



# Charge Asymmetry Dependence of Pion Azimuthal Anisotropy in Au + Au Collisions at STAR

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# CME & CSE

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- **Chiral Magnetic Effect (CME)**: nonzero axial charge density induces a vector (electric) current along external magnetic field.

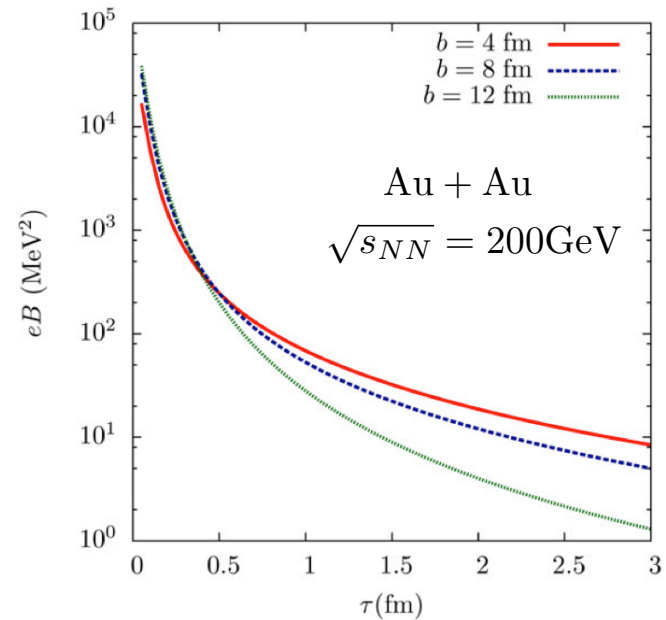
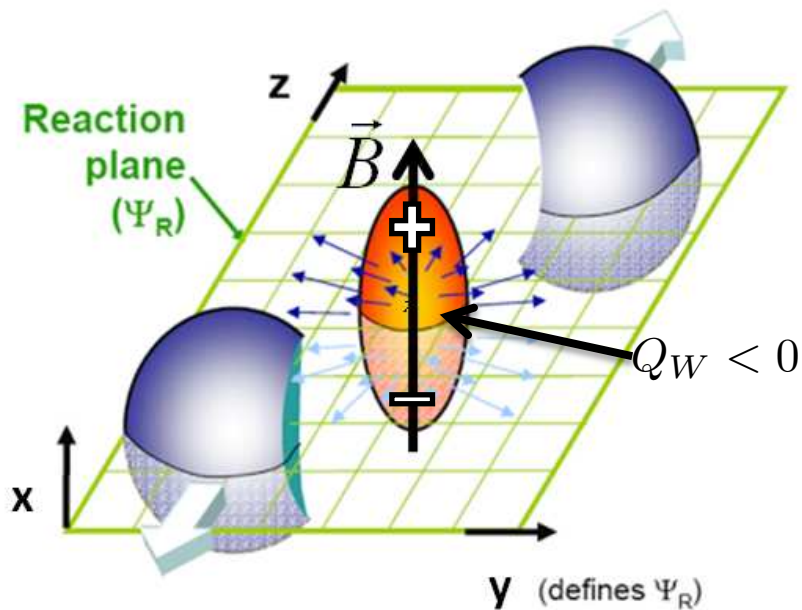
$$\mathbf{j}_V = \frac{N_c e}{2\pi^2} \mu_A \mathbf{B} \quad \rightarrow \quad \text{electric charge separation along } B \text{ field}$$

- **Chiral Separation Effect (CSE)**: nonzero vector charge density induces an axial current along external magnetic field.

$$\mathbf{j}_A = \frac{N_c e}{2\pi^2} \mu_V \mathbf{B} \quad \rightarrow \quad \text{chiral charge separation along } B \text{ field}$$

D. E. Kharzeev, L. D. McLerran, and H. J. Warringa, Nuclear Physics A **803**, 227 (2008)  
T. Son, D. and A. R. Zhitnitsky, Phys. Rev. D **70**, 074018 (2004)

