

Beam-energy dependence of transverse momentum and flow correlations in STAR

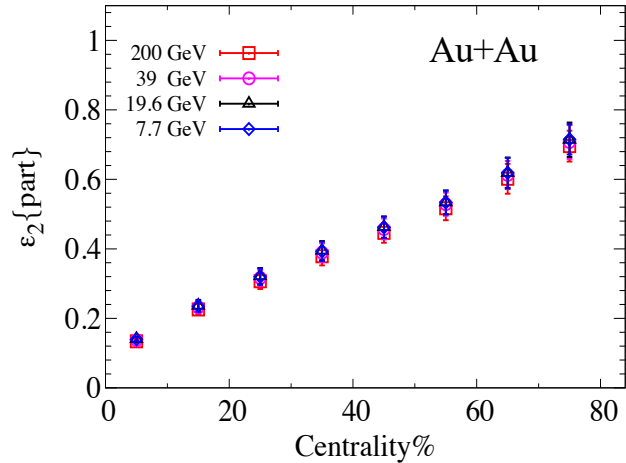


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Motivation:

- The beam-energy dependence of flow and p_T correlations will reflect the respective roles of ϵ_n , its fluctuations and $\frac{\eta}{s}$ as a function of T and μ_B

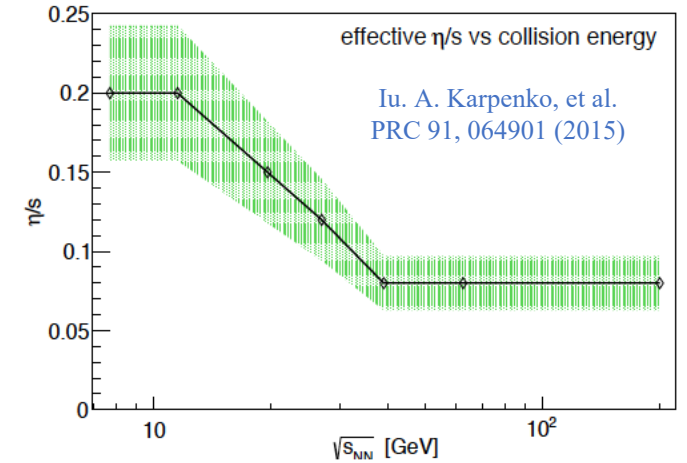
Beam energy dependence for a given collision system:



- Initial-state ϵ_2 is approximately energy independent
- Viscous attenuation ($\propto \frac{\eta}{s}(T)$) is beam energy dependent

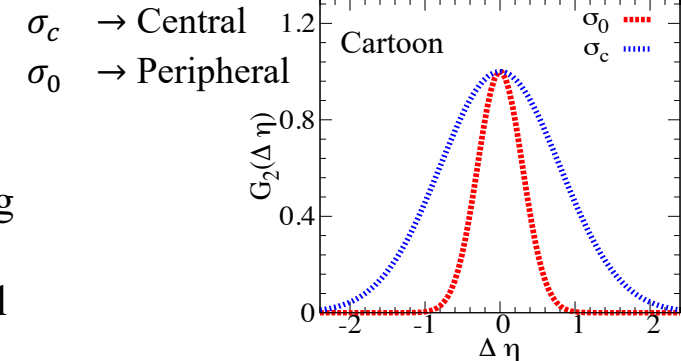
Piotr Bozek
PRC 93, 044908 (2016)

Niseem Magdy, Roy Lacey
PLB 821 136625 (2021)



- The Pearson correlation, $v_n - [p_T]$ correlation, coefficient (PCC) is expected to be more susceptible to the initial conditions of heavy-ion collisions.

S. Gavin and M. Abdel-Aziz
Phys.Rev.Lett. 97 (2006) 162302



The Gavin ansatz:

- The $p_T - p_T$ correlation function is sensitive to the dissipative viscous effects that are ensured during the transverse and longitudinal expansion of the medium created in the collisions.
- Because such dissipative effects are more prominent for long-lived systems, they lead to longitudinal broadening of p_T 2-P correlation function as collisions become more central.
- A proposed estimate of this broadening, $\Delta\sigma^2$, can be linked to η/s as:

$$\Delta\sigma^2 = \sigma_c^2 - \sigma_0^2 = \frac{4}{T_c} \frac{\eta}{s} \left(\frac{1}{\tau_0} - \frac{1}{\tau_{c,f}} \right)$$

Analysis procedure:

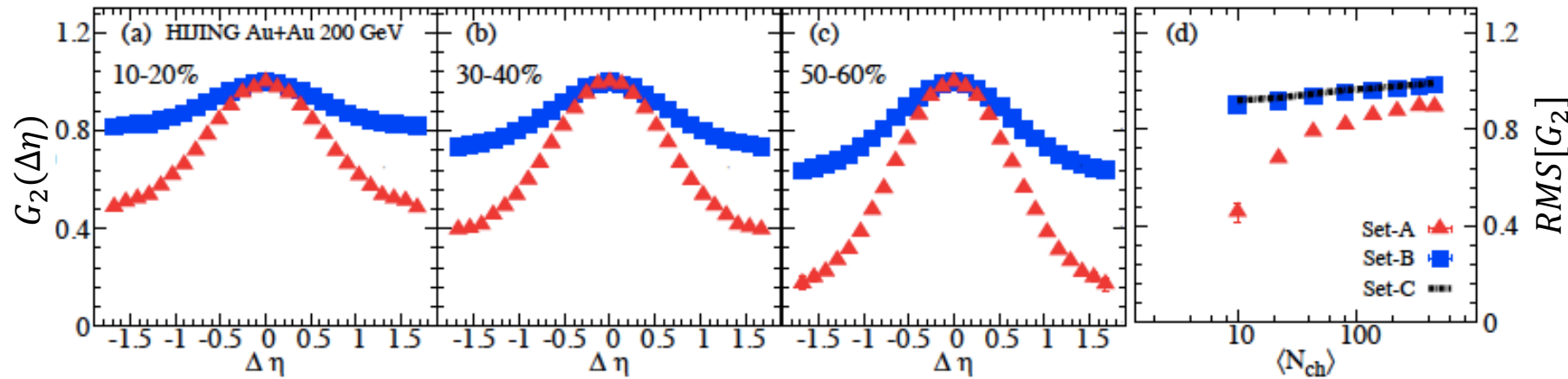
The $p_T - p_T$ correlator:

STAR Collaboration
PLB 704 (2011) 467–473

N. Magdy and R. Lacey
PRC 104, 014907 (2021)

Niseem Magdy, et al.
Eur.Phys.J.C 81 (2021) 8, 779

$$G_2(\eta_1, \varphi_1, \eta_2, \varphi_2) = \frac{\left\langle \sum_i^{n_1} \sum_{j \neq i}^{n_2} p_{T,i} p_{T,j} \right\rangle}{\langle n_1 \rangle \langle n_2 \rangle - \langle p_{T,1} \rangle_{\eta_1, \varphi_1} \langle p_{T,2} \rangle_{\eta_2, \varphi_2}} \rightarrow \frac{\left\langle \sum_i^{n_1} \sum_{j \neq i}^{n_2} p_{T,i} p_{T,j} \right\rangle}{\langle n_1 \rangle \langle n_2 \rangle} = \frac{\left\langle \sum_i^{n_1} \sum_{j \neq i}^{n_2} p_{T,i} p_{T,j} \right\rangle}{\left\langle \sum_i^{n_1} \sum_{j \neq i}^{n_2} n_i n_j \right\rangle} r_{1,2} \rightarrow r_{1,2} = \frac{\left\langle \sum_i^{n_1} \sum_{j \neq i}^{n_2} n_i n_j \right\rangle}{\langle n_1 \rangle \langle n_2 \rangle}$$



