



# Direct virtual photon measurements in Au+Au collisions with STAR BES-II data

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Shandong University

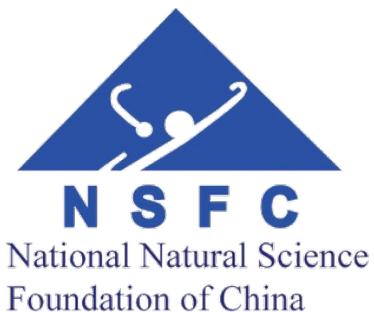
Hot Quark, May. 11-17, 2025, Hefei, China

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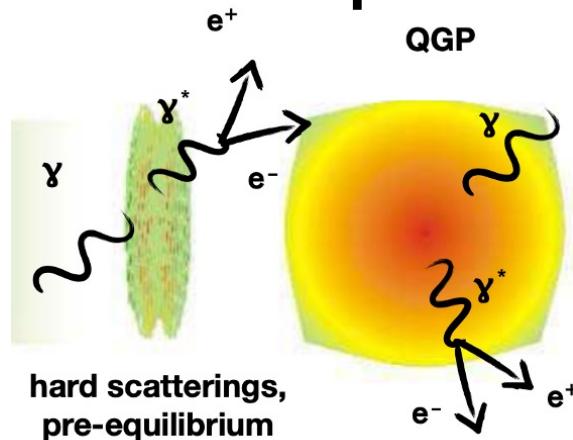
# Introduction



## Why direct virtual photon?

- No strong interaction with medium
- Emitted from initial stage to hadron gas phase stage
- Carry information on energy density, effective temperature and collective motion of QGP

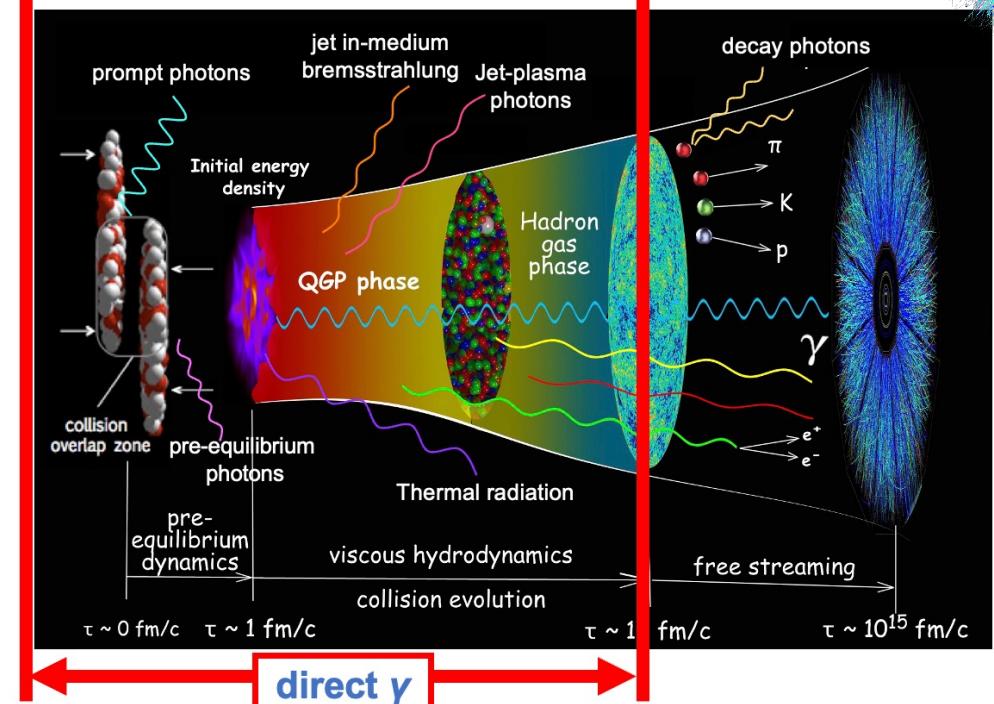
## Direct virtual photon production



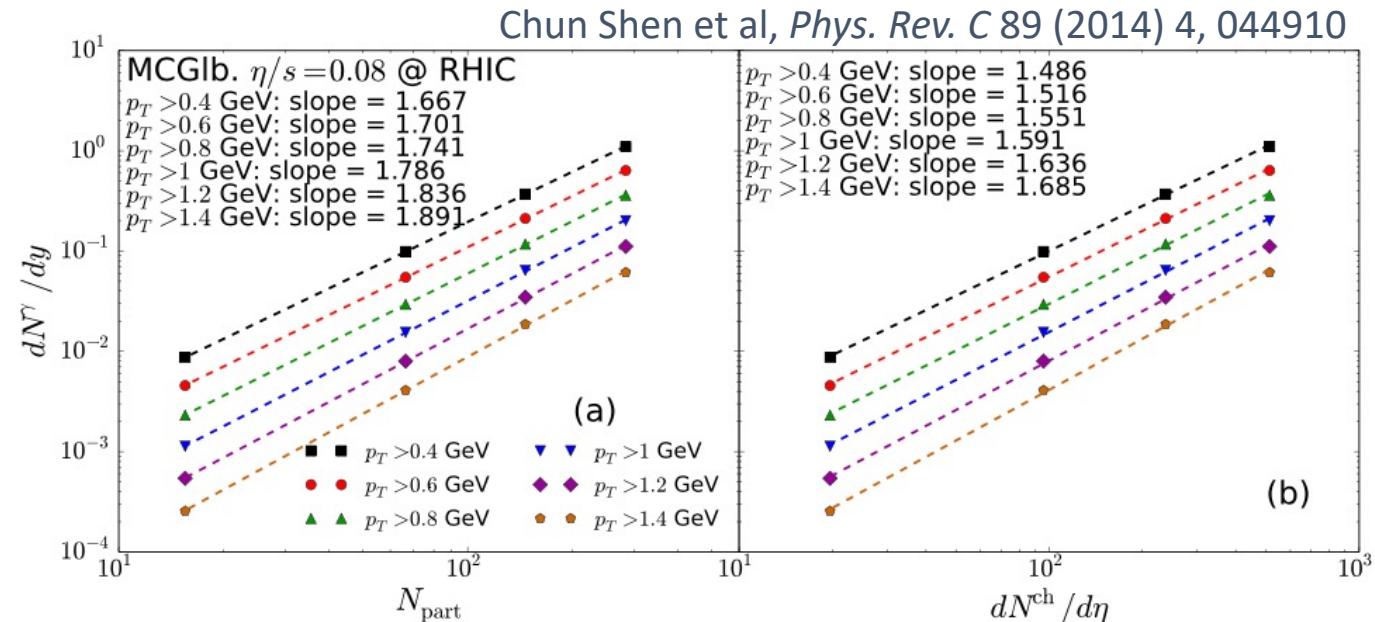
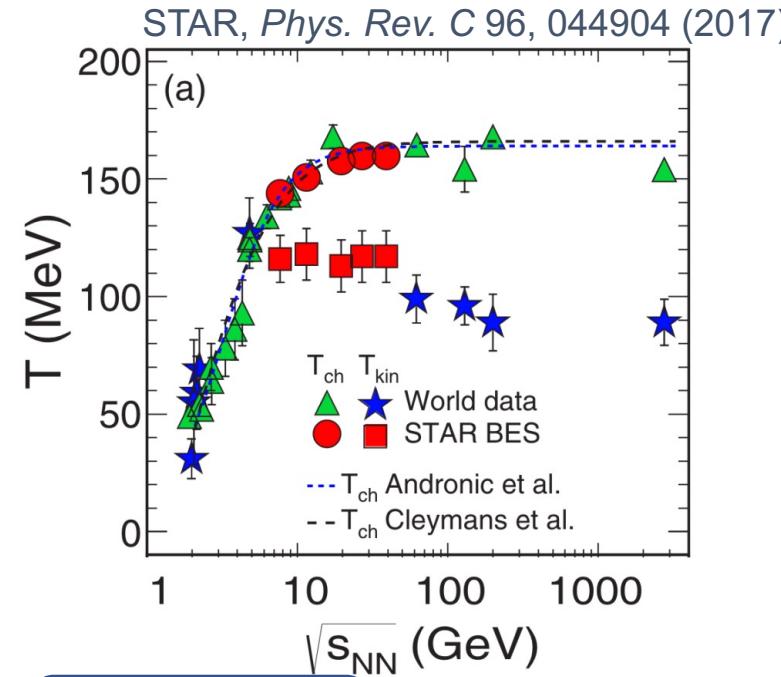
$$\frac{d^2N_{ee}}{dM^2} = \frac{\alpha}{3\pi} \frac{L(M)}{M^2} S(M, q) dN_\gamma \quad S(M, q) = dN_{\gamma^*}/dN_\gamma$$

- High  $p_T$ : mainly **prompt photon** from **earlier stage**
- Low  $p_T$ : mainly **thermal photon** from **later stage**

