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# Higher harmonic flow of $\phi$ meson in STAR at RHIC

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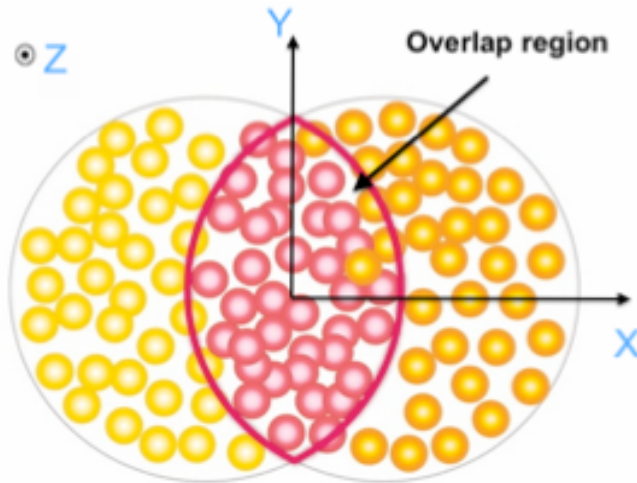
## OUTLINE:

- ✓ Introduction and Motivation
- ✓ STAR Detector and Data set
- ✓ Analysis Method
- ✓ Results
- ✓ Summary



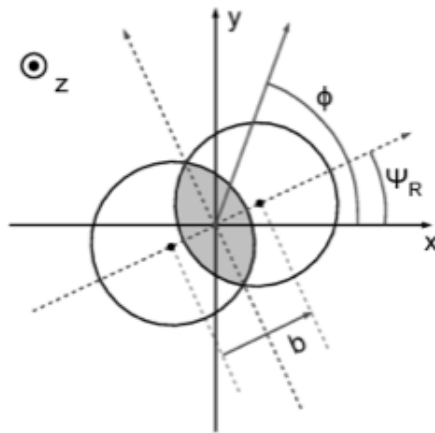
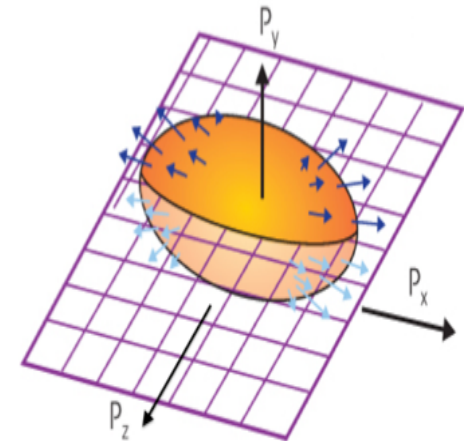
DAE HEP Symposium-2014, IIT Guwahati

# Introduction: Azimuthal anisotropy



Interactions  
↓  
Pressure (P)

$y > x \longrightarrow \frac{\partial P}{\partial x} > \frac{\partial P}{\partial y}$



$$\frac{dN}{d\phi} \propto \frac{1}{2\pi} \left[ 1 + \sum_{n=1}^{\infty} 2v_n \cos(n(\phi - \psi_R)) \right]$$

$$\phi = \tan^{-1} \left( \frac{P_y}{P_x} \right)$$

$$v_n = \langle \cos[n(\phi - \psi_R)] \rangle$$

- $\psi_R$  is the azimuthal angle of the reaction plane (spanned by impact parameter and beam direction)
- $v_2$ ,  $v_3$  and  $v_4$  are called elliptic, triangular and quadrangular flow

*A. M. Poskanzer and S. A. Voloshin, Phys. Rev. C 58, 1671(1998).*



# Motivation

- $\phi$  meson has small hadronic interaction cross section. Thus  $\phi$  meson  $v_n$  is less affected by later stage hadronic interaction. Hence  $\phi$  meson is a clean probe to study the medium created in the early stage of collisions.
- The ratios between various harmonics can be used to understand the properties of the system created in heavy-ion collisions.

## Coalescence Model

$$\frac{v_{4,M}(2p_T)}{v_{2,M}^2(2p_T)} \approx \frac{1}{4} + \frac{1}{2} \frac{v_{4,q}(p_T)}{v_{2,q}^2(p_T)}$$

