



Collision-system Dependence of the Charge Separation
Relative to the Event Plane:
Implications for Chiral Magnetic Effect Search in STAR

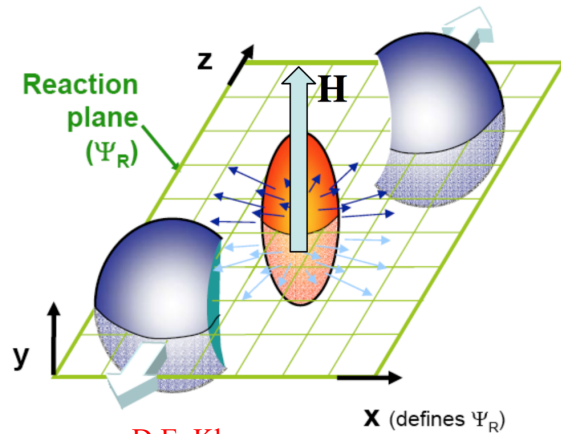
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Outline

- Introduction
- $R_{\Psi_m}(\Delta S)$ Correlator
- $\Delta \gamma$ vs $R_{\Psi_m}(\Delta S)$
- Data Analyses
- Results
- Conclusion

❖ Introduction

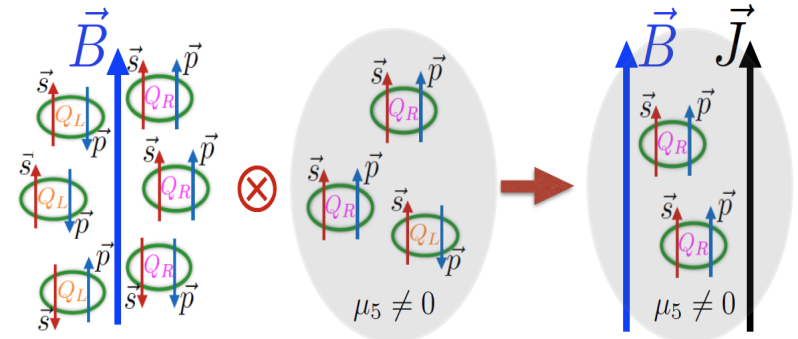
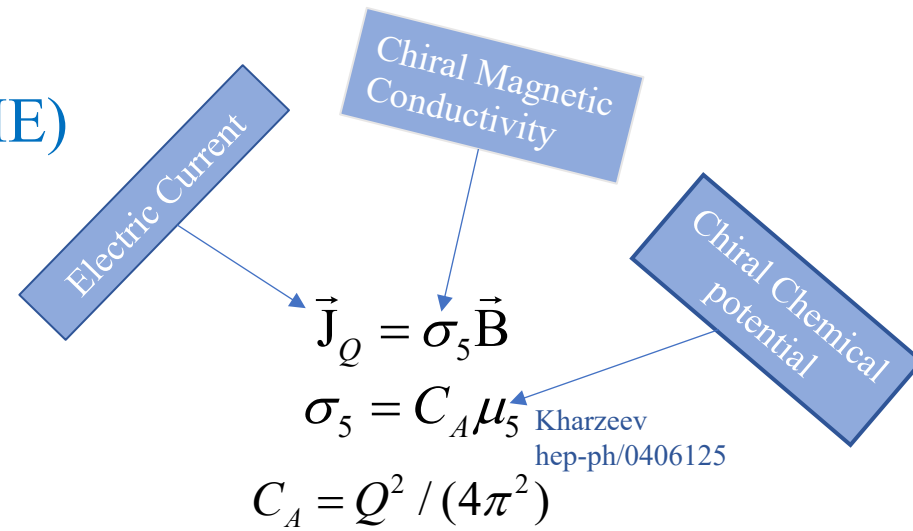
✓ Chiral Magnetic Effect (CME)



D.E. Kharzeev

Prog.Part.Nucl.Phys. 75 (2014) 133-151

- In non-central collisions, a strong magnetic field is created \perp to Ψ_{RP}



D.E. Kharzeev et al.

Prog.Part.Nucl.Phys. 88 (2016) 1-28

- Magnetic field acts on the chiral fermions with $\mu_5 \neq 0$ leading to an electric current along the magnetic field which leads to a charge separation

❖ Introduction

➤ Chiral Magnetic Effect (CME)

CME-driven charge separation leads to a dipole term in the azimuthal distribution of the produced charged hadrons:

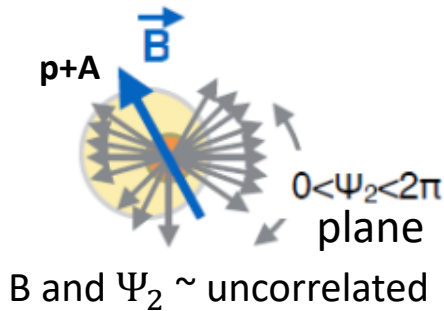
$$\frac{dN^{ch}}{d\phi} \propto 1 \pm 2 a_1^{ch} \sin(\phi) + \dots$$

$$a_1^{ch} \propto \mu_5 \vec{B}$$

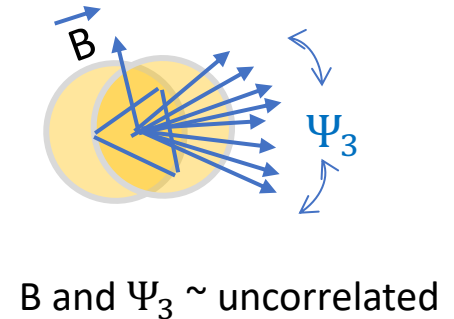
Can we identify & characterize this dipole moment?

➤ What a good correlator should establish?

✓ Leverage Small systems



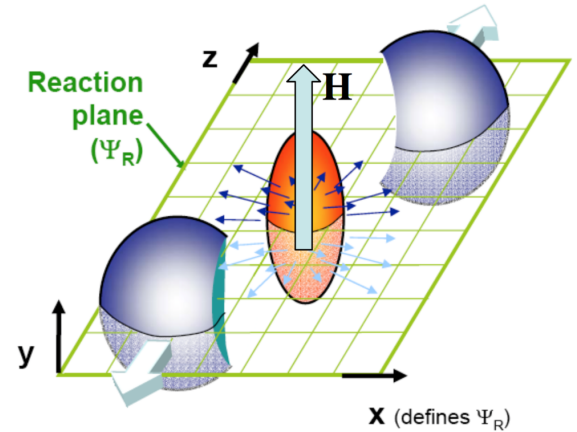
✓ Leverage Ψ_3 measurements



✓ Excellent benchmark

Strength : $eB \sim (m_\pi)^2 \sim 10^{18}$ Gauss

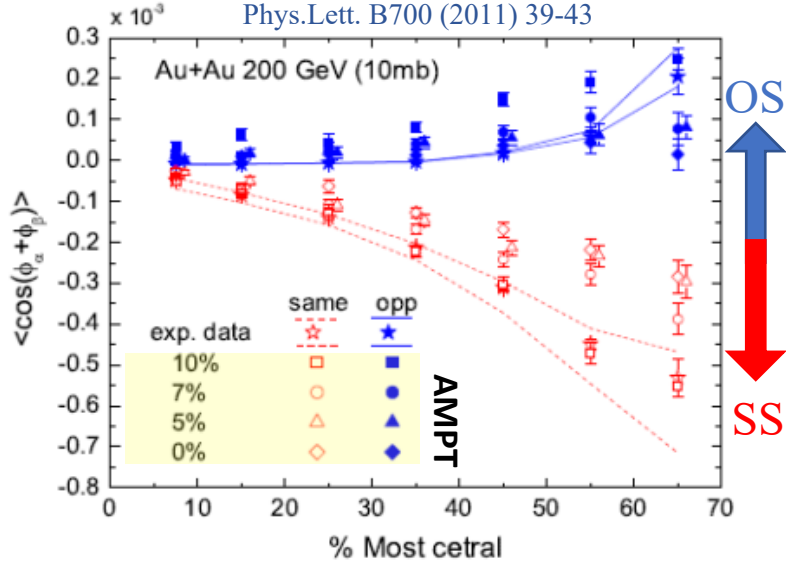
D.E. Kharzeev
Prog.Part.Nucl.Phys. 75 (2014) 133-151