Chun Shen et al, *Comput. Phys. Commun.*, 199:61–85, 2016



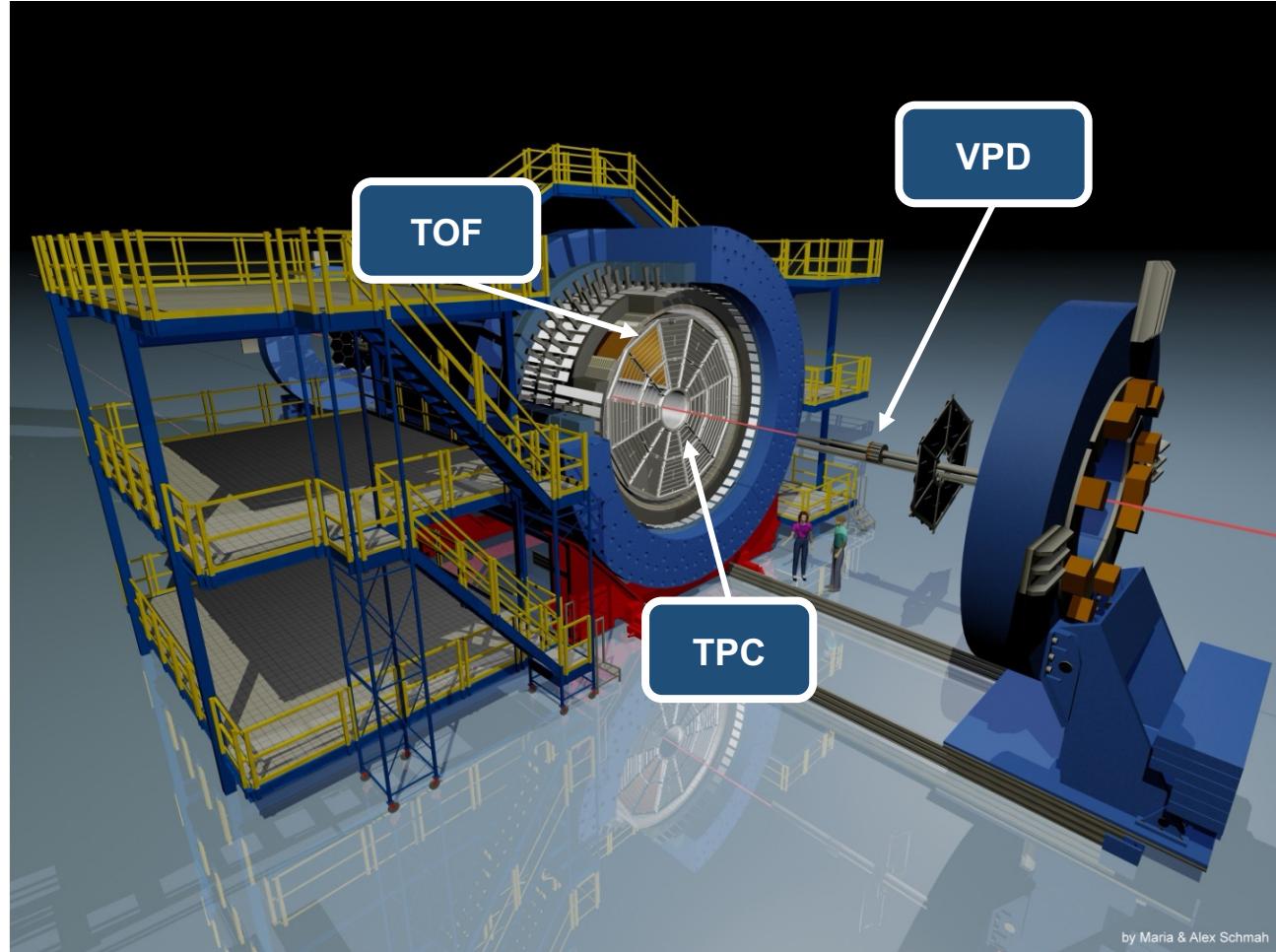
# Thermal radiation



Direct photon:

- **Thermometer**: extract  $T_{effective}$  from  $p_T$  spectra (affected by blue shift effect)
- **Chronometer**: integrated yield is sensitive to lifetime

# The Solenoidal Tracker At RHIC



## Time Projection Chamber

- Tracking
- Momentum and energy loss
- Acceptance:  $|\eta| < 1; 0 \leq \varphi \leq 2\pi$

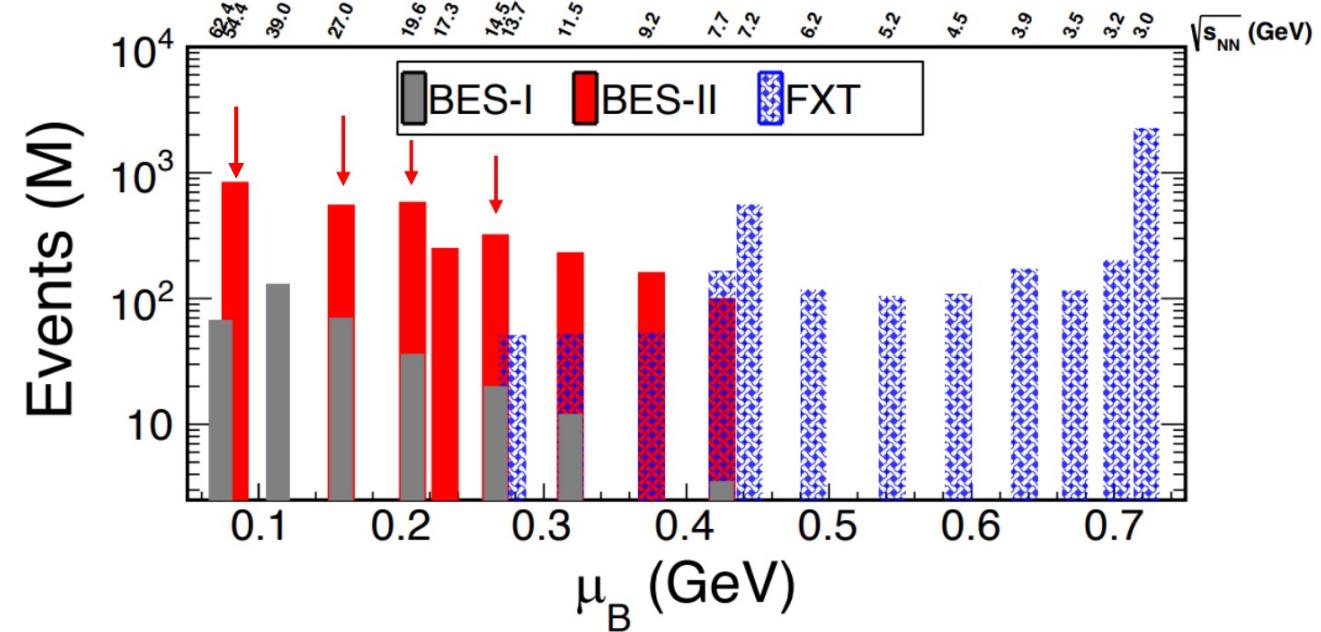
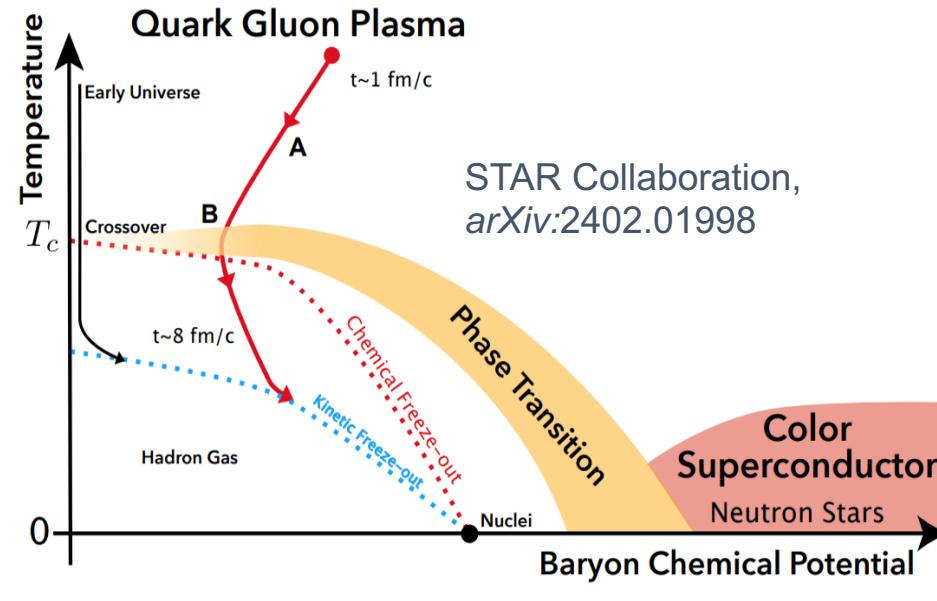
## Time Of Flight

- Time of flight
- Particle identification
- Acceptance:  $|\eta| < 1; 0 \leq \varphi \leq 2\pi$

## Collision energies (Au+Au)

- $\sqrt{s_{NN}} = 54.4 \text{ GeV}$  ( $\sim 430M$  events)
- $\sqrt{s_{NN}} = 27 \text{ GeV}$  ( $\sim 250M$  events)
- $\sqrt{s_{NN}} = 19.6 \text{ GeV}$  ( $\sim 213M$  events)
- $\sqrt{s_{NN}} = 14.6 \text{ GeV}$  ( $\sim 110M$  events)

