

1 Highlights from the STAR experiment

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4 **Abstract.** We present an overview of the recent results obtained by the STAR
5 Collaboration at RHIC. We report the beam energy dependence of directed flow
6 of strange particles. These results support the assumption, that directed flow
7 is formed before hadronization and observed particles are formed via coalescence
8 of constituent quarks. First evidence of a non-zero directed flow of D^0
9 meson is reported. We present the precision measurement of elliptic flow of D^0
10 meson and the first measurement of Λ_c baryon in Au+Au collisions at $\sqrt{s_{NN}}$
11 = 200 GeV, which suggests thermalization of charm quarks. We report results
12 of global Λ polarization measurements in Au+Au collisions, together with the
13 investigation of polarization dependence of centrality and transverse momentum
14 p_T . Results from the STAR fixed-target program show good agreement
15 with previously obtained results. A precise measurement of the ${}^3_\Lambda H$ lifetime is
16 presented. Mesonic decay modes are used to reconstruct the ${}^3_\Lambda H$ from Au+Au
17 collisions. The measured lifetime is about 50% shorter than the one of a free Λ ,
18 indicating strong hyperon-nucleon interaction in the hypernucleus system. Finally,
19 we give an outlook to detector upgrades for the Beam Energy Scan phase
20 II.

21 1 Azimuthal anisotropy measurements

22 One of the most important observables in heavy-ion experiments is the azimuthal anisotropic
23 flow [1] that is usually quantified by the Fourier coefficients of the azimuthal distribution
24 of the final state particles relative to the collision symmetry planes $Ed^3N/dp^3 =$
25 $1/2\pi \cdot d^2N/p_t dp_t d\eta \cdot (1 + \sum 2v_n \cos[n(\phi - \Psi_n)])$. The first-order coefficient, called the directed
26 flow, is argued to be sensitive to the equation of state of the matter, and could serve
27 as a possible signature of the QGP phase transition. STAR has recently performed directed
28 flow measurements at mid-rapidity for $\Lambda, \bar{\Lambda}, K^\pm, K_s^0$, and ϕ at $\sqrt{s_{NN}} = 7.7, 11.5, 14.5, 19.6,$
29 $27, 39, 62.4$ and 200 GeV in Au+Au collisions. Results show good sensitivity for testing a
30 picture where flow is assumed to be formed before hadron formation and the observed particles
31 are assumed to form via coalescence of constituent quarks. Figure 1 represents $v_1(y)$ for
32 indicated particles at two centralities. One can see that within errors, the plotted species have
33 a near linear $v_1(y)$ over the acceptance of the STAR detector.

34 Directed flow slope dv_1/dy versus beam energy for $\Lambda, \bar{\Lambda}, K^\pm, K_s^0, \phi$ was also studied [2]
35 and compared with the data for π^\pm, p and \bar{p} [3]. Combined data for the ten particle species
36 available allow a detailed investigation of constituent quark v_1 . Obtained results support
37 the assumption called the coalescence sum rule [4]. In this scenario the directed flow v_1 is

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38 formed at the pre-hadronic stage, specific types of quarks have the same directed flow and
 39 the detected hadrons are formed via coalescence.