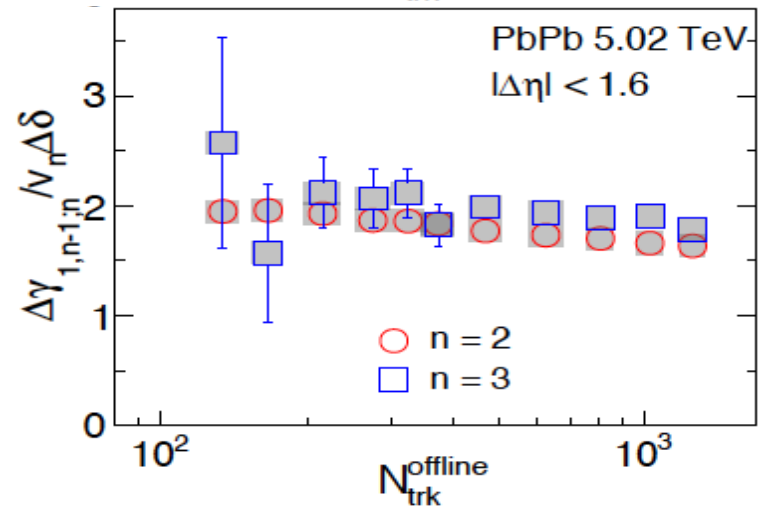
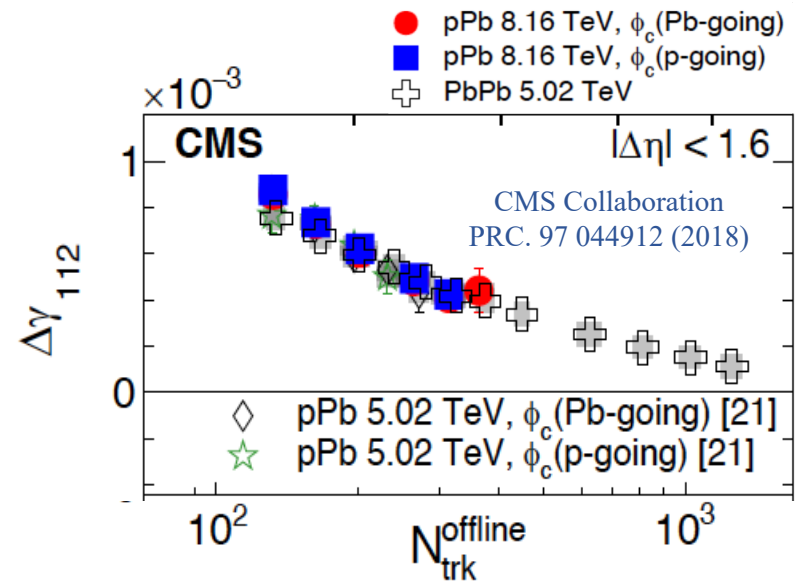
❖ Introduction

✓ The Gamma correlator

Ma, Zhang
Phys.Lett. B700 (2011) 39-43



- The Gamma Correlator's response is similar for signal and background
 - ✓ Background-driven correlations complicate CME-driven signal extraction
- Can background account for a **part**, or **all** of the observed charge separation signal?



- CMS at QM-18
 - ✓ pPb consistent with 100% BKG
 - ✓ Same for PbPb

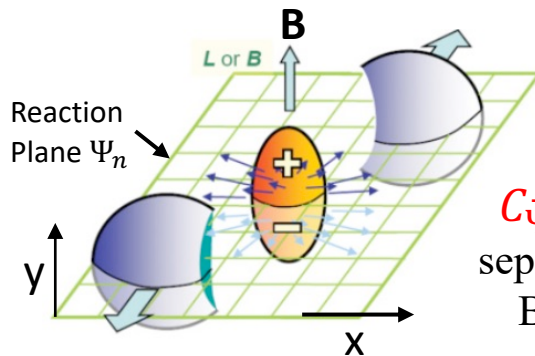
❖ $R_{\Psi_m}(\Delta S)$ Correlator

N. Magdy, et al.
PRC 97, 061901 (2018)

- The correlator is constructed for a given event plane Ψ_m via a ratio of two correlation functions

$C_{\Psi_2}(\Delta S)$ quantifies charge separation along the B-field

$$R_{\Psi_m}(\Delta S) = \frac{C_{\Psi_m}(\Delta S)}{C_{\Psi_m}^{\perp}(\Delta S)}, m = 2 \text{ and } 3$$



$C_{\Psi_2}^{\perp}(\Delta S)$ quantifies charge separation perpendicular to the B-field (only background)

The $R_{\Psi_2}(\Delta S)$ correlator measures the magnitude of charge separation parallel to the B-field, relative to that for charge separation perpendicular to the B-field

Note that both $C_{\Psi_3}(\Delta S)$ and $C_{\Psi_3}^{\perp}(\Delta S)$ are insensitive to the CME-driven charge separation (only background)

❖ $R_{\Psi_m}(\Delta S)$ Correlator

N. Magdy, et al.
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$$R_{\Psi_m}(\Delta S) = \frac{C_{\Psi_m}(\Delta S)}{C_{\Psi_m}^{\perp}(\Delta S)}$$

$$C_{\Psi_m}(\Delta S) = \frac{N(\Delta S)}{N(\Delta S_{sh})}$$

$$C_{\Psi_m}^{\perp}(\Delta S) = \frac{N(\Delta S^{\perp})}{N(\Delta S_{sh}^{\perp})}$$

$$\Delta\varphi = \varphi - \Psi_m$$

$N(\Delta S)$

$$\langle S_{\Psi_m}^+ \rangle = \frac{\sum_1^p w_p \sin(\frac{m}{2} \Delta\varphi)}{\sum_1^p w_p}$$

$$\langle S_{\Psi_m}^- \rangle = \frac{\sum_1^n w_n \sin(\frac{m}{2} \Delta\varphi)}{\sum_1^n w_n}$$

$$\Delta S = \langle S_{\Psi_m}^+ \rangle - \langle S_{\Psi_m}^- \rangle$$

Sensitive to charge separation
(CME and Background)

w_i : charge dependent
detector acceptance.

$N(\Delta S)$

$$\langle S_{\Psi_m}^+ \rangle^{\perp} = \frac{\sum_1^p w_p \cos(\frac{m}{2} \Delta\varphi)}{\sum_1^p w_p}$$

$$\langle S_{\Psi_m}^- \rangle^{\perp} = \frac{\sum_1^n w_n \cos(\frac{m}{2} \Delta\varphi)}{\sum_1^n w_n}$$

$$\Delta S^{\perp} = \langle S_{\Psi_m}^+ \rangle^{\perp} - \langle S_{\Psi_m}^- \rangle^{\perp}$$

$N(\Delta S_{sh})$

$$\Delta S_{sh} = \langle S_{\Psi_m}^+ \rangle_{sh} - \langle S_{\Psi_m}^- \rangle_{sh}$$

Shuffling of charges within an
event breaks the charge
separation sensitivity

$N(\Delta S_{sh})$

$$\Delta S_{sh}^{\perp} = \langle S_{\Psi_m}^+ \rangle_{sh}^{\perp} - \langle S_{\Psi_m}^- \rangle_{sh}^{\perp}$$

