



Three-dimensional Kaon Source Extraction from STAR experiment at RHIC

Michal Šumbera (Nuclear Physics Institute ASCR), for the STAR Collaboration



Abstract

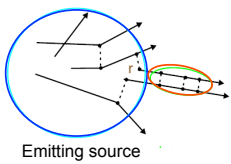
Three-dimensional source imaging techniques in conjunction with detailed model comparisons have shown the viability of disentangling the spatio-temporal information contained in two-pion interferometric measurements from ultra-relativistic heavy ion collisions. This has led to the observation of non-Gaussian tails in the 3D pion source function and the extraction of finite pion emission duration at RHIC energies [1]. The STAR Collaboration has recently extracted the 3D kaon source function from a high statistics two-kaon interferometric measurement from Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV. Such measurement offers a window into the fireball freeze-out dynamics with a much cleaner probe with smaller resonance decay contributions than for the pion case. The extracted space-time characteristics are compared with those obtained from pion analysis. Kaon emission is inferred to be instantaneous from a freeze-out hypersurface bearing no space-time correlation.

Motivation

Kaon 1D source shape results

Source imaging:

Model-independent femtoscopic technique of emission source shape extraction (beyond Gaussian assumption). D. Brown and P. Danielewicz [2].



Inversion of linear integral equation to obtain source function
1D Koonin-Pratt equation:

$$R(q) \equiv C(q) - 1 = 4\pi \int dr r^2 K(q,r) S(r)$$

Correlation function

Kernel (Encodes FS)

Source function (Distribution of pair separations in pair rest frame)

Freeze-out occurs after last scattering

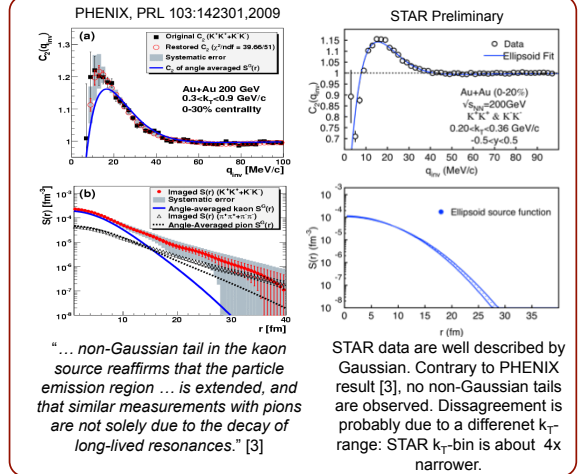
⇒ Only Coulomb & QS effect included in the kernel: $K(q,r) = \frac{1}{2} \int d\cos\theta_{q\vec{r}} \left[\phi(\vec{q},\vec{r})^2 - 1 \right]$

Inversion procedure: $S(r) = \sum_j S_j \cdot B_j(r)$ Expansion into B-spline basis

$$C^{Th}(q_i) = \sum_j K_{ij} \cdot S_j$$

$$K_{ij} = \int dr \cdot K(q_i,r) B_j(r)$$

$$\chi^2 = \frac{(C^{Exp}(q_i) - \sum_j K_{ij} \cdot S_j)^2}{(\Delta C^{Exp}(q_i))^2}$$



3D source shape analysis

Kaon and pion 3D source shape results

Expansion of $R(\vec{q})$ and $S(\vec{r})$ in Cartesian Harmonics basis [4]

$$R(\vec{q}) = \sum_{\alpha_1, \alpha_2, \alpha_3} R'_{\alpha_1, \alpha_2, \alpha_3}(\vec{q}) A'_{\alpha_1, \alpha_2, \alpha_3}(\Omega_{\vec{q}}) \quad (1)$$

$$S(\vec{r}) = \sum_{\alpha_1, \alpha_2, \alpha_3} S'_{\alpha_1, \alpha_2, \alpha_3}(\vec{r}) A'_{\alpha_1, \alpha_2, \alpha_3}(\Omega_{\vec{r}}) \quad (2)$$

$\alpha_i = x, y$ or z
 $x =$ out-direction
 $y =$ side-direction
 $z =$ long-direction

- Cartesian surface-spherical harmonic are based on the products of unit vector components, $n_{\alpha_1} n_{\alpha_2} \dots n_{\alpha_\ell}$.
- Unlike the spherical harmonics they are real.
- Due to the normalization $n_x^2 + n_y^2 + n_z^2 = 1$, at a given $\ell \geq 2$, the different component products are not linearly independent as functions of spherical angle.
- At a given ℓ , the products are spanned by spherical harmonics of rank $\ell' \leq \ell$, with ℓ' of the same evenness as ℓ .