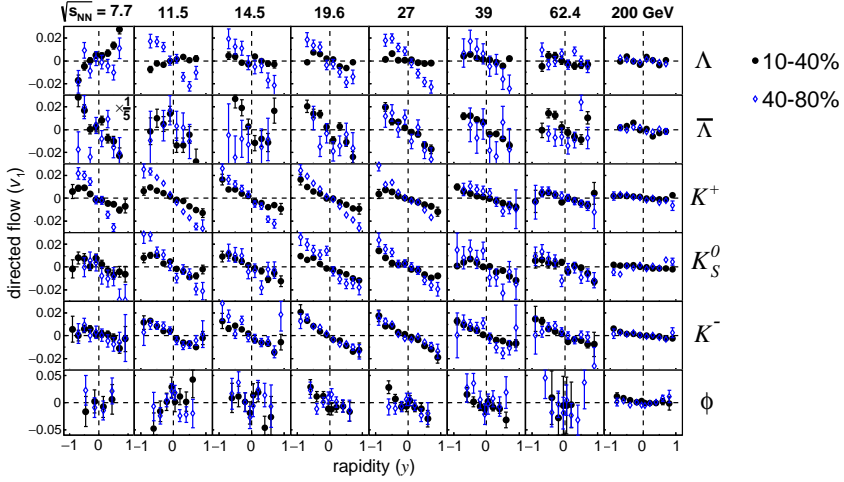


Figure 1. Directed flow as a function of rapidity for the six indicated particle species in 10-40% (black) and 40-80% (blue) central Au+Au collisions at $\sqrt{s_{NN}} = 7.7$ to 200 GeV. The error bars include statistical uncertainties only. All panels for species other than ϕ use the same v_1 scale with the exception of $\bar{\Lambda}$ at $\sqrt{s_{NN}} = 7.7$ GeV, where v_1 magnitudes are exceptionally large and require the measurements to be divided by five.

40 Direct measurements of charmed hadrons require very good vertex detector due to their
 41 very short lifetimes. STAR had installed the Heavy Flavor Tracker (HFT) [5], the MAPS-
 42 based vertex detector for run years 2014-2016. This allowed to perform precise measure-
 43 ments of flow for heavy flavor particles [6].

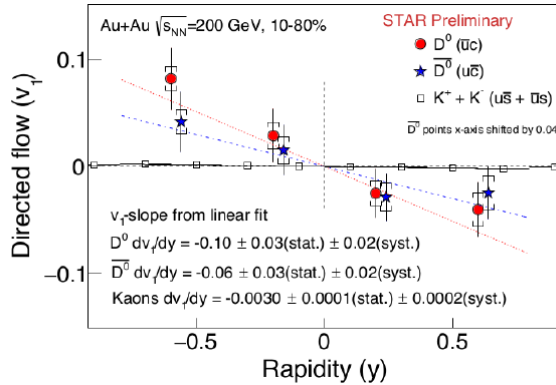


Figure 2. Directed flow as a function of rapidity for D^0 mesons in 10–80% centrality Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV compared with directed flow of $K^+ + K^-$. Data points for \bar{D}^0 are slightly shifted for visibility.

44 Figure 2 shows the first measurement of directed flow for charmed meson D^0 produced
 45 in Au-Au collisions for 10–80% centrality. It is clear that within the errors value of $v_1(y)$ is
 46 non-zero and follows the linear trend obtained previously for light particle sector. The second
 47 harmonic of the charmed meson D^0 was recently studied using data from run years 2014 and
 48 2016. The measured D^0 v_2 in 0–80% centrality Au+Au collisions can be described by a
 49 viscous hydrodynamic calculation for transverse momentum p_T less than 4GeV/c. The D^0 v_2
 50 as a function of transverse kinetic energy ($m_T - m_0$, where $m_T = \sqrt{p_T^2 + m^2}$) is consistent with
 51 that of light mesons in 10–40% centrality Au+Au collisions. These results suggest that charm
 52 quarks have achieved local thermal equilibrium with the medium created in such collisions.

53 The energy dependent measurements reported here will be enhanced after STAR acquires
 54 greatly increased statistics using upgraded detectors in Phase-II of the RHIC Beam Energy
 55 Scan (BES-II).

56 2 Heavy flavor production

57 The excellent HFT resolution enabled study of charmed hadrons with even smaller decay
 58 lengths than D^0 , in particular Λ_c . Figure 3 shows the reconstructed invariant mass of $pK\pi$
 59 in 10–60% Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV, with topological cuts optimized for Λ_c .
 60 A peak at the Λ_c mass is observed clearly. STAR has performed the first measurements of
 61 centrality dependence of Λ_c production in heavy ion-collisions. It was found that Λ_c/D^0
 62 ratio increases from peripheral to central, which indicates hot medium effects. The measured
 63 STAR data point was compared to model calculations with various levels of charm quark
 64 thermalization in the medium and different coalescence implementations of the coalescence
 65 mechanism.