❖ $R_{\Psi_m}(\Delta S)$ Correlator

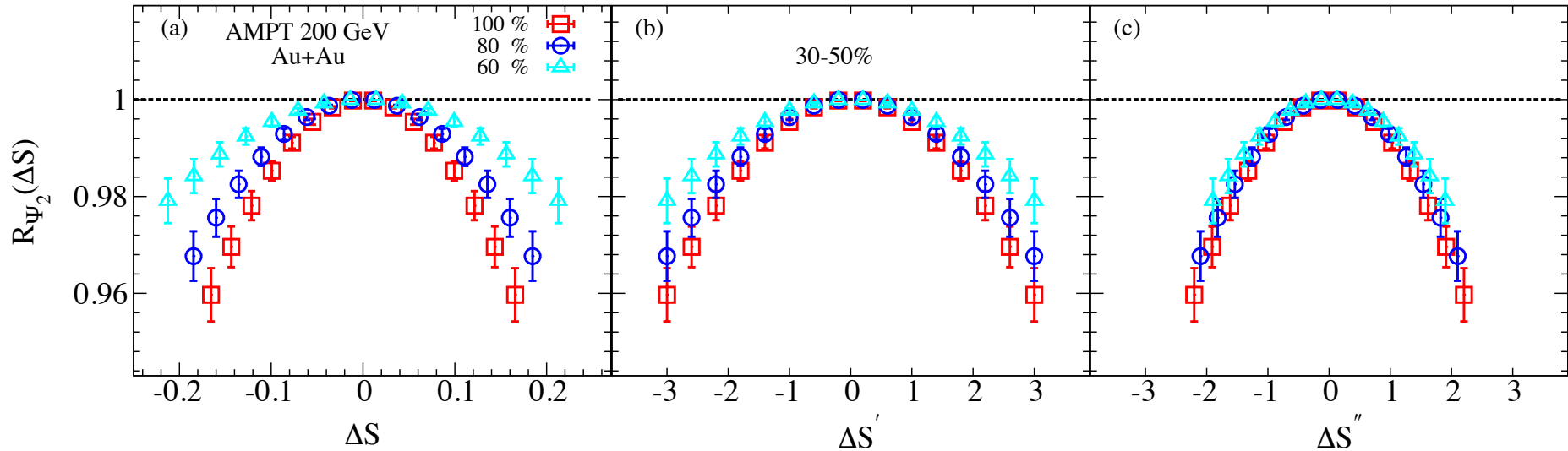
➤ Charge separation magnitude is reflected in the width of the $R_{\Psi_m}(\Delta S)$ distribution which is affected by:

✓ Number fluctuations

✓ Empirical EP-resolution

$$\Delta S' = \Delta S / \sigma_{\Delta S^{sh}}$$

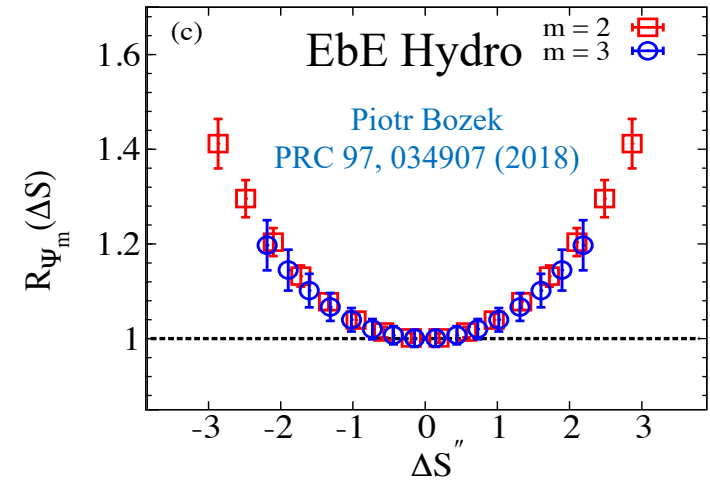
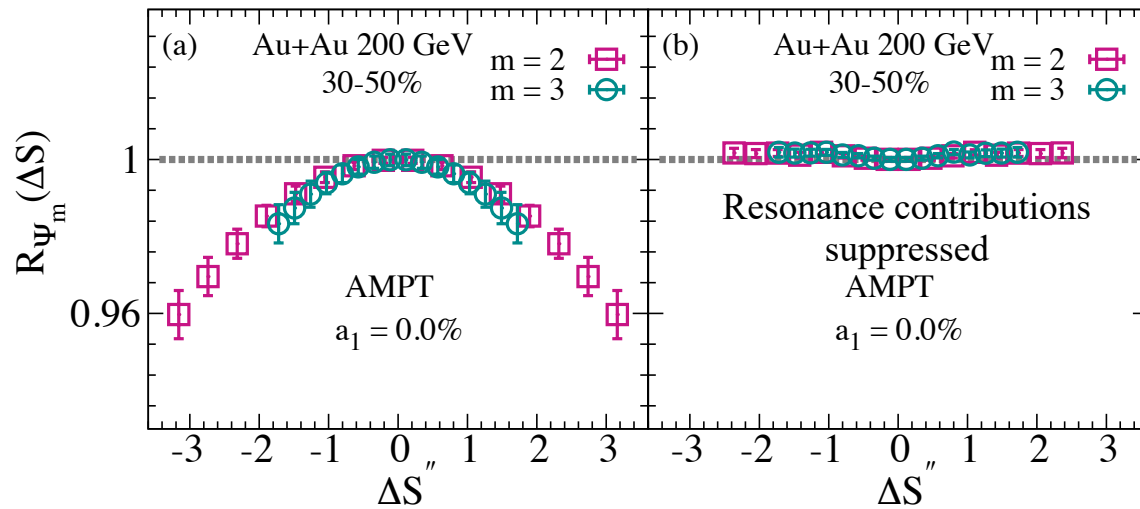
$$\Delta S'' = \Delta S' / e^{\left(\left(1 - \sqrt{\langle \cos(2\Delta\Psi_2) \rangle} \right)^2 \right)}$$



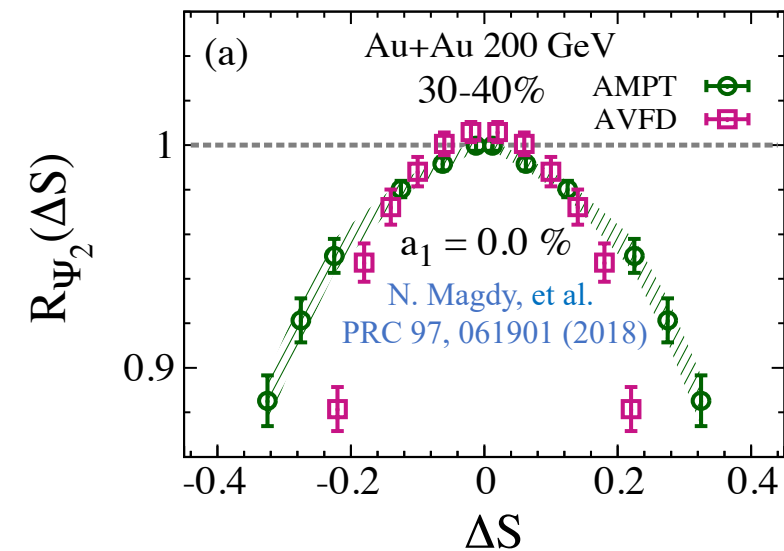
➤ We can account for both number fluctuations and EP-resolution effect on the width of the $R_{\Psi_m}(\Delta S)$

❖ $R_{\Psi_m}(\Delta S)$ Correlator

➤ $R_{\Psi_m}(\Delta S)$ response in background models



✓ Resonance suppression leads to flat R_{Ψ_m}



✓ The R_{Ψ_2} and R_{Ψ_3} give similar response to the background irrespective of the correlator shape

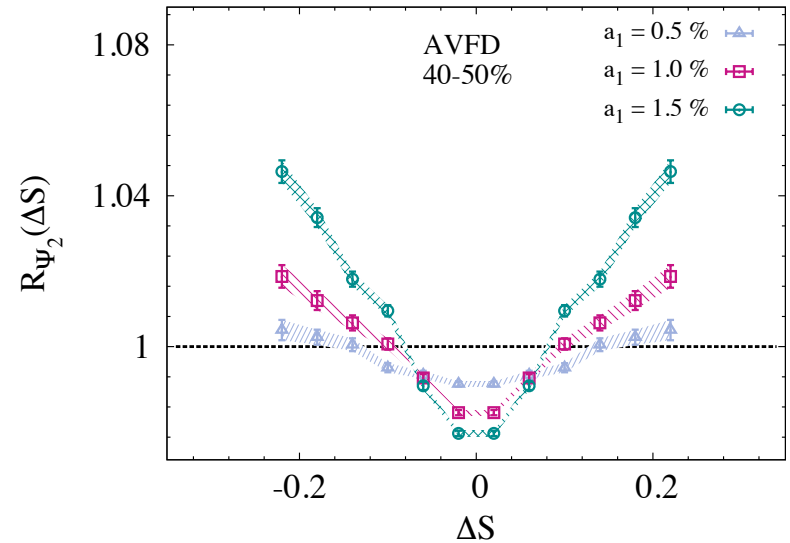
❖ $R_{\Psi_m}(\Delta S)$ Correlator

➤ $R_{\Psi_m}(\Delta S)$ response in CME models

- Signal magnitude is reflected in the widths of the distributions
 - ✓ Smaller widths for larger input signal
- Validation of the expected concave-shaped response of $R_{\Psi_2}(\Delta S)$ to the CME-driven charge separation input in CME-events.

