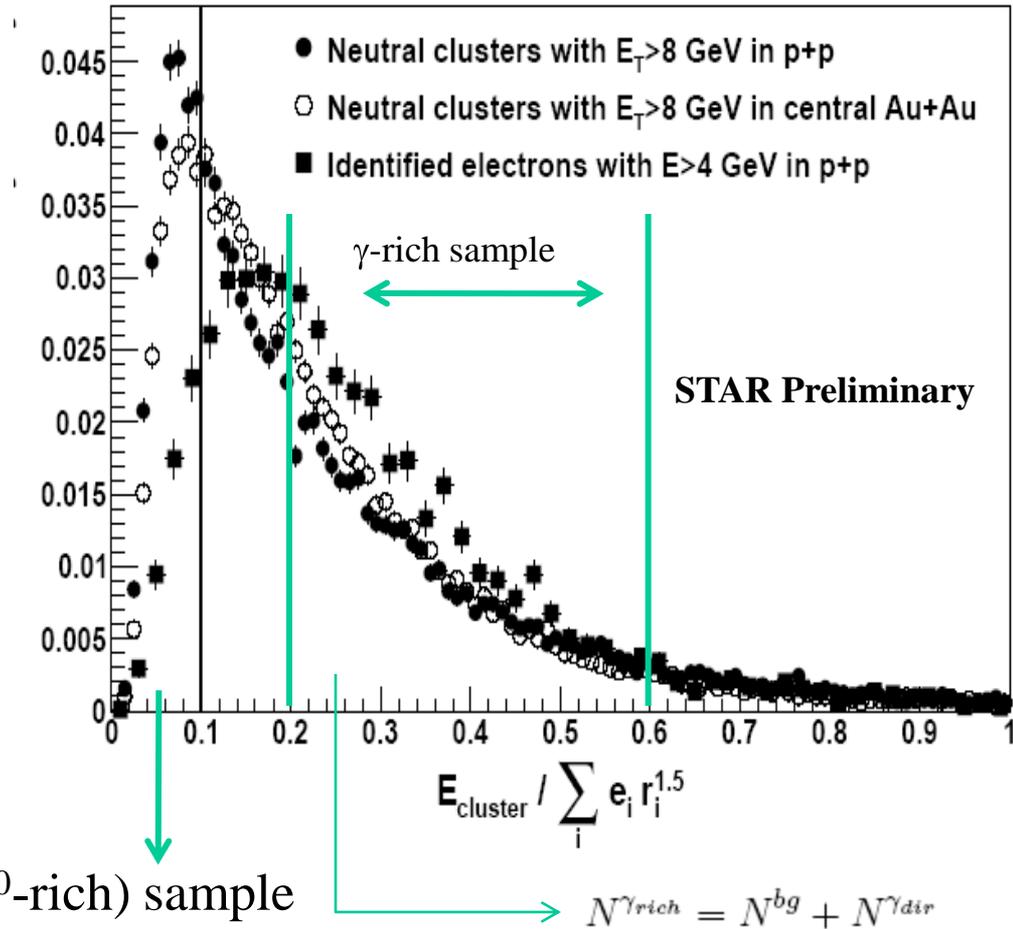
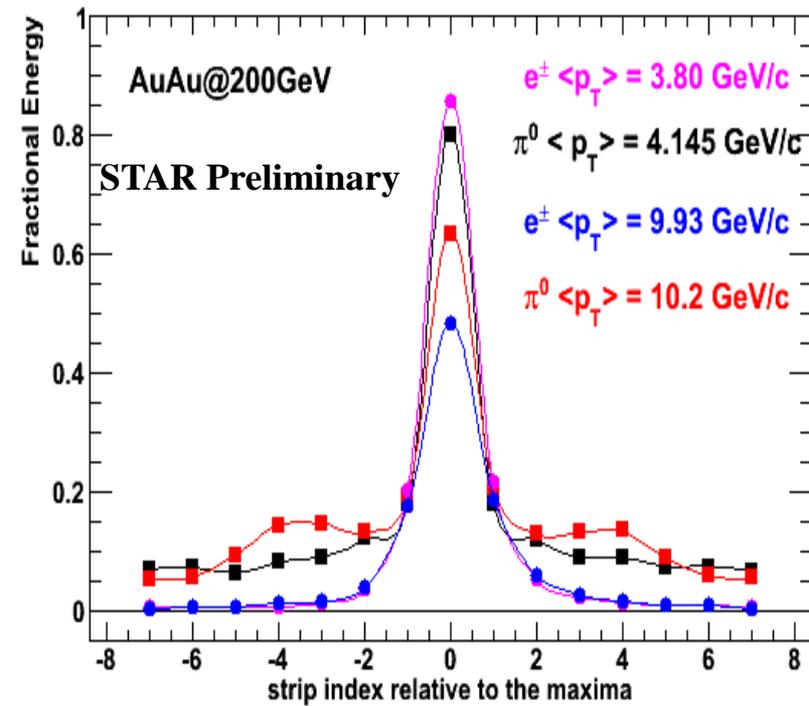
$$E_{\text{cluster}} / \sum_i e_i r_i^{1.5}$$



The two photons originated from  $\pi^0$  hit the same tower at  $p_T > 8 \text{ GeV}/c$

The shower shape is quantified with the cluster energy, measured by the BEMC, Normalized by the position-dependent energy moment, measured by the BSMD strips.