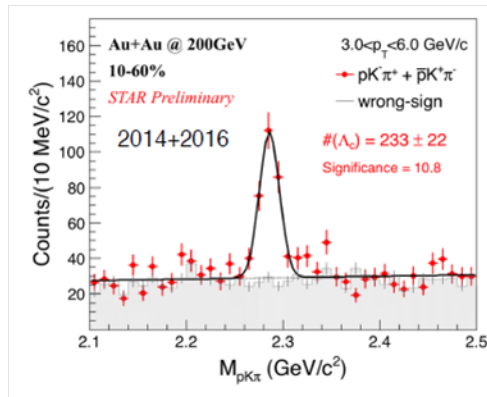


Figure 3. Reconstructed invariant mass of $pK\pi$ for 10-60% centrality Au+Au collisions from run years 2014 and 2016, representing signal of Λ_c baryons.

66 Energy loss of charm and bottom quarks is of great interest since quarks are expected
 67 to exhibit different radiative energy loss depending on their mass. The nuclear modification
 68 factor $R_{AA}(p_T) = \sigma_{in}^{pp} / \langle N_{coll}^{AA} \rangle \cdot [d^2N_{AA}/dp_T d\eta] / [d^2\sigma_{pp}/dp_T d\eta]$ of inclusive J/ψ via
 69 the di-muon channel at midrapidity in 0-40% Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV was
 70 compared to LHC results [7]. The strong suppression at high p_T obtained at RHIC indicates
 71 significant J/ψ dissociation.

72 STAR has shown ΥR_{AA} for the $2S + 3S$ states for Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV
 73 as a function of particles number N_{part} , compared to CMS data. There are indications that the
 74 suppression is weaker at $\sqrt{s_{NN}} = 200$ GeV than at $\sqrt{s_{NN}} = 2.76$ TeV, which might be due to
 75 the lower temperature at RHIC. Weaker suppression of the $1S$ state compared to the $2S+3S$
 76 states in central collisions is consistent with the picture of sequential melting.

77 3 Global polarization measurements

78 STAR has recently reported measurements of global Λ polarization in Au+Au collisions
 79 [8, 9]. The existence of an alignment between the angular momentum of a non-central collision
 80 and the spin of emitted particles was demonstrated for the first time, revealing that
 81 the fluid produced in heavy-ion collisions is by far the most vortical system ever observed.
 82 The measured effect for Λ and $\bar{\Lambda}$ hyperons show a positive polarization of the order of a
 83 few percent for energies of BES and of tens of percent for highest RHIC energy. The signal
 84 increases with decreasing collision energy, and a systematic splitting between Λ and $\bar{\Lambda}$ is
 85 observed which might indicate an additional magnetic component.

86 Spin-orbit coupling can generate a spin alignment, or polarization, along the direction of
 87 the vorticity which is on average parallel to the overall angular momentum of the system.
 88 Figure 4 shows the measured global polarization of Λ and $\bar{\Lambda}$ as a function of the collision
 89 energy for the 20–50% centrality bin in Au+Au collisions. A set of different theoretical
 90 models is argued to describe the measured polarization [10]-[12].