# STAR CME in Heavy Ion Collisions

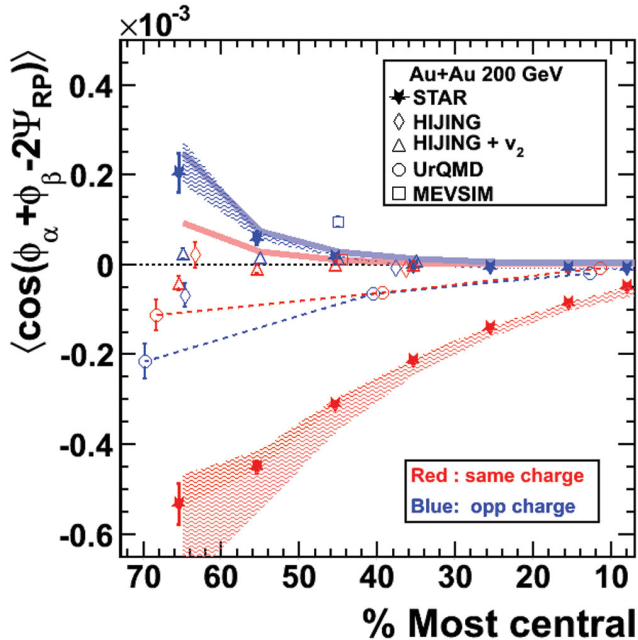
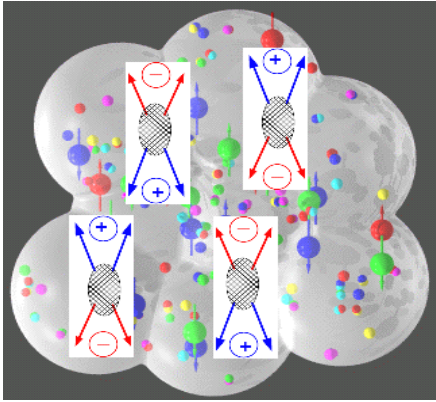


Nuclear Physics A **803**, 227 (2008).

- QCD vacuum transition in heavy ion collision  $\rightarrow Q_W \neq 0$
- Extremely strong magnetic field created in heavy ion collision
- CME causes out-of-plane electric charge separation
- Charge asymmetry w.r.t. reaction plane as a signature of LPV.



# CME in Heavy Ion Collisions



- Charge particle azimuthal angle distribution

$$\frac{dN_\alpha}{d\phi} \propto 1 + 2v_{1,\alpha} \cos(\Delta\phi) + 2v_{2,\alpha} \cos(2\Delta\phi) + \dots$$

$$+ 2a_{1,\alpha} \sin(\Delta\phi) + 2a_{2,\alpha} \sin(2\Delta\phi) + \dots$$

- Correlations  $\langle a_\alpha a_\beta \rangle$  can be measured

$$\alpha, \beta = +, -$$

- The three-particle correlations are directly sensitive to predicted local  $P$ -violation in heavy-ion collisions.

$$\langle \cos(\phi_\alpha + \phi_\beta - 2\Psi_{RP}) \rangle$$

- Out-of-plane charge separation  
 same charge correlation  $< 0$   
 opposite charge correlation  $> 0$

- The observed signal cannot be described by models (HIJING, HIJING +  $\nu_2$ , URQMD, MEVSIM)

S. A. Voloshin, Phys. Rev. C **70**, 057901 (2004)  
 STAR Phys. Rev. Lett. **103**, 251601 (2009), Phys. Rev. C **81**, 054908 (2010)



# Chiral Magnetic Wave

## Chiral Magnetic Effect

$$j_V = \frac{N_c e}{2\pi^2} \mu_A B \longrightarrow$$

**important at high beam energy:**  
enough energy for “jumping”  
over barriers to create initial  
axial charge density

