

# 1 Elliptic ( $v_2$ ) and triangular ( $v_3$ ) anisotropic flow of 2 identified hadrons from the STAR Beam Energy 3 Scan program

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## 8 **Abstract.**

9 Elliptic ( $v_2$ ) and triangular ( $v_3$ ) anisotropic flow coefficients for inclusive and identified  
10 charged hadrons ( $\pi^\pm$ ,  $K^\pm$ ,  $p$ ,  $\bar{p}$ ) at midrapidity in Au+Au collisions, measured by the STAR  
11 experiment in the Beam Energy Scan (BES) at the Relativistic Heavy Ion Collider at  $\sqrt{s_{NN}}$   
12 = 11.5 - 62.4 GeV, are presented. We observe that the triangular flow signal ( $v_3$ ) of identified  
13 hadrons exhibits similar trends as first observed for  $v_2$  in Au+Au collisions, i.e. (i) mass  
14 ordering at low transverse momenta,  $p_T < 2$  GeV/c, (ii) meson/baryon splitting at intermediate  
15  $p_T$ ,  $2 < p_T < 4$  GeV/c, and (iii) difference in flow signal of protons and antiprotons. New  
16 measurements of  $v_3$  excitation function could serve as constraints to test different models and  
17 to aid new information about the temperature dependence of the transport properties of the  
18 strongly interacting matter.

## 19 **1. Introduction**

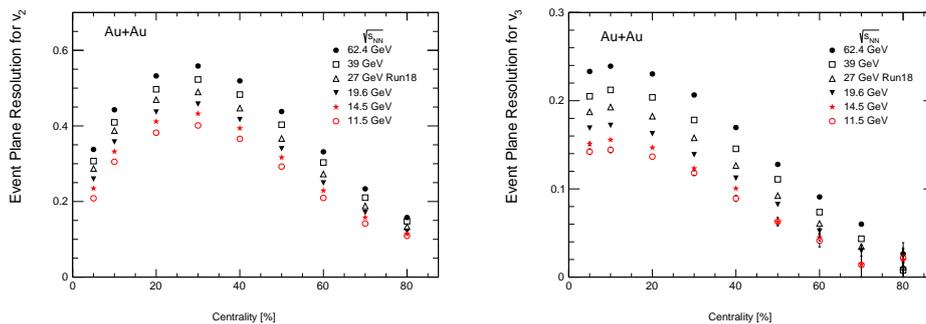
20 The heavy-ion experiments at the Relativistic Heavy Ion Collider (RHIC) and the Large Hadron  
21 Collider (LHC) have established the existence of a strongly coupled Quark Gluon Plasma [1, 2], a  
22 new state of QCD matter with partonic degrees of freedom and with low specific shear viscosity  
23  $\eta/s$  [3]. Lattice QCD calculations [4] indicate that the quark-hadron transition is a smooth  
24 crossover at top RHIC energy and above ( at high temperatures  $T$  and small values of baryonic  
25 chemical potential  $\mu_B$ ). A Beam Energy Scan (BES) program at RHIC plays a central role in  
26 the experimental study of the QCD phase diagram over a wide range in  $T$  and  $\mu_B$  [5, 6].

27 The anisotropic flow is one of the important observables sensitive to the equation of state (EOS)  
28 and transport properties of the strongly interacting matter such as the shear viscosity over  
29 entropy ratio  $\eta/s$  [3, 7, 8]. The azimuthal anisotropy of produced particles can be quantified  
30 by the Fourier coefficients  $v_n$  in the expansion of the particles azimuthal distribution as:  
31  $dN/d\phi \propto 1 + \sum_{n=1} 2v_n \cos(n(\phi - \Psi_n))$  [7, 9], where  $n$  is the order of the harmonic,  $\phi$  is the  
32 azimuthal angle of particles for a given type, and  $\Psi_n$  is the azimuthal angle of the  $n$ th-order  
33 event plane. The  $n^{\text{th}}$ -order flow coefficients  $v_n$  can be calculated as  $v_n = \langle \cos[n(\varphi - \Psi_n)] \rangle$ , where  
34 the brackets denote an average over particles and events. Elliptic ( $v_2$ ) and triangular ( $v_3$ ) flows  
35 are the dominant flow signals and have been studied very extensively both at top RHIC and  
36 LHC energies [10, 11, 12]. For low transverse momentum ( $p_T < 2$ -3 GeV/c), the  $p_T$  dependence  
37 of  $v_2$  and  $v_3$  for produced particles is well described by viscous hydrodynamic models and a good