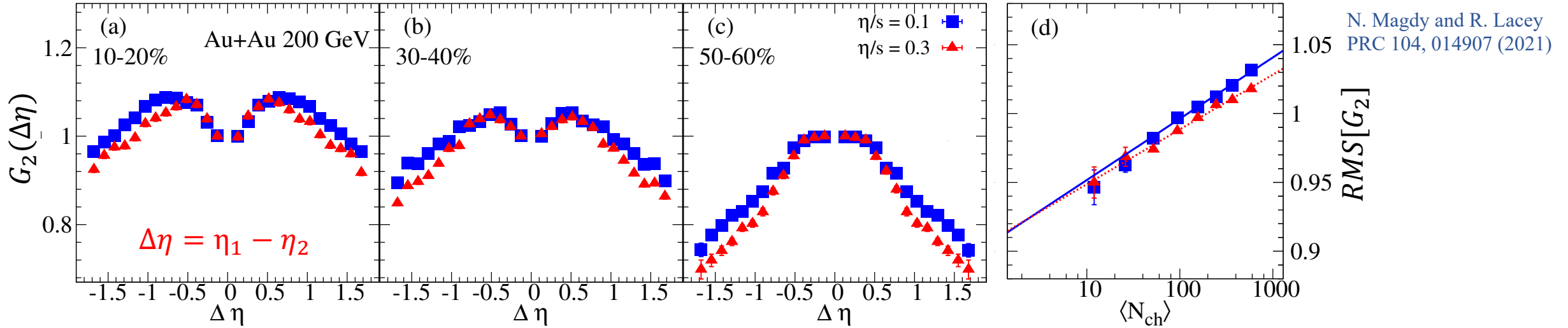
Comparison of the $G_2(\Delta\eta)$ correlators ($|\Delta\varphi| < 1$) obtained from 10-20%, 30-40% and 50-60% central HIJING events for Au+Au collisions at 200 GeV.

- (i) Set-A: with centrality defined using all charged particles in an event,
- (ii) Set-B: with centrality defined using random sampling of charged particles in an event
- (iii) Set-C: with centrality defined using the impact parameter distribution.

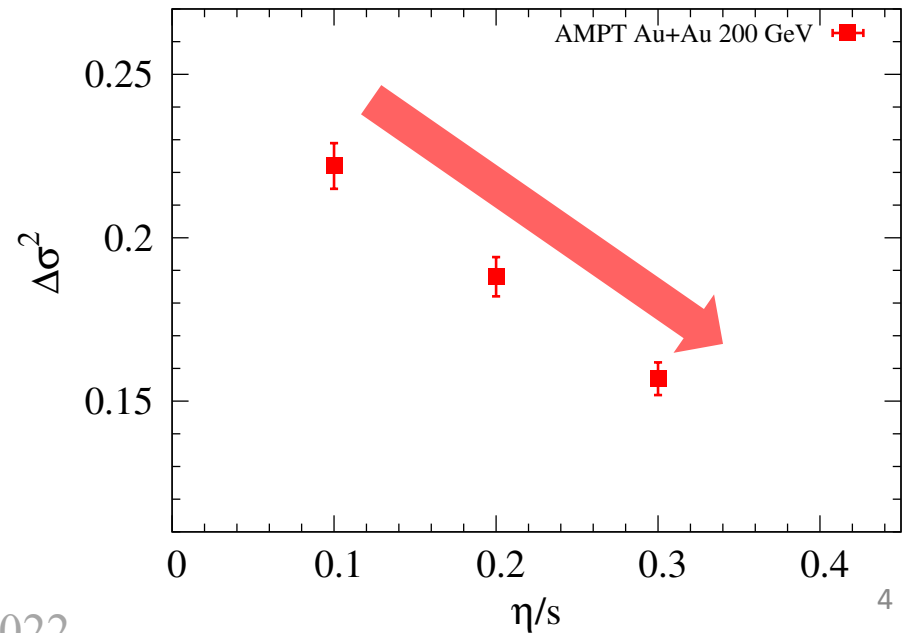
Excluding the POI from the collision centrality definition, serves to reduce the possible self-correlations.

Investigations of the $p_T - p_T$ correlations

➤ The longitudinal correlations using AMPT model for Au+Au at 200 GeV



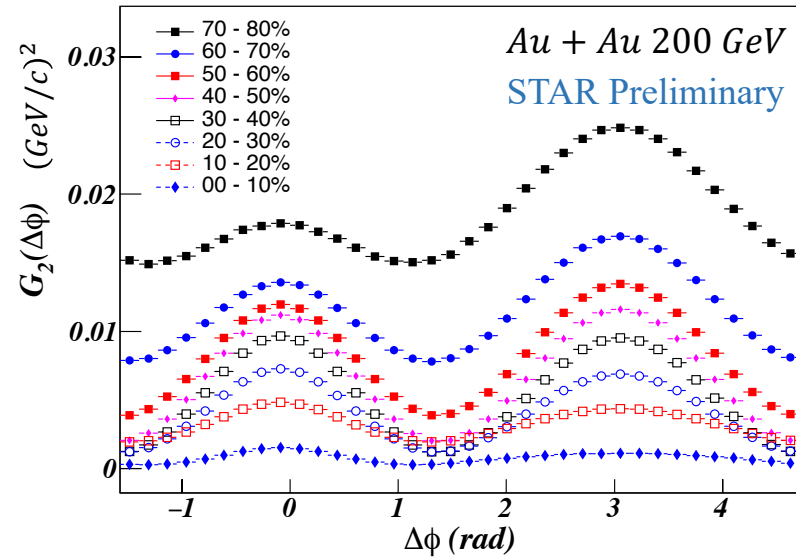
➤ The $\sigma_{\Delta\eta}(G_2)$ is sensitive to the η/s change