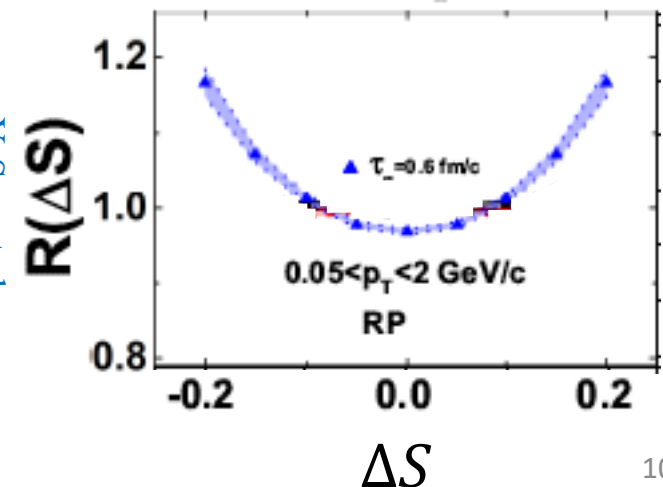
Concaved-shape distribution for input charge separation

N. Magdy, et al.
PRC 97, 061901 (2018)



Chiral kinetic approach
Ru+Ru @ 200 GeV, $b=7$ fm, $|\eta| \leq 1$

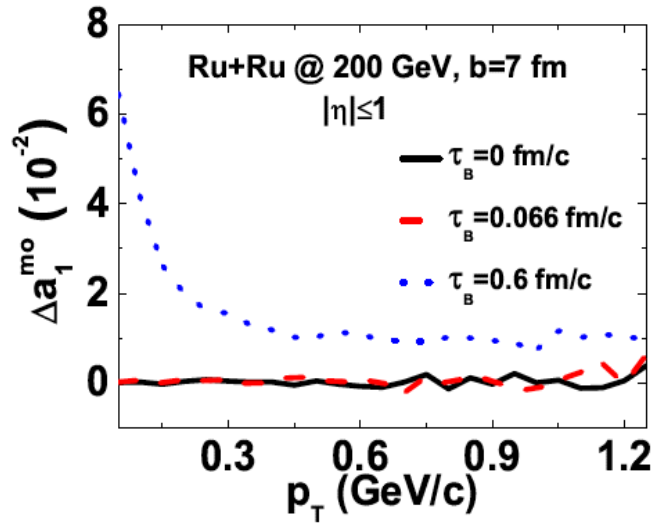
Y. Sun et al.
PRC 98, 014911 (2018)



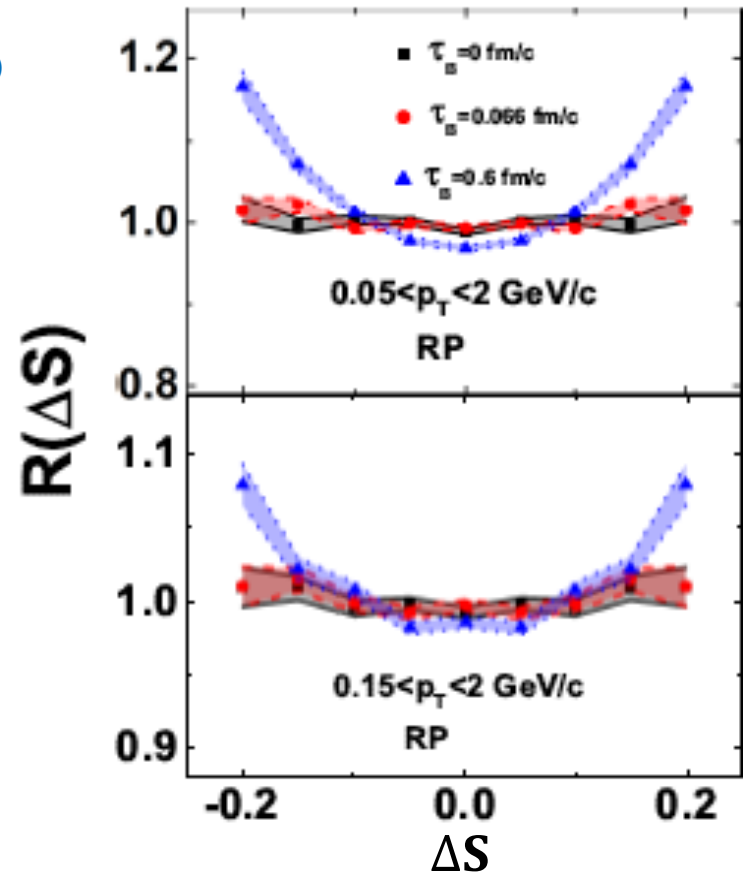
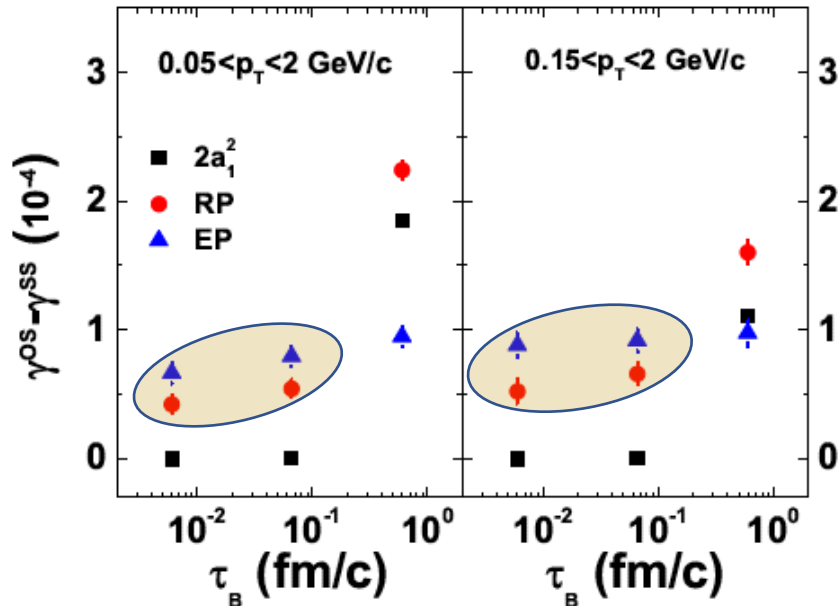
❖ $\Delta\gamma$ vs $R_{\Psi m}(\Delta S)$