38 agreement between data and model calculations can be reached for the small values of  $\eta/s$  close  
 39 to the lower conjectured bound of  $1/4\pi$  [3]. The shear viscosity suppresses triangular flow signal  
 40  $v_3$  more strongly than elliptic flow signal  $v_2$  [13, 14]. The data for top RHIC energy show that, for  
 41 a given collision centrality, the measured values of  $v_n$  ( $n = 2, 3$ ) for all hadrons scale to a single  
 42 curve when plotted as  $v_n/n_q^{n/2}$  versus scaled transverse kinetic energy,  $(m_T - m_0)/n_q$ , where  $n_q$  is  
 43 the number of constituent quarks in the hadron and  $m_0$  is mass [10, 11]. The observed empirical  
 44 Number-of-Constituent Quark (NCQ) scaling with transverse kinetic energy may indicate that  
 45 the bulk of the anisotropic flow at top RHIC energies is partonic, rather than hadronic [15].  
 46 The collision energy dependence of elliptic flow ( $v_2$ ) for inclusive and identified hadrons at  
 47 mid-rapidity in Au+Au collisions, has been studied very extensively by STAR experiment at  
 48  $\sqrt{s_{NN}} = 7.7 - 62.4$  GeV [16, 17, 18, 19]. The elliptic flow signal  $v_2(p_T)$  for inclusive charged  
 49 hadrons shows a very small change over such a wide range of collision energies [16]. Hybrid  
 50 model calculations show that the weak dependence of  $v_2(p_T)$  on the beam energy may result  
 51 from the interplay of the hydrodynamic and hadronic transport phase [14]. The triangular flow  
 52  $v_3$  is expected to be more sensitive to the viscous damping and might be an ideal observable  
 53 to probe the formation of a QGP at different collision energies [20]. However, a significant  
 54 difference in the  $v_2$  values between particles and the corresponding anti-particles was observed  
 55 [18, 19]. This difference increases with decreasing collision energy and is larger for baryons than  
 56 mesons. Several different theoretical models have been proposed for the possible physical reason  
 57 for this effect and new measurements of  $v_3$  for particles and anti-particles might be important for  
 58 distinguishing between them [18, 19]. In this work, we report new measurements of triangular  
 59 ( $v_3$ ) anisotropic flow coefficients for inclusive and identified charged hadrons ( $\pi^\pm, K^\pm, p, \bar{p}$ ) at  
 60 midrapidity in Au+Au collisions at  $\sqrt{s_{NN}} = 11.5 - 62.4$  GeV and compare them to  $v_2$  results.

## 61 2. Data Analysis

62 The data reported in this analysis are from Au+Au collisions at  $\sqrt{s_{NN}} = 11.5, 14.5, 19.6, 27, 39$   
 63 and  $62.4$  GeV, collected during the beam energy scan phase-I and II (BES-I & BES-II) programs  
 64 by the STAR detector using a minimum bias trigger. The collision vertices were reconstructed  
 65 using charged-particle tracks measured in the Time Projection Chamber (TPC). The TPC covers  
 the full azimuth and has a pseudorapidity range of  $|\eta| < 1.0$ .