Where  $v_{n,q}(p_T) = kv_{2,q}^{n/2}(p_T)$

If  $k=1$   $\frac{v_{4,M}(2p_T)}{v_{2,M}^2(2p_T)} \approx 0.75$

## Hydro Model

$$\frac{v_4}{v_2^2} = 0.5$$

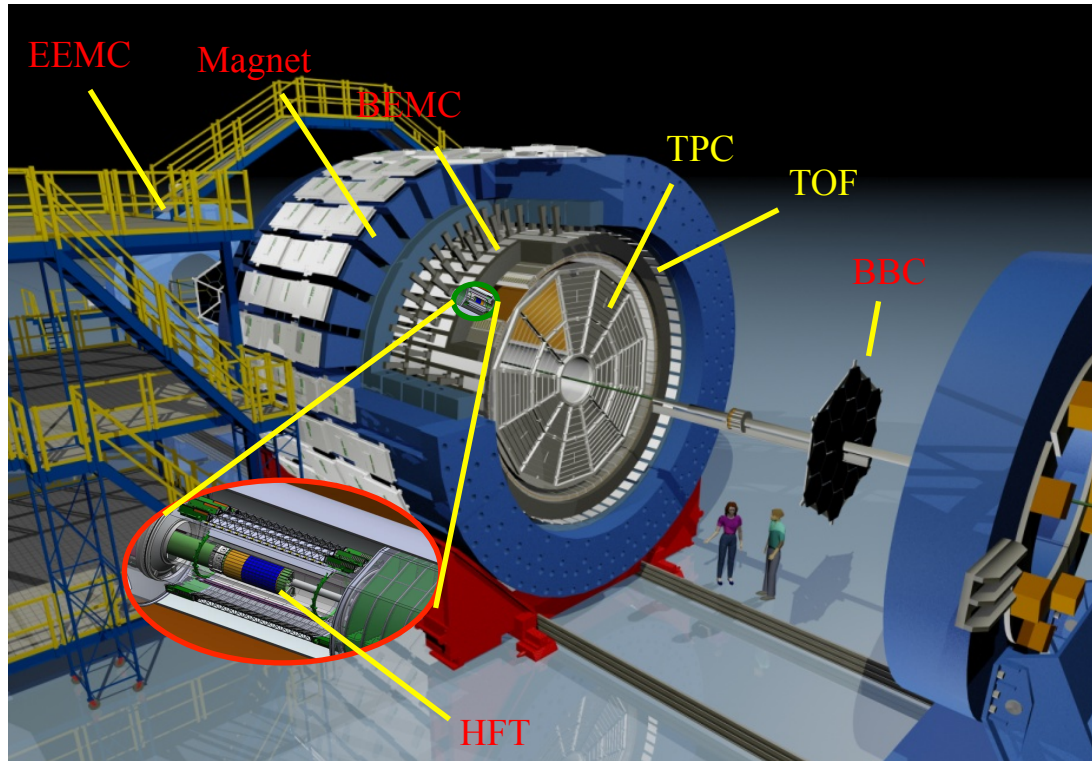
$$\frac{v_3}{v_2} = \text{Constant at high } p_T$$

*L. W. Chen et al., Phys. Rev. C 73, 044903 (2006).*

*J. Adams et al. (STAR Collaboration), Nucl. Phys. A 757, 102 (2005).*

*C. Lang et al., Eur. Phys. J. C 74 (2014) 2955.*





Magnetic field 0.5 Tesla  
 Full azimuthal coverage  
 $(0, 2\pi)$   
 $|\eta| < 1.0$  for TPC and  
 $|\eta| < 0.9$  for TOF

Data Set	Vertex Cut	Trigger	No. of events
AuAu 200 GeV (Run 11)	$ V_z  < 30$ cm $ V_r  < 2$ cm	MinBias	560 Million



# Particle Identification with STAR TPC and TOF

## ☐ TPC

- Full azimuthal coverage ( $0, 2\pi$ )
- Identifies kaon upto  $p = 0.65 \text{ GeV}/c$
- **Bethe Bloch Formula**

$$-\left\langle \frac{dE}{dx} \right\rangle \sim A \left( 1 + \frac{m^2}{p^2} \right)$$

- Particle identifies using

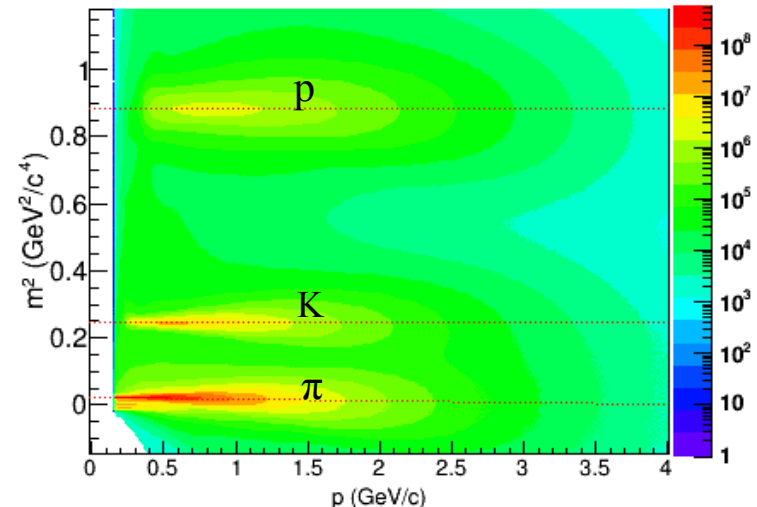
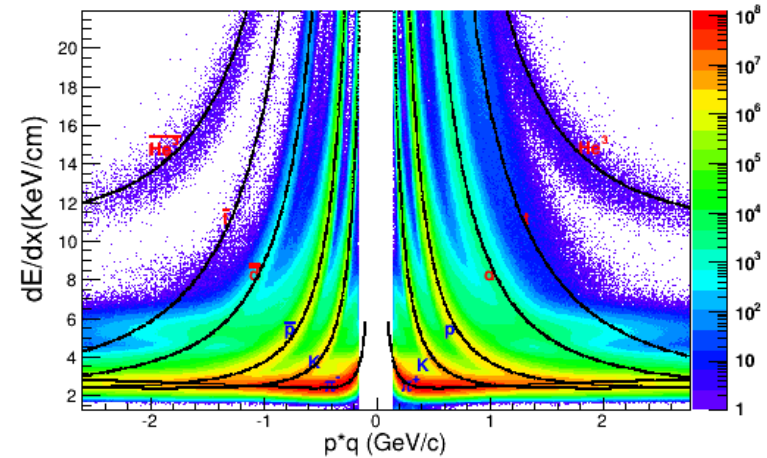
$$N\sigma = \frac{1}{R} \times \log \left( \frac{dE / dx_{\text{measured}}}{dE / dx_{\text{theory}}} \right)$$

