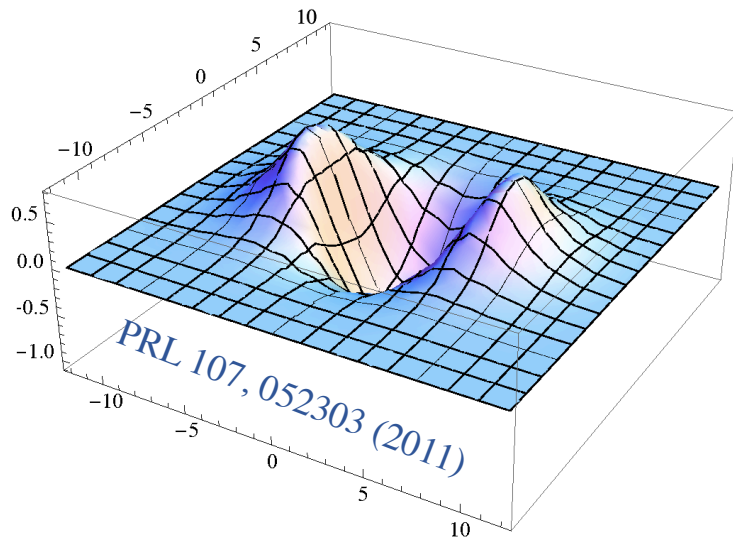


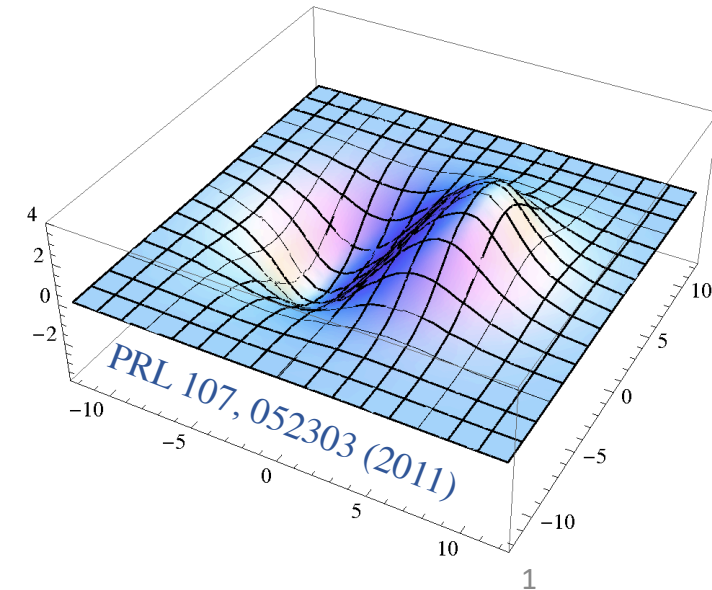
Characterization of Quadrupole charge separation at 200 GeV; Implications for the CMW



Niseem Magdy
University of Illinois at Chicago
niseem@uic.edu

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Outline



N. Magdy, et al, Phys.Lett.B 811 (2020) 135986
N. Magdy, et al, Phys.Rev.C 98 (2018) 6, 061902
N. Magdy, et al, Phys.Rev.C 97 (2018) 6, 061901

➤ Introduction

- ✓ Anomalous transport
- ✓ Prior correlators

➤ The $R_{\psi m}^{(d)}$ correlator

- ✓ $R_{\psi m}^{(d)}$ Sensitivity to Quadrupole charge separation

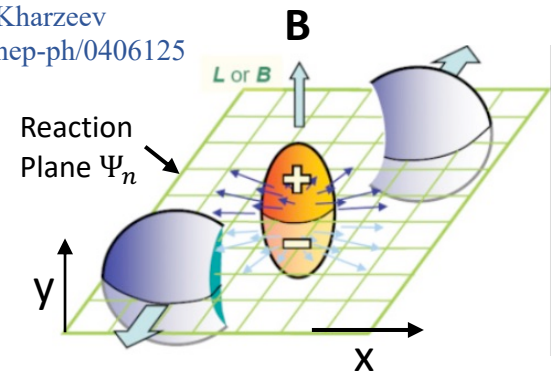
➤ Experimental Results

- ✓ Recent STAR measurements

➤ Conclusions

❖ Introduction:
 ✓ Anomalous Transport in the QGP

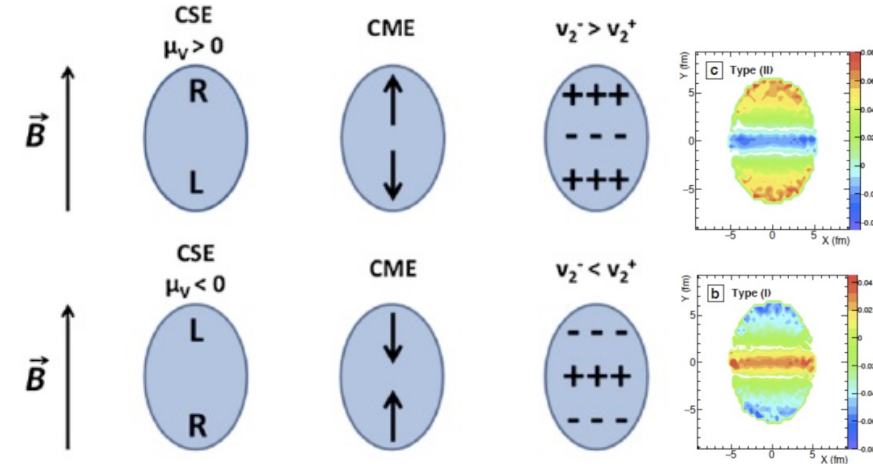
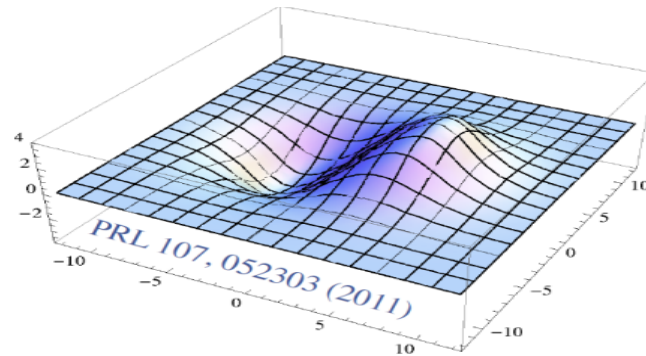
CME



CMW

The interplay between the CSE and CME can lead to the production of a gapless collective mode or Chiral Magnetic Wave (CMW)
 ✓ Stems from the coupling between the density waves of the electric and chiral charges

Dmitri E. Kharzeev and Ho-Ung Yee,
 Phys. Rev. D83, 085007 (2011)



The CMW transports positive (negative) charges out-of-plane and negative (positive) charges in-plane to form an electric quadrupole.