Investigations of the $p_T - p_T$ correlations from STAR

➤ The azimuthal correlations for Au+Au at 200 GeV

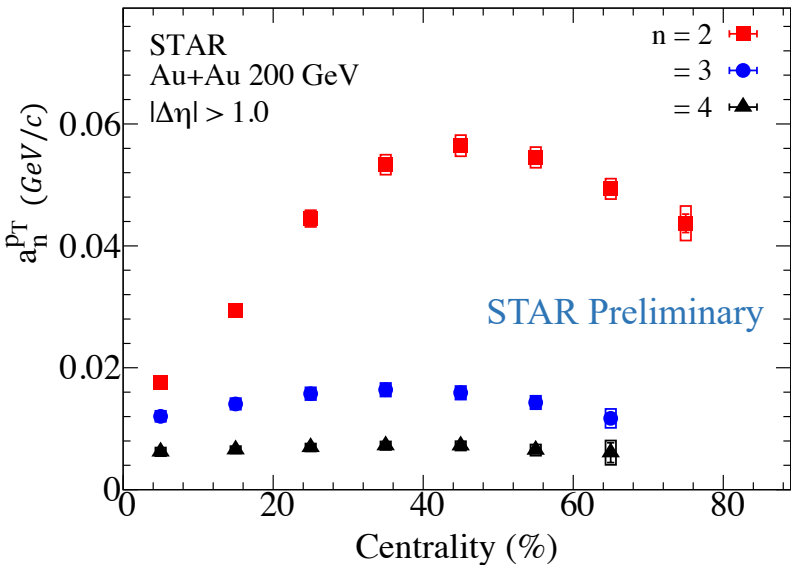
$$G_2(\Delta\phi) = A_0^{pT} + 2 \sum_{n=1}^6 A_n^{pT} \cos(n \Delta\phi)$$



$$a_n^{pT} = \sqrt{A_n^{pT}}$$

➤ The extracted a_2^{pT} :

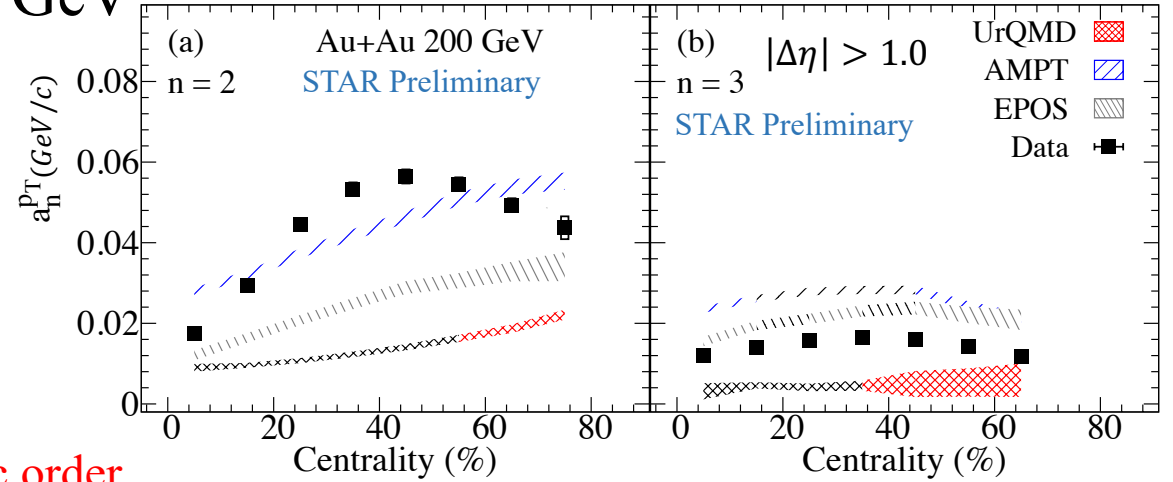
- ✓ Decrease with harmonic order
- ✓ Models do not describe the data
- ✓ Event shape dependent



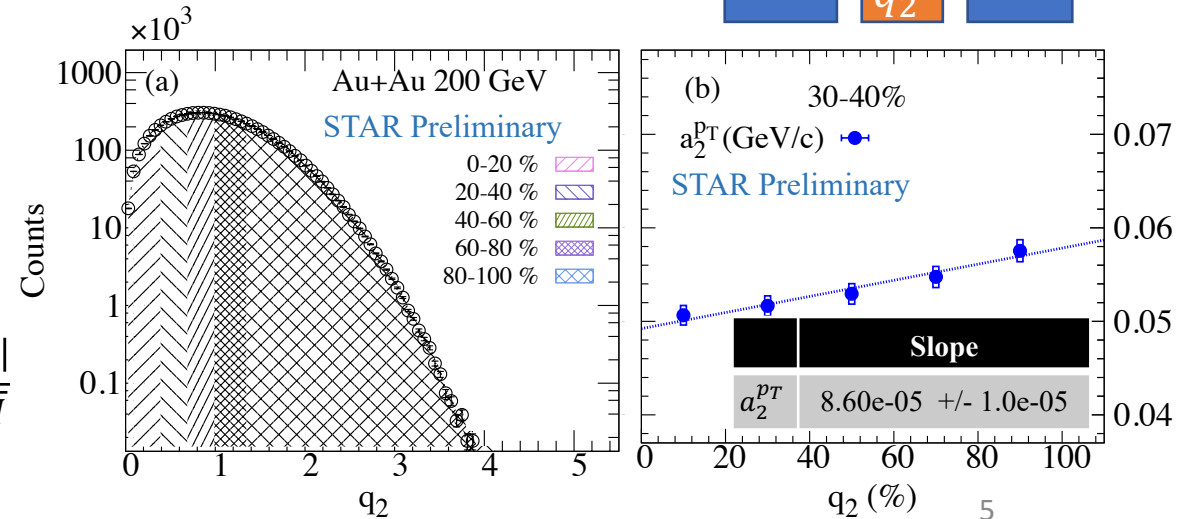
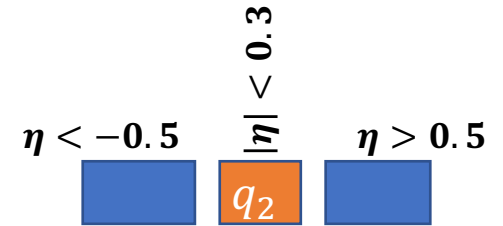
$$Q_{2,x} = \sum_{i=1}^M \cos(2 \varphi_i)$$