$$R(\vec{q}) = C(\vec{q}) - 1 = 4\pi \int dr r^2 K(\vec{q},\vec{r}) S(\vec{r}) \quad (3) \quad \text{3D Koonin-Pratt eq.}$$

$$\text{Plug (1) \& (2) into (3)} \Rightarrow R'_{\alpha_1, \alpha_2, \alpha_3}(\vec{q}) = 4\pi \int dr r^2 K_{\alpha_1, \alpha_2, \alpha_3}(\vec{q},\vec{r}) S'_{\alpha_1, \alpha_2, \alpha_3}(\vec{r}) \quad (4)$$

$$\text{Invert (1)} \Rightarrow R'_{\alpha_1, \alpha_2, \alpha_3}(\vec{q}) = \frac{(2\ell+1)!!}{\ell!} \int \frac{d\Omega_{\vec{q}}}{4\pi} A'_{\alpha_1, \alpha_2, \alpha_3}(\Omega_{\vec{q}}) R(\vec{q})$$

$$\text{Invert (2)} \Rightarrow S'_{\alpha_1, \alpha_2, \alpha_3}(\vec{r}) = \frac{(2\ell+1)!!}{\ell!} \int \frac{d\Omega_{\vec{r}}}{4\pi} A'_{\alpha_1, \alpha_2, \alpha_3}(\Omega_{\vec{r}}) S(\vec{r})$$

Fit to the 3D correlation function with a trial functional form for $S(\vec{r})$. Trial function: 4-parameter ellipsoid (3D Gaussian):

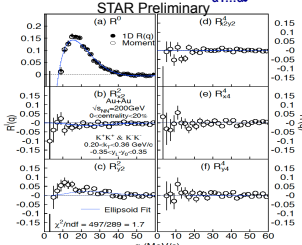
$$S^G(x,y,z) = \frac{\lambda}{(2\pi)^3 r_x r_y r_z} \exp\left[-\frac{x^2}{4r_x^2} - \frac{y^2}{4r_y^2} - \frac{z^2}{4r_z^2}\right]$$

Since the 3D correlation function has been decomposed into its independent moments this is equivalent to a simultaneous fit of 6 independent moments with the trial functional form.

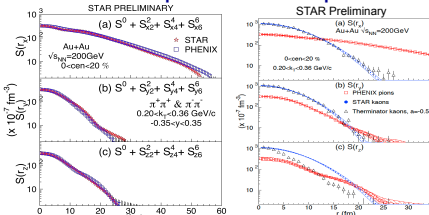
Reference

- [1] S. Afanasiev et al. (PHENIX Collaboration), Phys. Rev. Lett. 100:232301, 2008.
- [2] D. Brown and P. Danielewicz, Phys. Lett. B398:252, 1997; Phys. Rev. C57:2474, 1998.
- [3] S. Afanasiev et al. (PHENIX Collaboration) Phys. Rev. Lett. 103:142301, 2009.
- [4] P. Danielewicz and S. Pratt, Phys. Lett. B618:60, 2005.
- [5] A. Kisiel et al., Phys. Rev. C 73:064902 2006).
- [6] B. I. Abelev et al. (STAR Collaboration), Phys. Rev. C79:034909, 2009.

Correlation moments $R'_{\alpha_1, \dots, \alpha_\ell}$, $0 \leq \ell \leq 4$

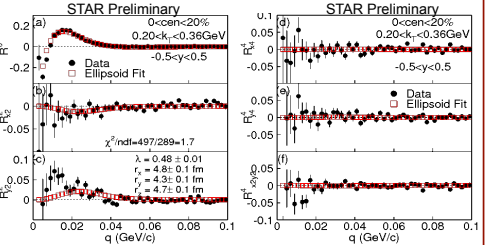


Kaon vs. pion 3D source shape



Good agreement on 3D pion source shape between PHENIX [1] and STAR. Pion [1] and kaon source shapes are very different: Is it due to different dynamics?

Ellipsoid fit



Comparison to the Therminator model [5]

Model main ingredients:

1. Boost invariance in z .
 2. Blast-wave expansion in transverse radius ρ : $v_t(\rho) = (\rho/D_{max})/(D/\rho_{max} + v_t)$, $v_t = 0.445$; $\langle v_t \rangle = 0.6$ [6].
 3. Thermal emission at proper time τ , from a cylinder of infinite length and finite transverse dimension D_{max} .
 4. Emission duration included via $\Delta\tau$.
- Freeze-out occurs at $\tau = \tau_0 + a\rho$. Particles which are emitted at (z, ρ) have LAB emission time $t^2 = (\tau_0 + a\rho)^2 + z^2$.

Summary

1. STAR high-statistics measurements allowed first model-independent extraction of kaon three dimensional source function.
2. The source function of mid-rapidity, low-momentum kaons from central Au+Au collisions at $\sqrt{s_{NN}}=200$ GeV obtained via Cartesian surface-spherical harmonic decomposition method is essentially ellipsoidal in shape. No long-range non-Gaussian tails are observed.
3. Comparison with Therminator model calculations indicates that kaons are emitted from a fireball whose transverse dimension and lifetime are consistent with those extracted with two-pion interferometry.
4. Contrary to an instantaneous emission duration and a negative ρ - τ correlation observed for pions [1], kaons are emitted instantaneously in the source element rest frame from a freeze-out hypersurface with no ρ - τ correlation.
5. Kaons and pions may be subjected to different dynamics owing to their emission over different timescales.