➤ The detection and characterization of both the dipole and quadrupole charge separations are paramount

$$\frac{dN}{d\Delta\phi} \propto [1 \pm 2a_1 \sin(\Delta\phi) + \dots]$$

$$\tilde{a}_1 = \langle a_1^2 \rangle^{1/2} \propto \mu_5 B$$

❖ Introduction:

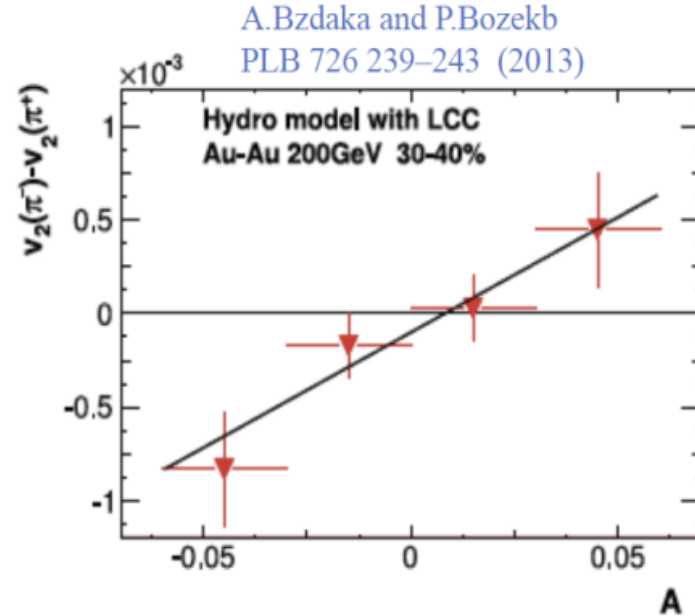
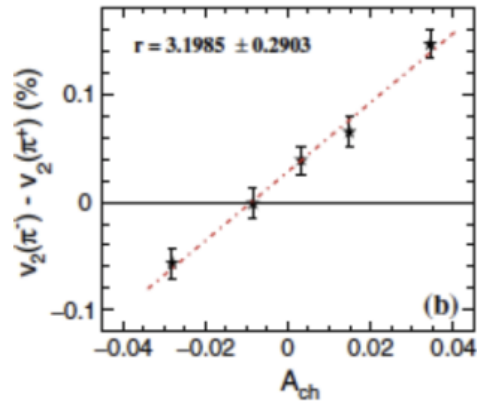
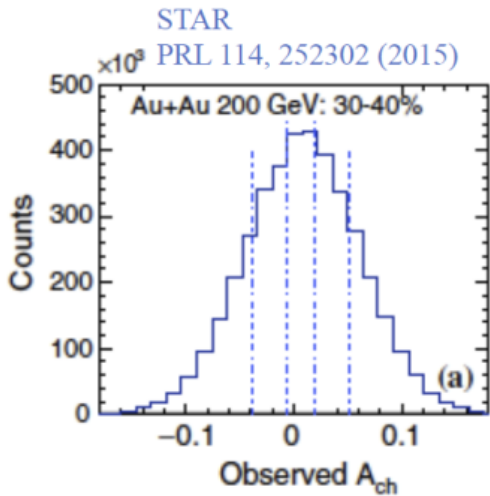


Prior quadrupole charge separation measurements

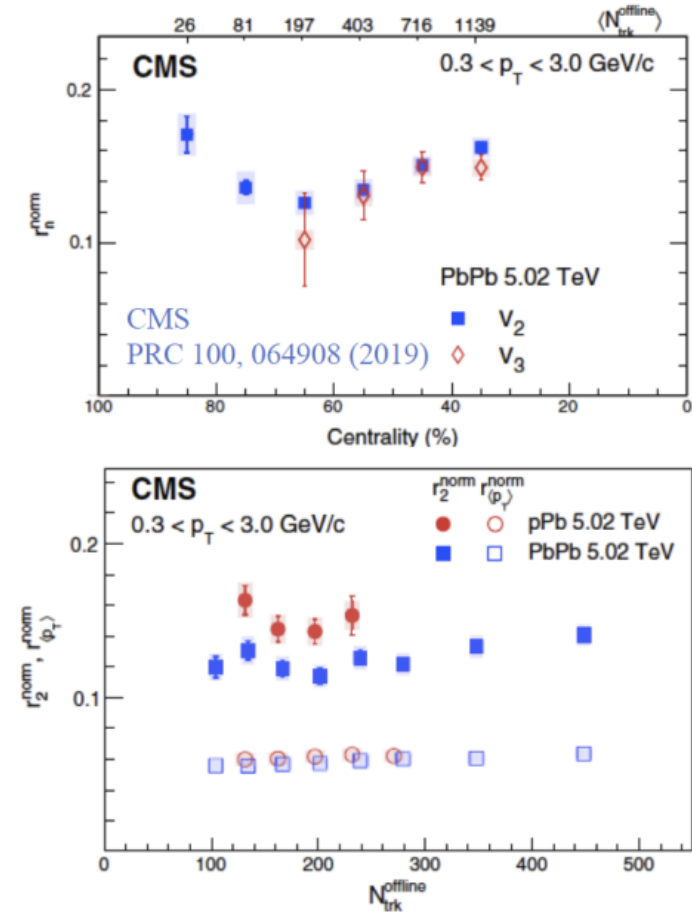
A pervasive approach is to measure the elliptic flow difference between negatively- and positively charged particles as a function of charge asymmetry

$$\Delta v_2 \equiv v_2^- - v_2^+ \simeq r A_{\text{ch}}$$

$$A_{\text{ch}} = \frac{(N^+ - N^-)}{(N^+ + N^-)}$$



A purely background scenario can give a similar dependence



- Similar slope parameter for:
 - ✓ v_2 and v_3
 - ✓ Small and large systems

Background can account for a part, if not all, of the observed charge separation signal with this correlator:

- ✓ Could one make a discerning measurements with a different correlator?

