

# Measurement of Intermittency for Charged Particles in Au + Au Collisions at $\sqrt{s_{NN}} = 7.7-200$ GeV from STAR

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## 1 Abstract

Local density fluctuations near the QCD critical point can be probed by intermittency analysis of scaled factorial moments in relativistic heavy-ion collisions. We report the first measurement of intermittency for charged particles in Au + Au collisions at  $\sqrt{s_{NN}} = 7.7-200$  GeV from the STAR experiment at RHIC. We observe scaling behaviors in central Au + Au collisions, with the extracted scaling exponent decreasing from mid-central to the most central Au + Au collisions. Furthermore, the scaling exponent exhibits a non-monotonic energy dependence with a minimum around  $\sqrt{s_{NN}} = 20-30$  GeV in central Au + Au collisions.

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## 19 1 Introduction

The major goal of the Beam Energy Scan (BES) at the Relativistic Heavy Ion Collider (RHIC) is to explore the phase diagram of quantum chromodynamics (QCD) [1, 2]. An important landmark of the QCD phase structure is the critical point (CP), which is the end point of first-order phase boundary between quark-gluon and hadronic phases [3]. In the thermodynamic limit, the correlation length diverges at the CP and the system becomes scale invariant and

25 fractal [4]. It is shown that the density fluctuations near the QCD critical point form a distinct  
 26 pattern of power-law or intermittent behavior in the matter produced in high energy heavy-ion  
 27 collisions [5].

28 In analogy to the critical opalescence observed in conventional matter near the critical  
 29 point, the related fractal and self-similar geometry of QCD matter will lead to local density  
 30 fluctuations that obey intermittent behavior [5]. Based on the effective action belonging to  
 31 three-dimensional Ising universality class, the intermittency of QCD matter is revealed in trans-  
 32 verse momentum spectra as a power-law (scaling) behavior of scaled factorial moment (SFM)  
 33 in heavy-ion collisions [5]. An intermittent behavior has observed in Si + Si collisions at 158A  
 34 GeV from the NA49 experiment [6]. Meanwhile, studies based on a critical Monte Carlo with  
 35 self-similar property [7] and transport model with hadronic potentials [8] demonstrate that  
 36 the intermittency could be visible in Au + Au collisions at RHIC energies.

## 37 2 Analysis Details

38 In high-energy experiments, local power-law fluctuations can be detectable through the mea-  
 39 surements of scaled factorial moment (SFM) which is defined as:

$$40 \quad F_q(M) = \frac{\langle \frac{1}{M^D} \sum_{i=1}^{M^D} n_i(n_i - 1) \cdots (n_i - q + 1) \rangle}{\langle \frac{1}{M^D} \sum_{i=1}^{M^D} n_i \rangle^q}, \quad (1)$$

41 where  $M^D$  is the number of cells in D-dimensional momentum space,  $n_i$  is the measured mul-  
 42 tiplicity in the  $i$ -th cell, and  $q$  is the order of moment.

43 Another expected power-law behavior that describes relationship between  $F_q(M)$  and  $F_2(M)$   
 44 is defined as [9, 10]:

$$45 \quad F_q(M) \propto F_2(M)^{\beta_q}. \quad (2)$$

46 Moreover, the scaling exponent  $\nu$  quantitatively describes the values of  $\beta_q$ :

$$47 \quad \beta_q \propto (q - 1)^\nu. \quad (3)$$

48 Here  $\nu$  specifies scaling (power-law) behavior of  $F_q(M)$ . According to Ginzburg-Landau (GL)  
 49 theory, the critical  $\nu$  is equal to 1.304 in entire space phase [9], while it is equal to 1.0 from  
 50 the two-dimensional Ising model [10].

51 The data reported here were obtained from Au + Au collisions at  $\sqrt{s_{NN}} = 7.7, 11.5, 14.5,$   
 52  $19.6, 27, 39, 54.4, 62.4$  and 200 GeV, which were recorded by the STAR experiment at RHIC  
 53 from 2010 to 2017. Protons ( $p$ ), antiprotons ( $\bar{p}$ ), kaons ( $K^\pm$ ) and pions ( $\pi^\pm$ ) are analyzed  
 54 as charged particles, and their identifications are carried out using the Time Projection Cham-  
 55 ber (TPC) and the Time-of-Flight (TOF) detectors. To avoid the self-correlation, the central-  
 56 ity was determined from uncorrected charged particles within a pseudo-rapidity window of  
 57  $0.5 < |\eta| < 1$ , which was chosen to be beyond the analysis window of  $|\eta| < 0.5$ .

58 To subtract the background at the level of SFM, a correlator  $\Delta F_q(M)$  is defined in terms  
 59 of original and mixed events, i.e.,  $\Delta F_q(M) = F_q(M)^{data} - F_q(M)^{mix}$  [6]. In addition, a cell-  
 60 by-cell method is proposed for efficiency correction on SFM [11]. The statistical uncertainties  
 are estimated by Bootstrap method, and the systematic uncertainties are estimated by varying  
 the experimental requirements for tracks in the TPC and TOF.