## Chiral Separation Effect

$$j_A = \frac{N_c e}{2\pi^2} \mu_V B \longrightarrow$$

**important at low beam energy:**  
there is sizable initial  
vector charge density to  
set off chiral magnetic waves

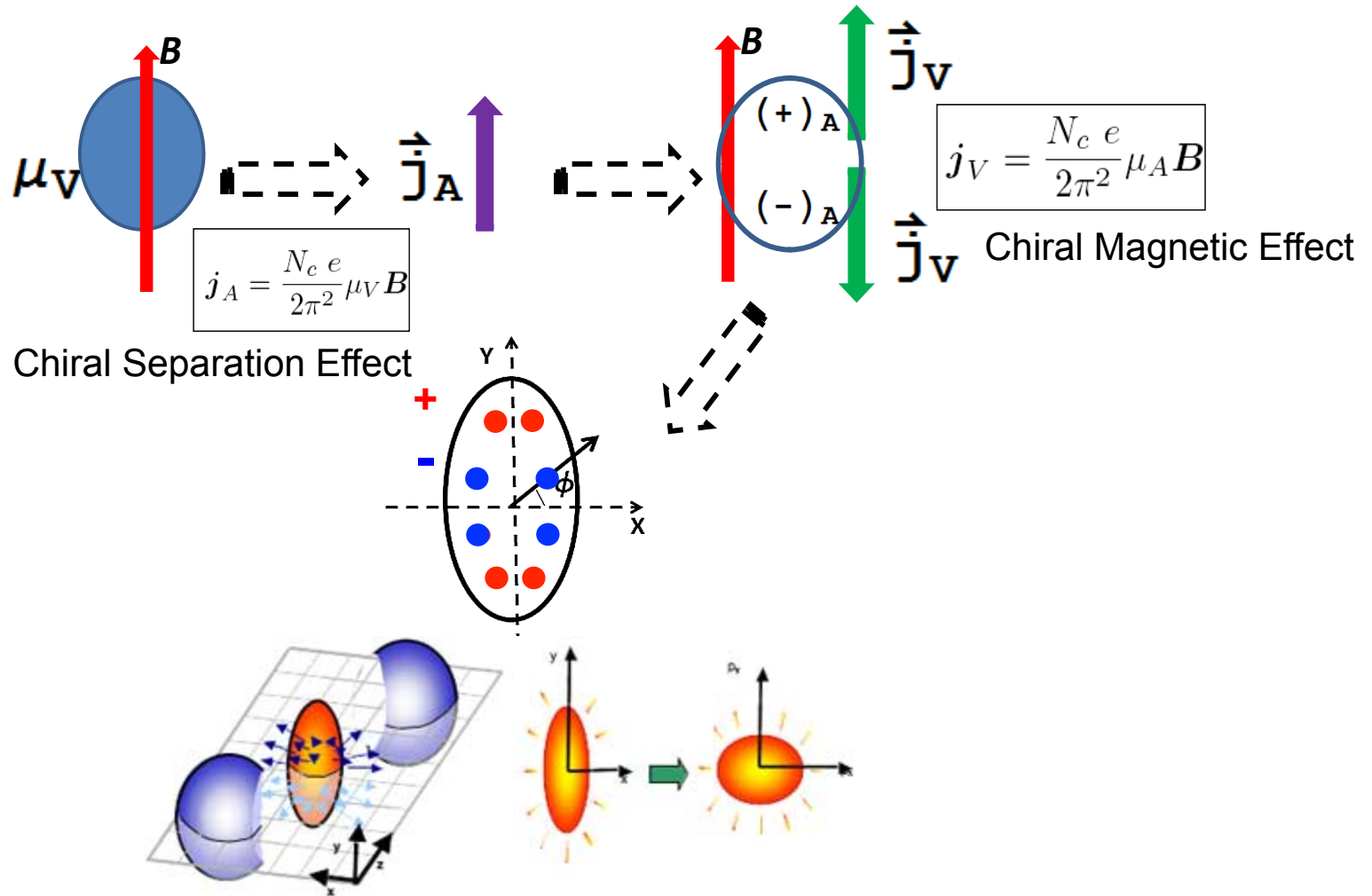
## Chiral Magnetic Wave: coupling two effects

$$\left( \partial_0 \mp \frac{N_c e B \alpha}{2\pi^2} \partial_1 - D_L \partial_1^2 \right) j_{L,R}^0 = 0$$