❖ The extended $R_{\Psi_2}^{(d)}$ correlator

$d = 2$ for quadrupole charge separation

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 N. Magdy, et al. Phys.Rev.C 98 (2018) 061902
 N. Magdy, et al. Phys.Rev.C 97 (2018) 061901
 D. Shen, et al. Phys.Rev.C 100 (2019) 064907

$$R_{\Psi_2}^{(d)}(\Delta S) = \frac{C_{\Psi_2}(\Delta S_d)}{C_{\Psi_2}^{\perp}(\Delta S_d)}$$

$$C_{\Psi_2}(\Delta S) = \frac{N(\Delta S_d)}{N(\Delta S_d)_{sh}}$$

$$\Delta\varphi = \varphi - \Psi_2$$

$$C_{\Psi_2}^{\perp}(\Delta S) = \frac{N(\Delta S_d)^{\perp}}{N(\Delta S_d)_{sh}^{\perp}}$$

Sensitive to charge separation
(Signal and Background):

$$N(\Delta S_d) = N(\langle S_{\Psi_2}^+ \rangle_d - \langle S_{\Psi_2}^- \rangle_d)$$

$$N(\Delta S_d)^{\perp} = N(\langle S_{\Psi_2}^+ \rangle_d^{\perp} - \langle S_{\Psi_2}^- \rangle_d^{\perp})$$

$$\langle S_{\Psi_2}^+ \rangle_d = \frac{\sum_1^p w_p \sin[d(\Delta\varphi)]}{w_p}$$

w_i : charge-dependent detector acceptance

$$\langle S_{\Psi_2}^+ \rangle_d^{\perp} = \frac{\sum_1^p w_p \cos[d(\Delta\varphi)]}{w_p}$$

$$\langle S_{\Psi_2}^- \rangle_d = \frac{\sum_1^n w_n \sin[d(\Delta\varphi)]}{w_n}$$

p/n : number of positive/negative hadrons per event

$$\langle S_{\Psi_2}^- \rangle_d^{\perp} = \frac{\sum_1^n w_n \cos[d(\Delta\varphi)]}{w_n}$$

Shuffling of charges within an event breaks the charge separation

sensitivity:

$$N(\Delta S_d)_{sh} = N(\langle S_{\Psi_2}^+ \rangle_d - \langle S_{\Psi_2}^- \rangle_d)_{sh}$$

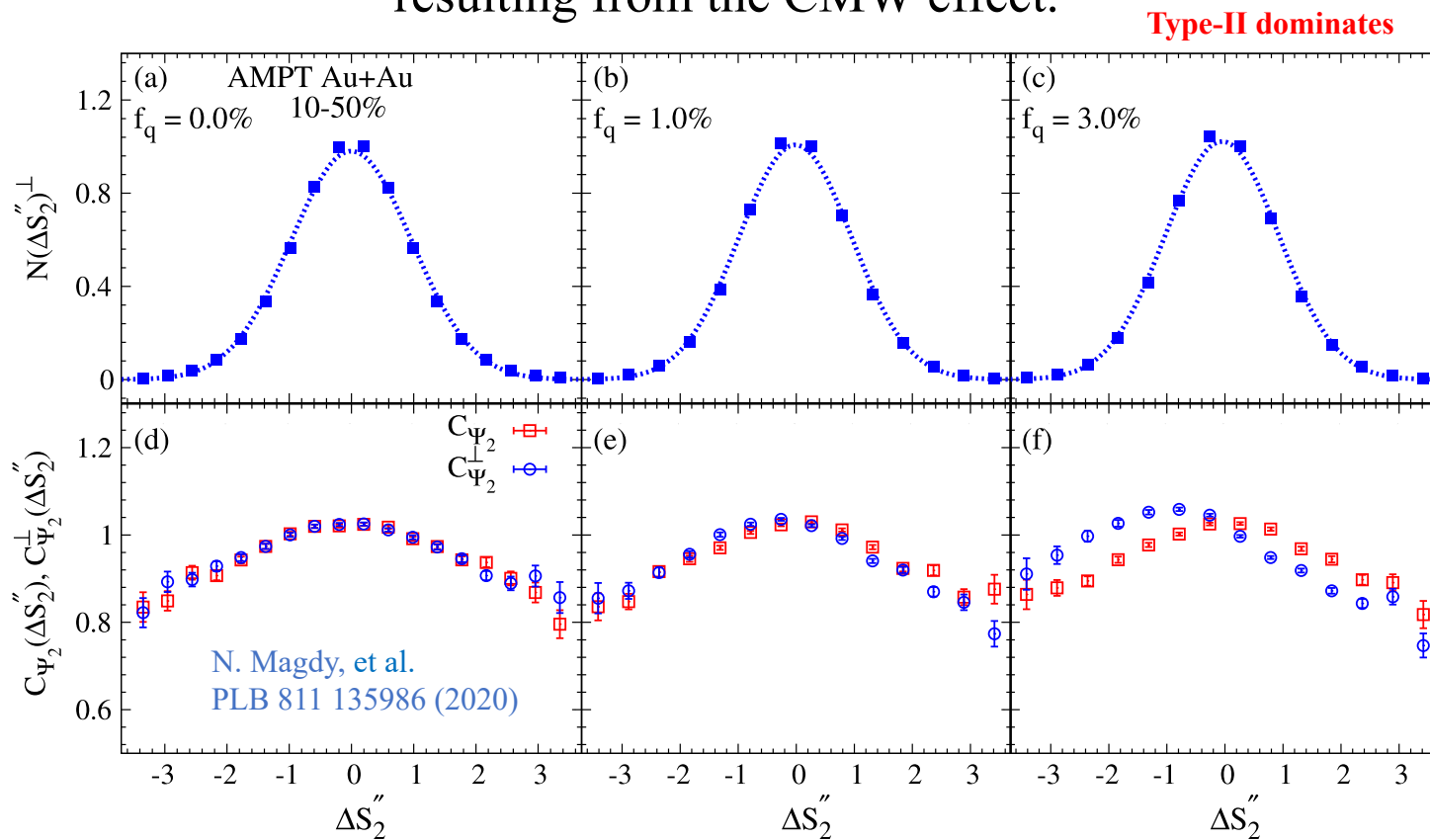
$$N(\Delta S_d)^{\perp} = N(\langle S_{\Psi_2}^+ \rangle_d^{\perp} - \langle S_{\Psi_2}^- \rangle_d^{\perp})_{sh}$$

❖ The $R_{\Psi_m}^{(2)}(\Delta S_2)$ [quadrupole correlator]



