

1 Identified hadron spectra and baryon stopping in $\gamma + \text{Au}$
2 collisions at STAR*

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6 Photonuclear collisions are one of the simplest processes possible in a
7 heavy-ion collision. They occur when one nucleus emits a quasi-real photon
8 which interacts with the other colliding nucleus, similar to an $e + A$ collision
9 except that the photon tends to have a much smaller virtuality. Results are
10 presented for identified π^\pm , K^\pm , and $p(\bar{p})$ spectra in photonuclear collisions
11 at STAR for Au+Au collisions at $\sqrt{s_{NN}} = 54.4$ GeV. Significant baryon
12 stopping and rapidity asymmetry are observed at low transverse momen-
13 tum, which could indicate the existence of a baryon junction within the
14 nucleon, a nonperturbative Y-shaped configuration of gluons which carries
15 the baryon number and is attached to all three valence quarks. Measure-
16 ments of identified particle spectra and their rapidity dependence in $\gamma + A$
17 events will give insights into the origin of baryon stopping and also inform
18 future measurements of identified particles at the Electron-Ion Collider.

19 **1. Introduction**

20 Baryon stopping is a well-documented phenomenon in heavy-ion physics
21 where more baryons than antibaryons are observed even at midrapidity [1].
22 Because the baryon number is strictly conserved, the excess baryons must
23 come from the colliding nuclei and their lost kinetic energy allows for the
24 large multiplicity of particles produced in heavy-ion collisions. The net
25 baryon yield is often estimated through the net proton yield, which in-
26 cludes contributions from the decays of heavier baryons. One of the pro-
27 posed mechanisms for inducing baryon stopping is the baryon junction: a
28 nonperturbative Y-shaped configuration of gluons connected to all three va-
29 lence quarks. In this model, the baryon junction, rather than the valence
30 quarks, carries the baryon number. It is theorized to be an effective mech-
31 anism to stop baryons in both proton-proton and heavy-ion collisions [2],
32 but this has yet to be confirmed experimentally.

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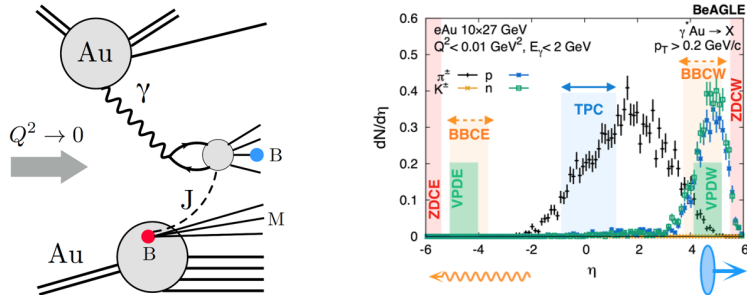


Fig. 1: (Left) A cartoon of a resolved photonuclear event [4], see text for details. (Right) The expected pseudorapidity distributions of particles in $\gamma + \text{Au}$ events simulated using the BeAGLE event generator [5, 6].

33 Heavy ions traveling at ultrarelativistic speeds produce a large flux of
 34 quasi-real photons. In a photonuclear collision, a photon from one ion col-
 35 lides with the other ion causing it to break up. As depicted in Fig.1 (Left) [4],
 36 in the majority of cases the photon fluctuates into a $q\bar{q}$ pair that scatters off
 37 of the partons in the target nucleus, also known as the resolved process [3].
 38 If the baryon number were carried by the three valence quarks, then this
 39 $q\bar{q}$ pair would not be able to stop the incoming target baryon (B) (pictured
 40 as the red dot). But if instead the junction (J) carried the baryon number,
 41 then the $q\bar{q}$ pair could slow the junction down and pull it to midrapidity.
 42 Because the junction is flavor blind, when it acquires new quarks from the
 43 vacuum this can result in a midrapidity baryon (B) (drawn as the blue dot)
 44 of a different flavor. The quarks of the initial target baryon then fragment
 45 into mesons (M), filling up the gap between midrapidity and the beam.