# Shower Profile of single $\gamma$ vs. two close $\gamma$ s

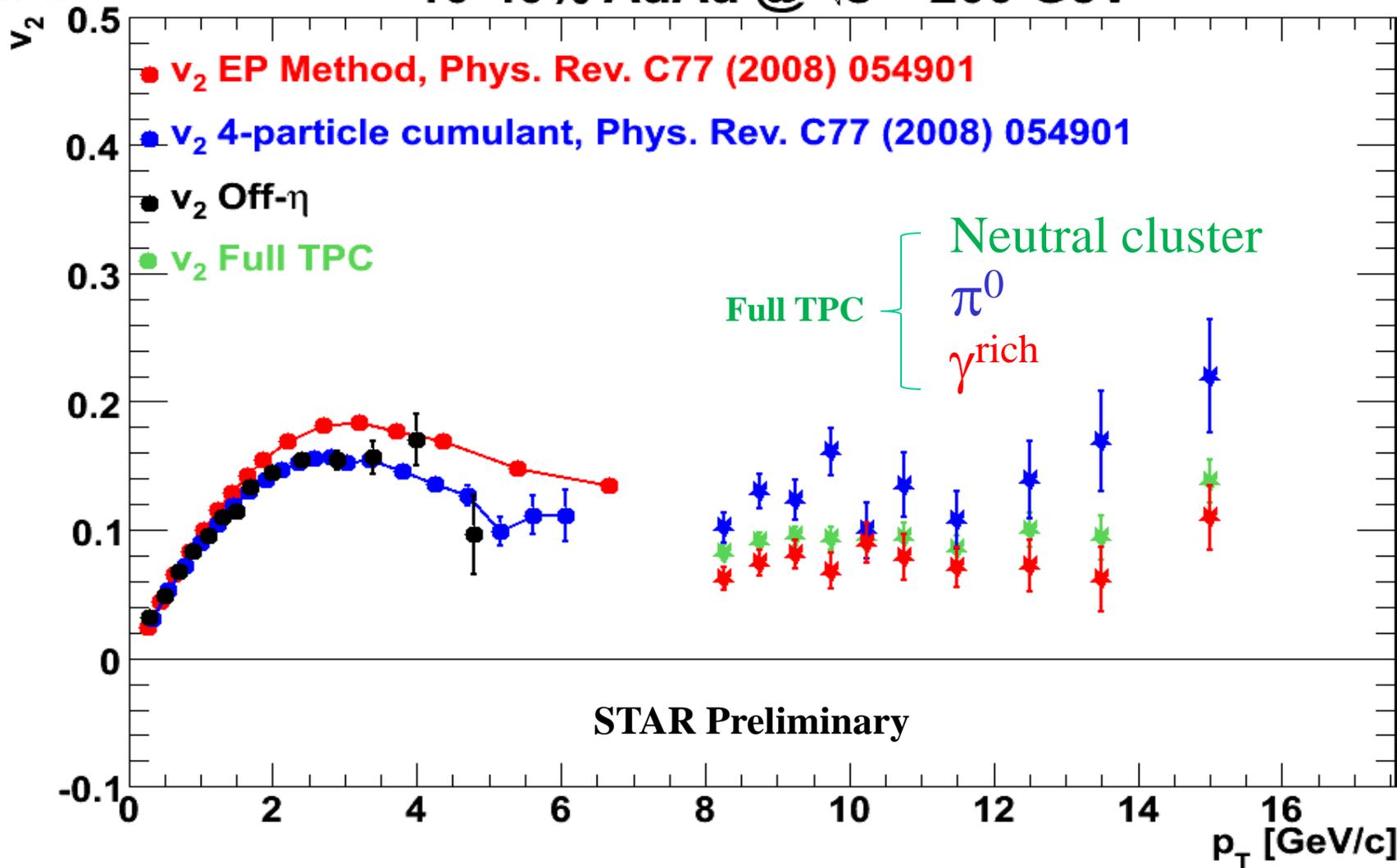


The probability distribution is peaked at smaller value in AuAu than in pp due to the larger relative fraction of  $\gamma^{\text{dir}}$ .