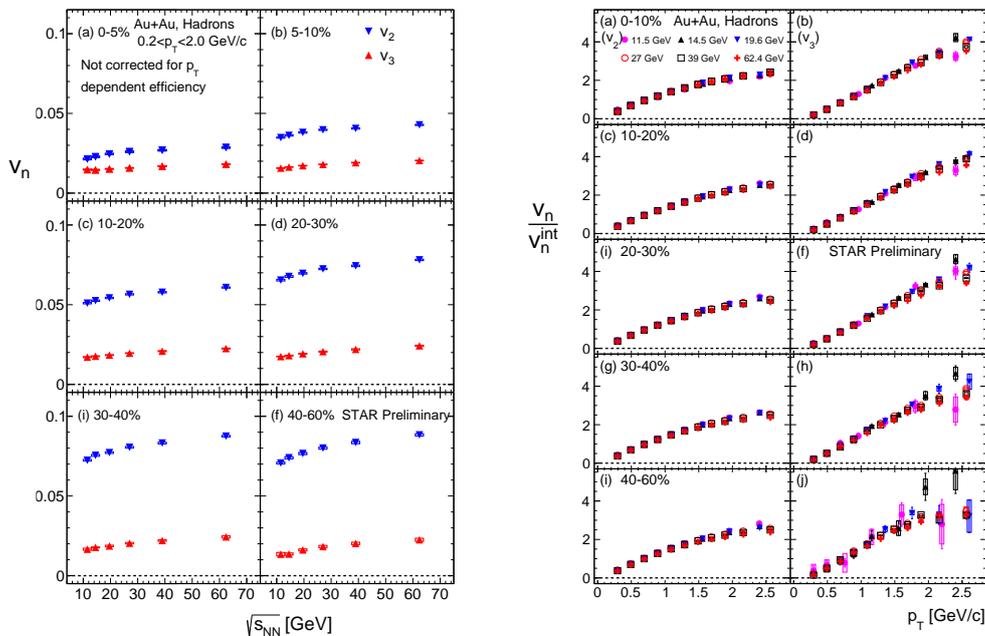
**Figure 1.** The centrality dependence of the event plane resolution for  $v_2$  (left panel) and  $v_3$  (right panel) for all six collision energies.

66 Events were selected to have a vertex position about the nominal center of the TPC in  
 67 the beam direction of  $\pm 40$  cm at  $\sqrt{s_{NN}} = 62, 39, 27, 19.6$  and  $14.5$  GeV,  $\pm 50$  cm at  
 68  $\sqrt{s_{NN}} = 11.5$  GeV, and to be within a radius of  $1 - 2$  cm with respect to the beam axis.  
 69 The centrality of each collisions was determined by measuring event-by-event multiplicity and  
 70 interpreting the measurement with a tuned Monte Carlo Glauber calculation [16, 18]. Analyzed  
 71

72 tracks were required to have a distance of closest approach to the primary vertex to be less  
 73 than 3 cm, and to have at least 15 TPC space points used in their reconstruction [16, 17]. The  
 74 particle identification for charged hadrons ( $\pi^\pm$ ,  $K^\pm$ ,  $p$ ,  $\bar{p}$ ) was based on a combination of the  
 75 ionization energy loss,  $dE/dx$ , in the TPC, and the squared mass,  $m^2$ , from the TOF detector  
 76 [17, 18, 19].

77 In this study, the event plane method with  $\eta$  sub-events, separated by an additional  $\eta$ -gap  
 78 of  $\Delta\eta > 0.1$ , was used to measure elliptic ( $v_2$ ) and triangular ( $v_3$ ) flow [16]. The  $\eta$  gap is  
 79 introduced to suppress the non-flow correlations between the two sub-events. The  $\eta$  sub-event  
 80 method was implemented using the procedure in Ref [16, 17, 18]. The centrality dependence of  
 81 event plane resolution for  $v_2$  and  $v_3$  for the six collision energies is shown in the Fig. 1. The  
 82 systematic uncertainty associated with the non-flow effects is estimated for each collision energy  
 83 by comparing  $v_2$  and  $v_3$  results obtained with different  $\Delta\eta$  gaps. Studies were performed for  $\Delta\eta$   
 84 values of 0.1, 0.3, 0.5, 0.7.