# Beam Energy Scan II

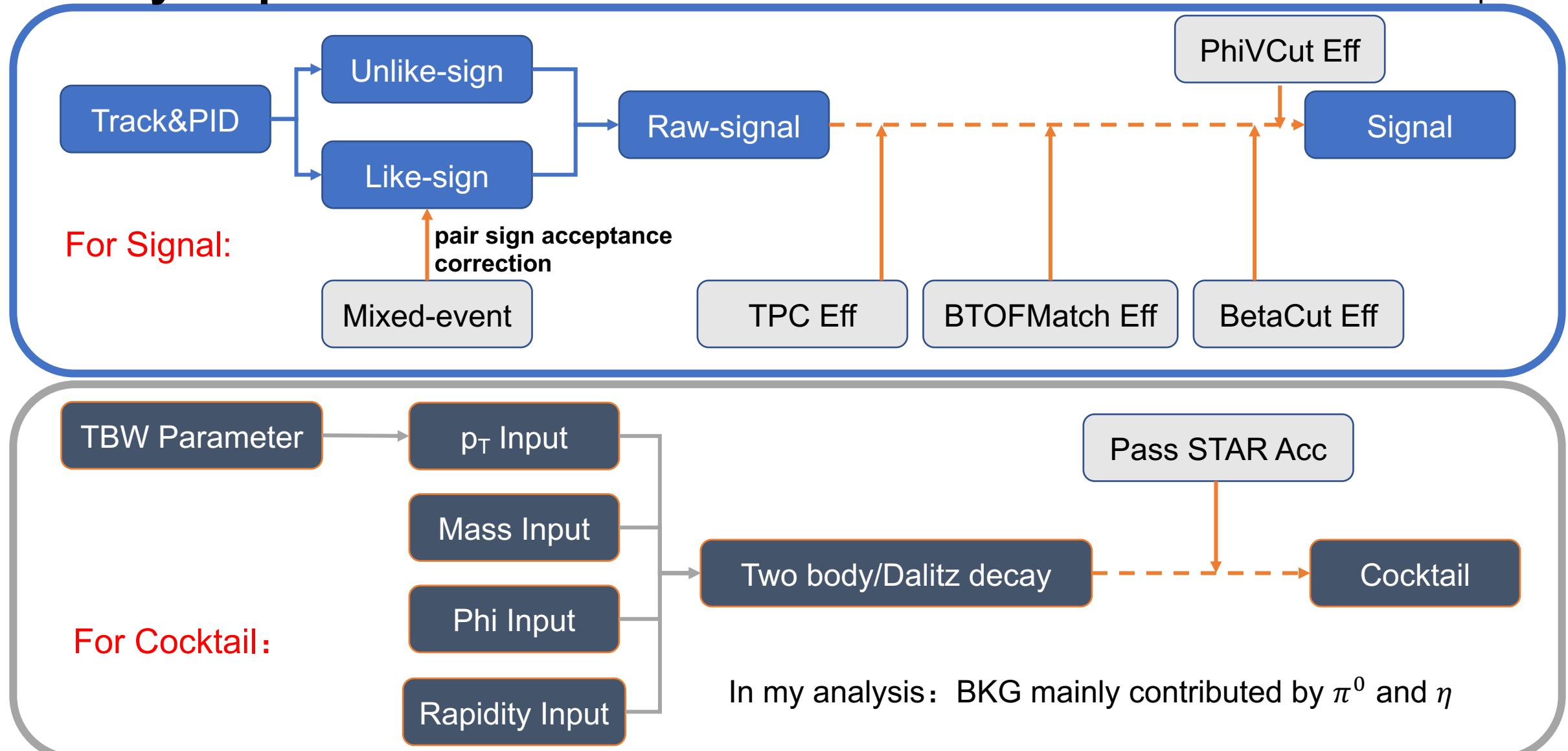


$\gamma_{direct}$ : a possible probe of Critical End Point (CEP)

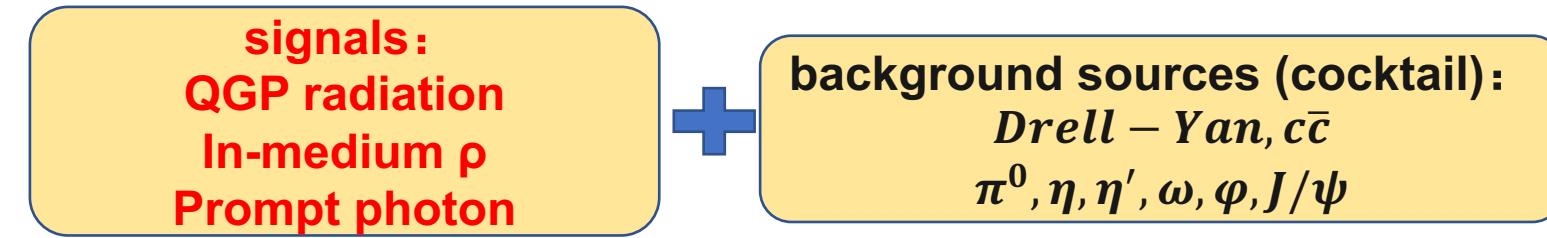
- Expect more thermal radiation close to CEP
- Need high statistics at lower collision energies

➤ BES-II: 10-20 times higher statistics than BES-I

# Analysis procedure



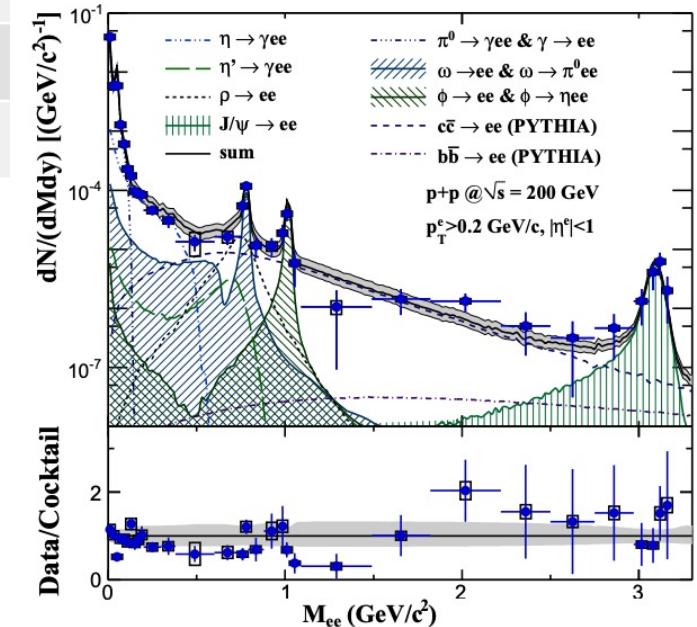
# Signal and Hadronic cocktail



**Hadronic cocktail :**

two-body decay	$\omega \rightarrow e^+e^-$ , $\phi \rightarrow e^+e^-$ , $J/\psi \rightarrow e^+e^-$ , $\psi' \rightarrow e^+e^-$
Dalitz decay	$\pi^0 \rightarrow \gamma e^+e^-$ , $\eta \rightarrow \gamma e^+e^-$ , $\eta' \rightarrow \gamma e^+e^-$ , $\omega \rightarrow \pi^0 e^+e^-$ , $\phi \rightarrow \eta e^+e^-$
heavy-flavor decay	$c\bar{c} \rightarrow e^+e^-X$
Drell-Yan process	$DY \rightarrow e^+e^-$

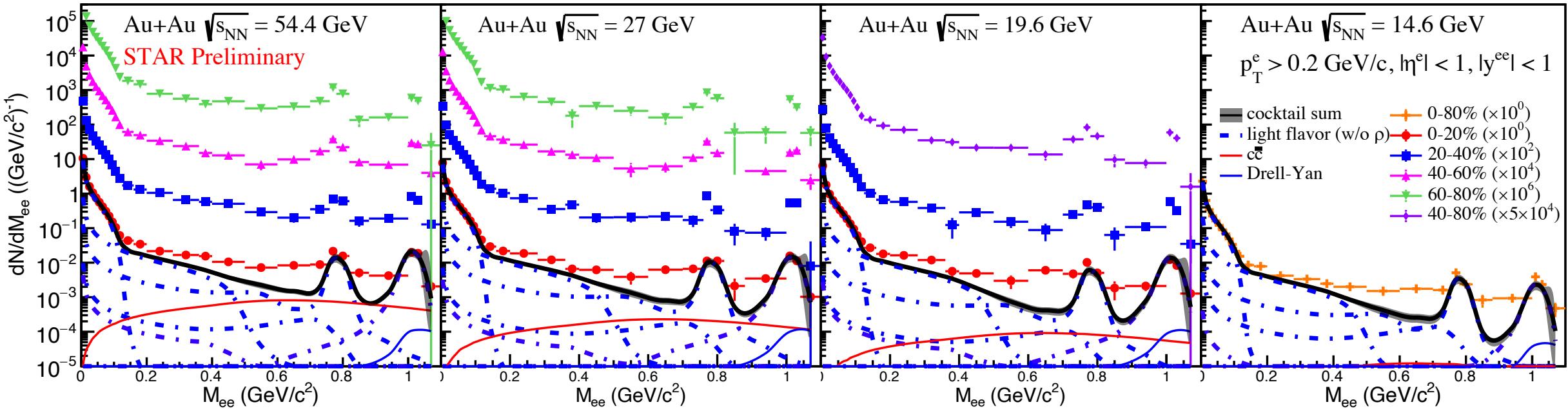
STAR, Phys. Rev. C 86 (2012) 024906



- Hadronic cocktail was examined in previous dielectron spectra studies in p+p collisions

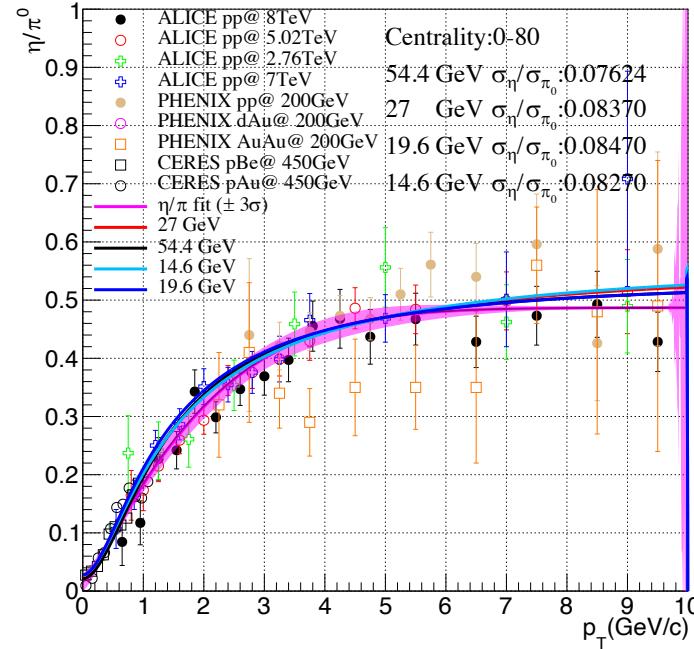
# Data vs. Cocktail

STAR, *Phys.Rev.C* 107 (2023) 6, L061901  
 STAR, *arXiv:2402.01998*



- Dielectron signal is consistent with cocktail in  $\pi^0$  mass region
- Observed significant excess yield contributed by  $\gamma_{dir}^*$ , in-medium  $\rho$  at low mass region  
**(LMR,  $0.1 < M_{ee} < 0.76 \text{ GeV}/c^2$ )**