The rejection power of direct photons is  $\sim 90\%$

# Elliptic flow of $\pi^0$ and $\gamma^{\text{rich}}$ at high $p_T$

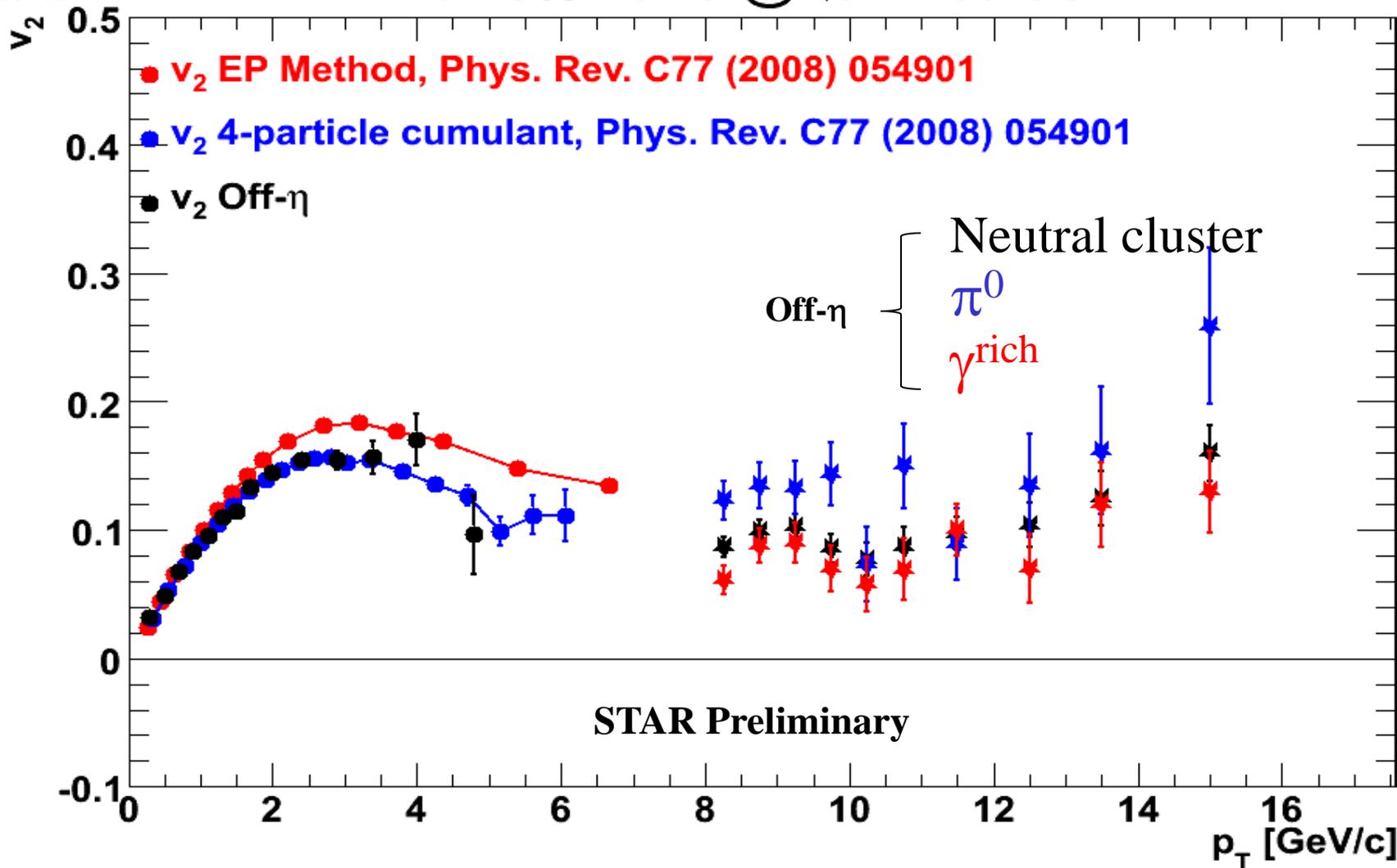
10-40% AuAu @  $\sqrt{s} = 200$  GeV



$v_2$  of  $\gamma^{\text{rich}} < v_2$  of  $\pi^0$  as expected

# Elliptic flow of $\pi^0$ and $\gamma^{\text{rich}}$ at high $p_T$

## 10-40% AuAu @ $\sqrt{s} = 200$ GeV



$v_2$  of  $\gamma^{\text{rich}} < v_2$  of  $\pi^0$  as expected

## Obtain $v_2$ of direct photons

- ⊕ Select EM neutral clusters
- ⊕ Use the transverse shower shape to select  $\gamma^{\text{dir}}$  free ( $\pi^0$ -rich) sample and  $\gamma^{\text{rich}}$  sample from the neutral clusters.

$$v_2^{\gamma^{\text{rich}}} N^{\gamma^{\text{rich}}} = v_2^{\text{bg}} N^{\text{bg}} + v_2^{\gamma^{\text{dir}}} N^{\gamma^{\text{dir}}}$$

$$\mathcal{R} = \frac{N^{\text{bg}}}{N^{\gamma^{\text{rich}}}} \simeq \frac{N^{\pi^0}}{N^{\gamma^{\text{rich}}}}$$