### 85 3. Results

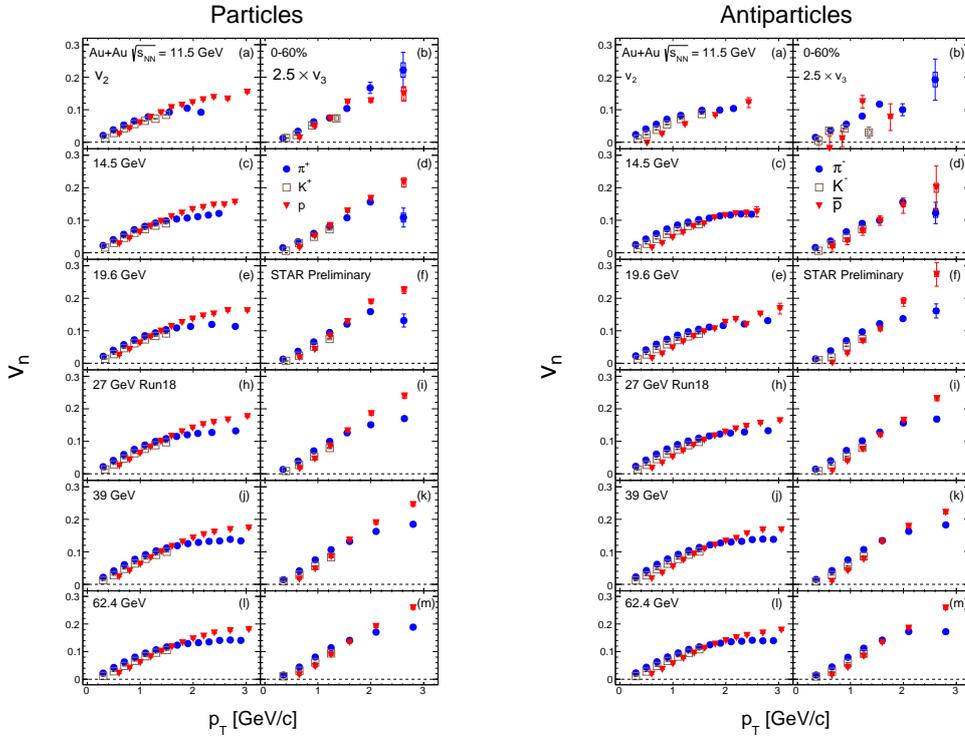


**Figure 2.** Left:  $p_T$ -integrated  $v_2^{\text{int}}$  and  $v_3^{\text{int}}$  of inclusive charged hadrons as a function of  $\sqrt{s_{NN}}$  for different bins in collision centrality. Right:  $p_T$ -dependence of  $v_2(p_T)/v_2^{\text{int}}$  and  $v_3(p_T)/v_3^{\text{int}}$  of charged hadrons for different bins in collision centrality. The measured  $v_n(p_T)$  values were divided by the corresponding  $v_n^{\text{int}}$  values from the left part of the figure. The results are presented for all 6 collision energies:  $\sqrt{s_{NN}} = 11.5, 14.5, 19.6, 27, 39$  and  $62.4$  GeV.

86 Preliminary results for the excitation functions of  $p_T$ -integrated ( $0.2 < p_T < 3.2$  GeV/c)  
 87 values of  $v_2^{\text{int}}$  and  $v_3^{\text{int}}$  of inclusive charged hadrons are presented in the left part of Fig. 2. The  
 88  $v_n^{\text{int}}$  results were not corrected for  $p_T$  dependent tracking efficiency, which will be explored in  
 89 future analysis. Although the efficiency is  $p_T$  dependent but is similar between different collision  
 90 energy, so it is not expected to influence the  $\sqrt{s_{NN}}$  trend. The results are presented for 6 bins  
 91 in collision centrality: 0 – 5%, 5 – 10%, 10 – 20%, 20 – 30%, 30 – 40% and 40 – 60%. The results  
 92 indicate an essentially monotonic increase for  $p_T$ -integrated  $v_2$  and  $v_3$  with  $\sqrt{s_{NN}}$ , as expected  
 93 from increase of the radial flow with collision energy which pushes the hadrons to larger  $p_T$

94 and renders the momentum spectra less anisotropic at low  $p_T$  [20]. The  $p_T$  dependence of  
 95  $v_2(p_T)/v_2^{\text{int}}$  and  $v_3(p_T)/v_3^{\text{int}}$  for inclusive charged hadrons is presented in the right part of the  
 96 Fig. 2 for different bins in collision centrality. The measured  $v_n(p_T)$  values were divided by the  
 97 corresponding  $p_T$ -integrated  $v_n^{\text{int}}$  values from the left part of the Fig. 2. The results in the figure  
 98 are presented for all 6 collision energies:  $\sqrt{s_{NN}} = 11.5, 14.5, 19.6, 27, 39$  and  $62.4$  GeV and  
 99 they show that  $v_n(p_T)/v_n^{\text{int}}$  has a very weak dependence on  $\sqrt{s_{NN}}$ . This is in agreement with  
 100 predictions from [21].