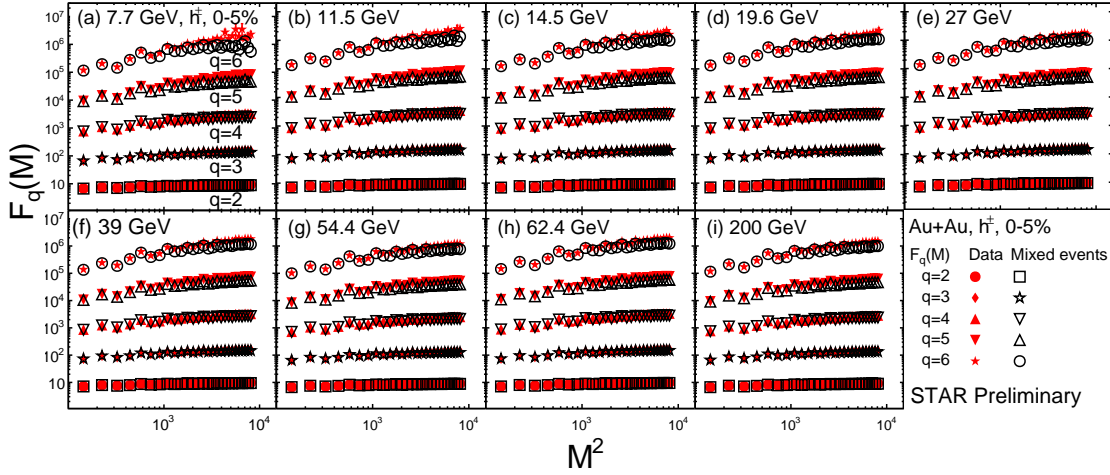


Figure 1:  $F_q(M)$  (up to sixth order) of charged particles in transverse momentum space for the most central (0-5%) Au + Au collisions at  $\sqrt{s_{NN}} = 7.7\text{-}200$  GeV in double-logarithmic scale.

### 61 3 Results and Discussion

62 Figure 1 shows  $F_q(M)^{data}$  and  $F_q(M)^{mix}$ , from the second order to the sixth order in the most  
 63 central (0-5%) collisions for various  $\sqrt{s_{NN}}$ . Based on the statistics of BES-I data,  $F_q(M)$  can be  
 64 calculated in the range of  $M^2$  from 1 to  $100^2$  and up to the sixth order ( $q=6$ ). It is observed  
 65 that  $F_q(M)^{data}$  is larger than  $F_q(M)^{mix}$  at large  $M^2$  region for various  $\sqrt{s_{NN}}$ , thus a deviation  
 66 of  $\Delta F_q(M)$  from zero is present in central Au + Au collisions.

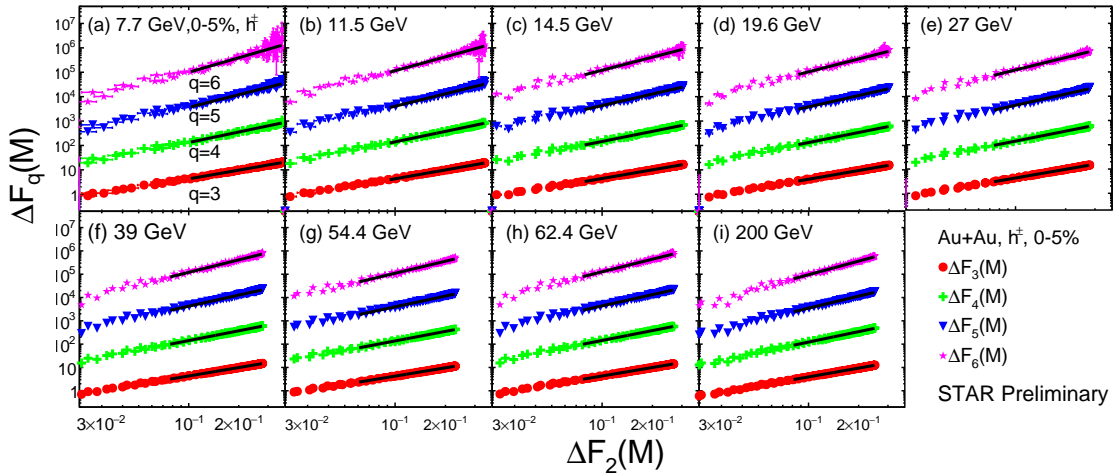


Figure 2:  $\Delta F_q(M)$  ( $q=3\text{-}6$ ) as a function of  $\Delta F_2(M)$  in the most central (0-5%) Au + Au collisions at  $\sqrt{s_{NN}} = 7.7\text{-}200$  GeV in double-logarithmic scale.