## ☐ TOF

- Full azimuthal coverage ( $0, 2\pi$ )
- Kaon can be identified upto  $p = 1.6 \text{ GeV}/c$

- **Time of Flight**

$$\langle t \rangle = \frac{L}{\beta} \quad \frac{1}{\beta} = \sqrt{1 + \frac{m^2}{p^2}}$$



Hans Bichsel, NIM Phys Research A 562 (2006) 154–197.



- Event Plane defined as:

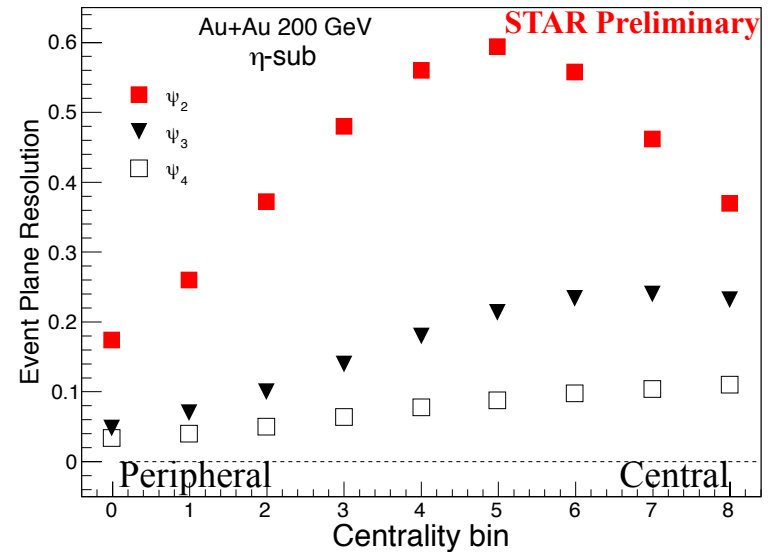
$$\Psi_n = \left( \tan^{-1} \left[ \frac{\sum_i w_i \sin(n\phi_i)}{\sum_i w_i \cos(n\phi_i)} \right] \right) / n$$

- Event Plane angle calculated in two different windows ‘west’ ( $\eta > 0.075$ ) and ‘east’ ( $\eta < -0.075$ )
- Event Plane Resolution then given by:

$$R = \sqrt{\langle \cos[n(\Psi_n^{west} - \Psi_n^{east})] \rangle}$$

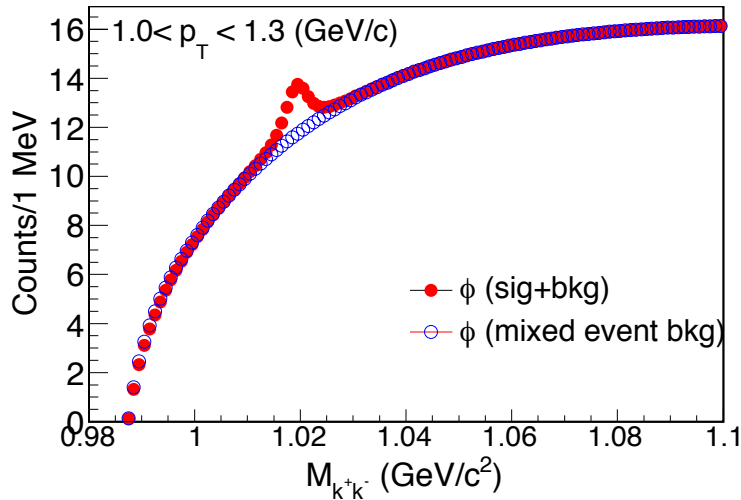
- event- by- event resolution correction