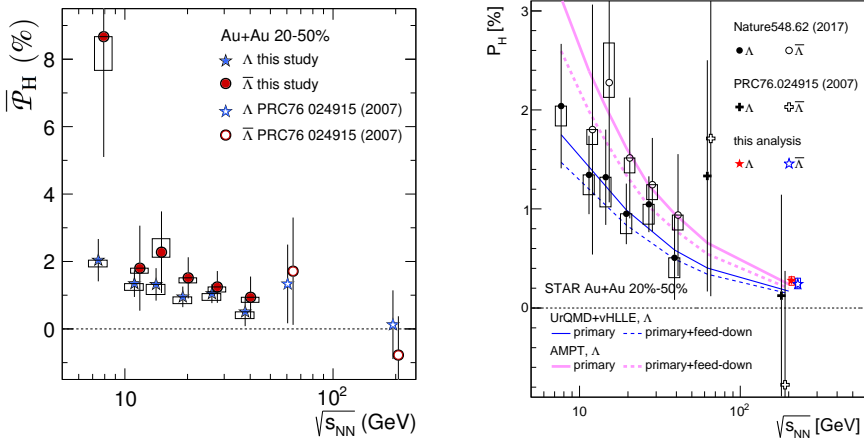


Figure 4. Global polarization \bar{P}_H (where $H = \Lambda$ or $\bar{\Lambda}$) of Λ and $\bar{\Lambda}$ as a function of the collision energy $\sqrt{s_{NN}}$ for 20-50% centrality Au+Au collisions. **Left.** The results of the measurements for BES energy range ($\sqrt{s_{NN}} < 40$ GeV) are compared with the data previously calculated for 62.4 and 200 GeV collisions, for which only statistical errors are plotted. Boxes indicate systematic uncertainties. **Right.** Thin lines show calculations from a (3+1)D cascade + viscous hydrodynamic model (UrQMD+vHLLLE) [10] and bold lines show the AMPT model calculations [11]. In the case of each model, primary Λ with and without the feed-down effect are indicated by dashed and solid lines, respectively. Open boxes and vertical lines show systematic and statistical uncertainties, respectively. Data points at 200 GeV and for $\bar{\Lambda}$ are slightly horizontally shifted for visibility.

91 Calculations for primary Λ and all Λ taking into account the effect of feed-down, from a
 92 (3+1)D viscous hydrodynamic model vHLLLE with the UrQMD initial state [10] are shown

93 for comparison. The model calculations agree with the data over a wide range of collision en-
 94 ergy, including $\sqrt{s_{NN}} = 200$ GeV within the current accuracy of experimental measurements.
 95 Calculations from a Multi-Phase Transport (AMPT) model [11] predict slightly higher pol-
 96 arization than the hydrodynamic model, but are also in good agreement with the data within
 97 uncertainties.

98 Figure 5 shows differential measurements of the polarization, versus the collision cen-
 99 trality, and hyperons transverse momentum. With the given large uncertainties, it is not clear
 100 if the polarization saturates or even starts to drop off in the most peripheral collisions (Fig.
 101 5, left). The polarization dependence on p_T is weak or absent (Fig. 5, right), considering
 102 the large uncertainties, which is consistent with the expectation that the polarization is gener-
 103 ated by a rotation of the system. No significant dependence on pseudorapidity or transverse
 104 momentum was observed. The statistical uncertainties need to be improved to reach a defini-
 105 tive conclusion on the event-by-event charge asymmetry, which is consistent with a possible
 106 contribution to the global polarization from the axial current induced by the initial magnetic
 107 field.

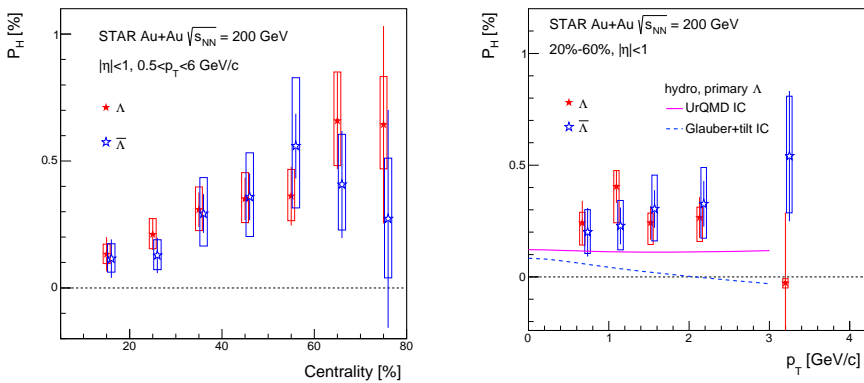


Figure 5. Left. Polarization of Λ and $\bar{\Lambda}$ as a function of the collision centrality in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV. **Right.** Polarization of Λ and $\bar{\Lambda}$ as a function of transverse momentum p_T for the 20 – 60% centrality bin in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV. Hydrodynamic model calculations for Λ with two different initial conditions are compared. Open boxes and vertical lines show systematic and statistical uncertainties. Data points for $\bar{\Lambda}$ are slightly shifted for visibility.

108 4 Results from fixed-target program

109 STAR has demonstrated the capability to operate at very low $\sqrt{s_{NN}} = 4.5$ GeV in the fixed-
 110 target mode. A test run with Au-target installed 2m from the TPC center within the beam
 111 pipe was done recently. During 3 minutes run 1M events of 0-30% centrality Au+Au were
 112 recorded. Figure 6 shows energy dependence of measured directed flow slope parameter
 113 $dv_1/dy|_{y=0}$, comparing the data from fixed-target mode to BES STAR [13] and AGS results
 114 [14]. Within the uncertainties the data are consistent with those measured in a collider exper-
 115 iments. Fig. 6, left, shows the slope parameter for lightest mesons (π^\pm , K^\pm and K_S^0) and Fig.
 116 6, right, the same for lightest baryons (p , Λ and \bar{p}).