$$Q_{2,y} = \sum_{i=1}^M \sin(2 \varphi_i)$$

$$|Q_2| = \sqrt{Q_{2,x}^2 + Q_{2,y}^2} \quad q_2 = \frac{|Q_2|}{\sqrt{M}}$$

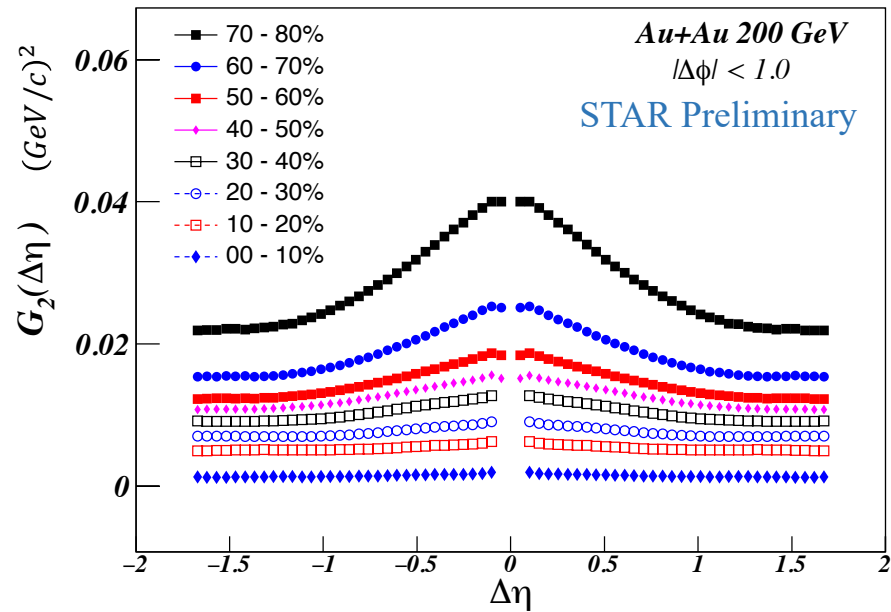


➤ q_2 is separated from particle of interest



Investigations of the $p_T - p_T$ correlations from STAR

➤ The longitudinal correlations for Au+Au at 200 GeV

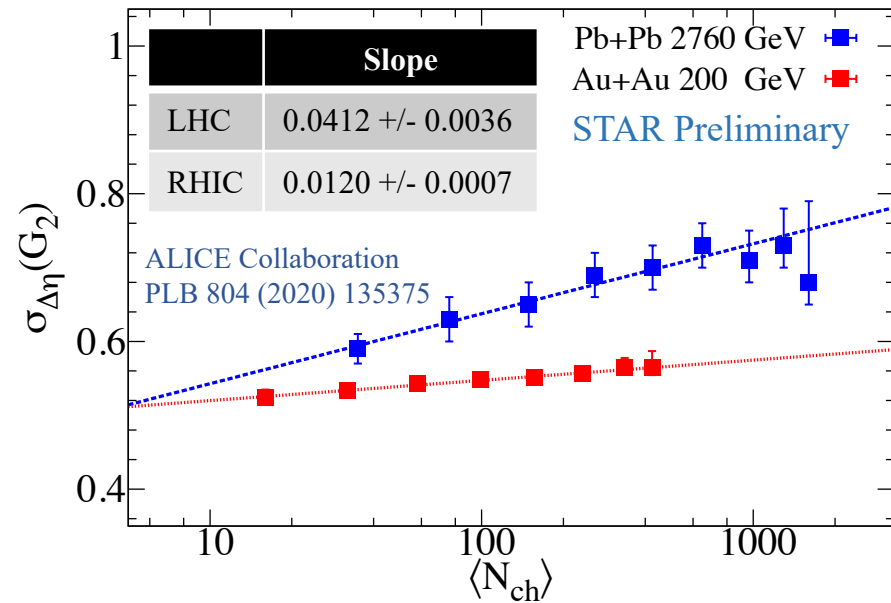
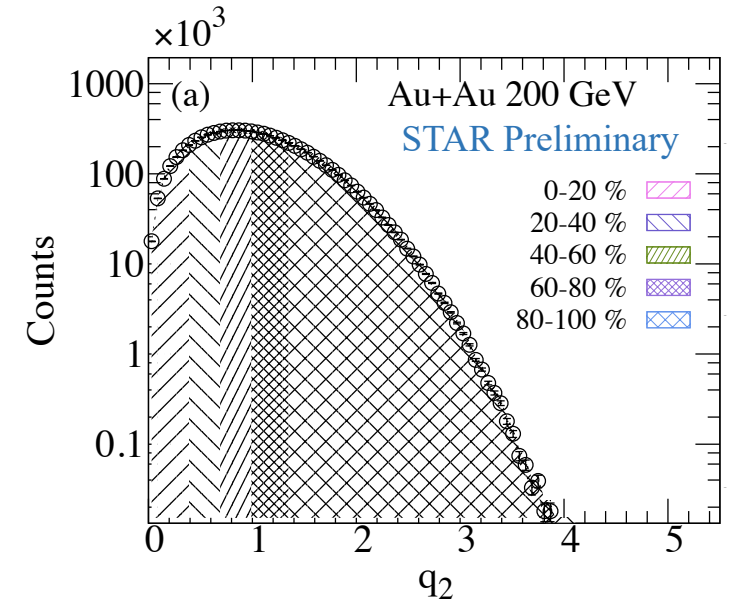


$$\sigma_{\Delta\eta}(G_2) = \text{RMS}[G_2(\Delta\eta)]$$

➤ The slope of $\sigma_{\Delta\eta}(G_2)$ vs. $\langle N_{ch} \rangle$ is softer for RHIC

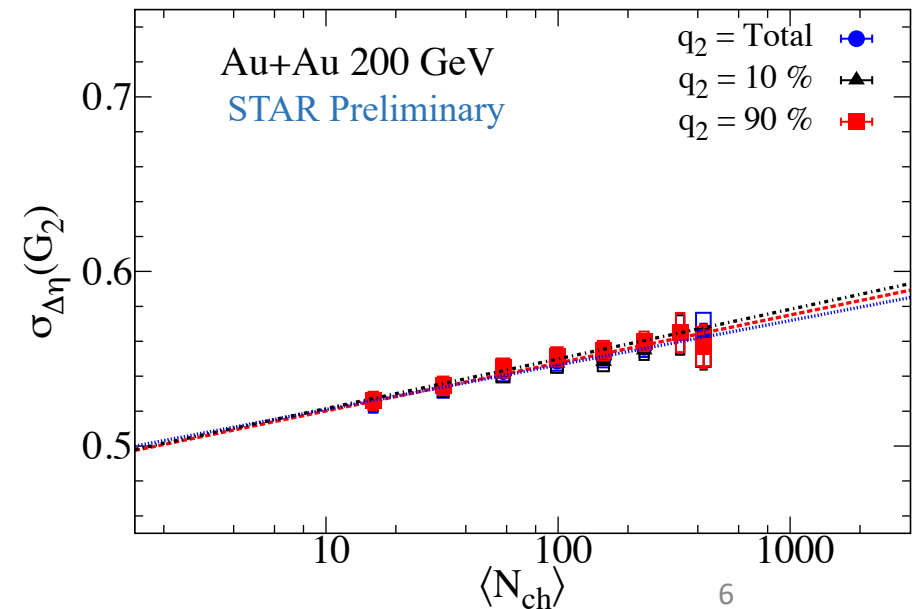
✓ Smaller η/s for RHIC

P. Alba et al.
PRC 98, 034909 (2018)

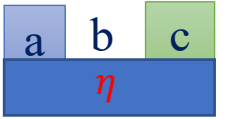


➤ The $\sigma_{\Delta\eta}(G_2)$ is event shape independent

	Slope
Total	0.0120 +/- 0.0007
10 %	0.0119 +/- 0.0009
90 %	0.0110 +/- 0.0012



$$|\Delta\eta| > 0.7$$



J. Jia, M. Zhou, A. Trzupek,
PRC 96 034906 (2017)

ATLAS Collaboration,
Eur. Phys. J. C 79, 985 (2019)

Piotr Bozek
PRC 93, 044908 (2016)

Niseem Magdy, Roy Lacey
PLB 821 136625 (2021)

Niseem Magdy, et al.
PRC 105 (2022) 4, 044901

Analysis procedure:

❖ Transverse momentum-flow correlations:

$$Var(v_n^2)_{dyn} = v_n^4\{2\} - v_n^4\{4\}$$

$$C_k = \left\langle \frac{\sum_b \sum_{b'} w_b w_{b'} (p_{T,b} - \langle [p_T] \rangle) (p_{T,b'} - \langle [p_T] \rangle)}{((\sum_b w_b)^2 - \sum_b (w_b)^2)} \right\rangle \Delta\eta_{b\hat{b}} > 0.2$$

$$cov(v_n^2, [p_T]) = Re \left(\left\langle \frac{\sum_{a,c} w_a w_c e^{in(\phi_a - \phi_c)} ([p_T] - \langle [p_T] \rangle)_b}{\sum_{a,c} w_a w_c} \right\rangle \right)$$

$$\rho(v_n^2, [p_T]) = \frac{cov(v_n^2, [p_T])}{\sqrt{Var(v_n^2)_{dyn} C_{\{k\}}}}$$

The Pearson correlation coefficient (PCC) measures the strength of the $v_n, [p_T]$ correlation.

