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# Beam-energy and collision-system dependence of the linear and mode-coupled flow harmonics from STAR

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## Abstract

Recent measurements and hydrodynamic model calculations suggest that the higher-order flow coefficients,  $v_4$  and  $v_5$ , have two contributions: a linear contribution driven by the initial-state eccentricities,  $\varepsilon_n$ , and a mode-coupled contribution derived from the lower-order eccentricity coefficients  $\varepsilon_2$  and  $\varepsilon_3$ . Measurements of these two contributions to  $v_4$  and  $v_5$  provide crucial insights to discern initial-state models and to constrain the temperature-dependent specific shear viscosity,  $\eta/s$ , of the plasma produced in heavy-ion collisions. In this work, we have employed the two-subevents cumulant technique to provide the first beam-energy and collision-system dependence of the linear and mode-coupled contributions to the higher-order flow harmonics. Our results are shown and discussed for several centrality intervals for U+U collisions at  $\sqrt{s_{NN}} = 193$  GeV, Au+Au collisions at  $\sqrt{s_{NN}} = 200$ , and 54 GeV and Cu+Au collisions at  $\sqrt{s_{NN}} = 200$  GeV. The results are compared with similar studies performed by ALICE experiments at LHC.

Keywords:

## 1. Introduction

<sup>2</sup> Ongoing investigations of the matter produced in heavy-ion collisions at the Relativistic Heavy Ion

<sup>3</sup> Collider (RHIC) and the Large Hadron Collider (LHC) indicate that an exotic state of matter called Quark-

<sup>4</sup> Gluon Plasma (QGP) is produced in these collisions. Many of these studies are aimed at understanding the

<sup>5</sup> dynamical evolution and the transport properties of the QGP [1, 2, 3].

The measurements of the azimuthal anisotropy of the particle production called anisotropic flow have been used in various studies to explain the viscous hydrodynamic response to the initial spatial distribution

<sup>8</sup> in energy density, created in the early stages of the collision [1, 2, 6].

The anisotropic flow can be described via the Fourier expansion [8] of the azimuthal angle distribution
 of the particle production,

$$\frac{dN}{d\phi} = \frac{N}{2\pi} \left( 1 + 2\sum_{n=1} V_n e^{-in\phi} \right),\tag{1}$$

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where  $V_n = v_n \exp(in\Psi_n)$  are the n<sup>th</sup> complex flow vectors,  $\Psi_n$  represent the flow vector direction, and  $v_n$  is the flow vector magnitude. The azimuthal anisotropic flow harmonic  $v_1$  is known as directed flow,  $v_2$  is the elliptic flow, and  $v_3$  is the triangular flow, etc.

To a good degree, the lower order flow harmonics  $v_2$  and  $v_3$  are linearly related to the initial-state anisotropies,  $\varepsilon_2$  and  $\varepsilon_3$ , respectively [9]. However, the higher-order flow harmonics  $v_{n>3}$  are arising from linear response to the same-order initial-state anisotropies along with nonlinear response to the lower-order eccentricities  $\varepsilon_2$  and/or  $\varepsilon_3$  [5]. Consequently, the full benefit of the higher-order flow harmonics for  $\eta/s$ extraction [12] benefits form a robust separation of their linear and nonlinear contributions.

<sup>19</sup> The higher-order flow harmonic  $V_4$  can be expressed as,

$$V_4 = V_4^{\text{Linear}} + V_4^{\text{Nonlinear}}, \qquad (2)$$

$$V_4^{\text{Nonlinear}} = \chi_{4,22} V_2 V_2, \tag{3}$$

where  $\chi_{4,22}$  is the nonlinear response coefficients. The value of  $\chi_{4,22}$  constrains the magnitude of  $V_4^{\text{Nonlinear}}$ , also the magnitude of  $V_4^{\text{Nonlinear}}$  encodes the correlations between the flow symmetry planes  $\Psi_2$  and  $\Psi_4$ .

In this work, we employ the multiparticle cumulant method [14] to measure the p<sub>T</sub>-integrated inclusive, nonlinear and linear higher-order flow harmonic  $v_4$ , in collisions of U+U at  $\sqrt{s_{NN}} = 193$  GeV, Cu+Au at  $\sqrt{s_{NN}} = 200$  GeV and Au+Au at at several beam energies.

#### 25 **2. Method**

The STAR data were analyzed with the multi-particle cumulant technique [13, 14]. The framework for the standard cumulant method is discussed in Ref. [13]; its extension to the subevents method is reported in Refs. [14]. In the two-subevent method and to minimize the non-flow correlations, the cumulants are constructed from two-subevents which are separated in  $\eta$ . Thus, the constructed two- and multi-particle correlations can be written as:

$$v_n = \langle \langle \cos(n(\varphi_1^A - \varphi_2^B)) \rangle \rangle^{1/2},$$

$$C_{n+m,n,m} = \langle \langle \cos((n+m)\varphi_1^A - n\varphi_2^B - m\varphi_3^B) \rangle \rangle,$$

$$\langle v_n^2 v_m^2 \rangle = \langle \langle \cos(n\varphi_1^A + m\varphi_2^A - n\varphi_3^B - m\varphi_4^B) \rangle \rangle,$$
(4)

where,  $\langle \langle \rangle \rangle$  represents the average over all particles then average over events, *k*, *n* and *m* are harmonic numbers and  $\varphi_i$  are the ith particles azimuthal angle. For the two-subevent method, subevent A and subevent B are required to have a minimum  $\Delta \eta > 0.6$  separation, i.e.  $\eta_A > 0.3$  and  $\eta_B < -0.3$ .