D. Kharzeev and H.-U. Yee, Physical Review D **83**, 085007 (2011)

# Chiral Magnetic Wave



Y. Burnier, D. E. Kharzeev, J. Liao and H-U Yee, Phys. Rev. Lett. **107**, 052303 (2011)



# Observable for Chiral Magnetic Wave

$$\frac{dN_{\pm}}{d\phi} = N_{\pm} [1 + 2v_2 \cos(2\phi)]$$

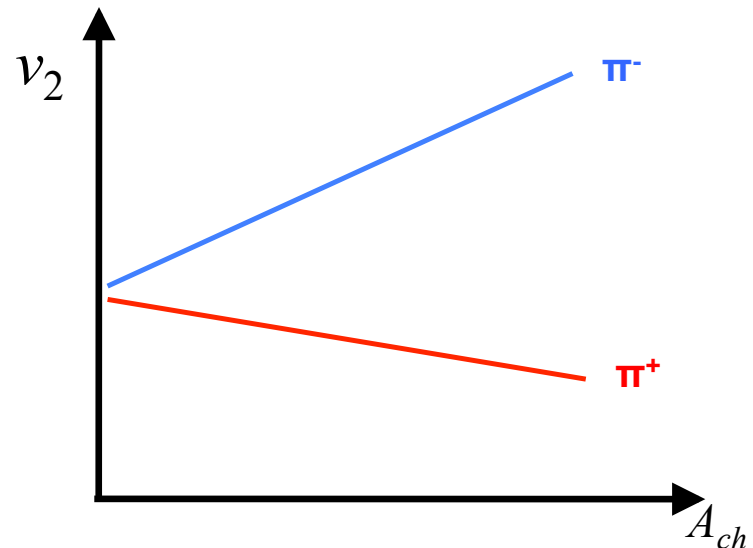
$$\approx \bar{N}_{\pm} [1 + 2v_2 \cos(2\phi) \mp A_{ch} r \cos(2\phi)]$$

$$A_{ch} = \frac{N_+ - N_-}{N_+ + N_-}$$

$$r = 2 \left( \frac{q_e}{\bar{\rho}_e} \right)$$

$$v_2^{\pm} = v_2 \mp \left( \frac{q_e}{\bar{\rho}_e} \right) A_{ch}$$

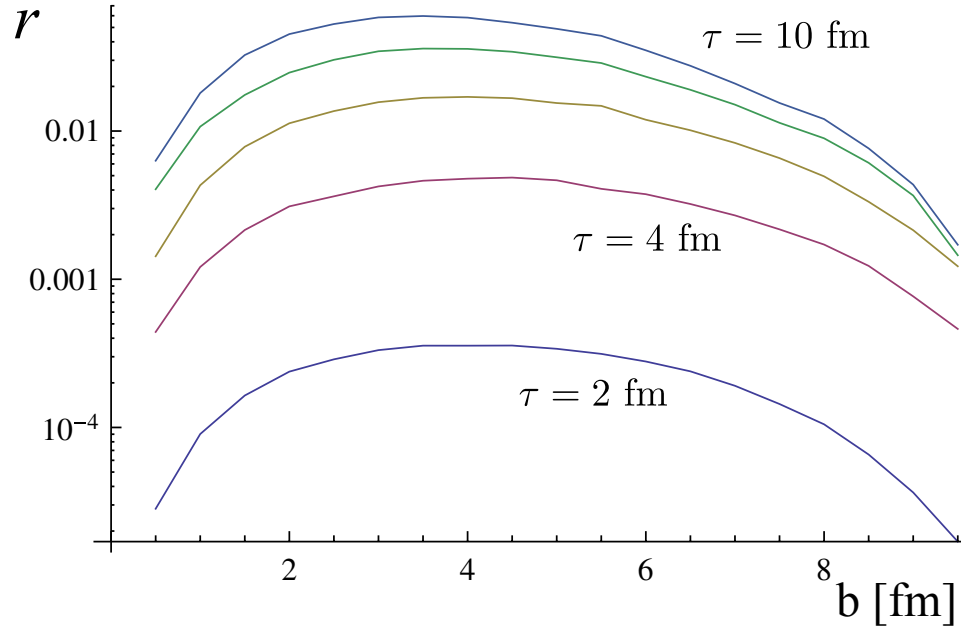
- $v_2(\pi^-) > v_2(\pi^+)$
- $v_2(\pi^-)$  and  $v_2(\pi^+)$  have opposite trend as a function of  $A_{ch}$
- difference between  $v_2(\pi^-)$  and  $v_2(\pi^+)$  has a linear relationship with  $A_{ch}$