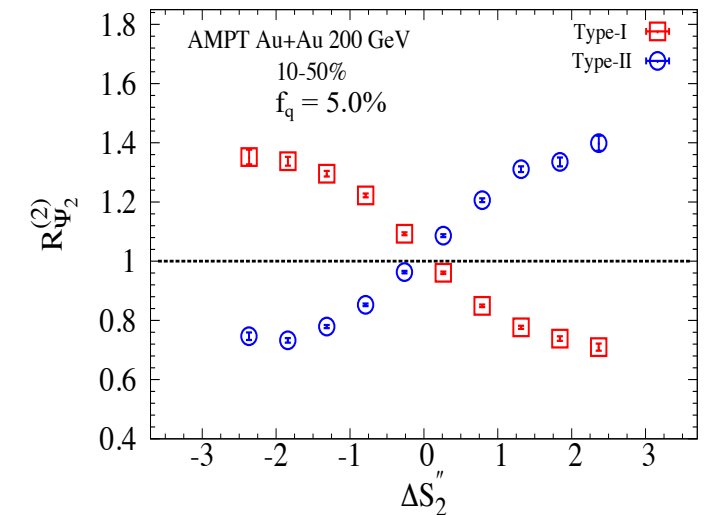
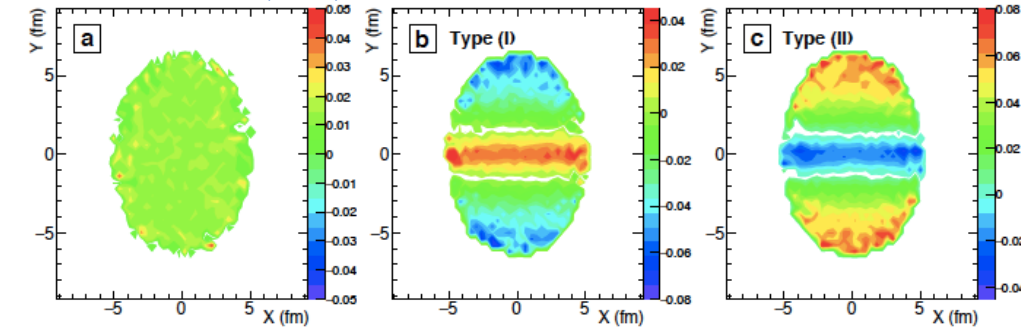
N. Magdy, et al. Phys.Lett.B 811 (2020) 135986
 D. Shen, et al. Phys.Rev.C 100 (2019) 064907

- The AMPT used to simulate the charge quadrupole moment resulting from the CMW effect.



The fraction f_q , characterize the strength of the quadrupole charge separation signal.

D. Shen, et al.
 PRC 100, 064907 (2019)

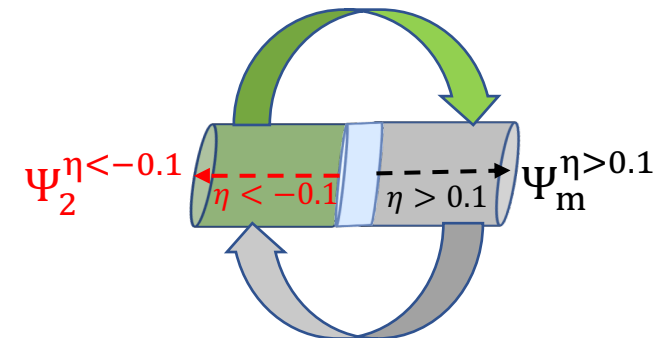
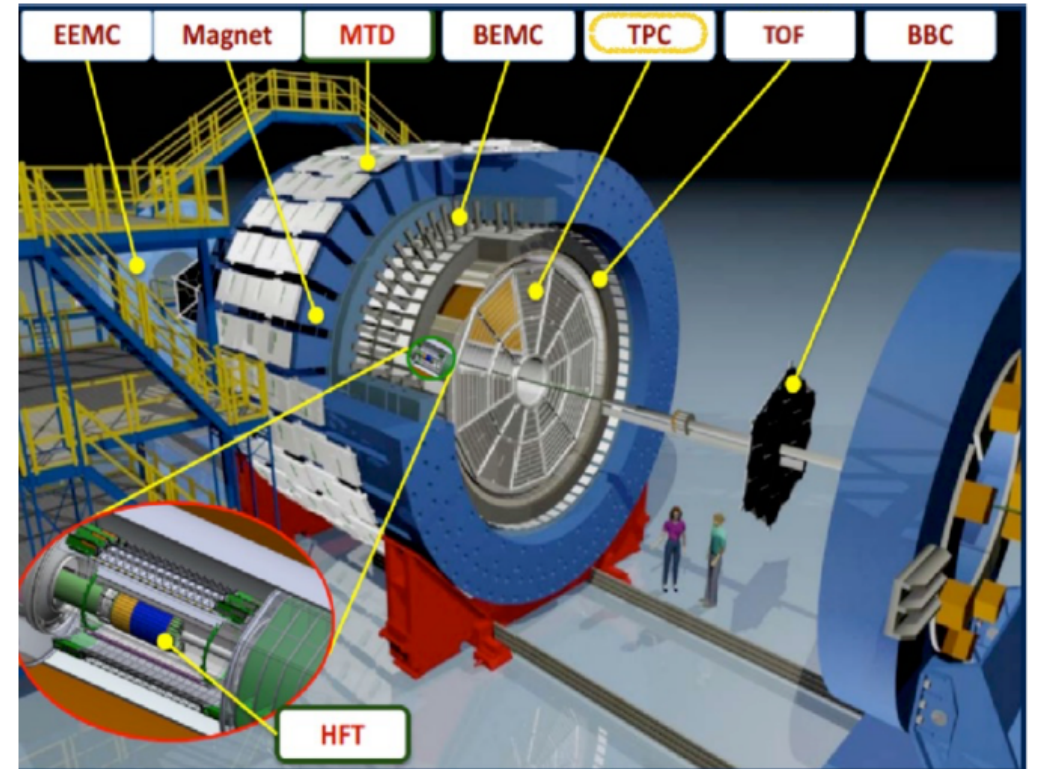


- The charge quadrupole shifts the distribution mean value

✓ A stronger shift for large input signal

❖ $R_{\Psi_m}^{(2)}$ in Data

- TPC detector used in the current analysis
- Au+Au at 200 GeV data from year 2011 is used
- Charged hadrons with $0.2 < p_T < 2.0$ GeV/c used to construct $\Psi_m^{\eta > 0.1}$ & $\Psi_m^{\eta < -0.1}$
- Particles with $0.35 < p_T < 2.0$ GeV/c and $-1 > \eta < -0.1$ analyzed using $\Psi_m^{\eta > 0.1}$
- Particles with $0.35 < p_T < 2.0$ GeV/c and $1 > \eta > 0.1$ analyzed using $\Psi_m^{\eta < -0.1}$