$$\langle v_n \rangle = \left\langle \frac{v_n^{obs.}}{R} \right\rangle$$



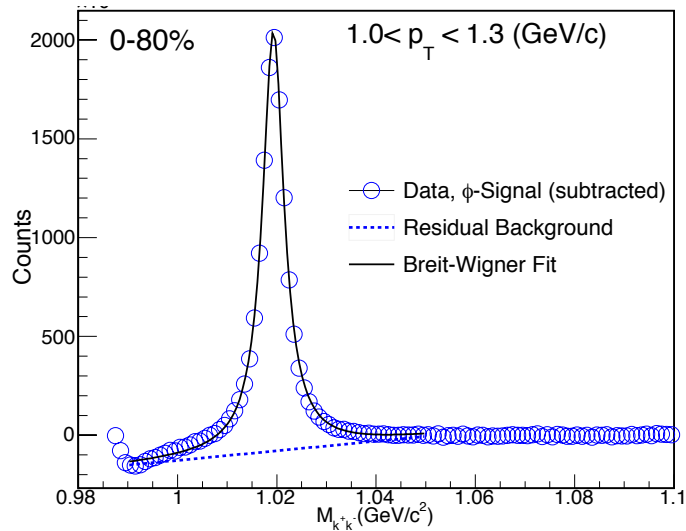
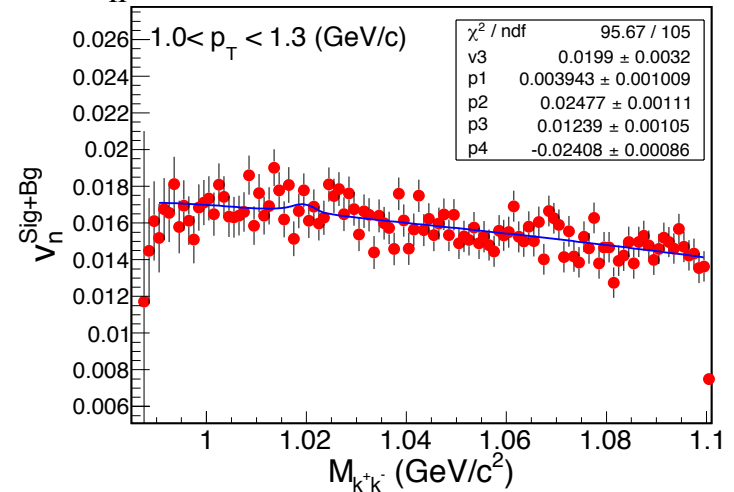
*A.M. Poskanzer & Voloshin, Phys.Rev. C58 (1998).*

# $\phi$ -meson signal extraction



- $\phi$  meson decay  $\rightarrow K^+K^-$  (B.R 48.9 %)
- Background reconstructed from mixed events
- $\phi$  signal is fitted with BW + 1<sup>st</sup> order polynomial

## $V_n$ vs Invariant mass method



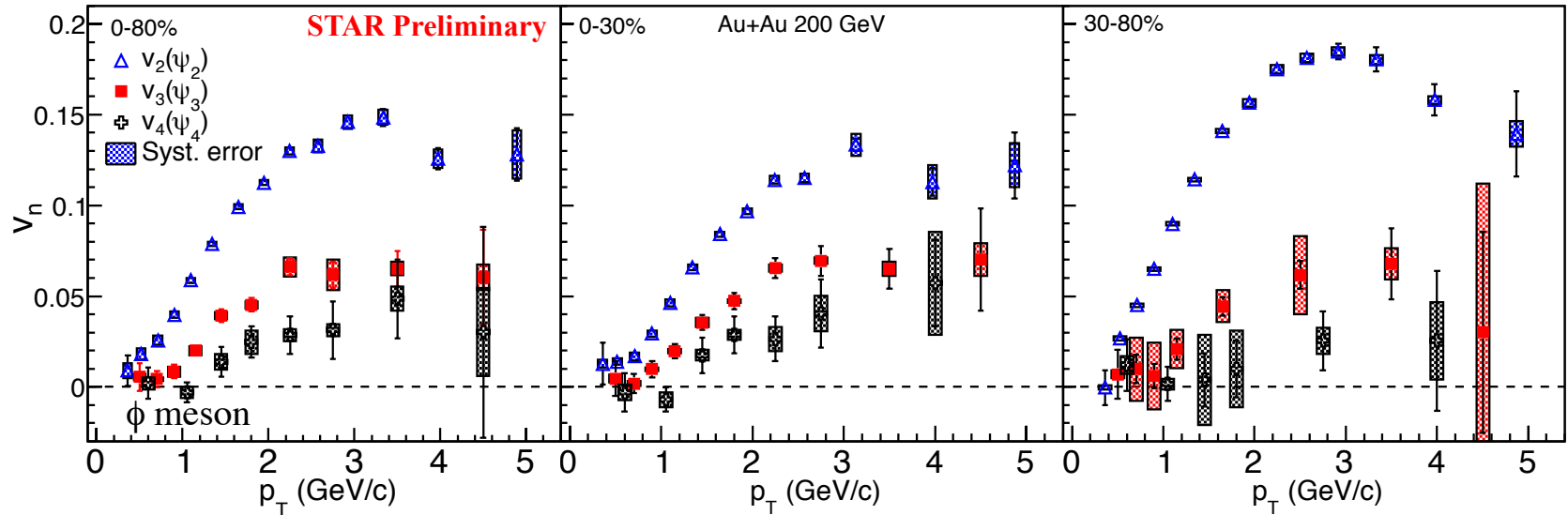
$$v_n^{\text{Sig+Bg}}(m_{\text{inv}}) = \langle \cos(n(\Phi - \Psi)) \rangle$$