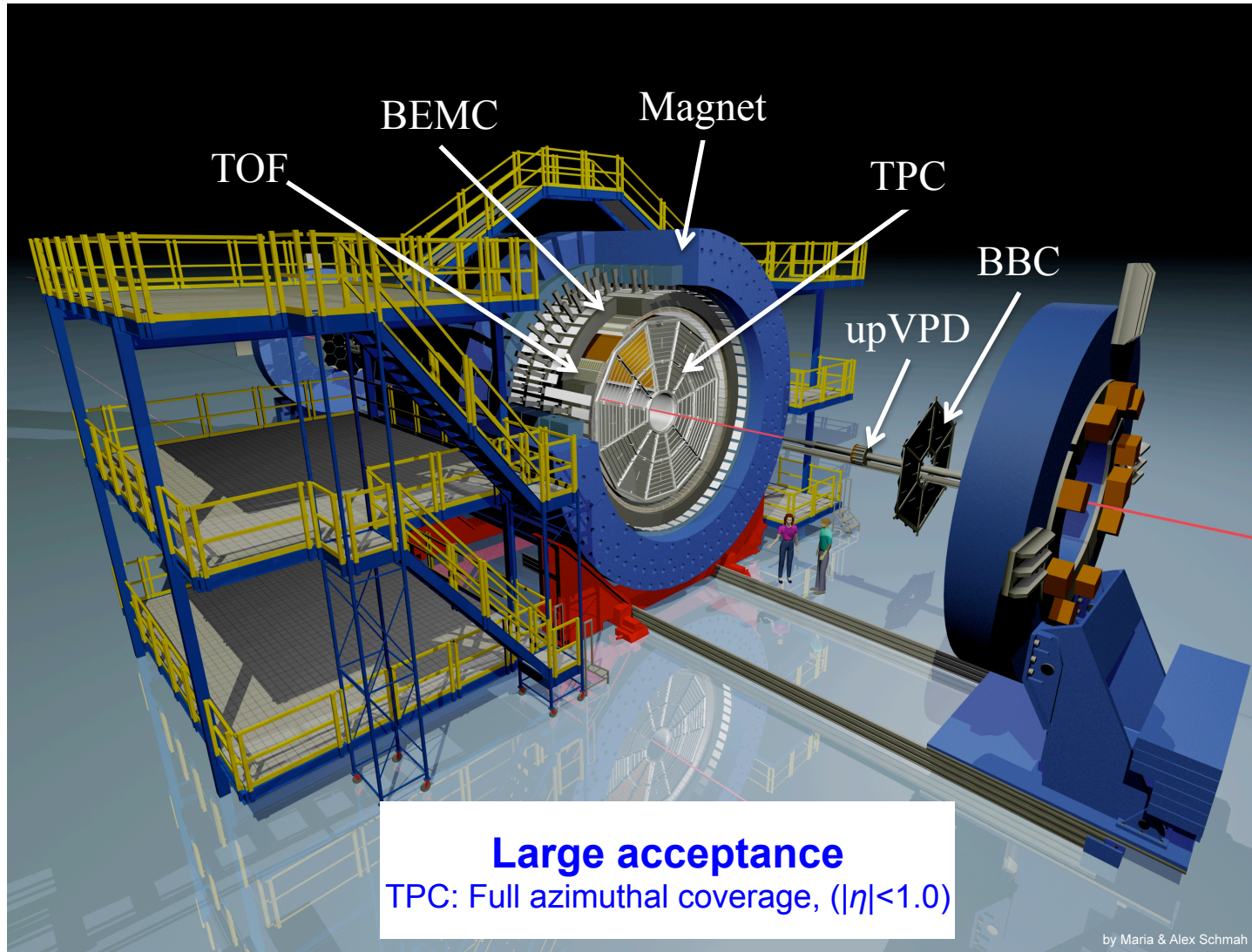
# Observable for Chiral Magnetic Wave

$$v_2^- - v_2^+ = 2 \left( \frac{q_e}{\bar{\rho}_e} \right) A_{ch} \quad r = 2q_e/\bar{\rho}_e$$