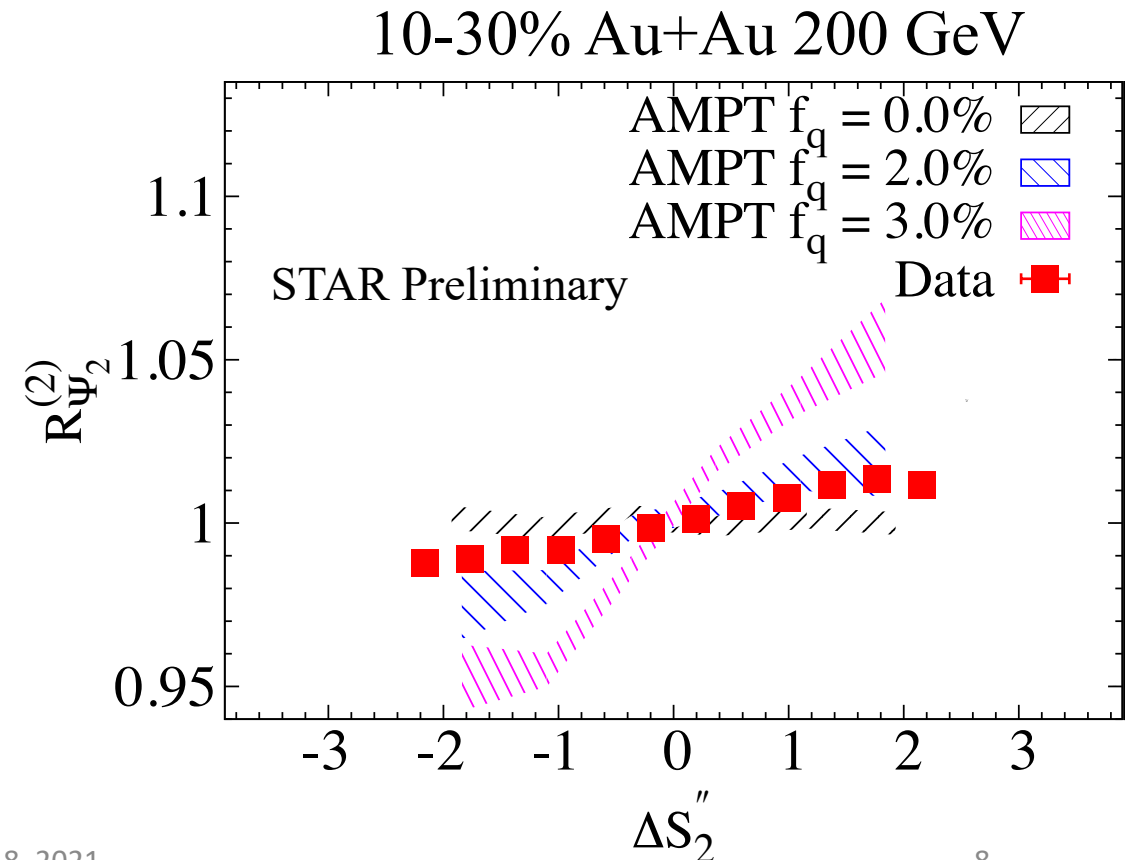
❖ The $R_{\Psi_m}^{(2)}(\Delta S_2)$ [quadrupole correlator]

➤ The correlator response to quadrupole charge separation:

The $R_{\Psi_2}^{(2)}$ distributions for charged particles for 10-30% central events in Au+Au collisions. Data are compared to AMPT with different quadrupole charge separation

➤ The charge quadrupole strength is reflected in the slope of $R_{\Psi_2}^{(2)}$ vs. $\Delta S_2''$:

- ✓ Linear dependence of $R_{\Psi_2}^{(2)}$ on $\Delta S_2''$ similar to AMPT with quadrupole charge separation signal



❖ The $R_{\Psi_m}^{(2)}(\Delta S_2)$ [quadrupole correlator]

➤ The correlator response in quadrupole charge separation:

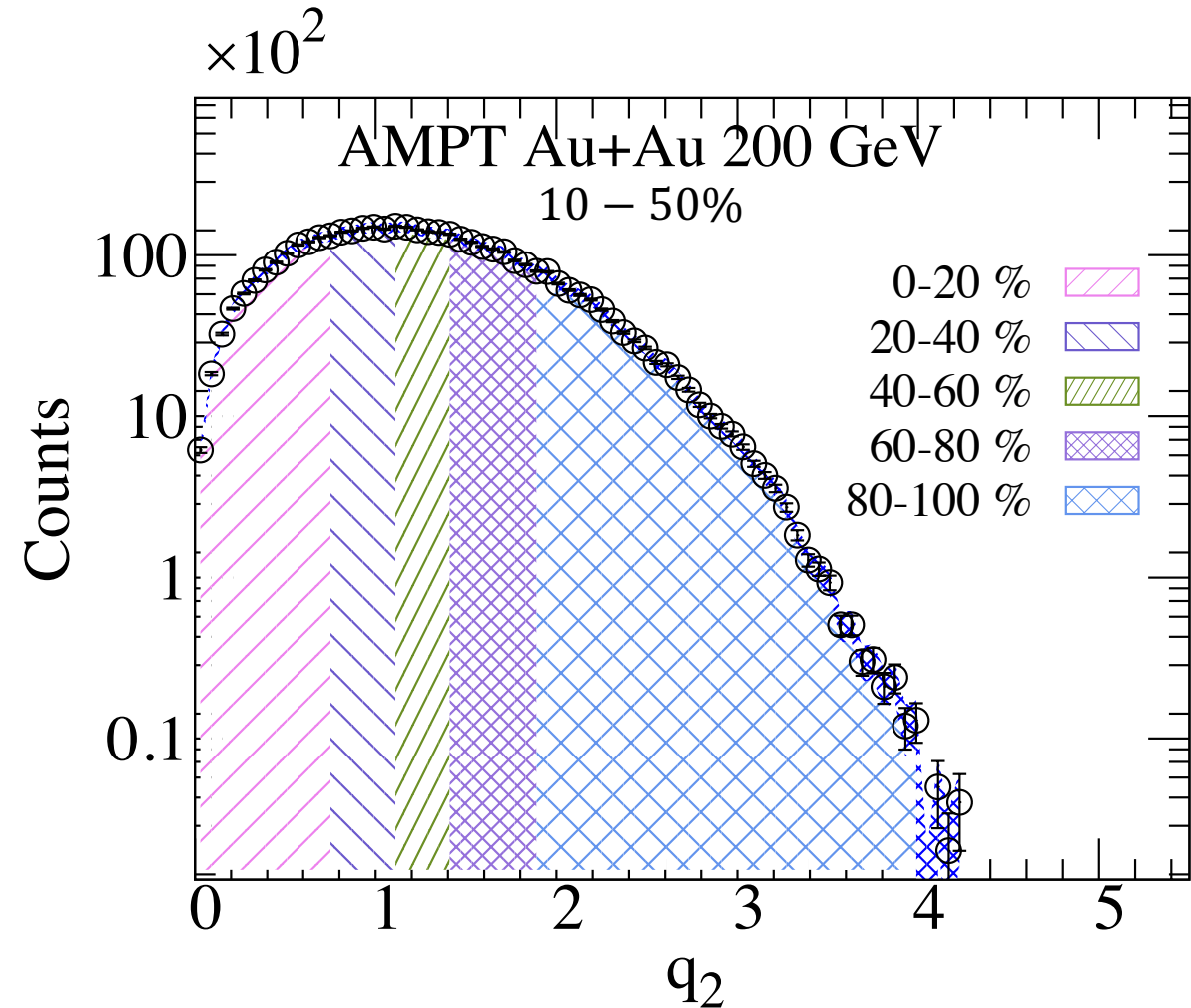
- ✓ The v_2 -driven background

$$Q_{2,x} = \sum_{i=1}^M \cos(2 \varphi_i)$$

$$Q_{2,y} = \sum_{i=1}^M \sin(2 \varphi_i)$$

$$|Q_2| = \sqrt{Q_{2,x}^2 + Q_{2,y}^2}$$

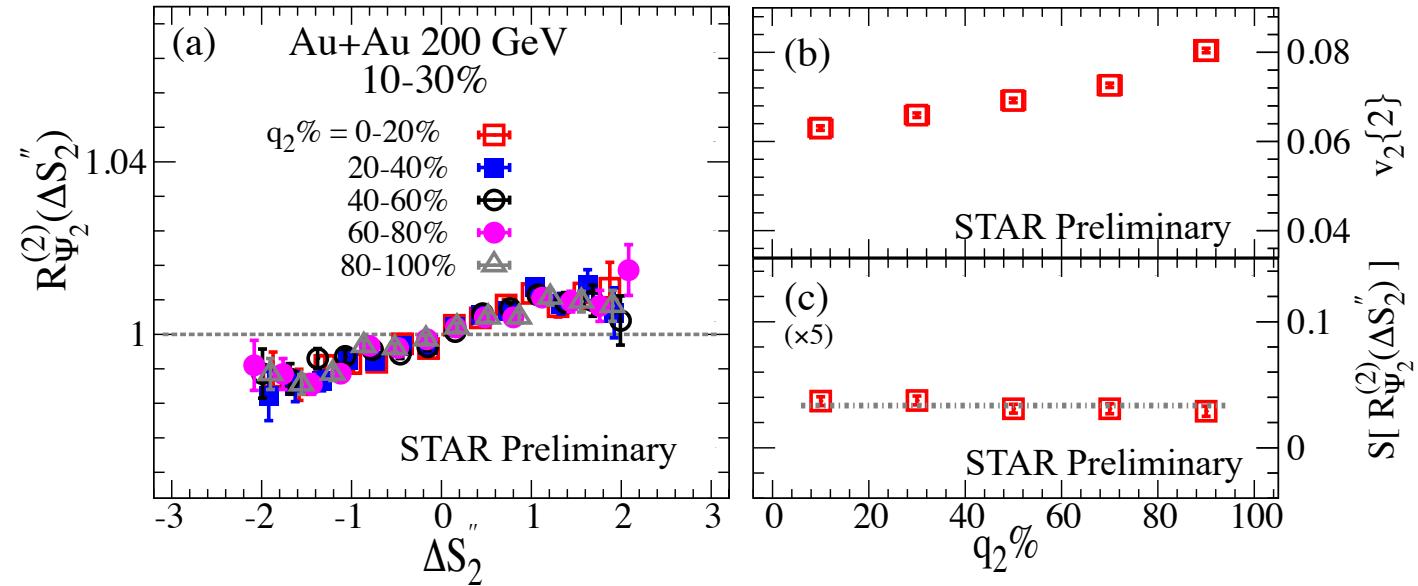
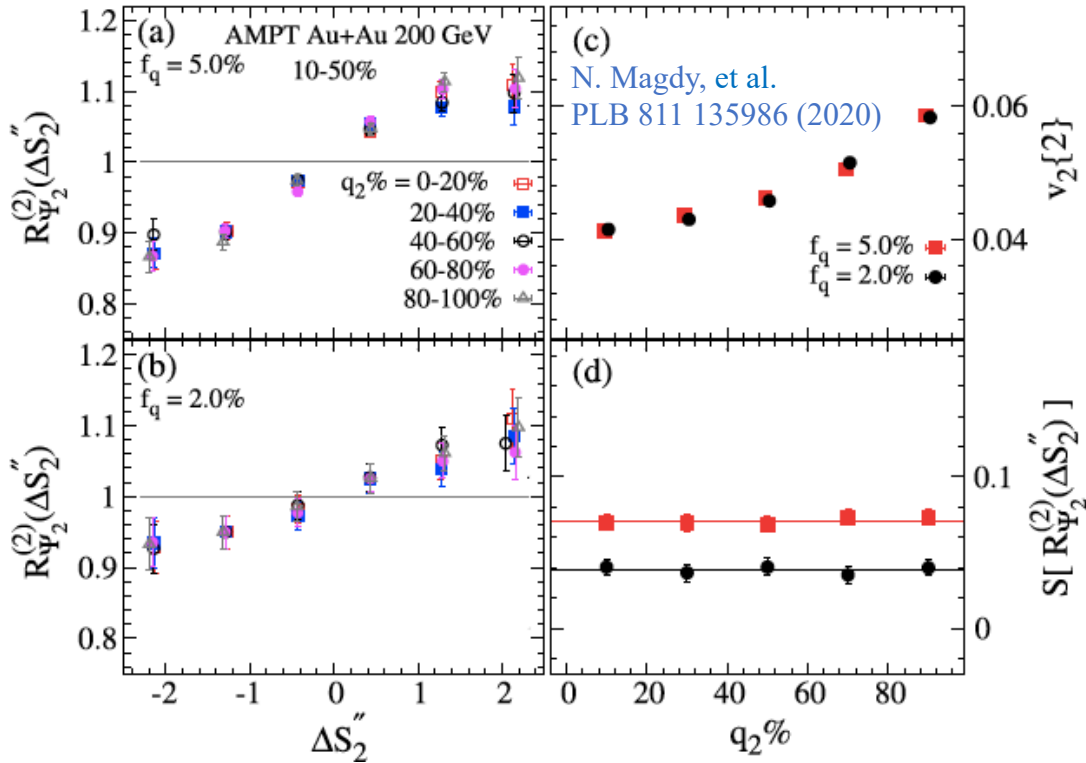
$$q_2 = \frac{|Q_2|}{\sqrt{M}}$$



❖ The $R_{\Psi_m}^{(2)}(\Delta S_2)$ [quadrupole correlator]

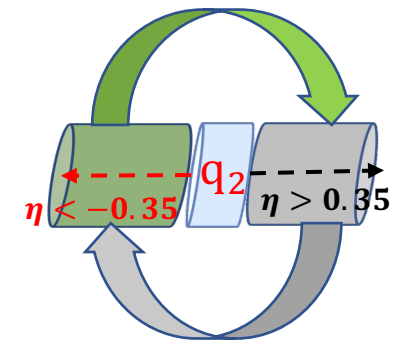
➤ The correlator response in quadrupole charge separation:

- ✓ The v_2 -driven background



- Variation of q_2 leading to $\sim 30\%$ variation of v_2 in 10-30% centrality does not lead to any significant variation of $R_{\Psi_2}^{(2)}$