101 Figure 3 shows the collision energy dependence in  $v_2(p_T)$  and  $v_3(p_T)$  for identified hadrons ( $\pi^\pm$ ,  
 102  $K^\pm$ ,  $p$ ,  $\bar{p}$ ) for 0-60% central Au+Au collisions. The results for particles (left panel) and anti-  
 103 particles (right panel) are presented separately. We observe that the  $v_3(p_T)$  signal of identified  
 104 charged hadrons exhibits similar trends as first observed for  $v_2$  in Au+Au collisions: mass  
 105 ordering at low transverse momenta,  $p_T < 2$  GeV/c, and meson/baryon splitting at intermediate  
 106  $p_T$ ,  $2 < p_T < 4$  GeV/c [17, 18, 19].

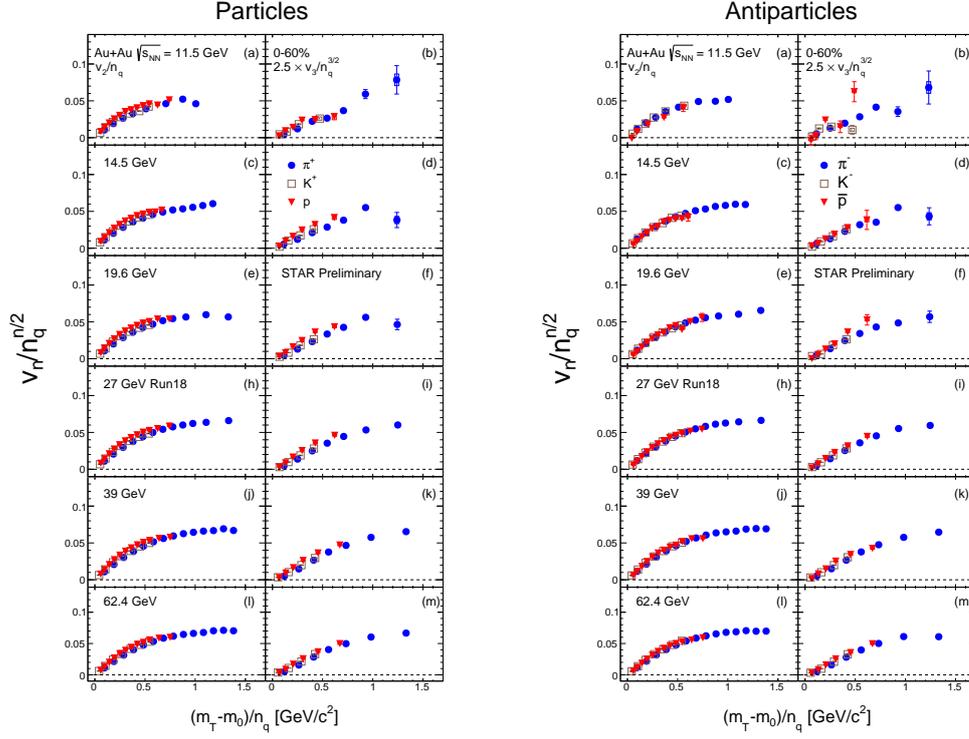


**Figure 3.**  $p_T$  dependence of  $v_2$  and  $v_3$  signals of  $\pi^+$ ,  $K^+$ ,  $p$  (left) and  $\pi^-$ ,  $K^-$ ,  $\bar{p}$  (right) for 0-60% central Au+Au collisions.

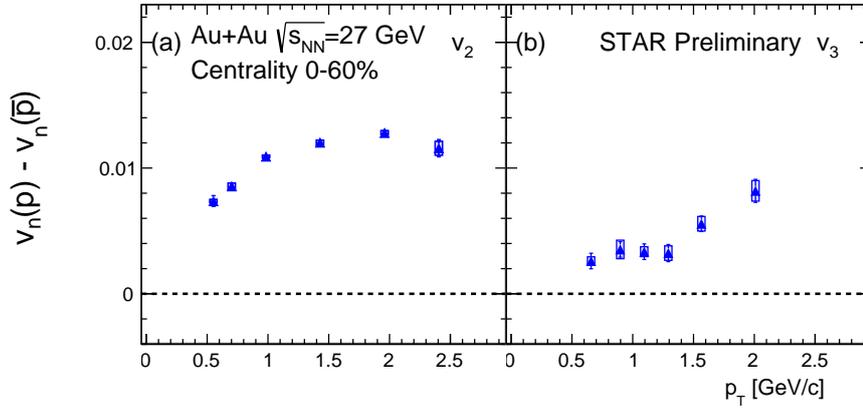
107 Figure. 4 shows that the measured  $v_3$  values of identified charged hadrons seems to follow  
 108 the NCQ scaling,  $v_n/n_q^{n/2}$  versus  $(m_T - m_0)/n_q$ , if we plot the results for particles (left part)  
 109 and anti-particles (right part) separately.