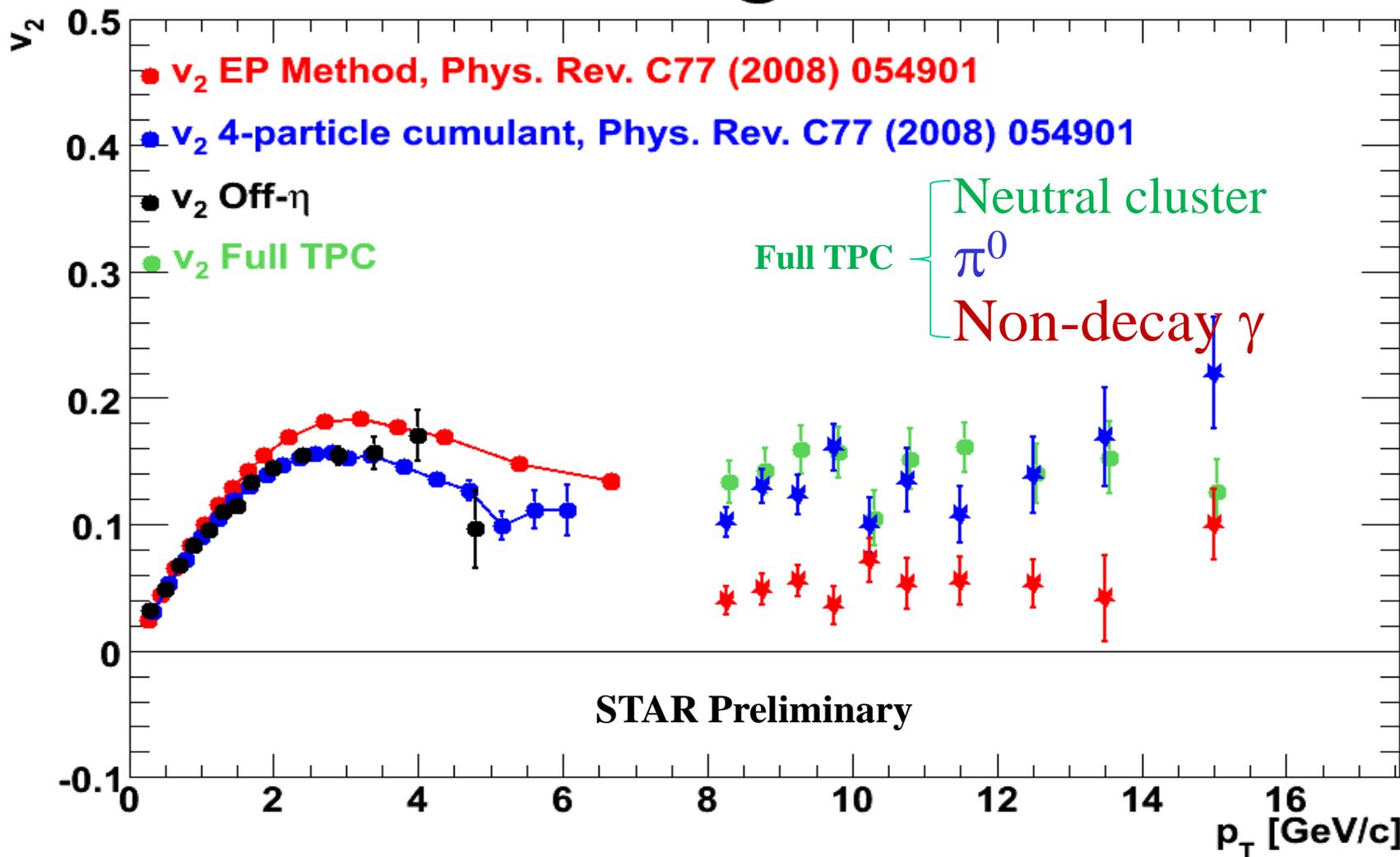
$$v_2^{\gamma^{\text{direct}}} = \frac{v_2^{\gamma^{\text{rich}}} - v_2^{\text{bg}} \mathcal{R}}{1 - \mathcal{R}}$$

$$v_2^{\gamma^{\text{direct}}} = \frac{v_2^{\gamma^{\text{rich}}} - v_2^{\pi^0} \mathcal{R}}{1 - \mathcal{R}}$$