117 Figure 7 demonstrates STAR capabilities for particle identification in fixed-target regime
 118 for strange particles. Spectra are plotted in dependence of transverse kinetic energy for K_S^0 ,
 119 left, and Λ , right.

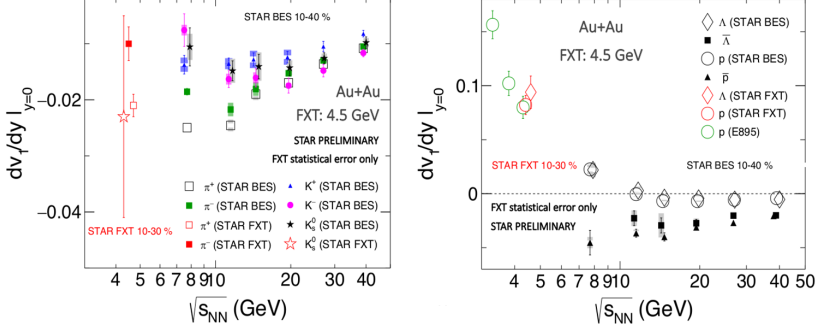


Figure 6. Directed flow slope at mid-rapidity as a function of the center-of-mass energy $\sqrt{s_{NN}}$ for the STAR fixed-target data. **Left:** $dv_1/dy|_{y=0}$ for light mesons, compared to results obtained at STAR BES-I. **Right:** $dv_1/dy|_{y=0}$ for light baryons, compared to results obtained at AGS and STAR BES-I.

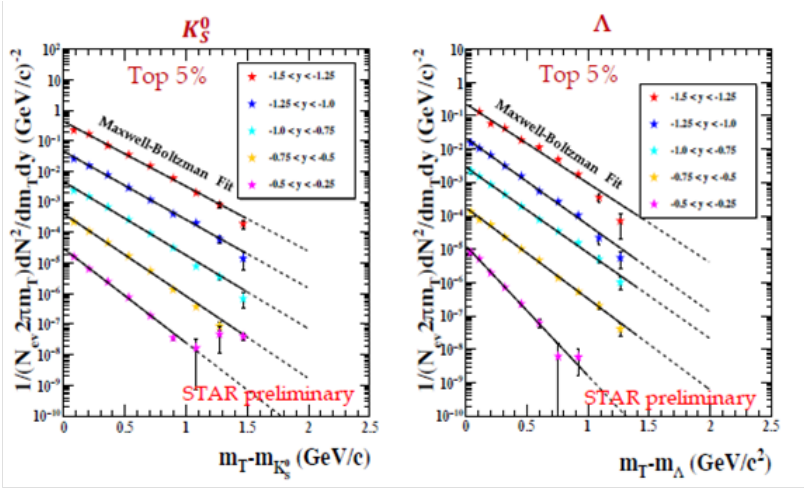


Figure 7. Spectra of strange particles obtained in fixed-target mode as a function of transverse kinetic energy. **Left:** Spectra of K_S^0 . **Right:** Spectra of Λ .

120 5 Hypertriton lifetime measurements

121 STAR has performed a precise measurement of the ${}^3_\Lambda H$ and ${}^3_\Lambda \bar{H}$ lifetime. Mesonic decay
 122 modes were used to reconstruct ${}^3_\Lambda H$ from Au+Au collisions data. All ${}^3_\Lambda H$ measurements, re-
 123 gardless of beam energy, were combined to increase the statistics. The hypertriton candidates
 124 were reconstructed from the invariant mass distributions of the daughters: ${}^3He + \pi^-$ for the
 125 2-body, and $d + p + \pi^-$ for the 3-body decay channel of ${}^3_\Lambda H$ and ${}^3\bar{He} + \pi^+$ for the 2-body, and
 126 $\bar{d} + \bar{p} + \pi^+$ for the 3-body decay channel for ${}^3_\Lambda H$ and ${}^3_\Lambda \bar{H}$, respectively.