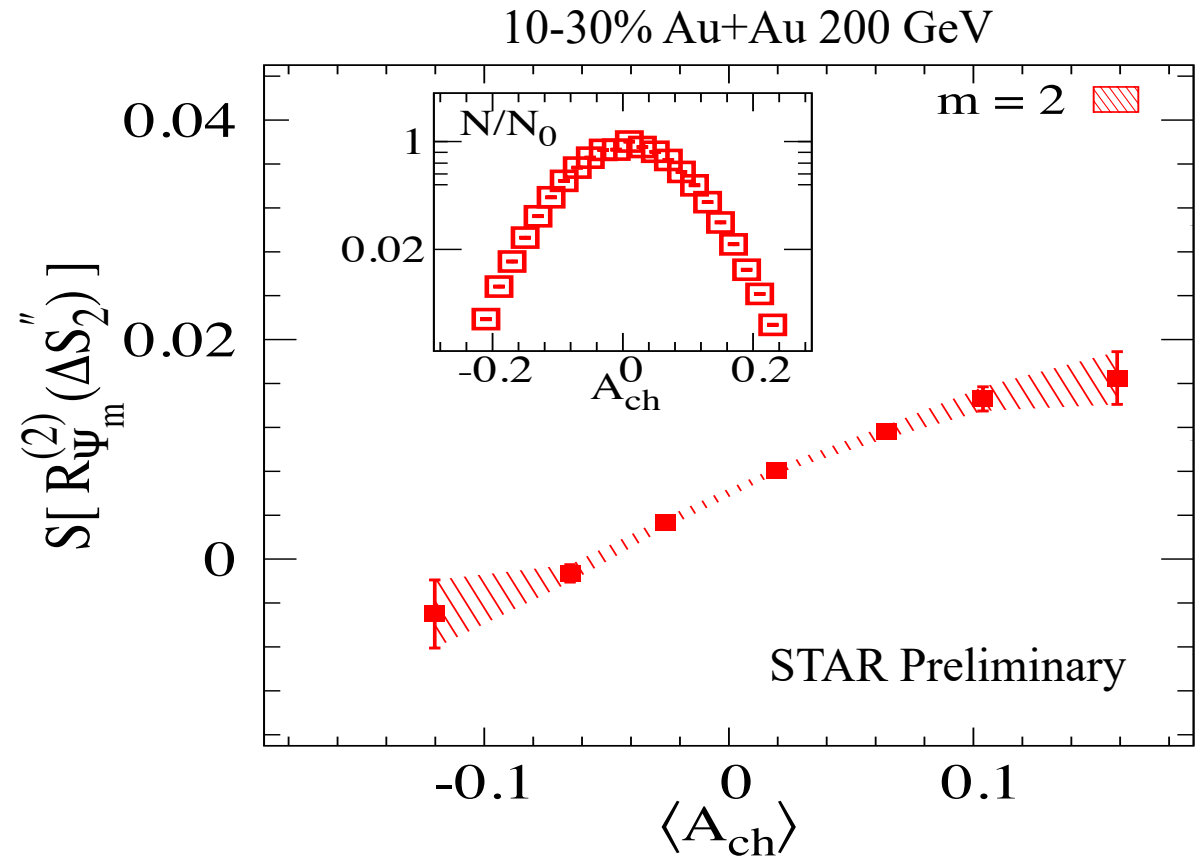
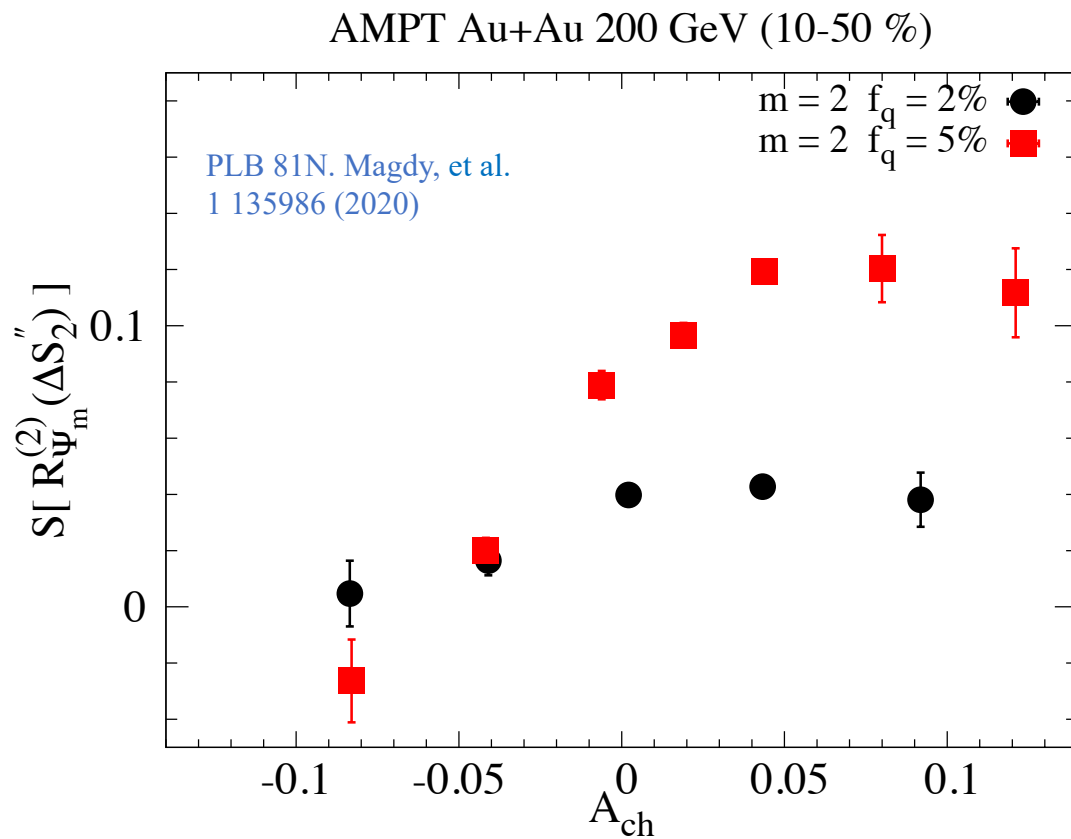
- $R_{\Psi_2}^{(2)}$ is not strongly influenced by v_2 background-driven charge separation



❖ The $R_{\Psi_m}^{(2)}(\Delta S_2)$ [quadrupole correlator]

➤ The correlator response in quadrupole charge separation:

The charge asymmetry dependence of the slope $S(R_{\Psi_2}^{(2)})$ from data and the AMPT model



➤ The $R_{\Psi_2}^{(2)}$ slope is sensitive to the charge asymmetry variations.