Chiral kinetic approach

Y. Sun et al.
PRC 98, 014911 (2018)



Ru+Ru @ 200 GeV, b=7 fm, $|\eta| \leq 1$

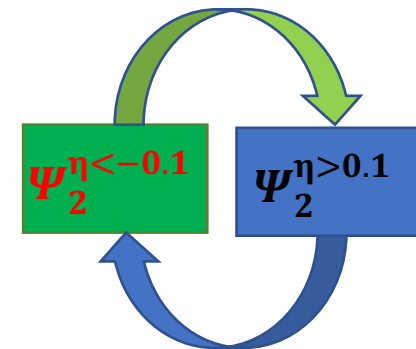
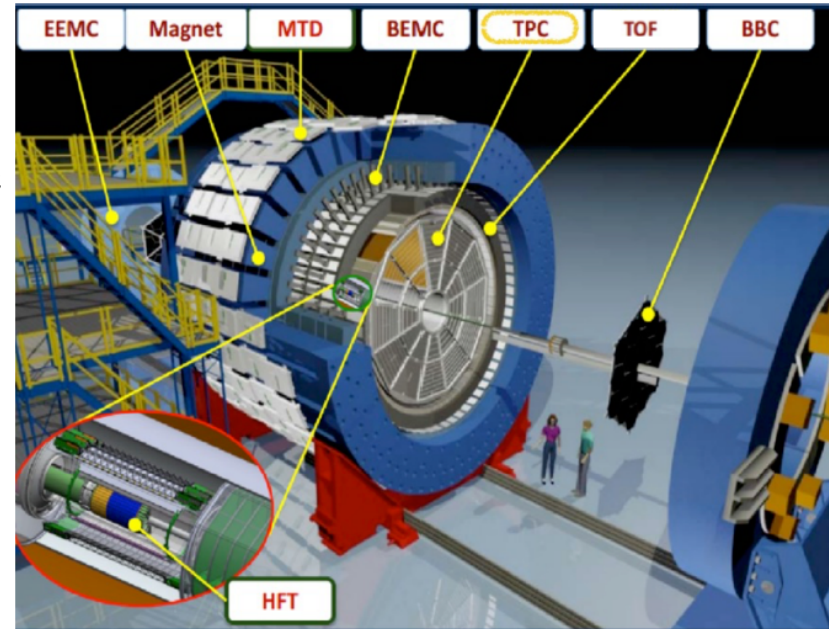


- $a_1 \sim 0$ for $\tau_B < 0.6$ fm/c
- ✓ $\Delta\gamma$ shows non-zero value
- ✓ $R_{\Psi 2}(\Delta S)$ shows a flat shape

❖ Data Analysis

The STAR Detector at RHIC

- The TPC detector is used in the current analysis
- Charged hadrons with $0.2 < p_T < 2.0 \text{ GeV}/c$ used to construct $\Psi_2^{\eta > 0.1}$ & $\Psi_2^{\eta < -0.1}$
- Particles with $0.35 < p_T < 2.0 \text{ GeV}/c$ and $\eta < 0$ analyzed using $\Psi_2^{\eta > 0.1}$
- Particles with $0.35 < p_T < 2.0 \text{ GeV}/c$ and $\eta > 0$ analyzed using $\Psi_2^{\eta < -0.1}$



❖ Data Analysis

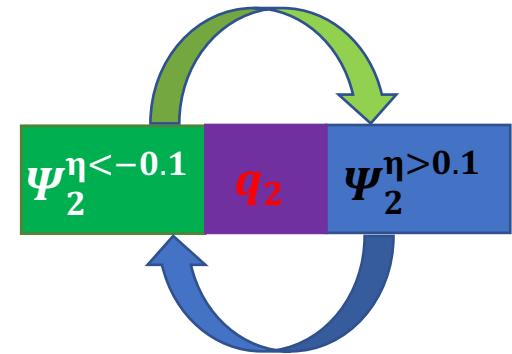
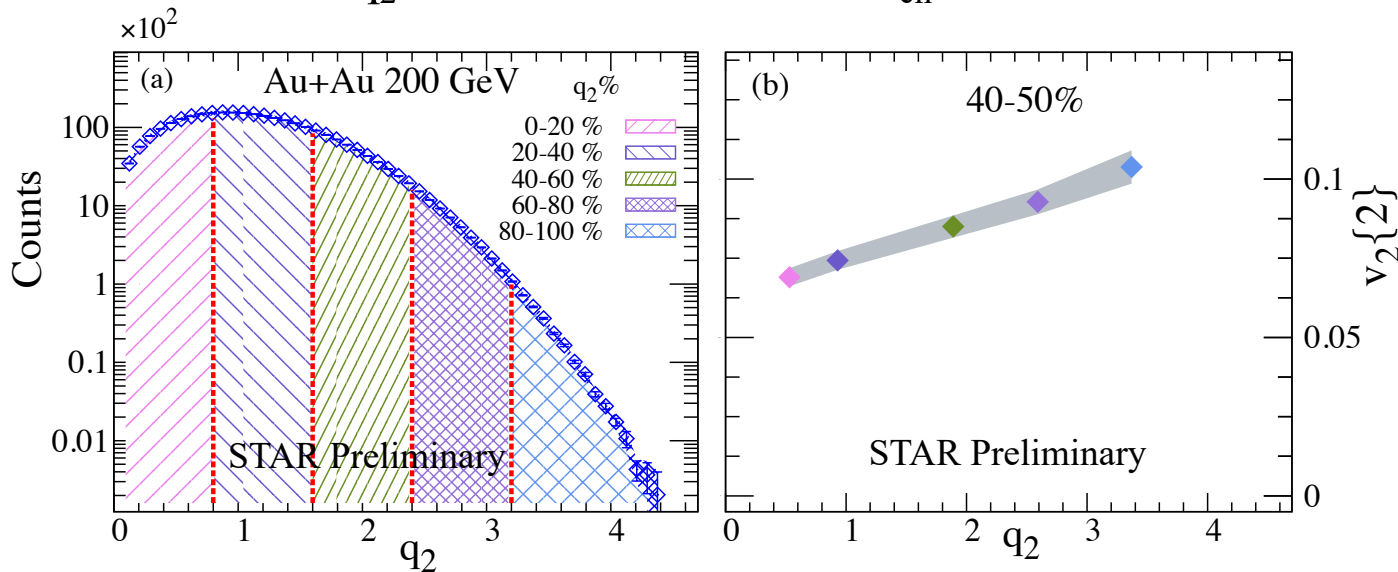
➤ Event-shape selections (Data)

- ✓ Events are further subdivided into groups with different q_2 magnitude:

$$Q_{2,x} = \sum_{i=1}^M \cos(2 \varphi_i) \quad Q_{2,y} = \sum_{i=1}^M \sin(2 \varphi_i)$$

$$|Q_2| = \sqrt{Q_{2,x}^2 + Q_{2,y}^2} \quad q_2 = \frac{|Q_2|}{\sqrt{M}}$$

- ✓ The q_2 was created for each N_{ch}



- $v_2\{2\}$ increases linearly with q_2
 - ✓ q_2 is good event-shape selector

Results

❖ Results

➤ $R_{\Psi_m}(\Delta S)$ measurements at RHIC

- ✓ Event plane dependence
- ✓ Small systems dependence

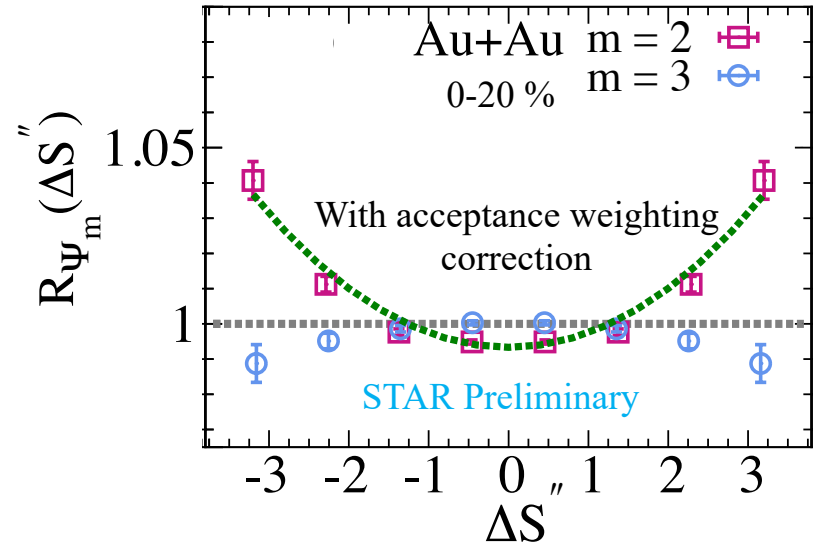
➤ Measurements for R_{Ψ_m} show:

- ✓ Different response for R_{Ψ_2} and R_{Ψ_3}
- ✓ Different response for small (p(d)+Au) and large (Au+Au) systems

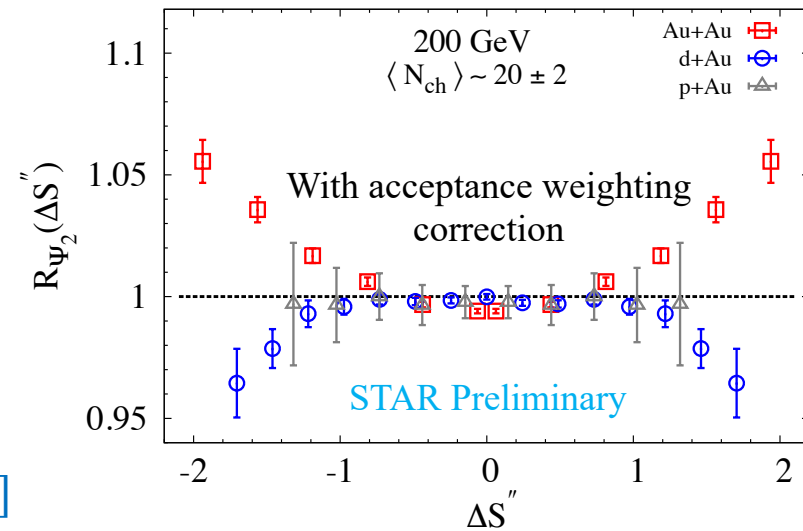
➤ R_{Ψ_m} results are consistent with the expectation for the CME-driven charge separation.

- ✓ Note that these observations contrast with those from the γ correlator.

[CMS Collaboration arXiv:1610.00263]

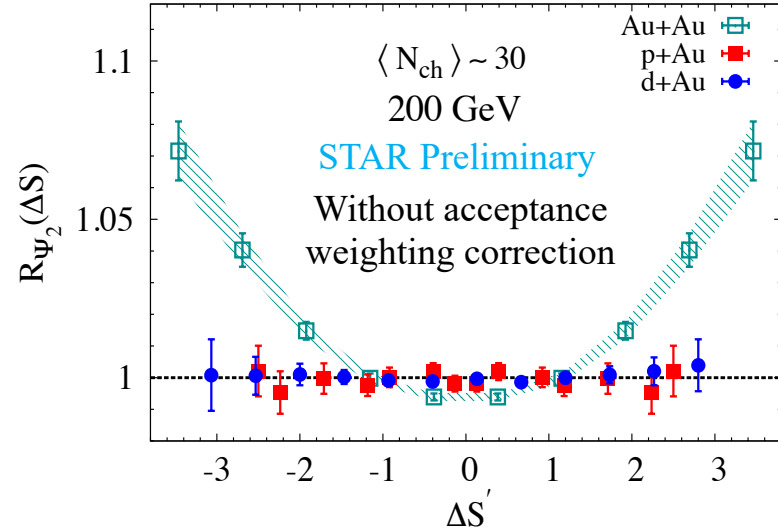
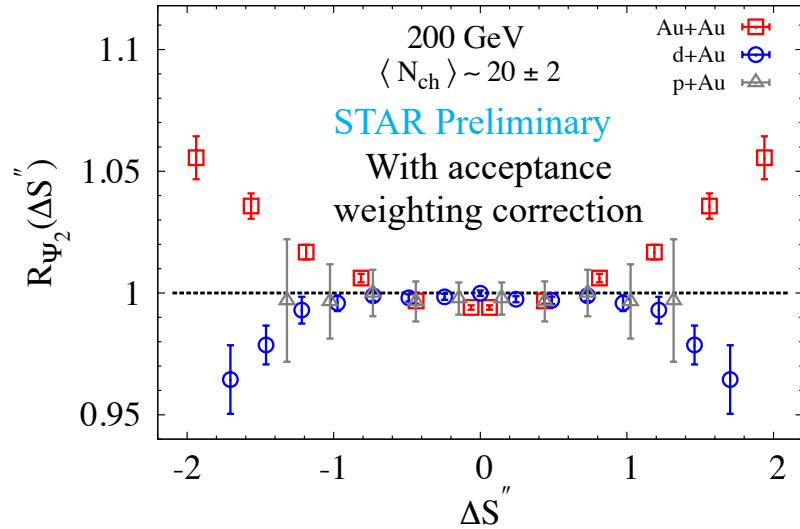


$$\frac{Res(p + Au)}{Res(Au + Au)} = 0.45 \quad \frac{Res(d + Au)}{Res(Au + Au)} = 0.85$$



❖ Results

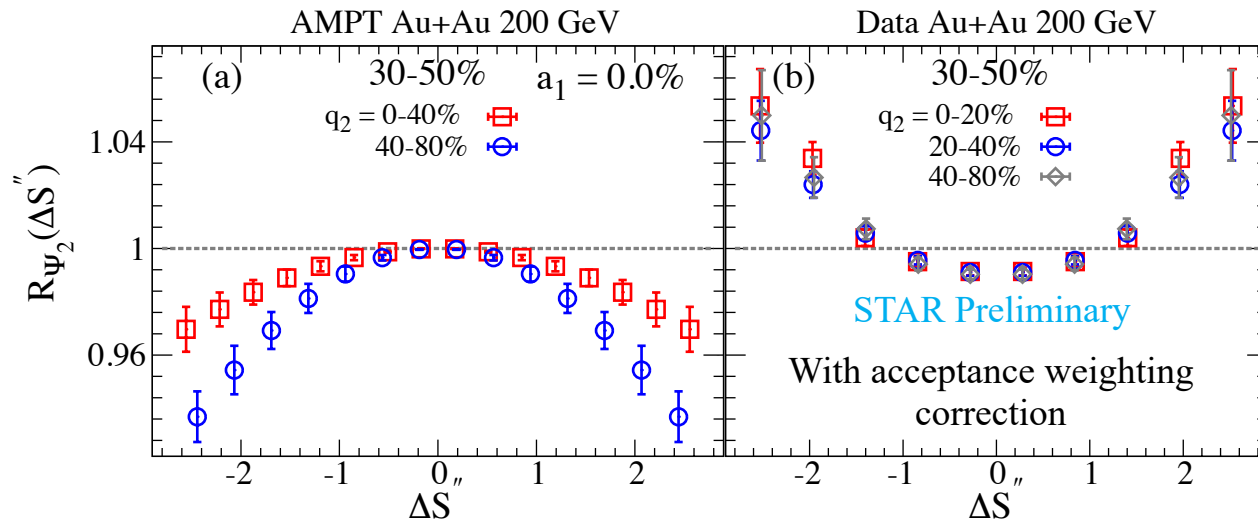
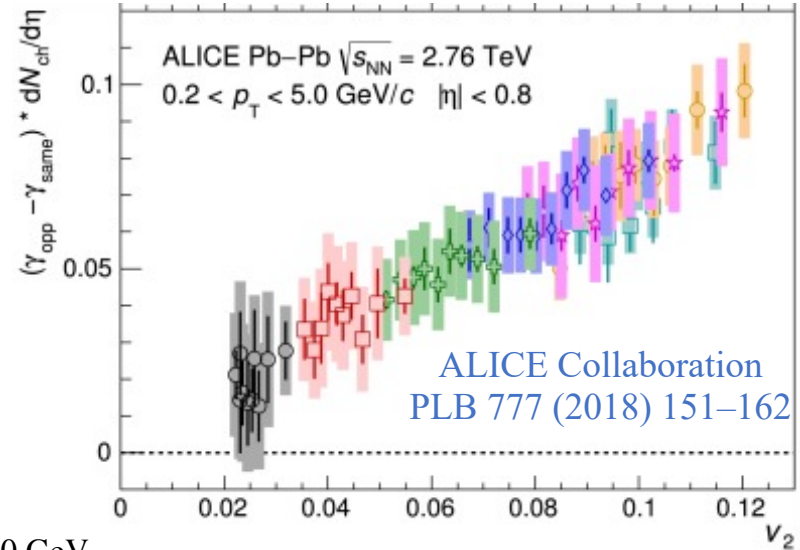
- Using acceptance correction (w_i) the charge dependent detector acceptance.