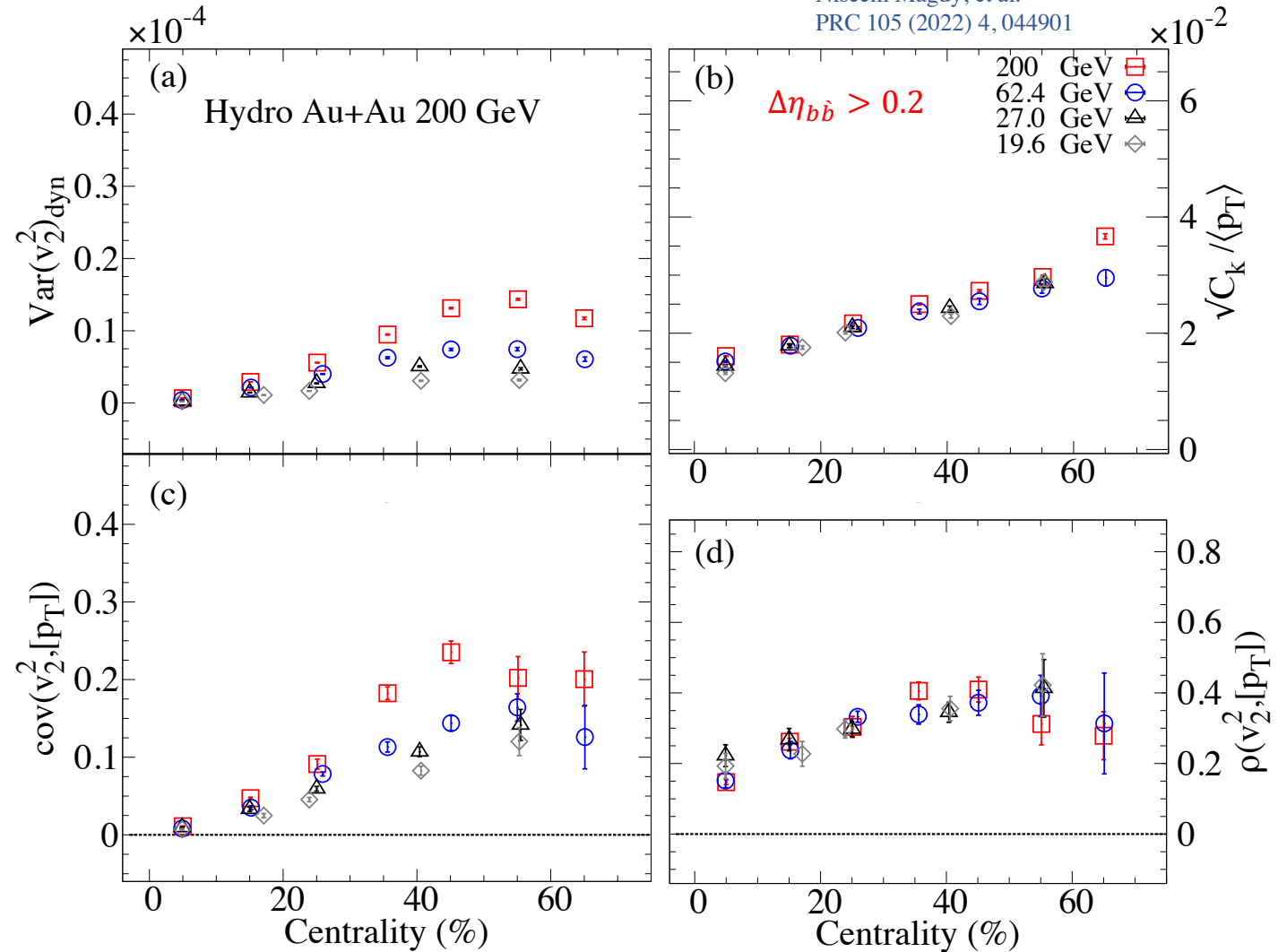
❖ Transverse momentum-flow correlations:

The beam-energy dependance of the transverse momentum-flow correlations
using hydro model

Niseem Magdy, et al.
PRC 105 (2022) 4, 044901

- $Var(v_2^2)_{dyn}$ decreases with beam-energy
- $\sqrt{C_k}/\langle p_T \rangle$ shows no change with beam energy
- $cov(v_2^2, [p_T])$ decreases with beam-energy