# Elliptic flow of $\pi^0, \gamma^{dir}$ , charged particles at high $p_T$

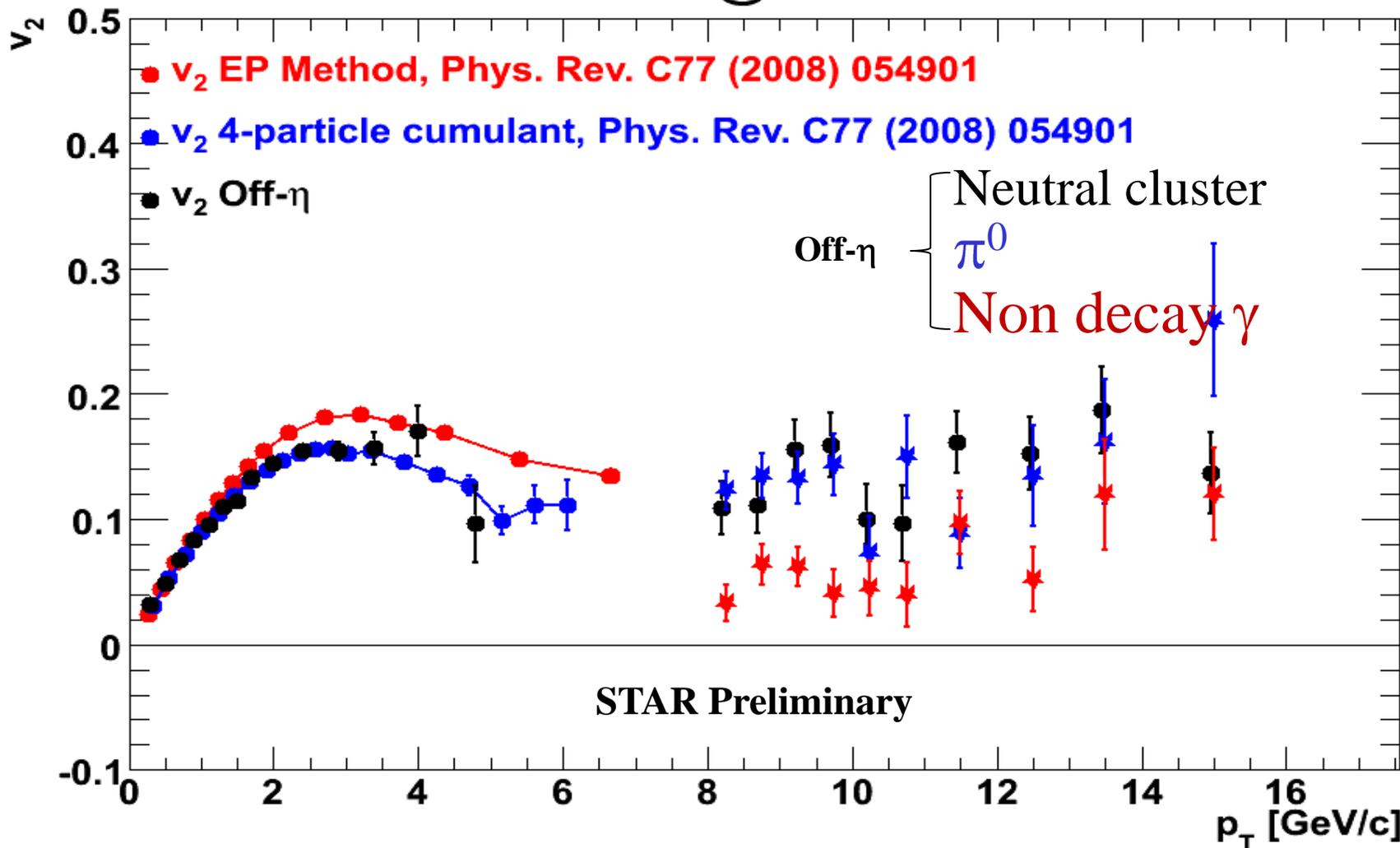
10-40% AuAu @  $\sqrt{s} = 200$  GeV



$v_2$  of non decay  $\gamma$  is  $\sim 1/3$  of  $v_2$  of  $\pi^0$  and charged particles

$v_2$  of non decay  $\gamma$  is not zero and not negative

10-40% AuAu @  $\sqrt{s} = 200$  GeV



$v_2$  of non decay  $\gamma$  is  $\sim 1/3$  of  $v_2$  of  $\pi^0$  and charged particles

$v_2$  of non decay  $\gamma$  is not zero and not negative

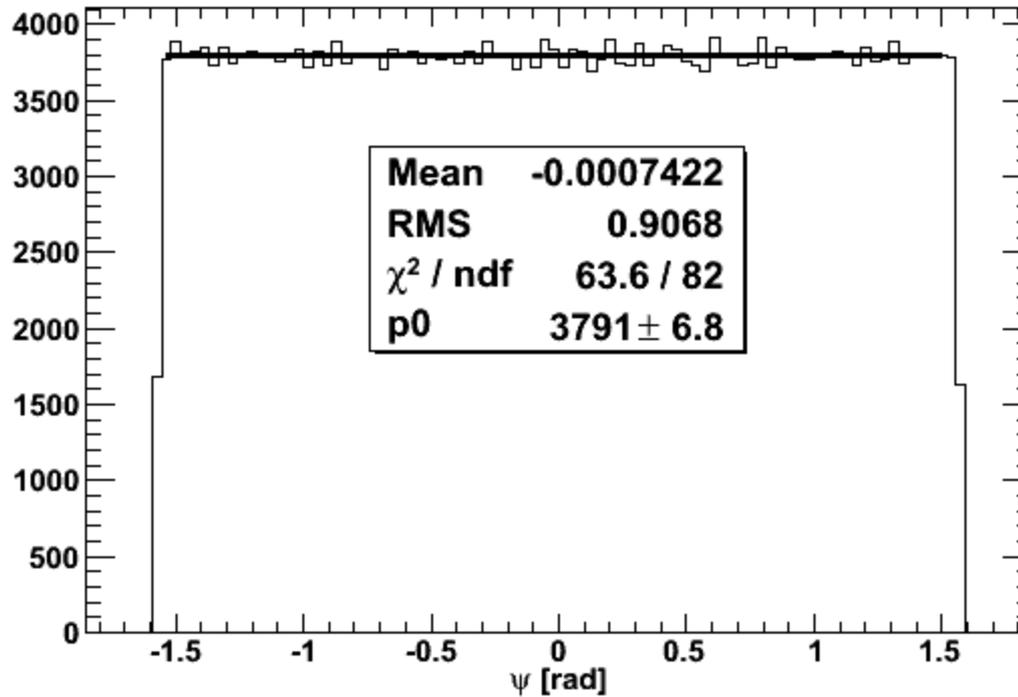


## Summary

- The geometrical effect of the medium can be probed by the elliptic flow measurement
- Finite and +ve value of  $v_2$  persist up to  $pt=16$  GeV/c for charged and neutral particles
- STAR has reported the first “preliminary” results of non decay photons elliptic flow at high  $pt$  at RHIC
  - No sign of negative  $v_2$  of non decay photons
  - Statistically significant value of +ve  $v_2$  of non decay photons
- The  $v_2$  at high  $pt$  can not be interpreted as path length dependence of energy loss