$$= v_n^{\text{Sig}} \cdot \frac{\text{Sig}}{\text{Sig} + \text{Bg}}(m_{\text{inv}}) + v_n^{\text{Bg}} \cdot \frac{\text{Bg}}{\text{Sig} + \text{Bg}}(m_{\text{inv}})$$

*N. Borghini and J. Ollitrault Phys.Rev. C 70, 064905(2004).*



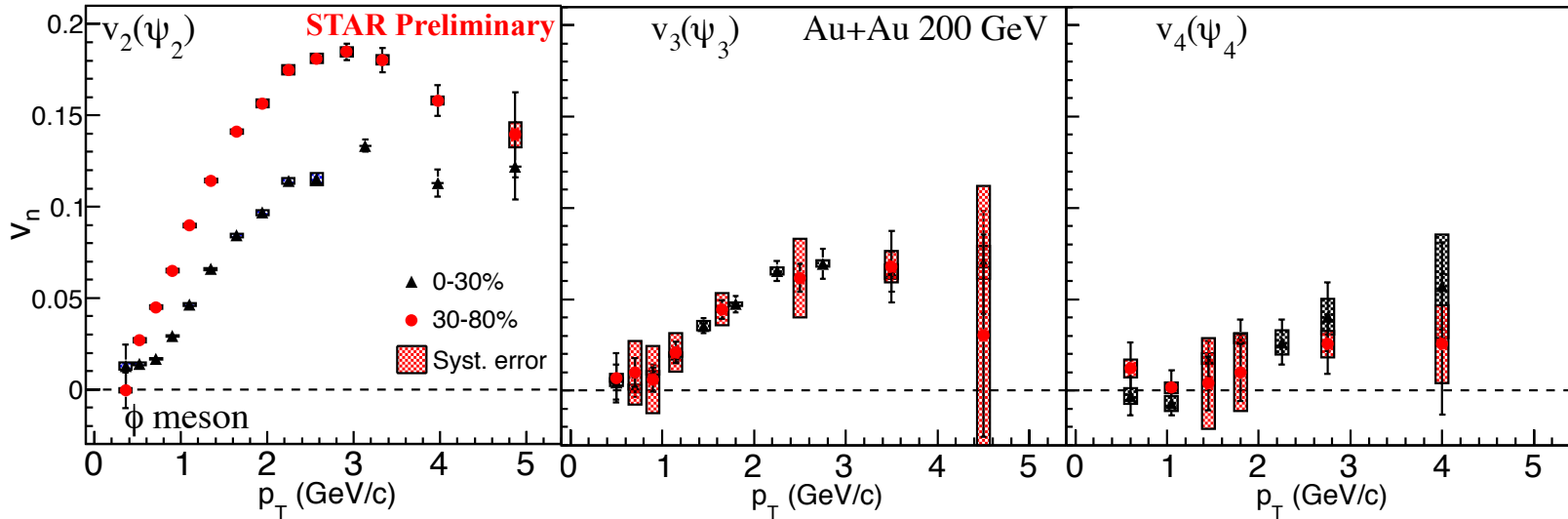
# $V_2$ , $V_3$ , $V_4$ of $\phi$ meson



- The magnitude of  $v_2(\psi_2)$  is greater than  $v_3(\psi_3)$  and  $v_4(\psi_4)$  for all centralities.
- $v_n$  increases with  $p_T$  and has a maximum value in 2-3 GeV/c