# Extract direct virtual photon



➤ Fit method:

$$R \frac{\eta}{\pi^0}(p_T) = A \frac{(e^{-a*p_T - b*p_T^2} + \left(\frac{R^\infty}{A}\right)^{-\frac{1}{n}} \frac{p_T}{p_0})^{-n}}{(e^{-a*p_T - b*p_T^2} + \frac{p_T}{p_0})^{-n}}$$

Yuanjin Ren et al, *Phys.Rev.C* 104 (2021) 5, 054902

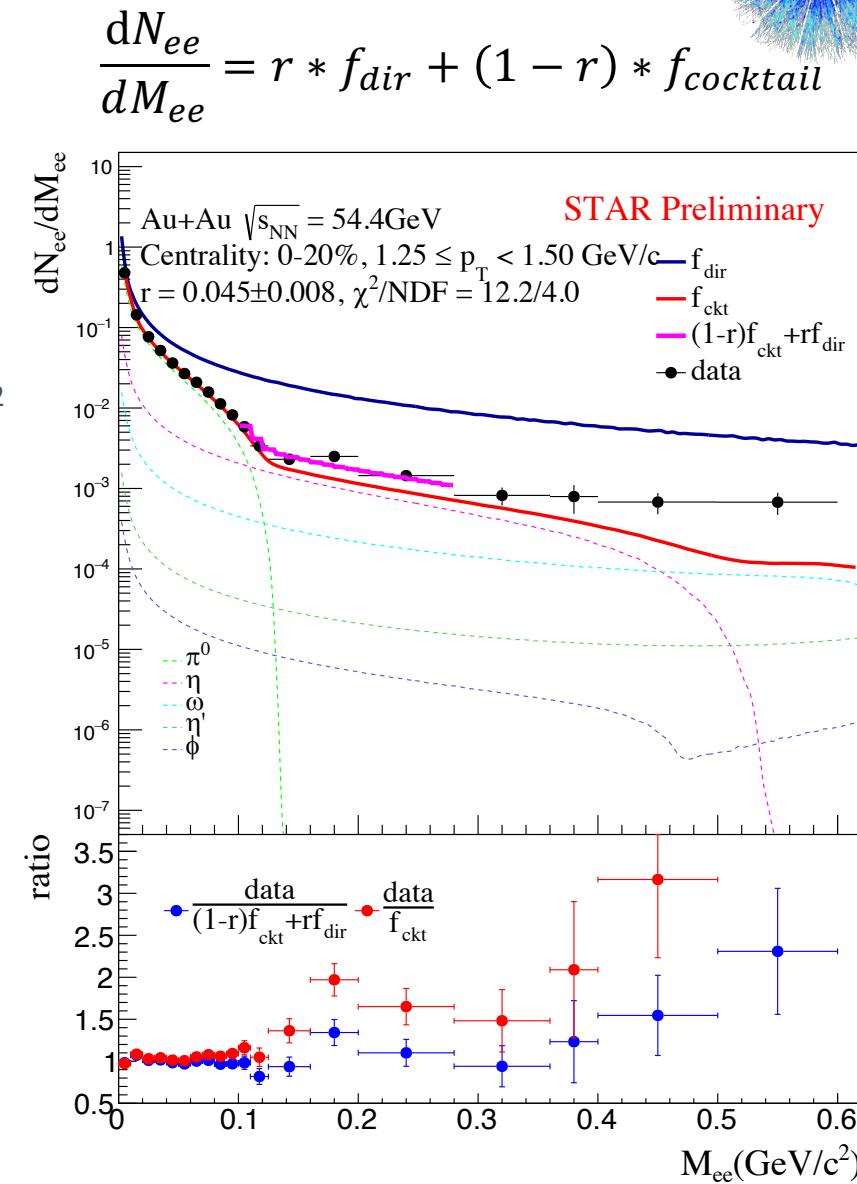
➤  $\eta/\pi^0$

- Parametrized using Tsallis Blast Wave (TBW) function
- Fixed to **0.470±0.017** at 5 GeV/c

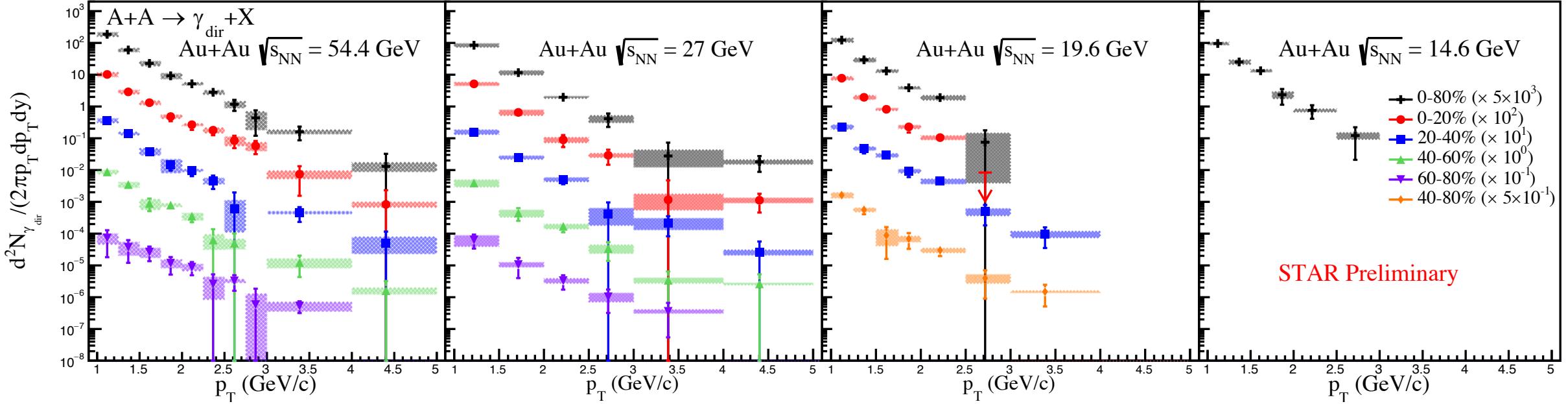
## Two component fit

PHENIX, *Phys.Rev.C* 81 (2010) 034911

- A clear enhancement at  $\eta$  mass region contributed by  $\gamma_{dir}^*$
- Extract direct virtual photon fraction  $r$  by fitting cocktail and  $\gamma_{dir}^*$  templates to the data in  $M_{ee}$  range [0.10,0.28] GeV/c<sup>2</sup>

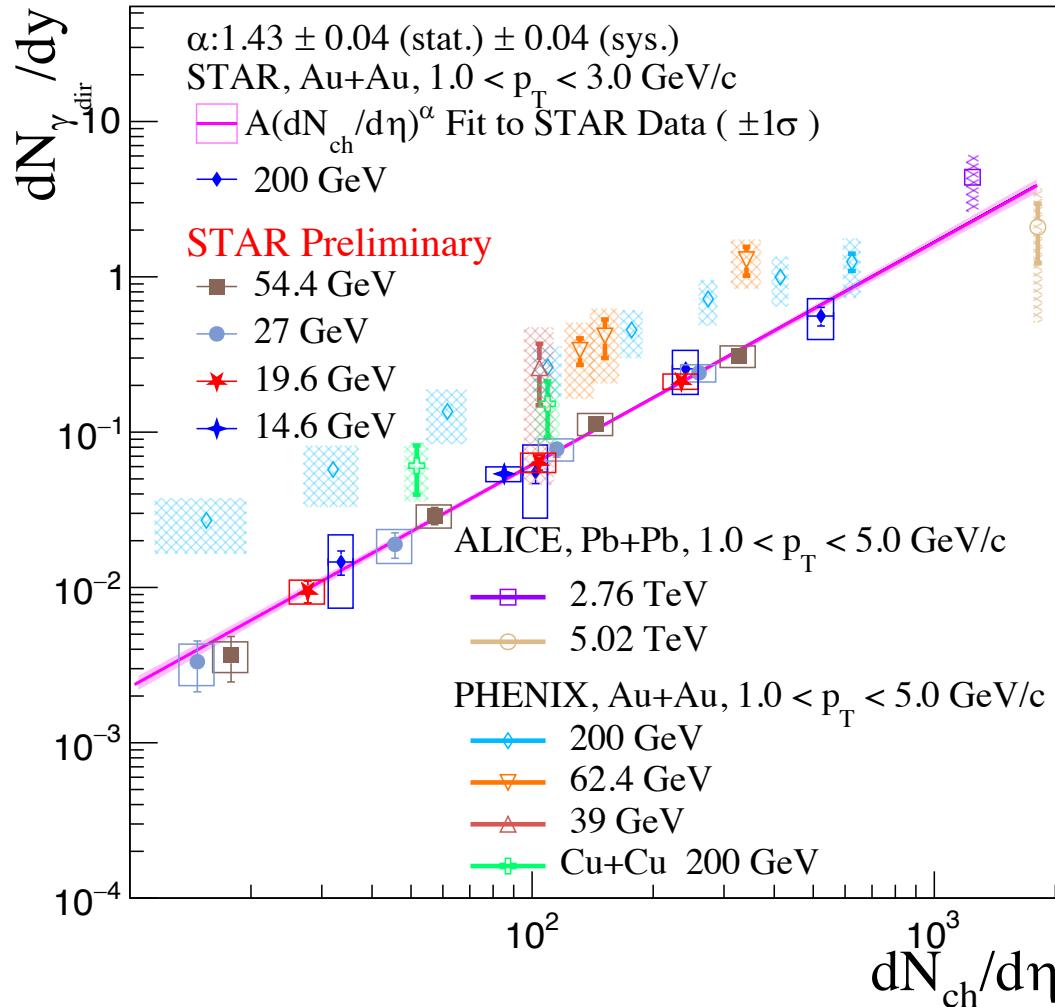


# Direct virtual photon $p_T$ spectrum



- First measurement of  $\gamma_{dir}^*$  in Au+Au collisions at BES-II in different centrality intervals