***Backup slides***

# Reaction plane



Only shift method is used to flatten the RP up to the 20<sup>th</sup> harmonic:

$$\Psi' = \Psi + \sum_n \frac{1}{n} [-\langle \sin(2n\Psi) \rangle \cos(2n\Psi) + \langle \cos(2n\Psi) \rangle \sin(2n\Psi)]$$

Reaction plane is flat “cos and sin  $\sim 0$ ”

$$v_2^{inclusive \gamma} = \frac{v_2^{direct \gamma} N_{direct \gamma} + v_2^{b.g.} N_{b.g.}}{N_{direct \gamma} + N_{b.g.}}$$

$\sigma_{RP} = \langle \cos(2(\Phi_{measured} - \Phi_{RP})) \rangle$  is 0.3

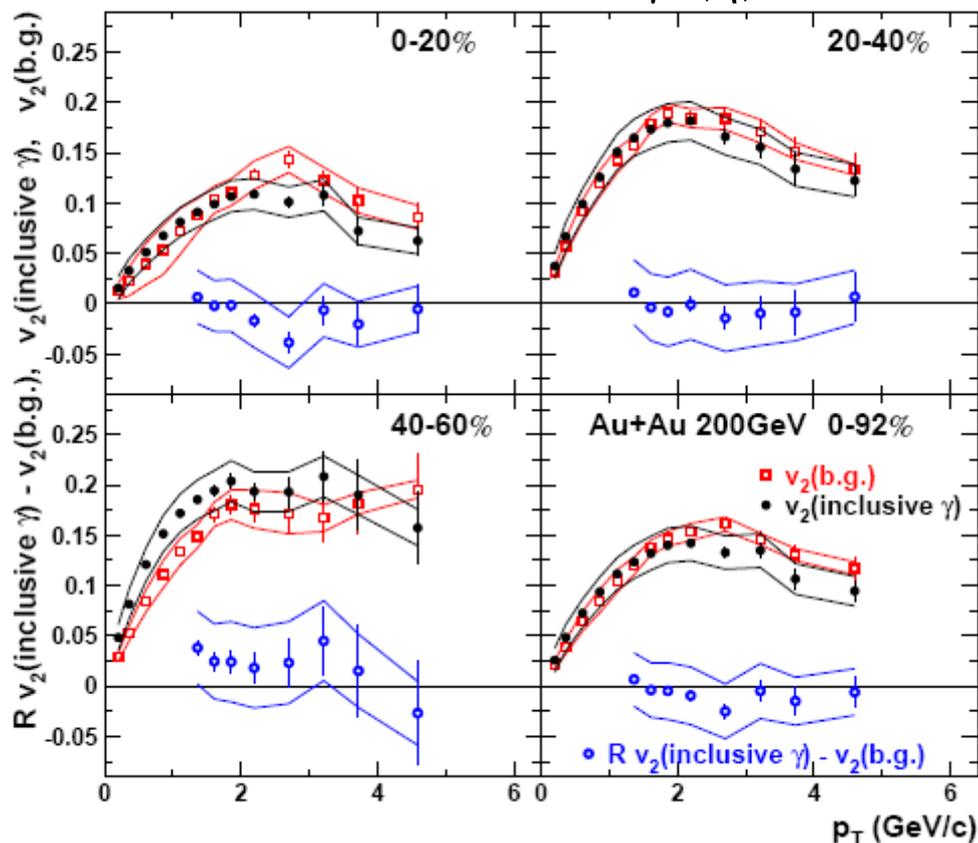
PHENIX BBC:  $3.1 < |\eta| < 3.9$

$$R = (N_{direct \gamma} + N_{b.g.}) / N_{b.g.}$$

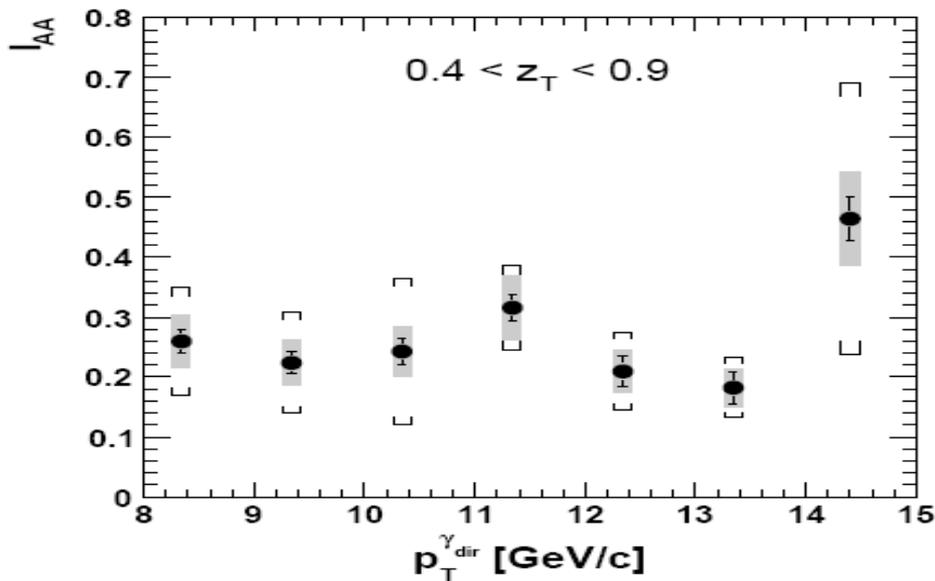
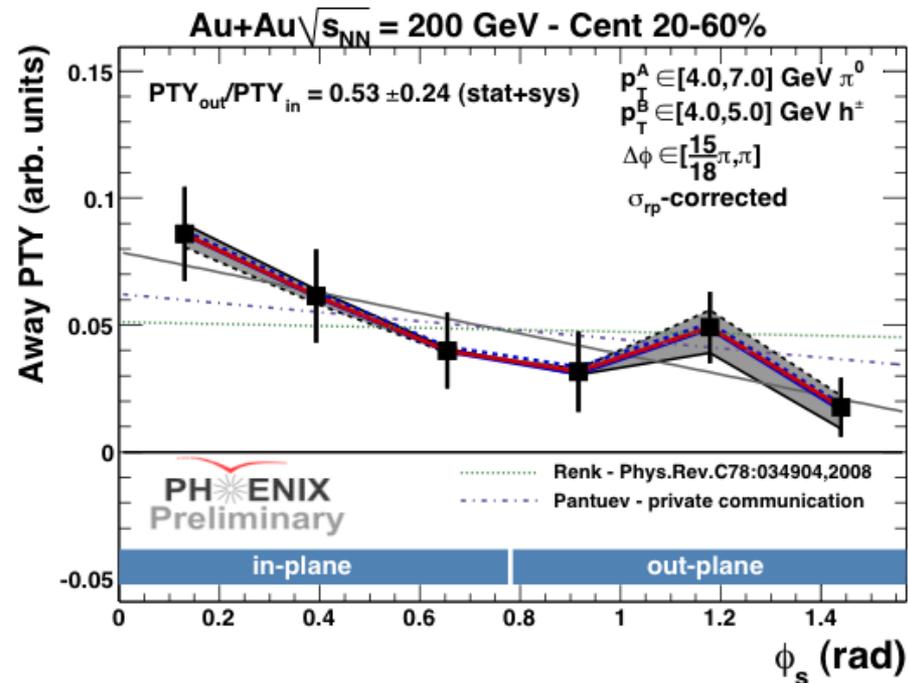
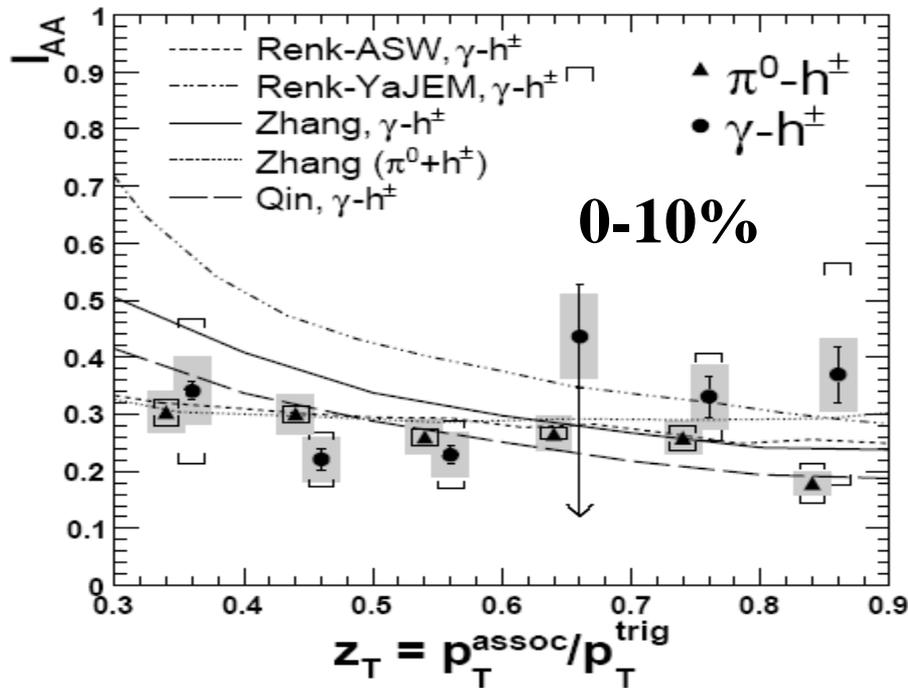
Measured  $\pi^0$  and inclusive  $\gamma$  :  $|\eta| < 0.35$