➤ The Pearson correlation, $\rho(v_2^2, [p_T])$ shows no change with beam energy



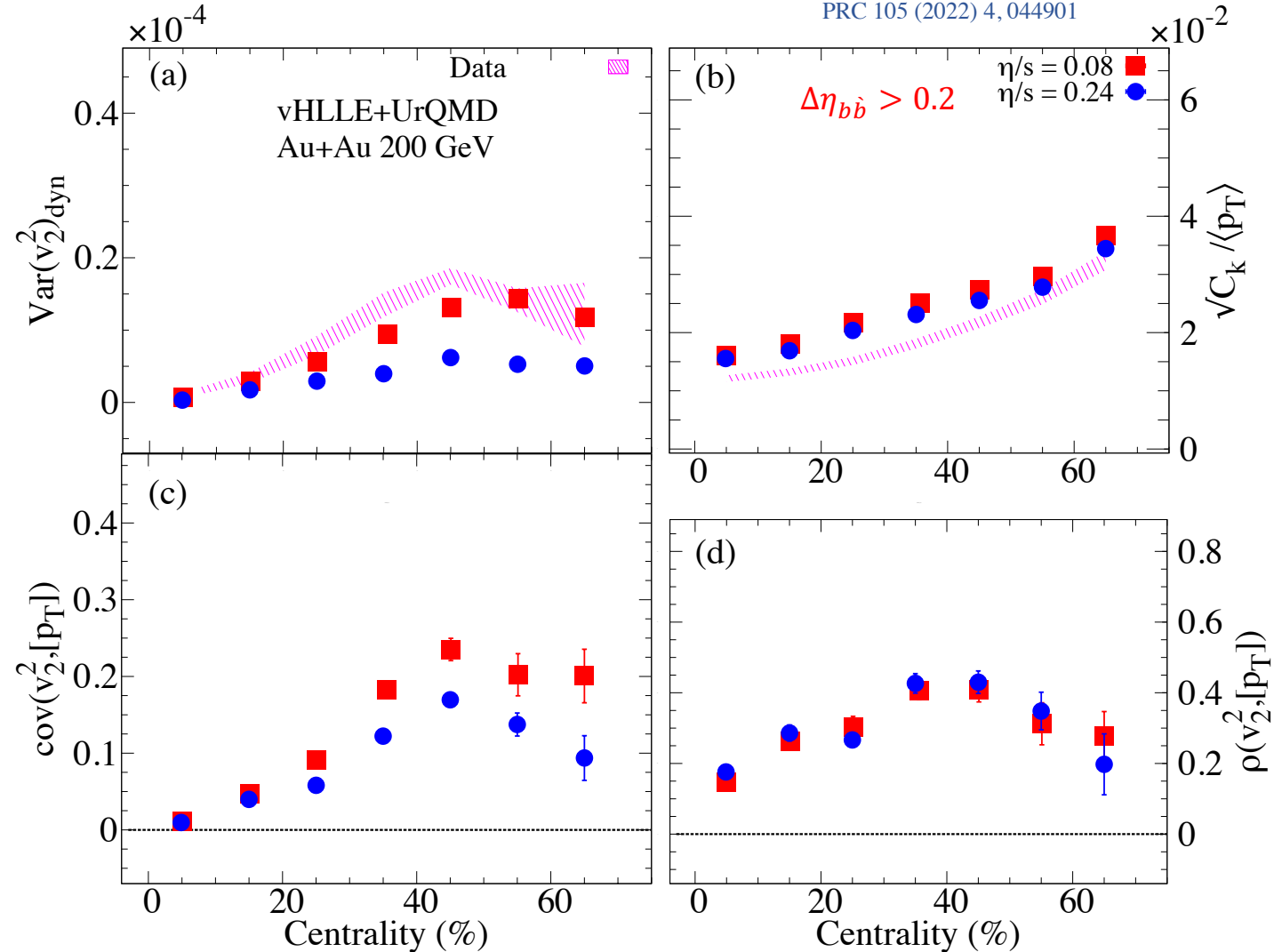
❖ Transverse momentum-flow correlations:

The beam-energy dependance of the transverse momentum-flow correlations using hydro model

Niseem Magdy, et al.
PRC 105 (2022) 4, 044901

- $Var(v_2^2)_{dyn}$ decreases with increasing η/s
- $\sqrt{C_k}/\langle p_T \rangle$ shows no change with η/s
- $cov(v_2^2, [p_T])$ decreases with increasing η/s

➤ The Pearson correlation, $\rho(v_2^2, [p_T])$ shows little change with η/s



Transverse momentum-flow correlations measurements:

❖ Hydro comparisons

➤ (A) B.Schenke, C.Shen, and P.Tribedy
PRC 99, 044908 (2019)

➤ (B) P. Alba, et al.
PRC 98, 034909 (2018)

	Hydro-A	Hydro-B
η/s	0.12	0.05
Initial conditions	IP-Glasma	TRENTO
Contributions	Hydro + Hadronic cascade	Hydro + Direct decay

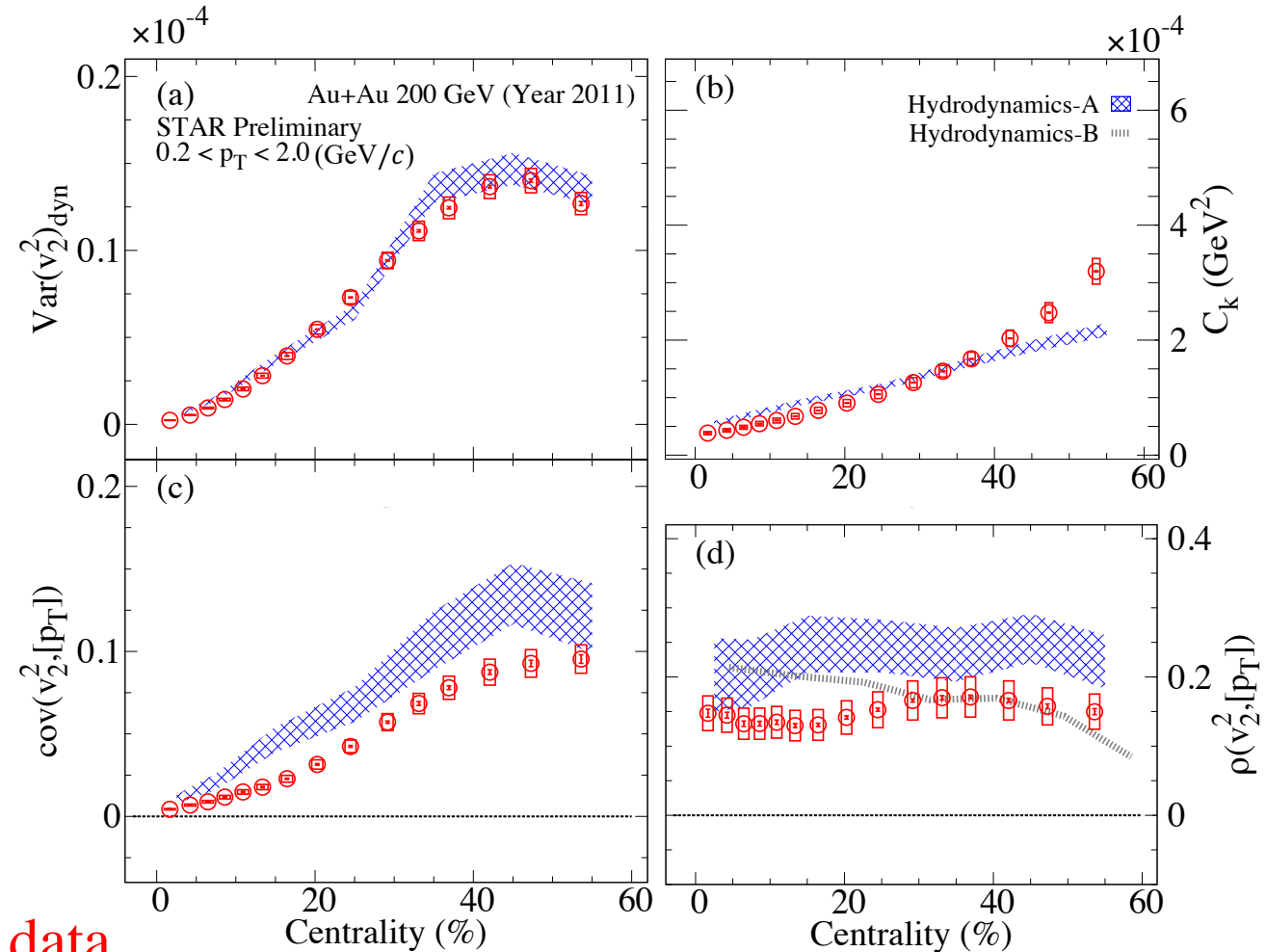
➤ $Var(v_2^2)_{dyn}$ shows a good agreement with Hydro-A