# Direct virtual photon $dN/dy$ vs. $dN_{ch}/d\eta$



- New measurements of  $dN_{\gamma_{dir}}/dy$  at STAR
- Strong  $dN_{ch}/d\eta$  dependence
- The yields at  $\sqrt{s_{NN}} = 14.6, 19.6, 27, 54.4, 200 \text{ GeV}$  measured by STAR follow a common scaling, with

$$\alpha = 1.43 \pm 0.04 \pm 0.04$$

STAR, *Phys. Lett. B* 770 (2017) 451-45

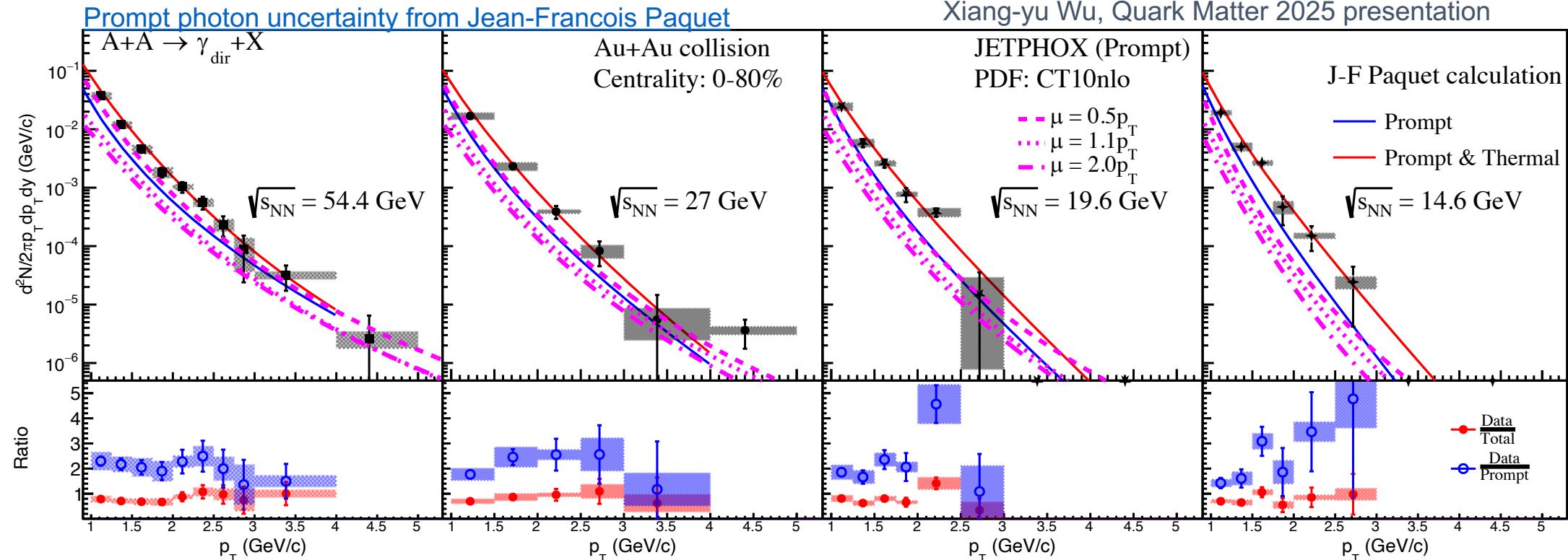
PHENIX, *Phys. Rev. C* 109 (2024) 044912

ALICE, *arXiv*: 2308.16704

ALICE, Jerome Jung, Hard Probe 2024 presentation

# Thermal photon enhancement

Patrick Aurebche et al, *Phys.Rev. D73:094007,2006*  
 PHENIX, *Phys.Rev.C 107 (2023) 2,024914*  
 Xiang-yu Wu, Quark Matter 2025 presentation



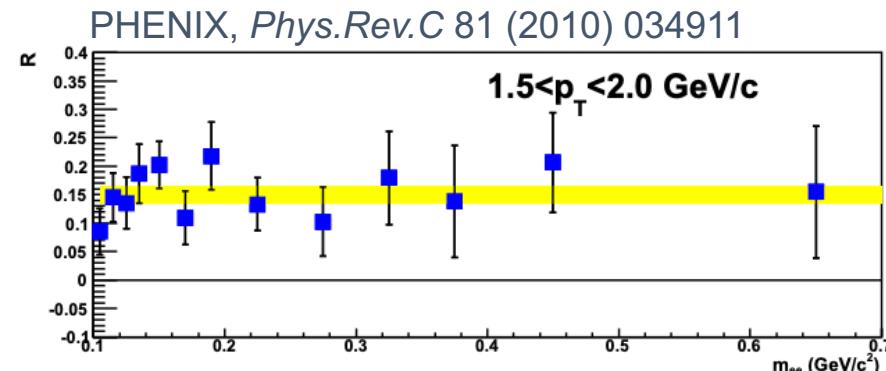
- Enhancement observed compared with prompt photon yield in 0-80% Au+Au at  $\sqrt{s_{NN}} = 14.6 - 54.4 \text{ GeV}$
- Indication of **thermal photon** contribution
- **Measured direct photon yield** can be well described by **theoretical calculation**

# Direct virtual photon mass shape

STAR, *Phys.Lett.B* 770 (2017) 451-458  
Peter Lichard, *Phys.Rev.D* 51 (1995) 6017-6035

For  $p_T > 1.5 \text{ GeV}/c$  the fit gives good  $\chi^2/\text{NDF}$ , demonstrating that the shape of the excess is consistent with  $1/m_{ee}$  as expected for internal conversion.

Since the shape of  $f_{\text{dir}}(m_{ee})$  is  $1/m_{ee}$  smeared by the detector effects, a fit of  $R = (\text{data} - \text{cocktail})/f_{\text{dir}}(m_{ee})$  to a constant can be used to test that the excess has the shape expected for internal conversion of direct photons.



- Mass Shape:
  - Power-law ( $1/M$ ) ← dominated by internal conversion from **earlier stage**
  - Exponential ( $e^{-M}$ ) ← dominated by thermal radiation from **later stage**
- Are the mass shapes in different  $p_T$  ranges sensitive to these contributions from different stages of evolution?

## Strategy:

For  $\gamma_{\text{dir}}^*$  mass:

$$\frac{dN}{dM} = C * [1 + (q - 1) \frac{M_{ee}}{\beta}]^{-\frac{1}{q-1}}$$

- Apply one **unified** mass shape to fit dielectron mass spectrum **in different  $p_T$  ranges**
- **Different values of q parameter:**
  - **Power-law ( $1/M$ ) :  $q \sim 2$**
  - **Exponential ( $e^{-M}$ ):  $q \sim 1$**

# Direct virtual photon production stage

$q \in [1, 2.5]$

$\beta \in (0, 0.25]$

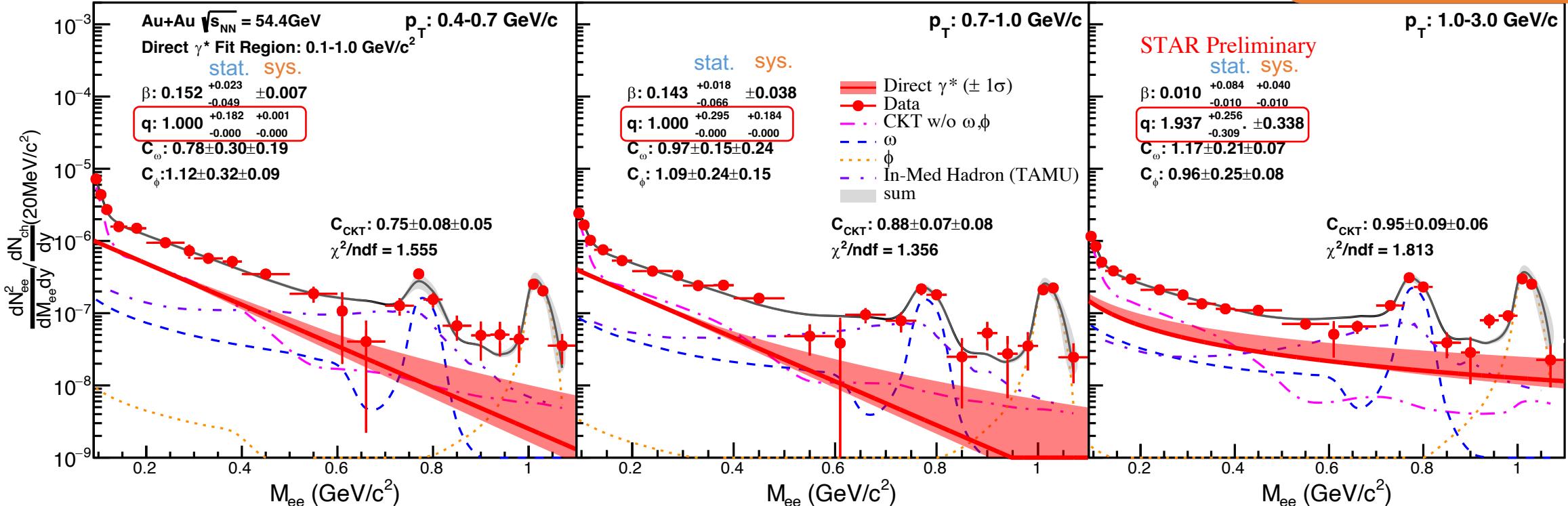


For  $\gamma_{dir}^*$  mass:

$$\frac{dN}{dM} = C * [1 + (q - 1) \frac{M_{ee}}{\beta}]^{-\frac{1}{q-1}}$$

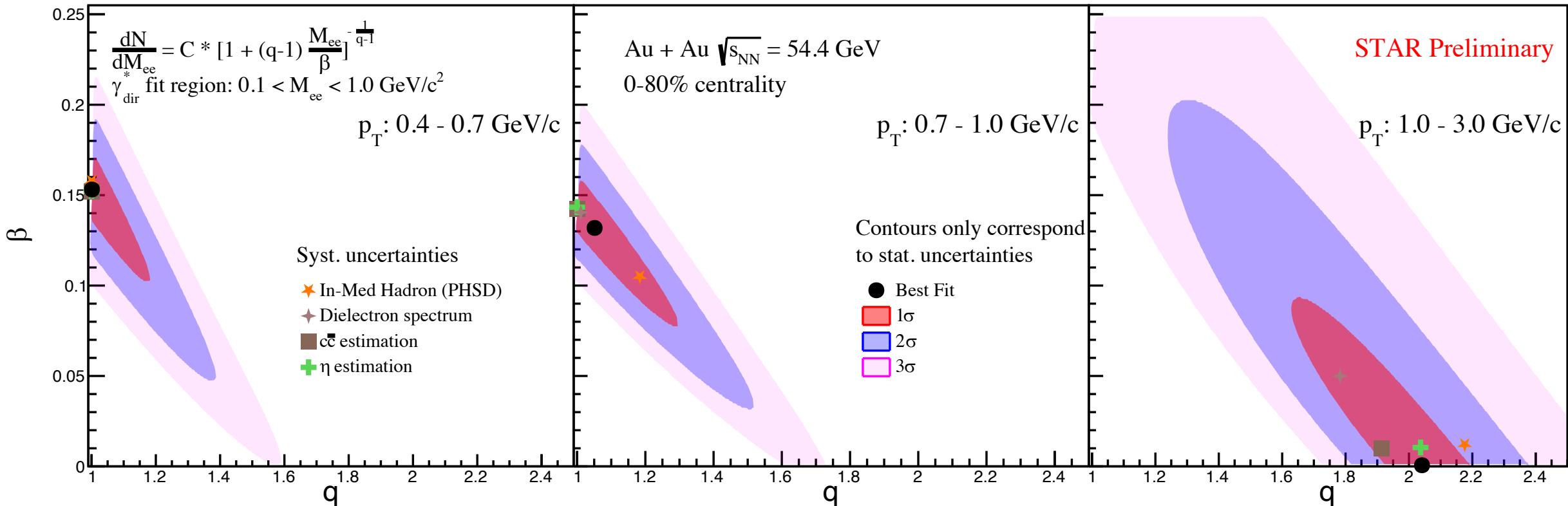
Exponential  $q \sim 1$   
Power-law  $q \sim 2$

$$\frac{dN_{ee}}{dM_{ee}} = \gamma^{direct} + C_{ckt} * f_{ckt(w/o \omega, \phi)} + H_{In-medium} + C_\omega * f_\omega + C_\phi * f_\phi$$



- As the transverse momentum increases, its shape changes from **exponential** to **power-law**, corresponding to the direct photon produced in **later** to **earlier** stages respectively

# Direct virtual photon production stage



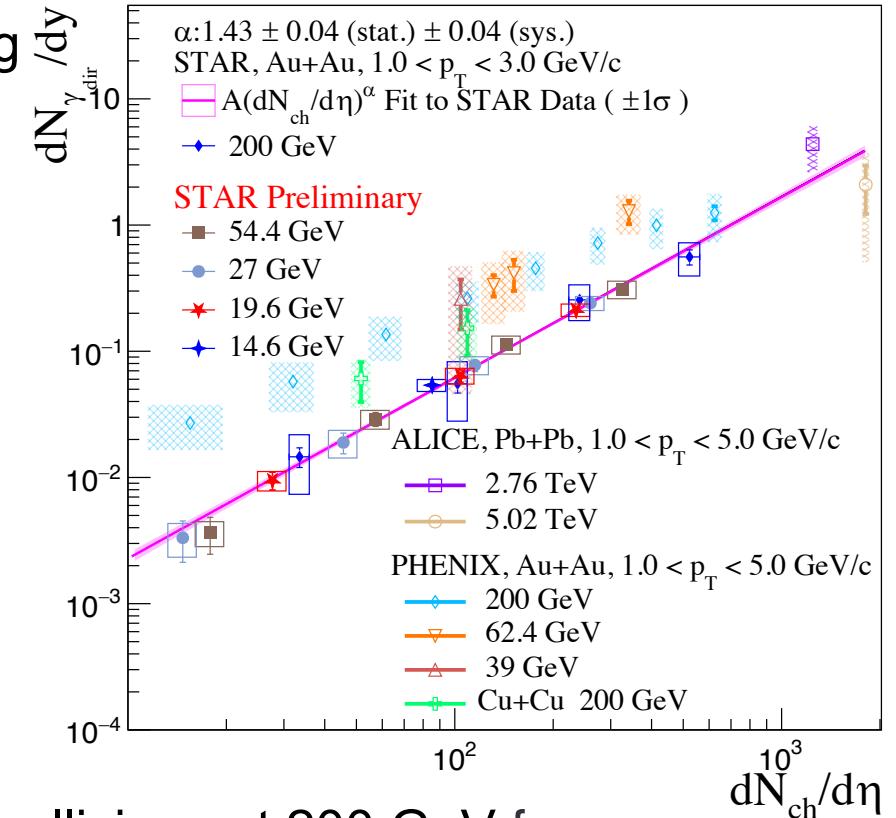
- Acquire 1, 2, 3 $\sigma$  contours by scanning different  $q$  and  $\beta$ 
  - $1 \leq q \leq 2.5$
  - $0 < \beta \leq 0.25$
- Transition from exponential to power-law is observed

# Summary

- First measurements of direct virtual photons in Au+Au collisions with BES-II data at RHIC
- The measured yields from STAR follow a common scaling
  - Strong  $dN_{ch}/d\eta$  dependence
  - Scaling power  $\alpha = 1.43 \pm 0.04$  (stat.)  $\pm 0.04$  (sys.)
- Direct virtual photon mass spectrum shape:
  - High  $p_T$  ( $>1$  GeV/c) → **power-law** → earlier production
  - Low  $p_T$  ( $<1$  GeV/c) → **exponential** → later production

**Outlook:**

- Fixed-target data at lower collision energies
  - High  $\eta$  region
- Direct virtual photon polarization/ $v_2$  in Au+Au collisions at 200 GeV from Run23 and Run25



# Summary

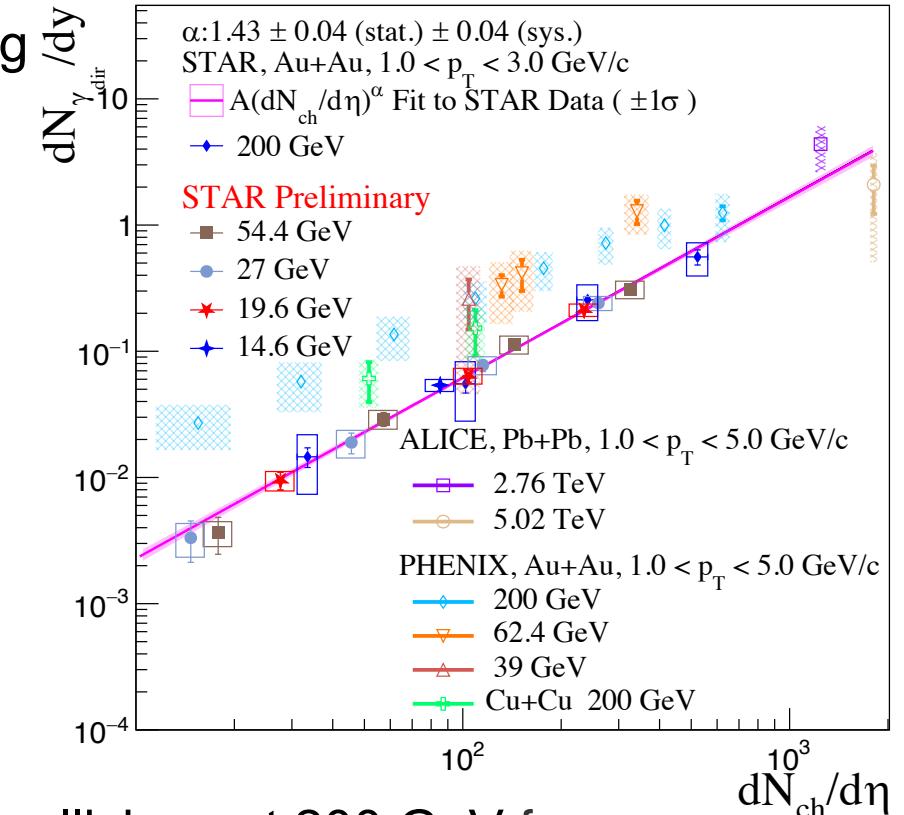
## Thanks for your attention!