A novel correlator $R_{\psi_2}^{(2)}(\Delta S_2)$, has been used to study the quadrupole charge separation in Au+Au at 200 GeV from data and AMPT model:

❖ Au+Au at 200 GeV AMPT simulations:

➤ The $R_{\psi_2}^{(2)}$ shows a sensitivity to small input quadrupole charge separation signals

❖ Au+Au at 200 GeV measurements:

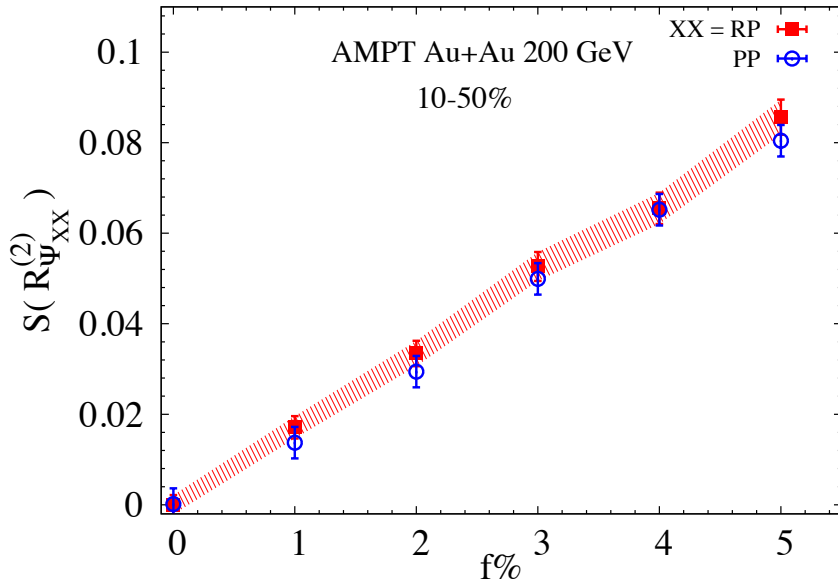
➤ The measurements show patterns and trends compatible with those expected for small quadrupole charge separation signal in the Au+Au data.

We observed a qualitative similarity between data and AMPT with quadrupole charge separation signal. Therefore, providing a comprehensive set of measurements as well as simulations will lead more to a better understanding of the new observable.

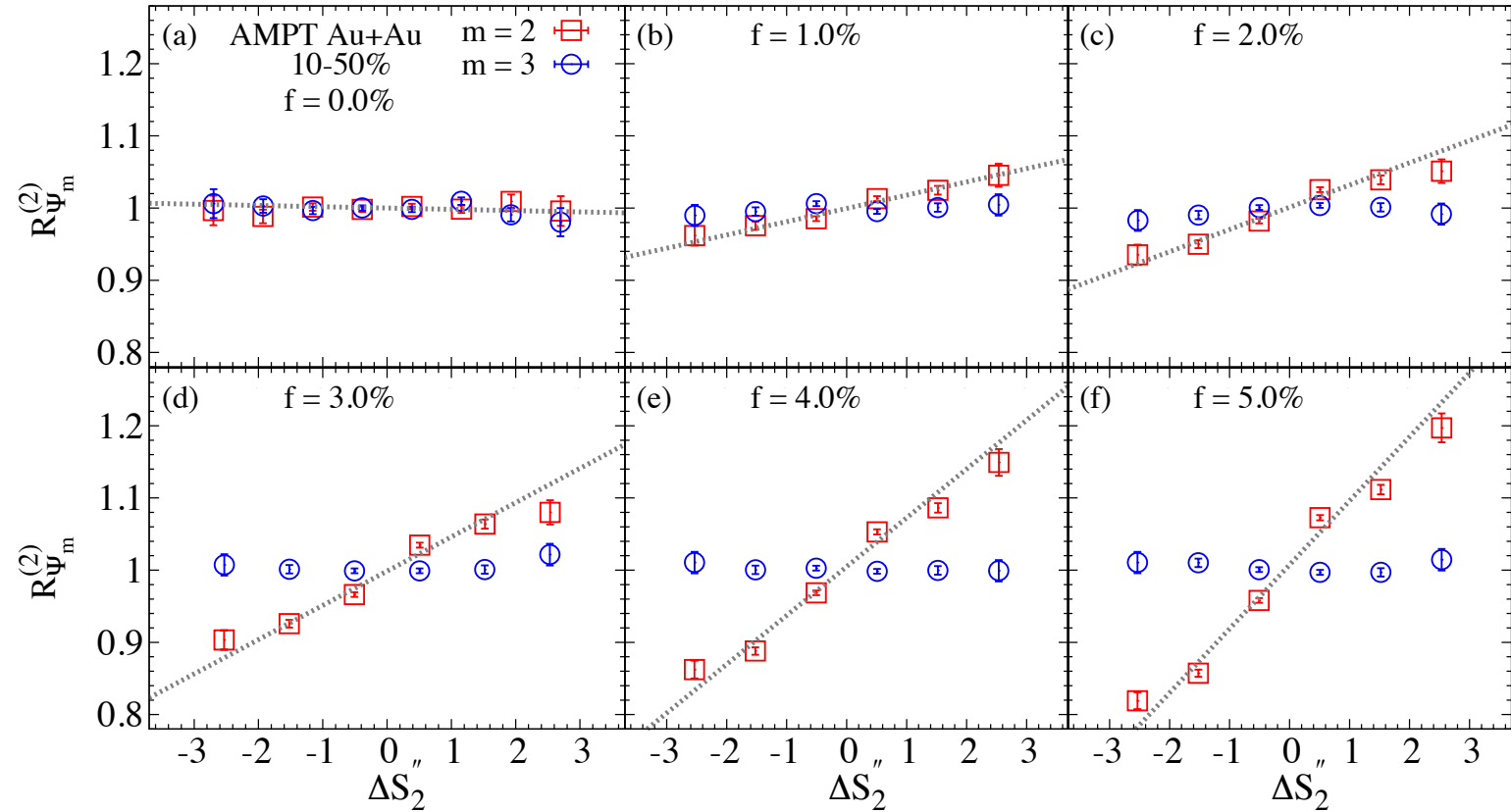
Thank You

❖ The $R_{\Psi_m}^{(2)}(\Delta S_2)$ [quadrupole correlator]

➤ The correlator response in CMW scenarios:



The $R_{\Psi_m}^{(2)}(m = 2,3)$ correlators from the AMPT model for several input charge quadrupole signal.



➤ The $R_{\Psi_{XX}}^{(2)}$ slope is sensitive to all input signals

- ✓ $R_{\Psi_{XX}}$ is essentially event plane independent
- ✓ The slopes are linearly related to the input signal

➤ The charge quadrupole strength captured in the $R_{\Psi_2}^{(2)}$ slope

- ✓ $R_{\Psi_3}^{(2)}$ show no sensitivity to the input signal