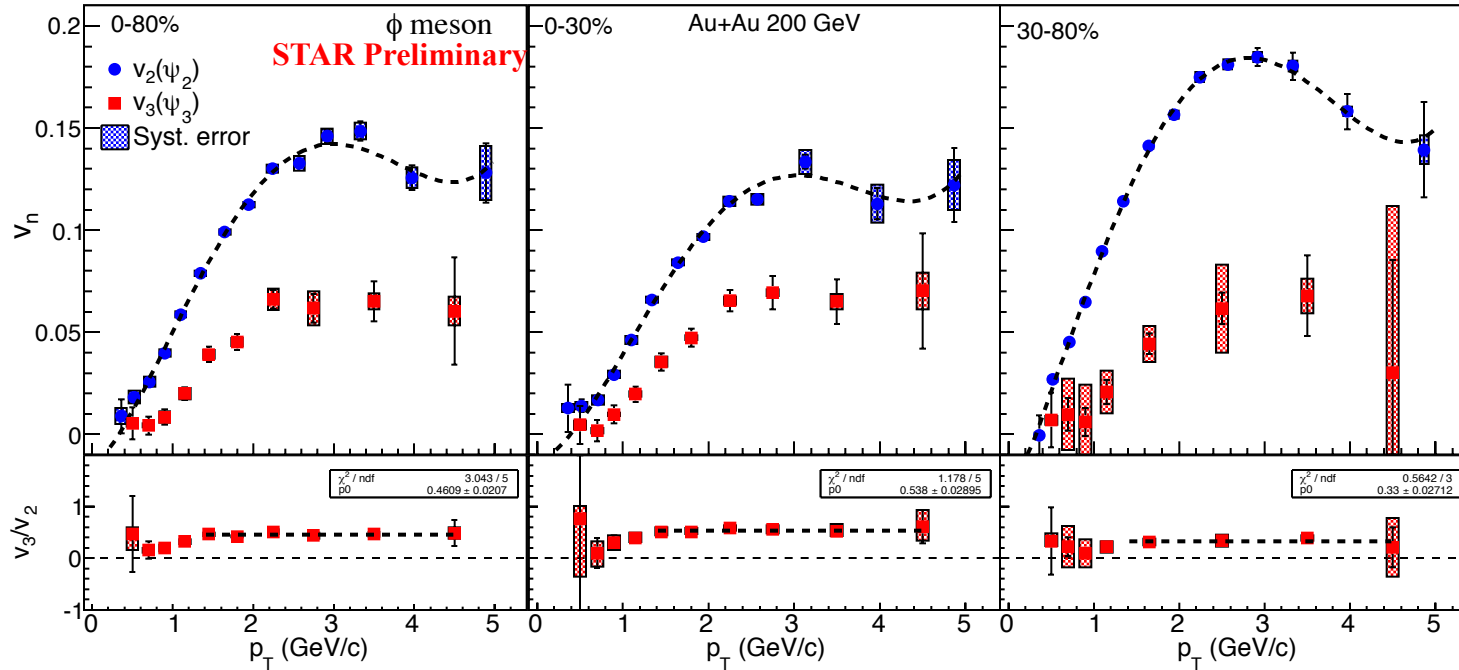


# $v_n$ : Centrality dependence



- $v_2(\psi_2)$  shows strong centrality dependence
- No centrality dependence for  $v_3(\psi_3)$  and  $v_4(\psi_4)$  within statistical uncertainties