➤ C_k shows a good agreement with Hydro-A from central to mid central

➤ Hydro-A overestimate $cov(v_2^2, [p_T])$

➤ Hydro models can qualitatively describe the data

✓ Both Hydro-A and -B overestimates $\rho(v_2^2, [p_T])$



Transverse momentum-flow correlations measurements:

❖ Hydro comparisons

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PRC 99, 044908 (2019)

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PRC 98, 034909 (2018)

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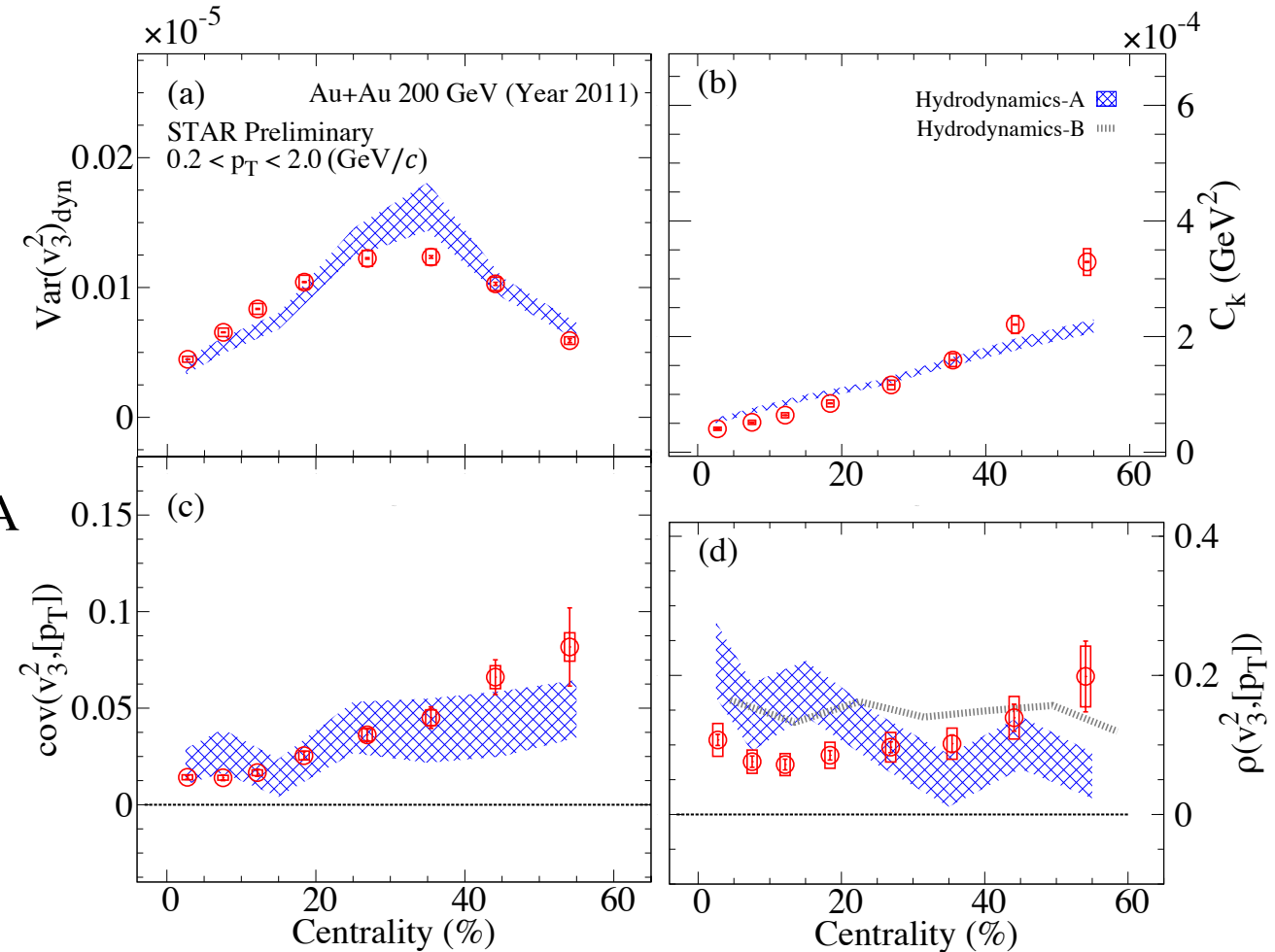
➤ $Var(v_3^2)_{dyn}$ shows a good agreement with Hydro-A

➤ C_k shows a good agreement with Hydro-A from central to mid central

➤ Hydro-A within the uncertainty shows a good agreement with $cov(v_3^2, [p_T])$

➤ Hydro models can qualitatively describe the data

✓ Both Hydro-A and -B overestimates $\rho(v_3^2, [p_T])$ in more central collisions