*Thanks for theoretical contribution by X. Wu, A. Alaoui, C. Gale, S. Jeon, J-F. Paquet, B. Schenke and C. Shen*

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- Direct virtual photon mass spectrum shape:
  - High  $p_T$  ( $>1\text{GeV}/c$ ) → **power-law** → **earlier production**
  - Low  $p_T$  ( $<1\text{GeV}/c$ ) → **exponential** → **later production**

**Outlook:**

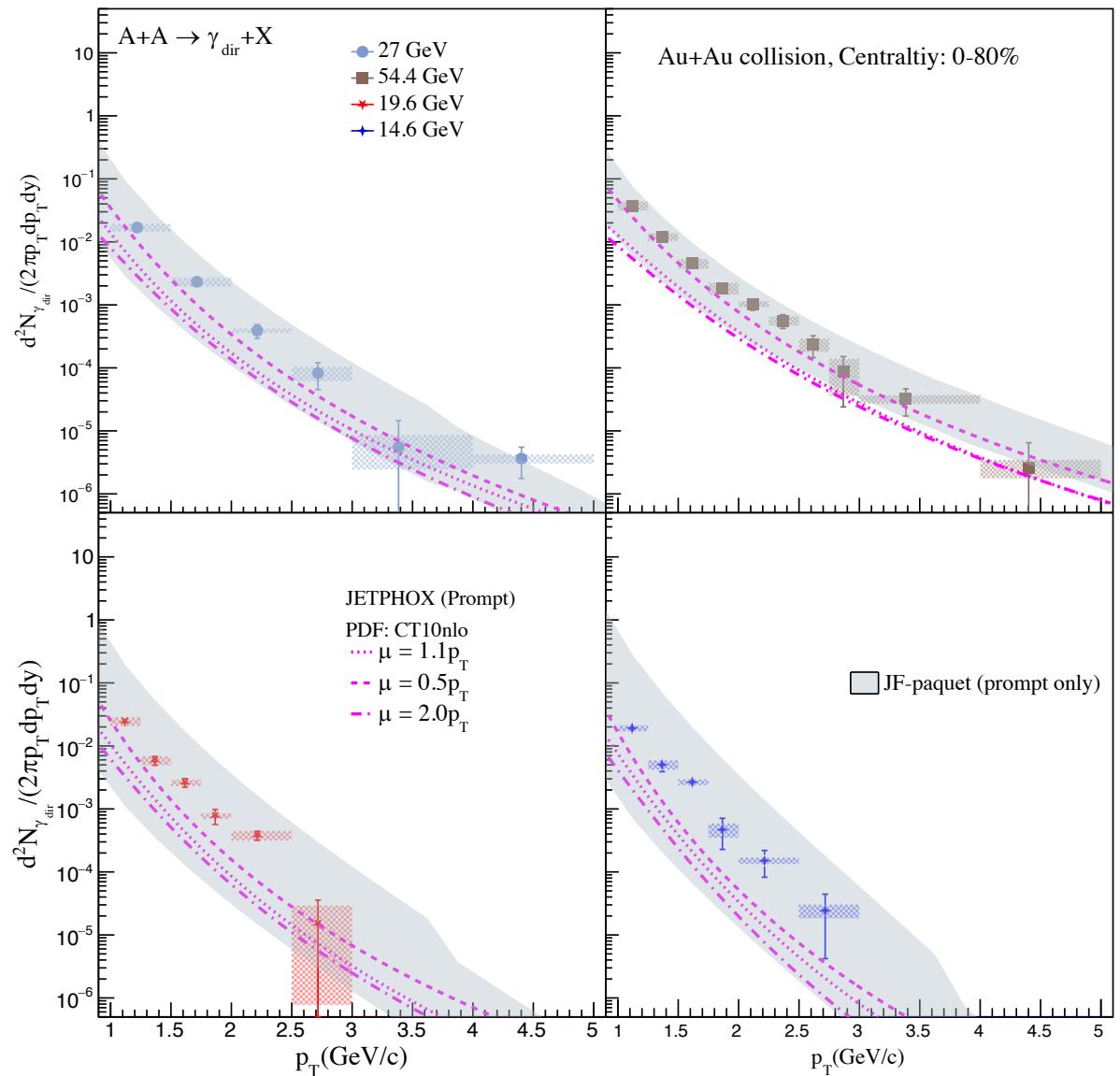
- Fixed-target data at lower collision energies
  - High  $\eta$  region
- Direct virtual photon polarization/ $v_2$  in Au+Au collisions at 200 GeV from Run23 and Run25



# Backup

# Compared with theoretical calculation

The direct virtual photon  $p_T$  spectra  
compared with Jean-Francois Paquet  
theoretical calculation



# Effective temperature based on Jetphox

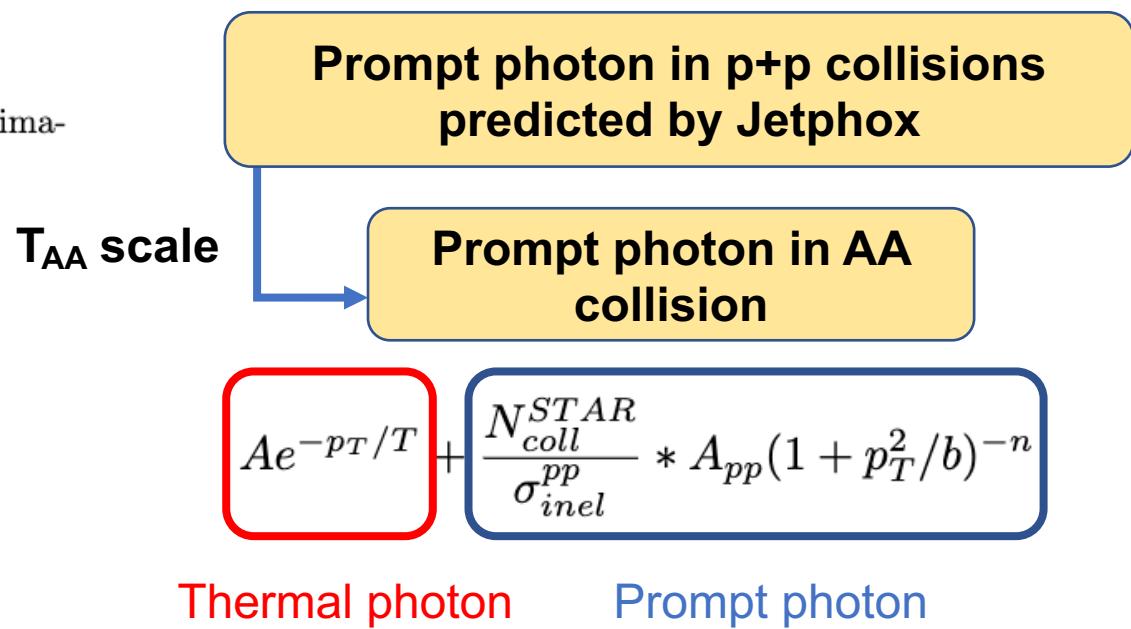
Prompt photon distribution can be comparable with data in

p+p collision at 200 GeV and 63 GeV when  $\mu = 1.1 p_T$

Table 1: Effective temperature extraction after different prompt photon estimation by jetphox

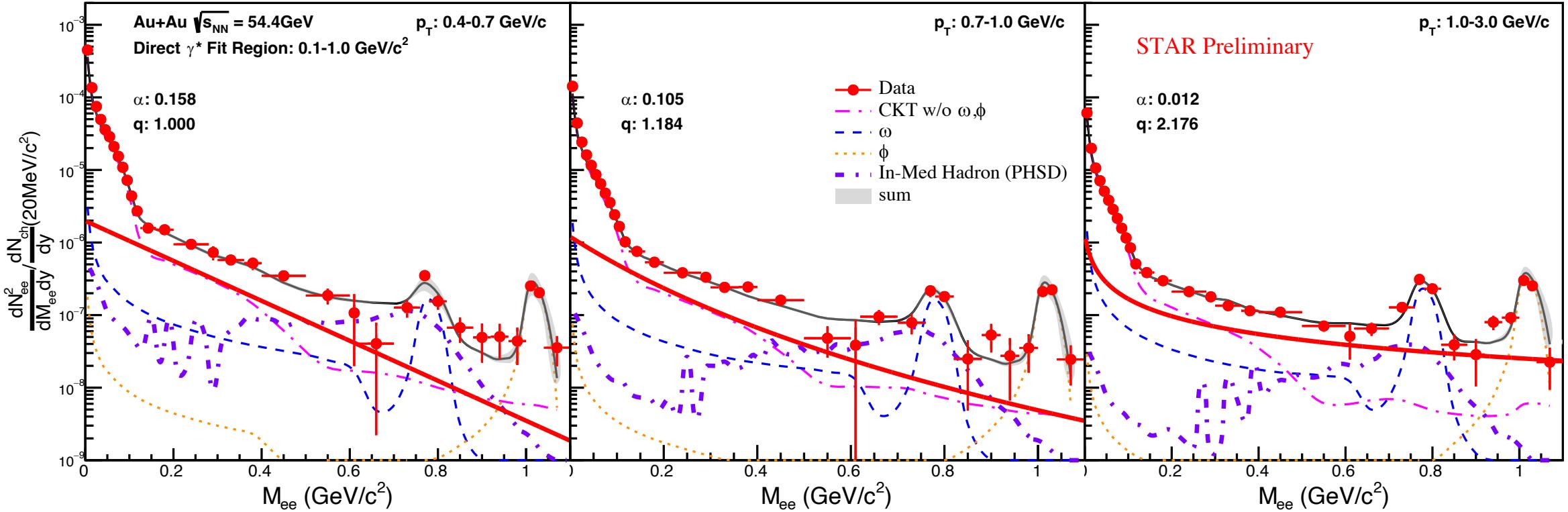
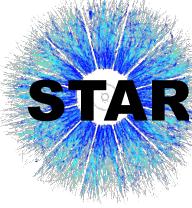
$\mu$	54.4 GeV	27 GeV	19.6 GeV	14.6 GeV
0.5 $p_T$	316±44±11	298±78±26	259±42±5	240±46±9
0.6 $p_T$	301±36±11	282±71±25	255±40±5	234±44±9
0.7 $p_T$	300±33±11	278±68±25	248±38±5	233±42±9
0.8 $p_T$	291±31±11	275±67±19	248±38±2	233±42±8
0.9 $p_T$	288±30±11	282±66±20	246±37±2	233±41±8
1.0 $p_T$	291±30±11	280±65±20	245±37±2	234±41±8
1.1 $p_T$	289±29±11	279±64±16	244±36±2	233±41±8
1.2 $p_T$	289±29±11	277±63±16	244±36±2	233±41±8
1.3 $p_T$	289±29±11	278±63±16	242±36±2	232±41±8
1.4 $p_T$	288±29±11	276±63±16	242±36±2	232±41±8
1.5 $p_T$	286±28±11	277±63±16	242±36±2	232±40±8
1.6 $p_T$	287±28±11	277±63±16	242±36±2	232±40±8
1.7 $p_T$	289±28±11	277±62±15	241±36±2	232±40±8
1.8 $p_T$	289±28±11	278±63±15	241±35±2	232±40±8
1.9 $p_T$	288±28±10	278±63±15	242±36±2	232±40±8
2.0 $p_T$	289±28±11	278±62±16	240±35±2	232±40±8