$$v_2^{direct \gamma} = \frac{R v_2^{inclusive \gamma} - v_2^{b.g.}}{R - 1}$$

**This measurements implies that  $v_2$  of direct photons is  $\sim 0$**



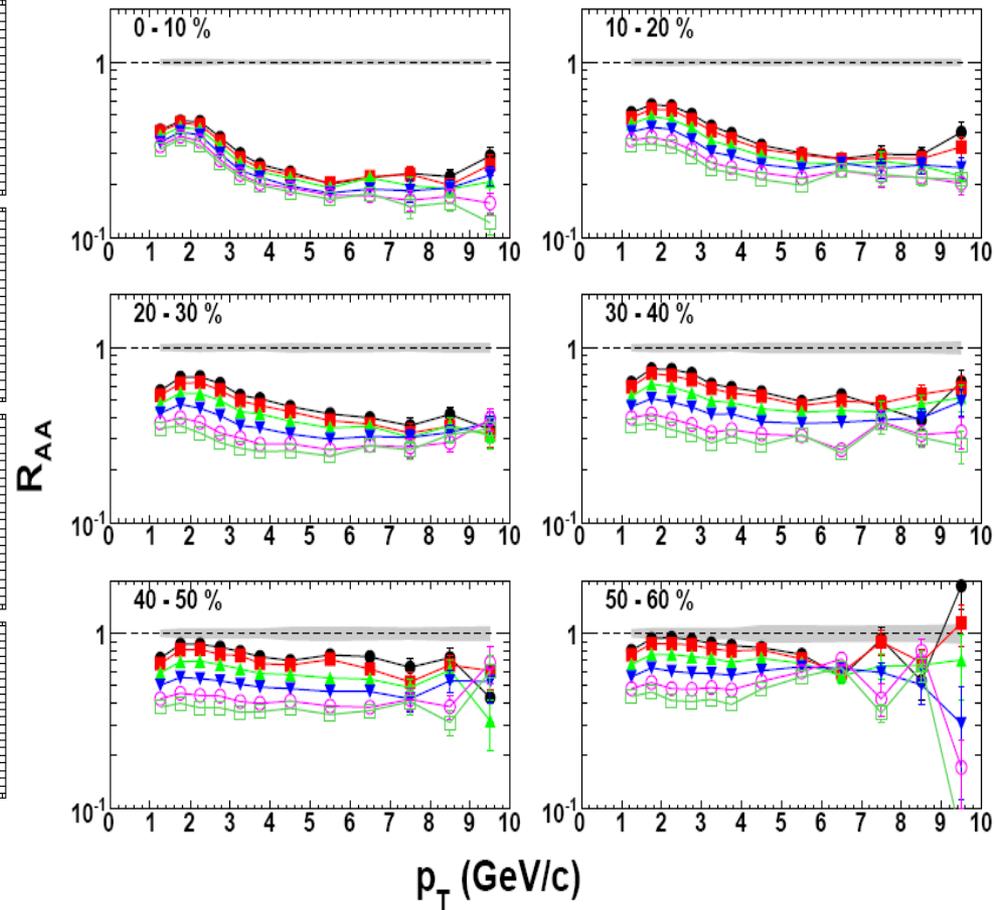
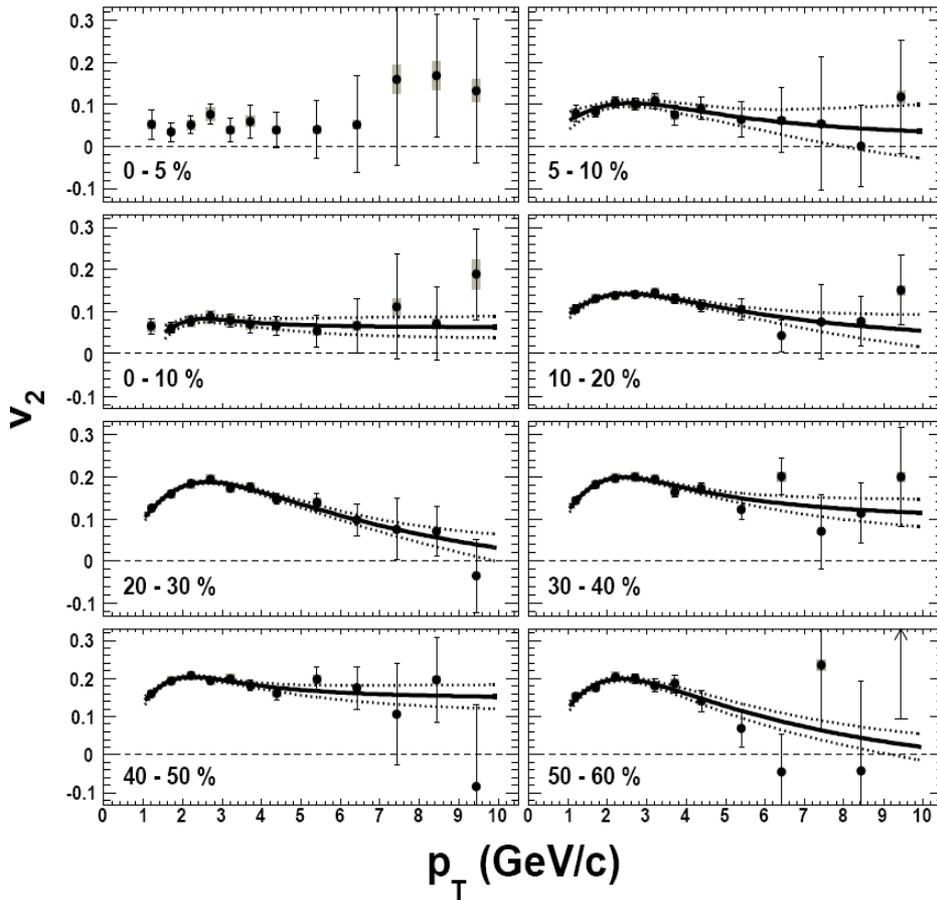
# Path length dependence of energy loss



☀ STAR measurement does not show path-length dependence.

☀ PHENIX measurement show path-length dependence.

# Path length dependence of energy loss



☀  $\pi^0 v_2(p_T)$  and  $R_{AA}(\Delta\phi)$  show statistically significant dependence on the path length particularly at  $p_T < 6 \text{ GeV}$