# $v_3/v_2$ ratio



0-80%

$$v_3/v_2 = 0.46 \pm 0.02$$

0-30%

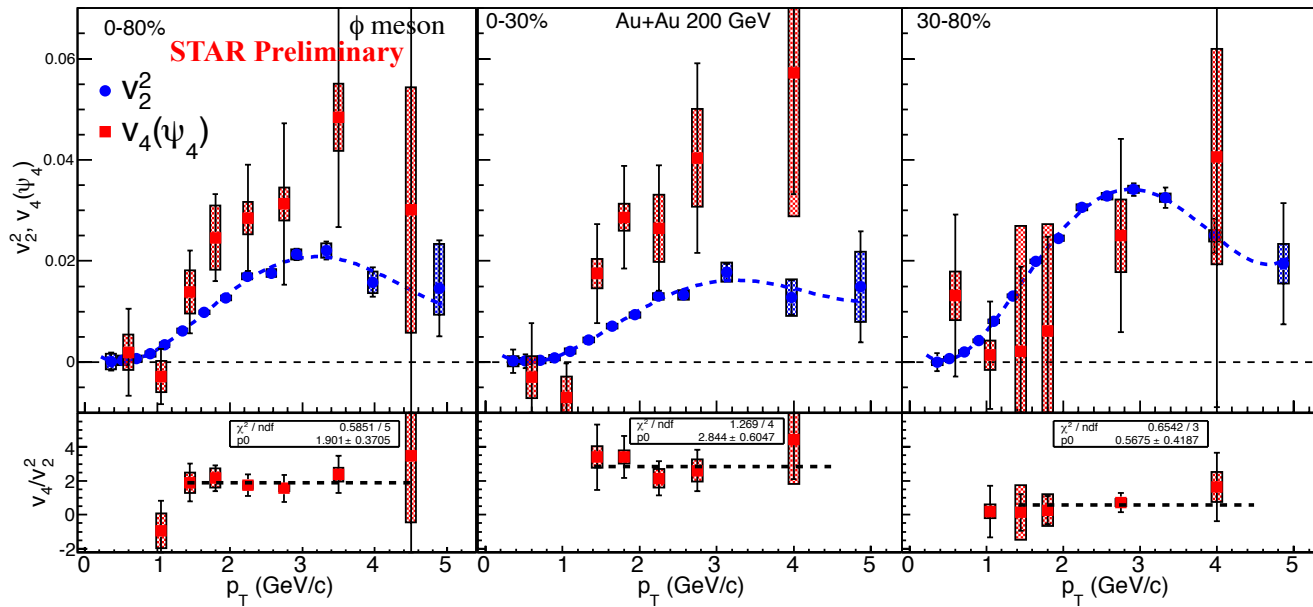
$$v_3/v_2 = 0.53 \pm 0.02$$

30-80%

$$v_3/v_2 = 0.33 \pm 0.02$$

- $v_3/v_2$  ratio is constant for  $p_T > 1.5$  GeV/c

# $v_4/v_2^2$ vs $p_T$



0-80%

$$v_4(\psi_4)/v_2^2 = 1.90 \pm 0.37$$

0-30%

$$v_4(\psi_4)/v_2^2 = 2.84 \pm 0.60$$

30-80%

$$v_4(\psi_4)/v_2^2 = 0.56 \pm 0.42$$



# Summary

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- We have presented  $v_3(p_T)$  and  $v_4(p_T)$  of  $\phi$  meson in Au+Au collisions at  $\sqrt{s_{NN}}=200$  GeV
- $v_n$  increases with  $p_T$  and has a maximum value in 2-3 GeV/c
- No centrality dependence for  $v_3(\psi_3)$  and  $v_4(\psi_4)$  within statistical uncertainties
- $v_3/v_2$  and  $v_4(\psi_4)/v_2^2$  ratios are calculated in Au+Au collisions at  $\sqrt{s_{NN}}=200$  GeV
- $v_3/v_2$  ratio is constant for  $p_T > 1.5$  GeV/c

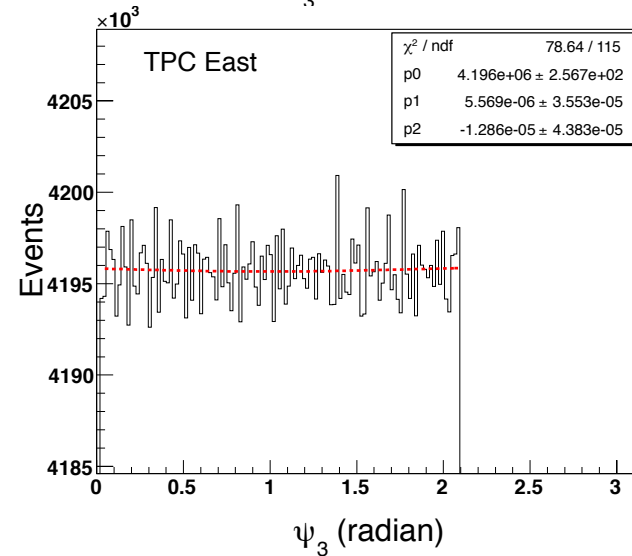
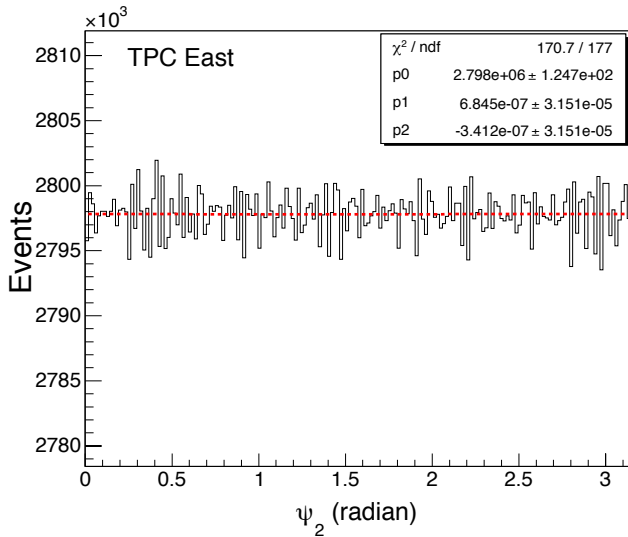
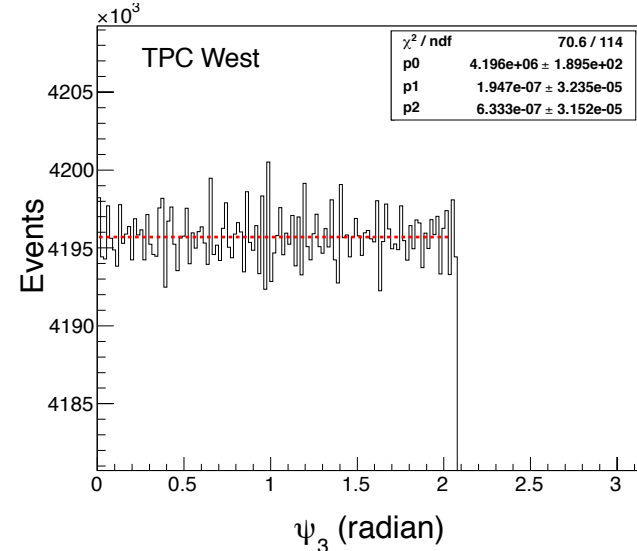
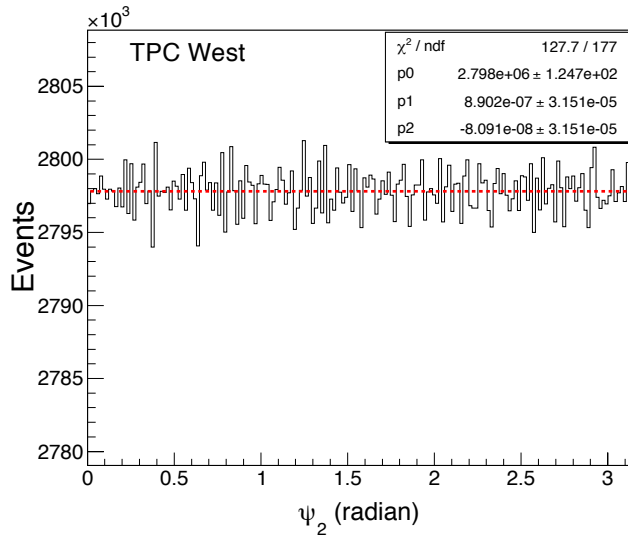