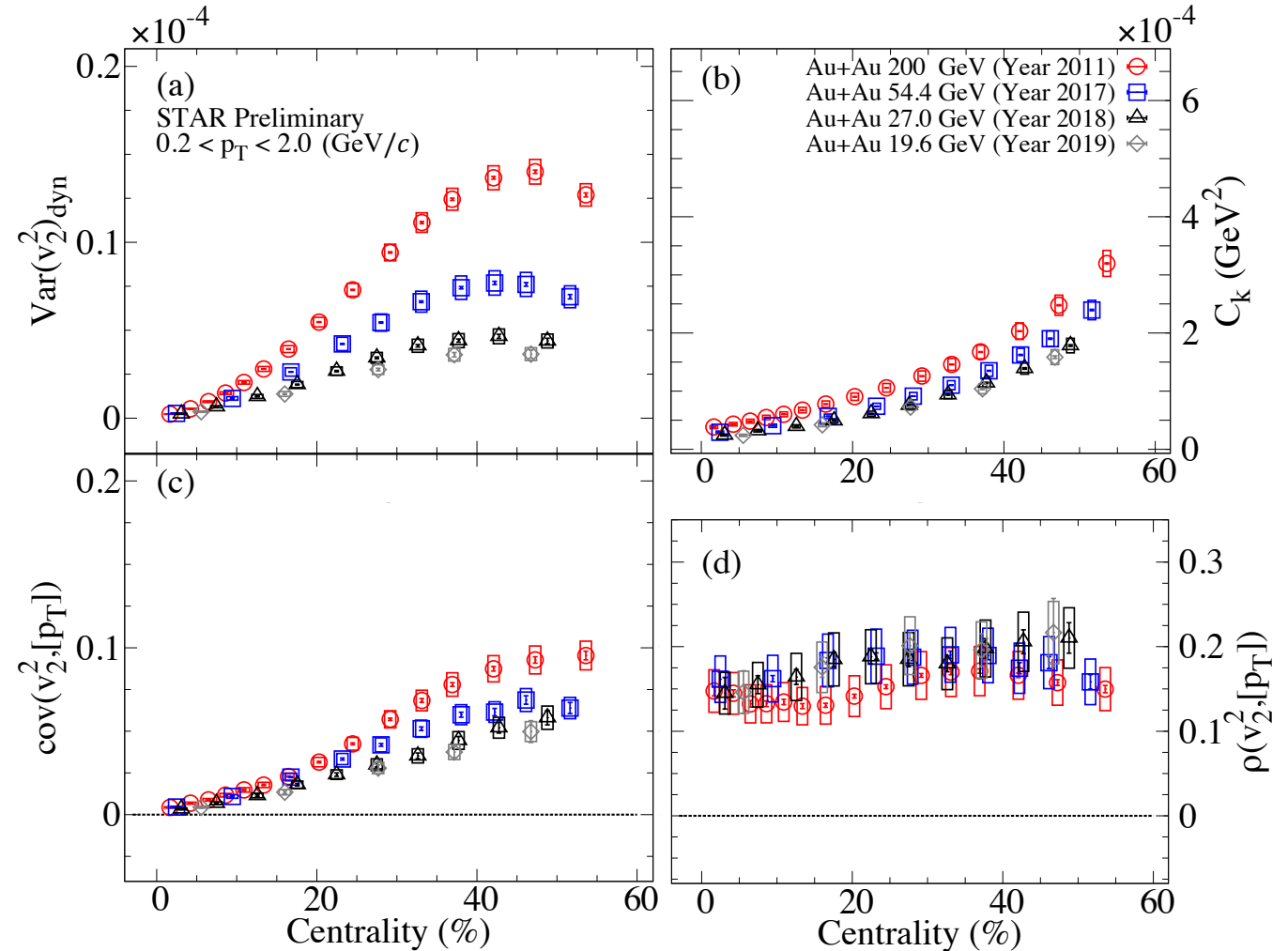


Transverse momentum-flow correlations measurements:

❖ The beam-energy dependence of the transverse momentum-flow correlations

- $Var(v_2^2)_{dyn}$ decreases with beam-energy
- C_k decreases with beam-energy
- $cov(v_2^2, [p_T])$ decreases with beam-energy

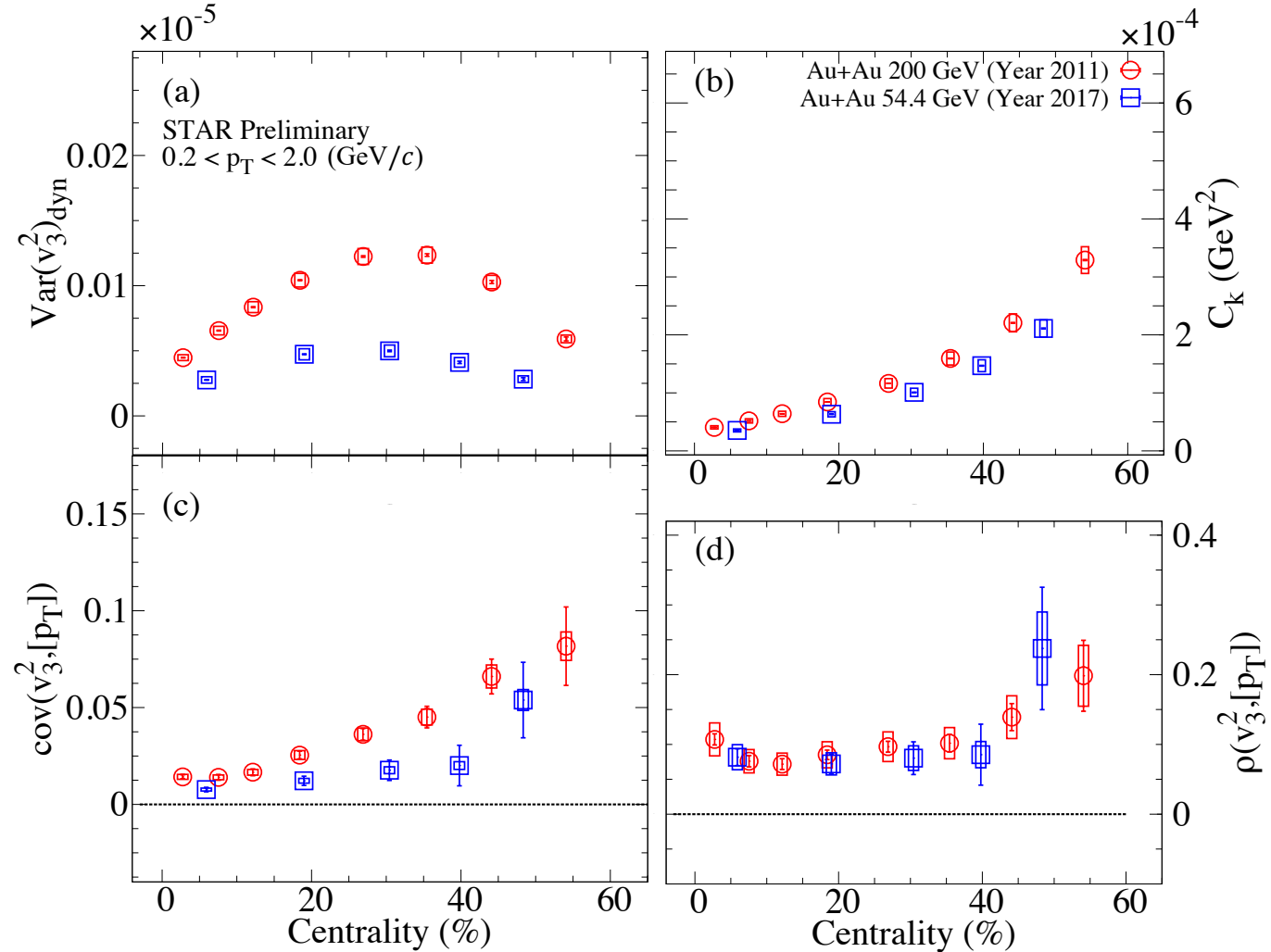
➤ The Pearson correlation, $\rho(v_2^2, [p_T])$, shows no significant energy dependence within the systematic uncertainties



Transverse momentum-flow correlations measurements:

❖ The beam-energy dependance of the transverse momentum-flow correlations

- $Var(v_3^2)_{dyn}$ decreases with beam-energy
- C_k decreases with beam-energy
- $cov(v_3^2, [p_T])$ decreases with beam-energy
- The Pearson correlation, $\rho(v_3^2, [p_T])$, show no significant energy dependence within the systematic uncertainties



Conclusions



We studied the transverse momentum and the transverse momentum-flow correlations as a function centrality for different beam energies

- The extracted $a_2^{p_T}$:
 - ✓ Decrease with harmonic order
 - ✓ Models don't describe the $a_2^{p_T}$ data
 - ✓ Event shape dependent
- The slope of $\sigma_{\Delta\eta}(G_2)$ vs multiplicity is:
 - ✓ Softer for RHIC (indicating smaller η/s for RHIC) than LHC
 - ✓ Event shape independent
- Transverse momentum-flow correlations:
 - ✓ The $cov(v_n^2, [p_T])$ increases with beam energy
 - ✓ The normalized $\rho(v_n^2, [p_T])$:
Show little, if any, change with beam energy

These comparisons are reflecting the efficacy of the $G_2(\Delta\eta, \Delta\phi)$ correlator to differentiate among theoretical models as well as to constrain the η/s .

The $\rho(v_n^2, [p_T])$ measurements show little, if any, change with beam energy, suggesting that $\rho(v_n^2, [p_T])$ is dominated by initial state effects.

Thank You