❖ Results

➤ Event shape selection

Comparison of the $R_{\Psi_2}(\Delta S'')$ correlators for q_2 selected events for 30–50% central, Au+Au collisions at 200 GeV for (a) AMPT simulations and (b) experimental measurements.

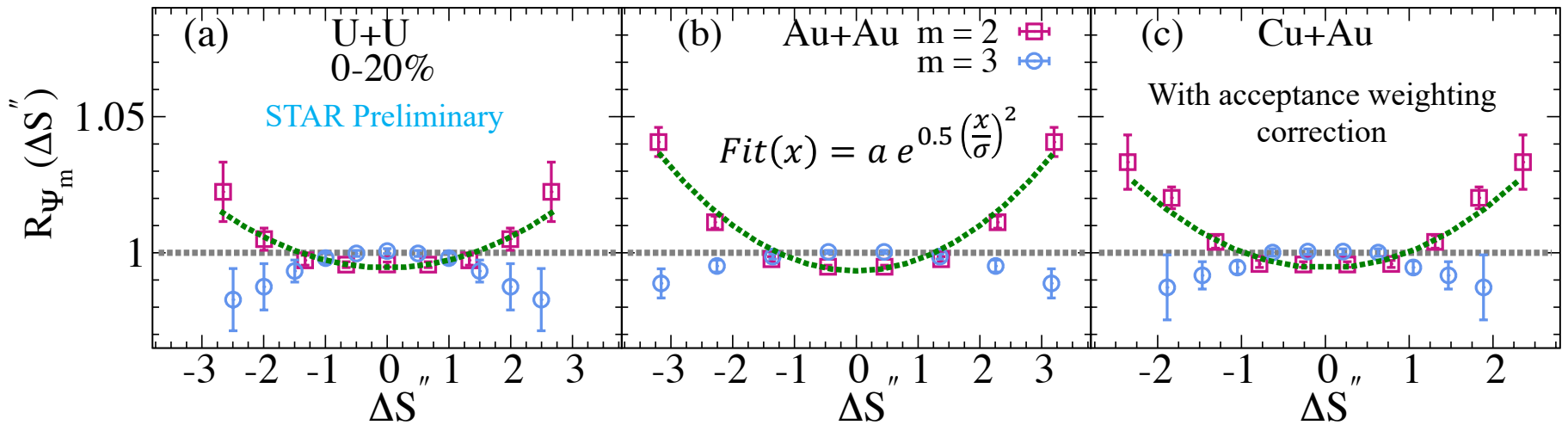


➤ Different q_2 selections (right panel) suggests that $R_{\Psi_2}(\Delta S'')$ is not strongly influenced by the v_2 background-driven charge separation.

❖ Results

- Different collision systems $R_{\Psi_2}(\Delta S'')$ and $R_{\Psi_3}(\Delta S'')$

$R_{\Psi_2}(\Delta S'')$ and $R_{\Psi_3}(\Delta S'')$ for 0-20% centrality selection in different collision systems.

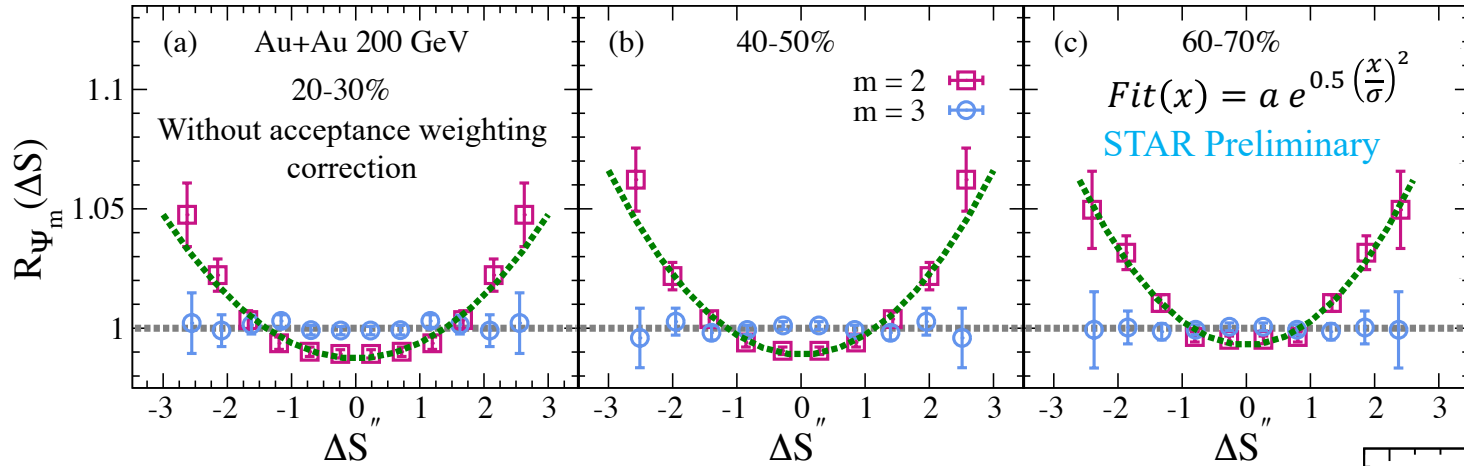


- The $R_{\Psi_2}(\Delta S'')$ correlators for different collision systems is strikingly different from those for $R_{\Psi_3}(\Delta S'')$ correlators.
- $R_{\Psi_2}(\Delta S'')$ decidedly concave-shaped, as would be expected for CME-driven charge separation with limited influence from background-driven charge separation.

❖ Results

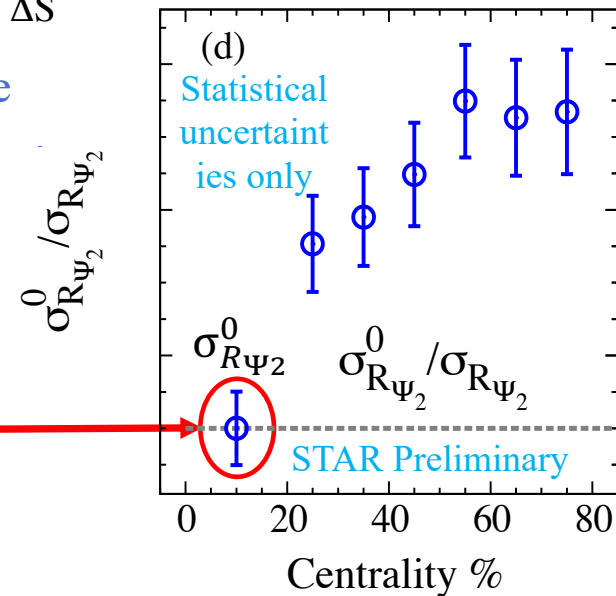
➤ Au+Au collision system $R_{\Psi_2}(\Delta S)$

$R_{\Psi_2}(\Delta S)$ for different centrality selections in Au+Au at 200 GeV.



➤ $\sigma_{R_{\Psi_2}}^{-1}$ indicates a sizable centrality dependence, suggestive of the expected increase in the magnitude of the CME-driven charge separation when the \vec{B} -field increases as collisions become more peripheral.