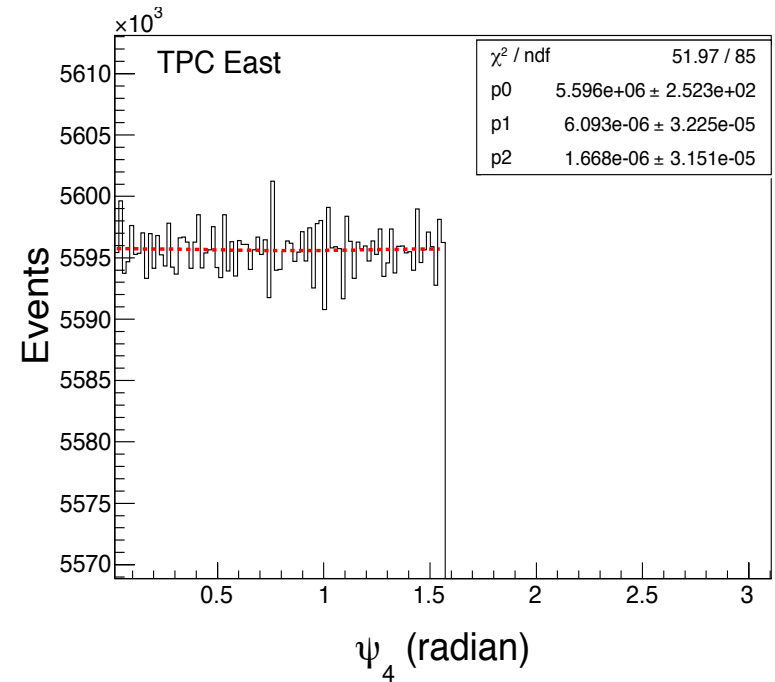
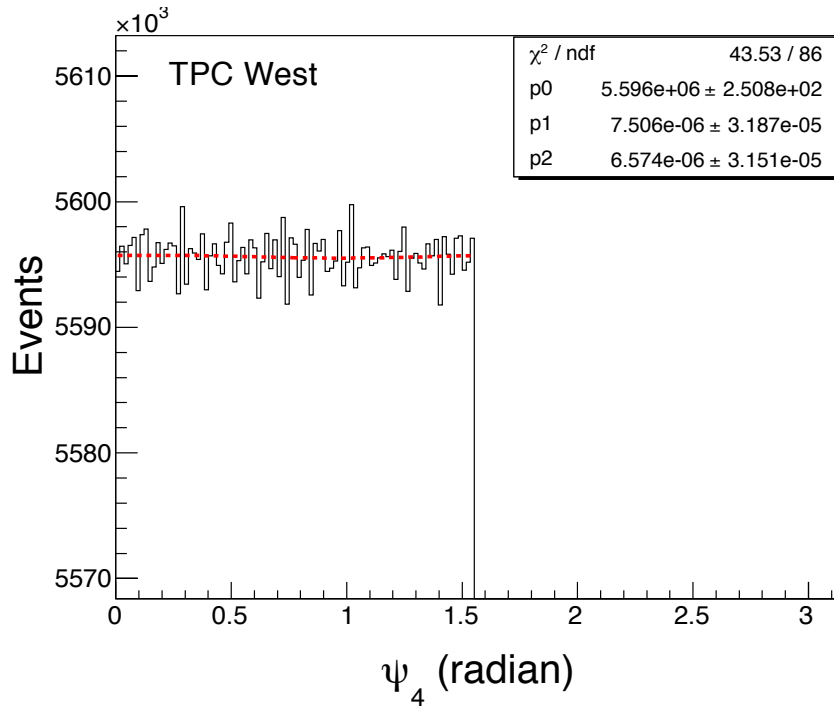


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*Thank you*







- Corrected by Recentre + Shift method
- Fitted with  $p0*(1+p1*\cos[n\Psi_n] + p2*\sin[n\Psi_n] )$
- $\eta$  gap between east & west event plane is 0.1