110 The analysis of the new dataset of Au+Au collisions at  $\sqrt{s_{NN}} = 27$  GeV, collected by STAR  
 111 experiment in 2018, allows us to observe the difference in the triangular  $v_3$  flow between protons  
 112 and anti-protons, see Fig. 5. It shows that, similar to elliptic flow  $v_2$ , the  $v_3$  flow signal of protons  
 113 is larger than  $v_3$  of antiprotons and the difference has  $p_T$  dependence.

114 For other collision energies we can estimate the difference in  $v_3$  values between particles and  
 115 corresponding anti-particles for  $p_T$ -integrated  $v_3$  values. The right part of Fig. 6 shows the  
 116 difference in  $v_3$  between particles (X) and their corresponding anti-particles ( $\bar{X}$ ) as a function



**Figure 4.** The Number-of-Constituent Quark (NCQ) scaled elliptic and triangular flow,  $v_n/n_q^{n/2}$  versus  $(m_T - m_0)/n_q$ , for 0-60% central Au+Au collisions for selected particles (left part) and corresponding anti-particles (right part).

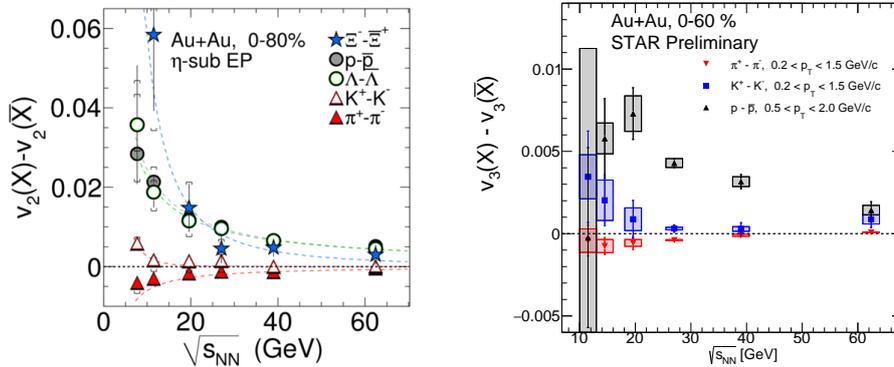


**Figure 5.** The difference between proton and antiproton  $v_n$  as a function of the transverse momentum  $p_T$  for 0-60% central Au+Au collisions at  $\sqrt{s_{NN}} = 27$  GeV.

117 of  $\sqrt{s_{NN}}$  for 0 – 60% central Au+Au collisions. Similar to  $v_2$ , the  $v_3(X) - v_3(\bar{X})$  difference  
 118 increases with decreasing collision energy and it is larger for baryons than mesons [17, 18, 19].

#### 119 4. Summary

120 In summary, we have employed the event plane method with  $\eta$  sub-events to carry out new  
 121 measurements of the triangular ( $v_3$ ) anisotropic flow coefficients for inclusive and identified  
 122 charged hadrons ( $\pi^\pm$ ,  $K^\pm$ ,  $p$ ,  $\bar{p}$ ) at midrapidity in Au+Au collisions, spanning the collision



**Figure 6.** (left) The difference in  $v_2$  between particles ( $X$ ) and their corresponding anti-particles ( $\bar{X}$ ) (see legend) as a function of  $\sqrt{s_{NN}}$  for 0–80% central Au+Au collisions. The figure is taken from [18]. (right) The preliminary results for the difference in the  $v_3$  values between particles ( $X$ ) and its corresponding anti-particles ( $\bar{X}$ ) as a function of  $\sqrt{s_{NN}}$  for 0-60% central Au+Au collisions.

123 energy range  $\sqrt{s_{NN}} = 11.5 - 62.4$  GeV. We observe that the triangular flow signal ( $v_3$ ) of  
 124 identified hadrons exhibits similar trends as first observed for  $v_2$  [16, 17, 18]. New measurements  
 125 of  $v_3$  excitation function could serve as constraints to test different models and to aid new  
 126 information about the temperature dependence of the transport properties of the strongly  
 127 interacting matter.

## 128 5. Acknowledgments

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