67 Figure 2 shows  $\Delta F_q(M)$  ( $q=2\text{-}6$ ), as a function of  $\Delta F_2(M)$  in the most central (0-5%)  
 68 collisions for various  $\sqrt{s_{NN}}$ . We clearly observe that the correlators  $\Delta F_q(M)$  ( $q=3\text{-}6$ ) exhibit  
 69 scaling behavior with  $\Delta F_2(M)$ .

70 The value of  $\beta_q$  is obtained through a power-law fit of Eq. (2) as shown in Figure 2, and its  
 71 statistical error is determined by the fit. Figure 3(a) shows  $\beta_q$  as a function of  $q-1$  in the most

72 central Au + Au collisions for  $\sqrt{s_{NN}} = 7.7-200$  GeV. Consistent with theoretical expectation,  
 73  $\beta_q$  also obeys a good scaling behavior with  $q$ , thus  $\nu$  can be obtained through a power-law  
 74 fit of Eq. (3). Figure 3(b) shows the extracted  $\nu$  as a function of  $\langle N_{part} \rangle$  in central Au + Au  
 75 collisions at various  $\sqrt{s_{NN}}$ . We find that  $\nu$  decreases from mid-central (30-40%) to the most  
 76 central (0-5%) Au + Au collisions.

77 Figure 4 shows the energy dependence of  $\nu$  of charged particles in central Au + Au colli-  
 78 sions at  $\sqrt{s_{NN}} = 7.7-200$  GeV. It is observed that the  $\nu$  exhibits a non-monotonic behavior on  
 79 collision energy and seems to reach a minimum around  $\sqrt{s_{NN}} = 20-30$  GeV. Higher statistics  
 80 data from BES-II will help to confirm the trend of energy dependence of  $\nu$ .

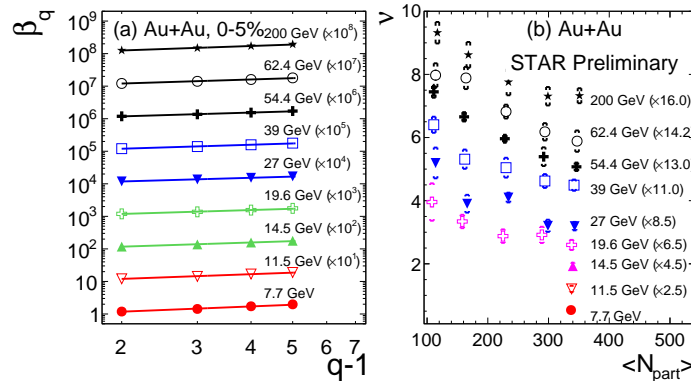


Figure 3: (a)  $\beta_q$  ( $q=3-6$ ) as a function of  $q-1$  in most central Au + Au collisions at  $\sqrt{s_{NN}} = 7.7-200$  GeV. (b)  $\nu$  as a function of  $\langle N_{part} \rangle$  in central Au + Au collisions.

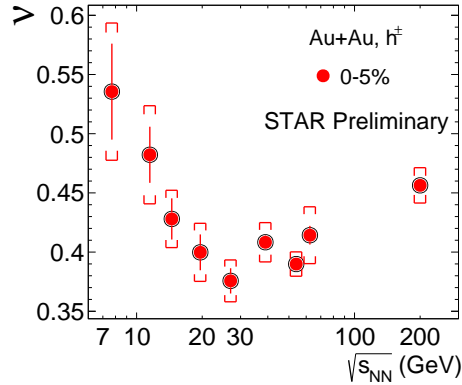


Figure 4: Energy dependence of  $\nu$  for charged particles in Au + Au collisions at  $\sqrt{s_{NN}} = 7.7-200$  GeV. The statistical and systematic errors are shown in bars and brackets, respectively.

## 81 4 Summary

82 In summary, we report the first measurements of intermittency for charged particles in Au +  
 83 Au collisions at  $\sqrt{s_{NN}} = 7.7-200$  GeV from the STAR experiment. Scaled factorial moments  
 84 (up to the sixth order) for  $p$ ,  $\bar{p}$ ,  $K^\pm$  and  $\pi^\pm$  within  $|\eta| < 0.5$ , have been measured in available  
 85 transverse momentum space. Scaling behavior is clearly visible in Au + Au collisions which

86 is consistent with theoretical predictions. The scaling exponent is related to the critical com-  
87 ponent, and we observe that it shows a non-monotonic behavior on  $\sqrt{s_{NN}}$  with a dip around  
88 20-30 GeV in the most central (0-5%) Au + Au collisions. This non-monotonic behavior needs  
89 to be understood with more theoretical inputs. With significantly improved statistics, the RHIC  
90 BES Phase-II program will allow for a more precise measurement of intermittency in heavy-ion  
91 collisions.

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