46 2. Selecting Photonuclear Events

47 This measurement uses data from Au+Au collisions with $\sqrt{s_{NN}} = 54.4 \text{ GeV}$
 48 which was taken at STAR in 2017. Figure 1 (Right) shows simulated
 49 $e + A$ collisions generated by BeAGLE [5, 6] where the photon virtually
 50 ($Q^2 < 0.01 \text{ GeV}^2/c^2$) and energy ($E_\gamma < 2 \text{ GeV}$) are restricted and the ion
 51 energy is set to 27 GeV in order to reproduce the kinematics of photonuclear
 52 events in this data set. Photonuclear events are selected using the minimum
 53 bias trigger, which yielded a total of 700 million events and crucially did not
 54 require events to be symmetric across the forward detectors since the $\gamma + A$
 55 collisions are highly asymmetric. The Zero Degree Calorimeters (ZDCs)
 56 select for events in which one neutron travels in the photon-going direction
 57 and multiple neutrons travel in the ion-going direction. The Beam-Beam
 58 Counters (BBCs) and Vertex Plain Detectors (VPDs) require no event ac-
 59 tivity on the photon-going side and some activity on the ion-going side.

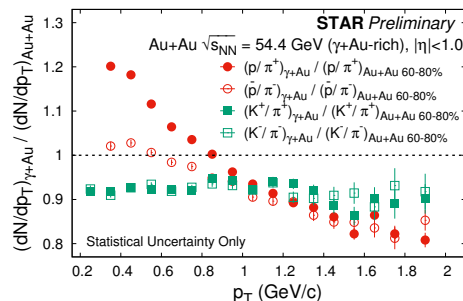


Fig. 2: Double ratios of K/π and p/π in $\gamma + A$ enriched events divided by the same ratio in peripheral $A + A$ events for $|\eta| < 1.0$.

60 Once these cuts are applied, the distribution of charged particle tracks in
 61 the Time Projection Chamber (TPC) is highly asymmetric similar to that
 62 shown in Fig. 1 (Right). As expected, the photonuclear events have very low
 63 multiplicity and so for a baseline comparison we study the particle spectra
 64 in peripheral events with a centrality of 60-80%.

65

3. Results

66 The identification of pions, kaons, and protons is done with the TPC and
 67 Time of Flight (TOF) detectors. Figure 2 shows the double ratios of K/π
 68 and p/π in the photonuclear enriched event sample divided by the same
 69 ratio in peripheral Au+Au events. This double ratio is constructed such
 70 that the tracking efficiency cancels out. The K/π double ratios are less
 71 than 1, which is consistent with there being less access to strangeness in
 72 photonuclear collisions compared to hadronic events. The p/π double ratios
 73 having a steeper slope in p_T compared to K/π is consistent with there being
 74 more radial flow in peripheral hadronic events. The striking feature of this
 75 plot is that the p/π^+ ratio is larger than the \bar{p}/π^- ratio for $p_T < 1$ GeV/c,
 76 indicating the presence of soft baryon stopping in photonuclear events.

77 Figure 3 shows the double ratios of antiparticles to particles in photonuclear
 78 events divided by the same ratio in peripheral events. The pion and
 79 kaon ratios are flat with p_T and consistent with 1, while the \bar{p}/p double ratio
 80 decreases at $p_T < 1$ GeV/c. As a function of rapidity, \bar{p}/p double ratio is
 81 slightly bigger than 1 in the backward or photon-going direction, and less
 82 than 1 and progressively smaller with increasing rapidity at low p_T . It means
 83 that not only is there soft baryon stopping present in $\gamma + A$ events, it is a
 84 stronger effect than in hadronic collisions with similar multiplicities. This
 85 is possible evidence of a baryon junction existing within the nucleon.

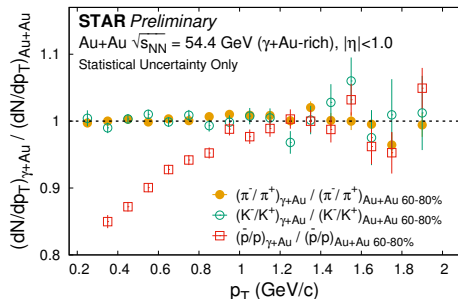


Fig. 3: The antiparticle to particle double ratio in $\gamma + A$ enriched events divided by the same ratio in peripheral $A + A$ events for $|\eta| < 1.0$.

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4. Summary

87 Identified particle spectra are measured in photonuclear collisions at the
 88 STAR experiment using Au+Au collisions at $\sqrt{s_{NN}} = 54.4$ GeV. These
 89 photonuclear events are selected by a combination of cutting on the num-
 90 ber of neutrons in the ZDC and the event activity close to the beam line.
 91 Significant baryon stopping is measured at low p_T , indicating the existence
 92 of a baryon junction within the nucleon. These results can inform future
 93 measurements of baryon stopping in electron-ion collisions at the upcoming
 94 Electron-Ion Collider, assuming the ability to cleanly identify protons and
 95 antiprotons at low momentum.

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