$\mu$  : factorisation/renormalisation/fragmentation scale

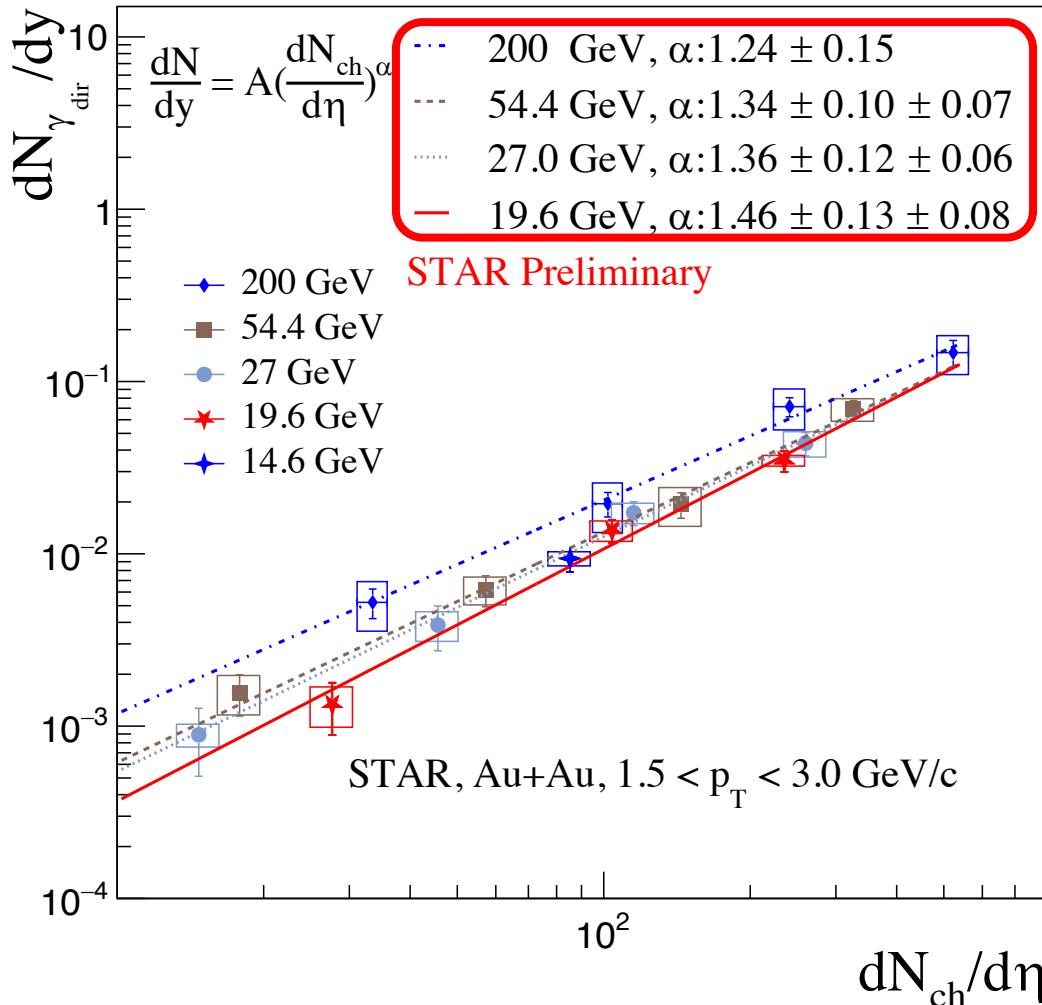


Observed that the effective temperature **decreases** as the energy **decreases** in different estimation of prompt photon by Jetphox

# Direct virtual photon mass shape (In-Med Hadron, PHSD)



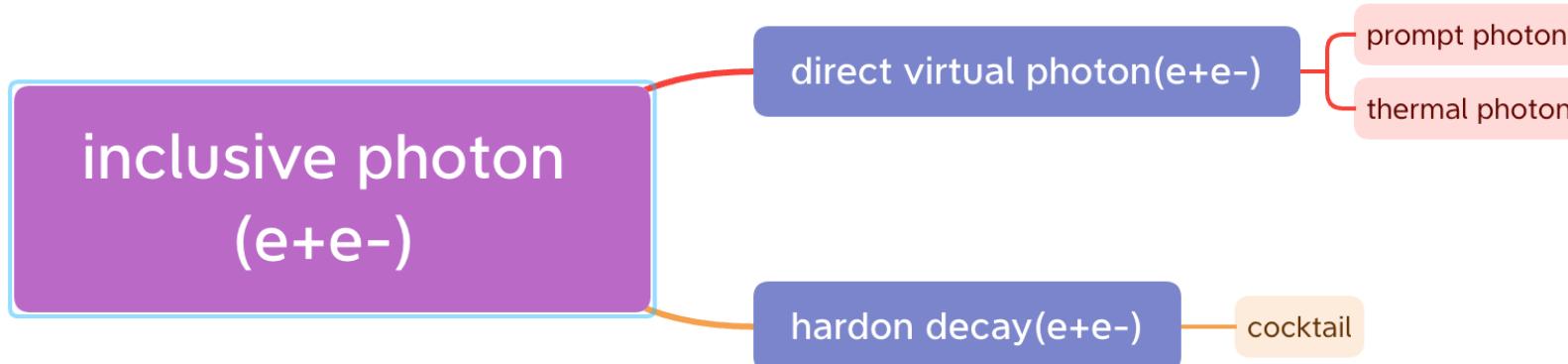
# $dN/dy$ vs. $dN_{ch}/d\eta$



Energy(GeV)	200	54.4	27	19.6
A	6.81e-5	2.83e-5	2.37e-5	1.26e-5
error	5.45e-5	1.49e-5	1.41e-5	8.39e-6

- Obvious  $dN_{ch}/d\eta$  dependence in energies from 200 to 19.6 GeV
- Hint of increasing  $\alpha$  with decreasing collision energies from 200 to 19.6 GeV
  
- ✓ **Hint that increasing proportion of prompt photon at lower collision energies at high  $p_T$**
- ✓ **Thermal photons emitted by QGP rarely extend to higher  $p_T$  at lower collision energies**

# Internal conversion



**Direct photon invariant yield:**

$$d_{direct \gamma} = r * F * \frac{3\pi}{2\alpha}$$

$$\frac{d^2N}{2\pi p_T dp_T dy} = \frac{3 * r * F}{4\alpha p_T dp_T dy}$$

inclusive dielectron consist of:

- From dalitz decay
- From direct photon decay

Parameter  $r$  to measure the directphoton weight in inclusive photon

$$if: p_T \gg M_{ee} \\ S \rightarrow 1, L(M_{ee}) \rightarrow 1$$

$$\begin{aligned} \frac{d^2N_{ee}}{dM} &= \frac{d^2N_{ee}^{direct \gamma*}}{dM} + \frac{d^2N_{ee}^{dalitz \gamma*}}{dM} \\ &= \frac{2\alpha}{3\pi} \frac{L(M)}{M} dN_{direct \gamma*} + \frac{2\alpha}{3\pi} \frac{L(M)}{M} dN_{dalitz \gamma*} \\ &= \frac{2\alpha}{3\pi} \frac{L(M)*S_{direct \gamma}}{M} dN_{direct \gamma} + \frac{2\alpha}{3\pi} \frac{L(M)*S_{dalitz \gamma}}{M} dN_{dalitz \gamma} \\ &= \frac{2\alpha r}{3\pi M} dN_{inclusive \gamma} + \frac{2\alpha (1-r)}{3\pi M} dN_{inclusive \gamma} \\ &= r * f_{dir} + (1 - r) * f_{cocktail} \end{aligned}$$

two-body decay process or Kroll-Wada

$$L(M) = \sqrt{1 - \frac{4M_e^2}{M_{ee}^2}}(1 + \frac{2M_e^2}{M_{ee}^2})$$

$$S = \frac{dN_{\gamma^*}}{dN_\gamma}$$

we interested inclusive background

$$= r * F * \frac{1}{M} + (1 - r) * f_{cocktail}$$

two component fit