127 A minimum χ^2 estimation is used to determine the lifetime of $\tau = 142 + 24 - 21(stat.) \pm$
 128 $31(sys.)$ ps. This lifetime is about 50% shorter than the lifetime $\tau = 263 \pm 2$ ps of a free Λ ,

129 indicating strong hyperon-nucleon interaction in the hypernucleus system. Figure 8 shows
 130 the experimentally measured hypertriton lifetime.

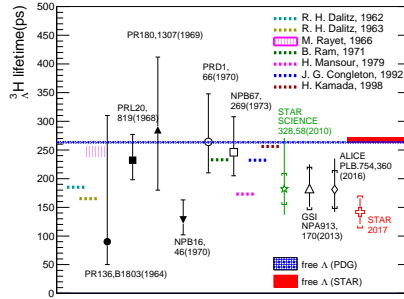


Figure 8. A summary of worldwide ${}^3\Lambda H$ lifetime experimental measurements and theoretical calculations. The two star markers are the STAR results published in 2010 and in 2018.

131 The STAR experiment will collect large datasets for Au+Au collisions during BES phase
 132 II in years 2019-2020, which will further reduce experimental uncertainties on the ${}^3\Lambda H$ life-
 133 time and will likely provide new insight into its structure.

134 6 Detector upgrades for BES-II

135 BES phase II, planned to run in years 2019 and 2020, covers the energy range $\sqrt{s_{NN}} < 20$
 136 GeV in both collider and fixed-target mods. The primary goals of BES-II are the search for
 137 the QCD critical point and for signatures of a phase transition between hadronic gas and
 138 QGP phases. In order to increase detector acceptance and particle identification capabilities,
 139 three major subsystem upgrades were proposed and to be ready till the end of 2018: Event
 140 Plane Detector (EPD), the inner TPC upgrade (iTTPC) and the endcap Time-Of-Flight detector
 141 (eTOF). The EPD was fully installed and became operational for run year 2018, replacing the
 142 Beam-Beam Counter (BBC) as a minimum-bias trigger detector. EPD allows forward mea-
 143 surements of both centrality and event plane determination reducing systematic uncertainty
 144 due to autocorrelations for midrapidity analysis. This new detector consists of two disks that
 145 are placed on both sides of the STAR detector. EPD has pseudorapidity acceptance of $2.1 <$
 146 $|\eta| < 5.1$ with 16 radial segments and 24 azimuthal segments.

147 The iTTPC will increase the acceptance of the TPC up to $|\eta| = 1.5$ increasing number of
 148 readout pad rows from 13 to 40 and maximum number of hits per track from 45 to 72. It will
 149 improve the dE/dx resolution, and will allow tracks to be reconstructed down to p_T of 60
 150 MeV/c. A single inner sector was installed for tests for run year 2018 and the full complement
 151 of 24 will be installed for run year 2019.

152 The eTOF will be installed on one side of STAR, which will extend PID capabilities at
 153 forward rapidity. Three modules of eTOF were installed behind one of the TPC sectors for
 154 run year 2018 and the full complement of 36 will be installed for run year 2019. Combin-
 155 ing all three detector systems will reduce significantly the systematic uncertainties of many
 156 observables and improve the statistical precision of STAR BES measurements.

157 7 Summary

158 STAR has presented a variety of experimental results in wide energy range $\sqrt{s_{NN}}$ from 4.5
 159 GeV to 200 GeV obtained in the fixed-target and collider mode. Systematic study of col-

160 lision energy dependence of strange particles indicates the constituent quark coalescence.
161 Results on heavy flavor production measured with HFT include directed and elliptic flow of
162 D^0 mesons and reconstruction of Λ_c baryons, which supports the idea of charm thermaliza-
163 tion at RHIC energies. Global Λ polarization measurements in Au+Au collisions previously
164 made for the energies of the BES now has been presented revealing a non-zero effect for the
165 highest RHIC energy 200 GeV. These observations can provide new insight on the vorticity
166 in heavy-ion collisions. Three detector upgrades are expected to be ready prior to the run
167 year 2019 for the upcoming Beam Energy Scan phase II. The iTPC, eTOF and EPD detec-
168 tors will improve the acceptance, particle identification capabilities, and reduce systematic
169 uncertainties for many observables.

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