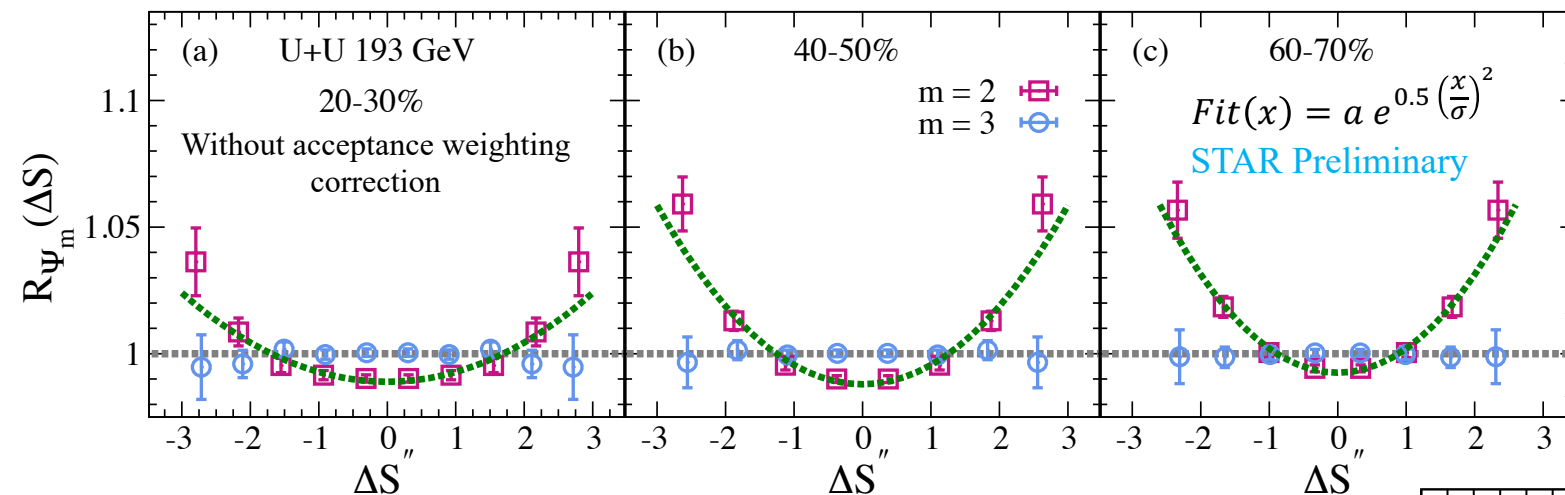
All widths are normalized to the width of 0-20% (smallest a_1)



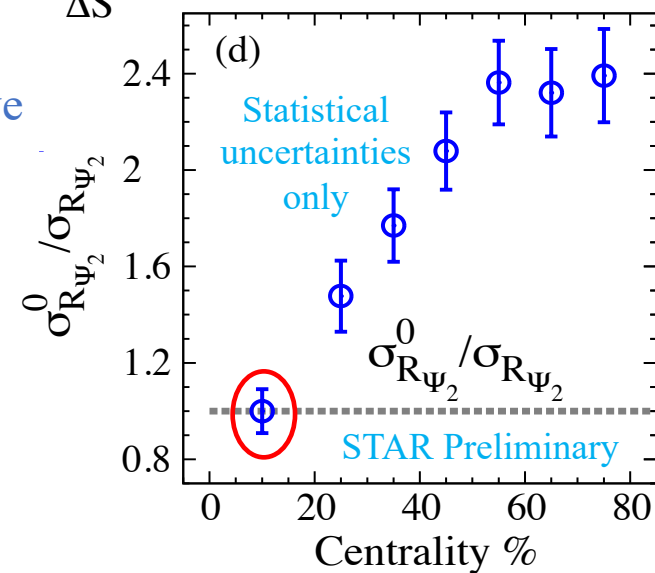
❖ Results

➤ U+U collision system $R_{\Psi_2}(\Delta S)$

$R_{\Psi_2}(\Delta S)$ for different centrality selections in U+U at 200 GeV.



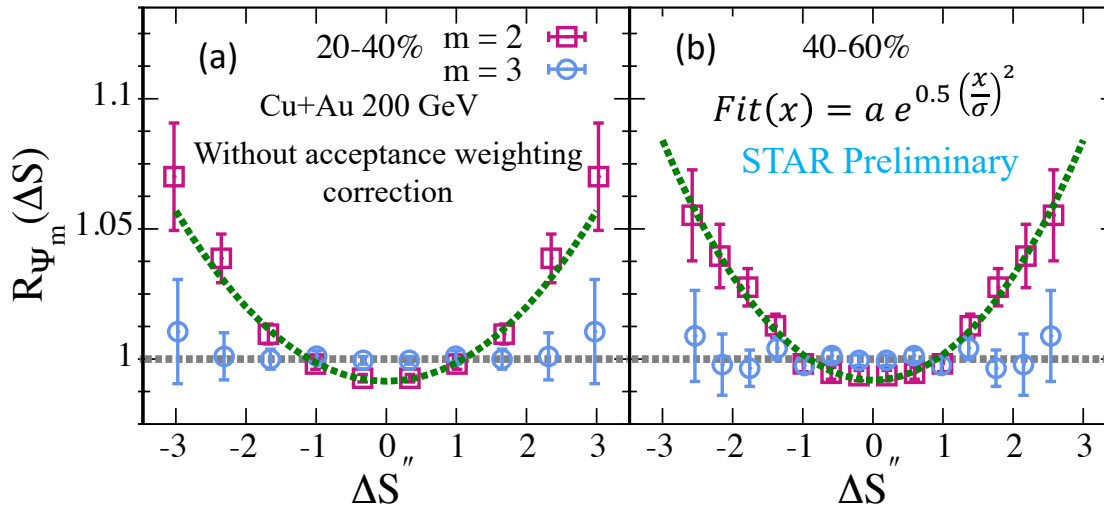
➤ $\sigma_{R_{\Psi_2}}^{-1}$ indicates a sizable centrality dependence, suggestive of the expected increase in the magnitude of the CME-driven charge separation when the \vec{B} -field increases as collisions become more peripheral.



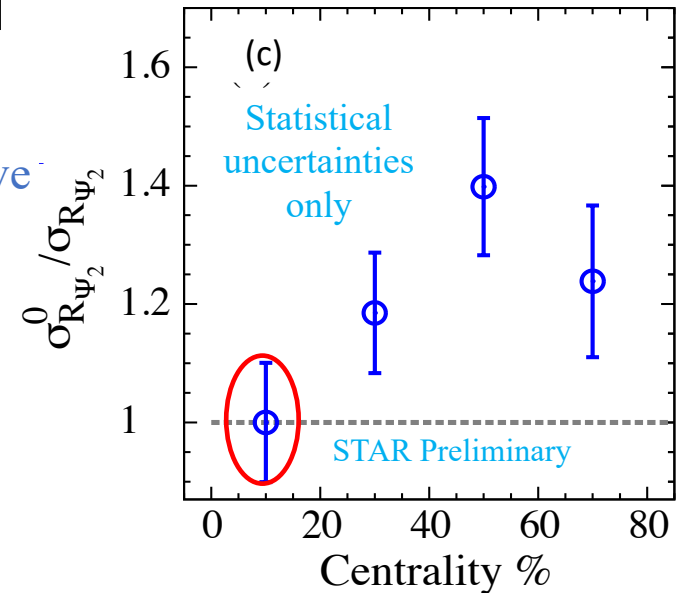
❖ Results

➤ Cu+Au collision system $R_{\Psi_2}(\Delta S)$

$R_{\Psi_2}(\Delta S)$ for different centrality selections in Cu+Au at 200 GeV.



➤ $\sigma_{R_{\Psi_2}}^{-1}$ indicates a sizable centrality dependence, suggestive of the expected increase in the magnitude of the CME-driven charge separation when the \vec{B} -field increases as collisions become more peripheral.



❖ Conclusions

Charge separation measurements performed with R_{Ψ_m} (for $m = 2,3$) correlator, for U+U (193 GeV), p(d)+Au, Au+Au and Cu+Au at 200 GeV collisions with the STAR detector.

➤ R_{Ψ_m} measurements show:

- ✓ Difference in the response for Ψ_2 and Ψ_3 for different collision systems
- ✓ Difference in the response for small (p(d)+Au) and large systems (Au+Au)

➤ R_{Ψ_2} width:

- ✓ Is q_2 -independent (weak sensitivity to background)
- ✓ Is centrality-dependent

➤ R_{Ψ_m} results are consistent with the expectation for CME-driven charge separation.