Using Eq.(4) the linear and nonlinear modes in the higher order anisotropic flow harmonic,  $v_4$ , can be expressed as,

$$v_4^{\text{Nonlinear}} = C_{4,22} / \sqrt{\langle v_2^2 v_2^2 \rangle},$$

$$v_4^{\text{Linear}} = \sqrt{(v_4^{\text{Inclusive}})^2 - (v_4^{\text{Nonlinear}})^2}.$$
(5)

Equation (5) assumes that the linear and nonlinear contributions in  $v_4$  are independent, which is a correct approach if the correlation between the lower  $v_n$  (n = 2, 3) and higher order flow coefficients (n > 3) is weak.

#### 38 3. Results

The centrality dependencies of the inclusive, linear and nonlinear  $v_4$  in the  $p_T$  range from 0.2 to 4.0 GeV/*c* for Au+Au collisions at  $\sqrt{s_{NN}} = 200$  GeV are shown in Fig. 1. Our study indicates that the  $v_4^{\text{Linear}}$ depends weakly on the collision centrality and it dominates over the nonlinear contribution to the inclusive  $v_4$  in central collisions.

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Fig. 1. The inclusive, nonlinear and linear higher-order flow harmonic  $v_4$  using the two subevent cumulant method as a function of centrality in the  $p_T$  range from 0.2 to 4.0 GeV/care shown. The respective systematic uncertainties are shown as open boxes.



Fig. 2. The inclusive, nonlinear and linear higher-order flow harmonic  $v_4$  using the two subevent cumulant method as a function of centrality in the  $p_T$  range from 0.2 to 4.0 GeV/*c* are shown. The respective systematic uncertainties are shown as open boxes. The results are compared with the LHC measurements in the  $p_T$  range from 0.2 to 5.0 GeV/*c* for Pb+Pb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV [17].



**Fig.** 3. The inclusive, nonlinear and linear higher-order flow harmonic  $v_4$  for U+U collisions at  $\sqrt{s_{NN}} = 193$  GeV, and Au+Au and Cu+Au collisions at  $\sqrt{s_{NN}} = 200$  GeV are shown. The presented results are measured using the two subevent cumulant method as a function of centrality in the p<sub>T</sub> range from 0.2 to 4.0 GeV/*c*. The respective systematic uncertainties are shown as open boxes.

Figure 2 compares the centrality dependence of the inclusive, linear and nonlinear  $v_4$  in the  $p_T$  range from 0.2 to 4.0 GeV/*c* for Au+Au collisions at  $\sqrt{s_{NN}} = 200$  and 54 GeV. At both presented energies we observe that the linear mode of  $v_4$  has a weak centrality dependence and it's the dominant contribution to the inclusive  $v_4$  in central collisions. The preliminary results are compared with similar LHC measurements in the  $p_T$  range from 0.2 to 5.0 GeV/*c* for Pb+Pb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV [17]. The observed difference in the magnitude of  $v_4$  between Au+Au collisions at  $\sqrt{s_{NN}} = 200$  and 45 GeV and Pb+Pb collisions at  $\sqrt{s_{NN}}$ = 2.76 TeV could be driven by the difference in the viscous effects between those energies.

The preliminary results for U+U collisions at  $\sqrt{s_{NN}} = 193$  GeV, and Au+Au and Cu+Au collisions at

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 $\sqrt{s_{NN}} = 200 \text{ GeV}$  are shown in Fig. 3. The magnitudes and trends for both inclusive and nonlinear  $v_4$  show

a weak system dependence, albeit with more visible differences between Cu+Au and Au+Au than between
 U+U and Au+Au.

# 54 4. Summary

In summary, we have used the cumulant method to measure the inclusive, linear and nonlinear  $v_4$  as a function of collision centrality in U+U collisions at  $\sqrt{s_{NN}} = 193$  GeV, Cu+Au at  $\sqrt{s_{NN}} = 200$  GeV and Au+Au at several beam energies. The measurements show the expected characteristic dependencies of the inclusive, linear and nonlinear  $v_4$  on centrality, system size and beam energy. Our study indicates that the linear contribution to the inclusive  $v_4$  dominates over the nonlinear contribution in central collisions for all presented energies and systems. These newly presented measurements may give extra constraints to test different initial-state models and to assist an accurate extraction of the QGP specific shear viscosity.

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