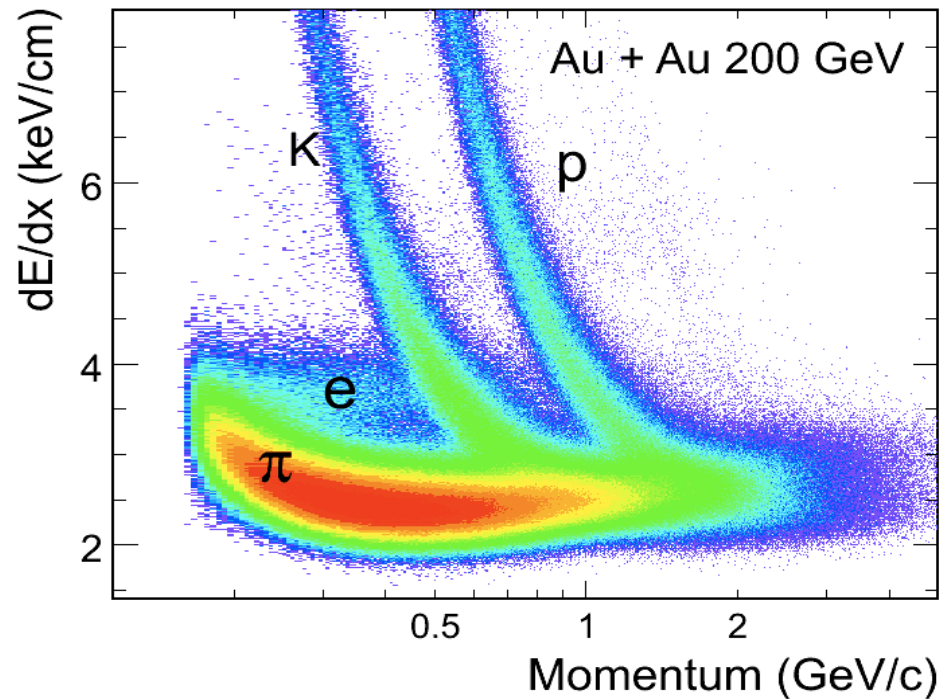


Y. Burnier, D. E. Kharzeev, J. Liao and H-U Yee, Phys. Rev. Lett. **107**, 052303 (2011) and arXiv: 1208.2537





## Ionization energy loss in the Time Projection Chamber



$$n\sigma_{\pi} = \frac{1}{R} \log \frac{dE/dx_{measured}}{\langle dE/dx \rangle_{\pi}}$$

**dE/dx Particle Identification:**

( $\pi$ , K):  $p \sim 0.6$  GeV/c; ( $\pi$ /K, p):  $p \sim 1.0$  GeV/c



# Analysis Details

## ➤ Data Set

- Au + Au
- Minimum Bias trigger
- 0 – 80% centrality

$\sqrt{s_{NN}}$ (GeV)	Events (M)
200	238
62.4	63
39	104
27	46
19.6	23
11.5	10
7.7	4

## ➤ Pion Selection

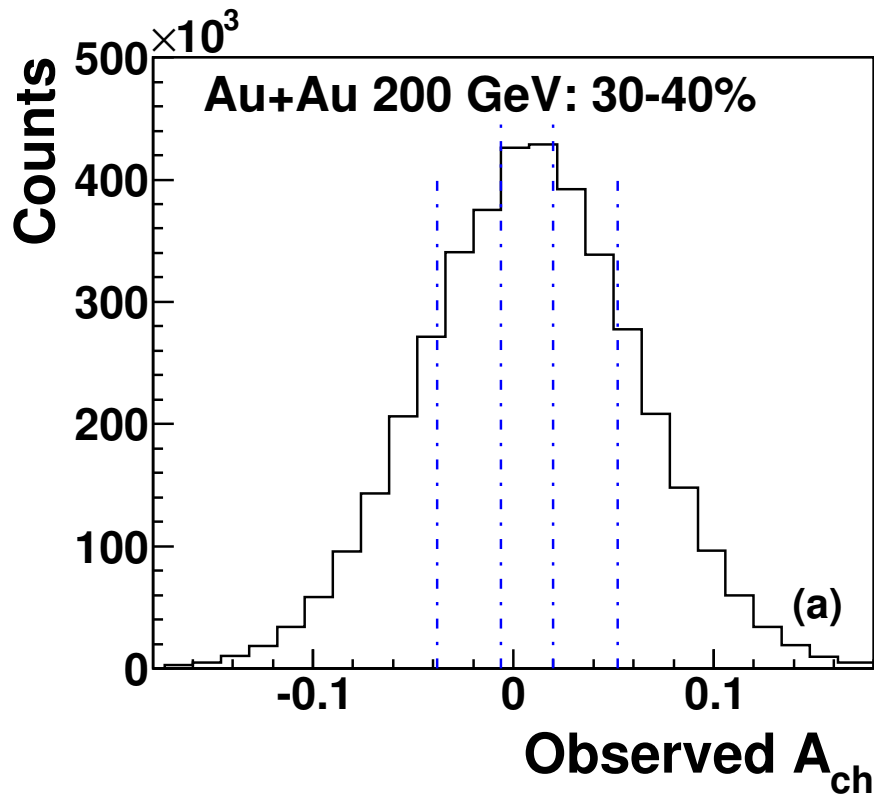
- PID:  $|n\sigma_{\pi}| < 2$
- $0.15 < p_T < 0.5$  GeV/c
- $|\eta| < 1.0$

## ➤ Particles for charge asymmetry

- charged particle
- $0.15 < p_T < 12$  GeV/c
- $|\eta| < 1.0$
- exclude (anti-)protons with  $p_T < 0.4$  GeV/c



# Observed Charge Asymmetry



- $N_+$  and  $N_-$  : number of positive and negative particles

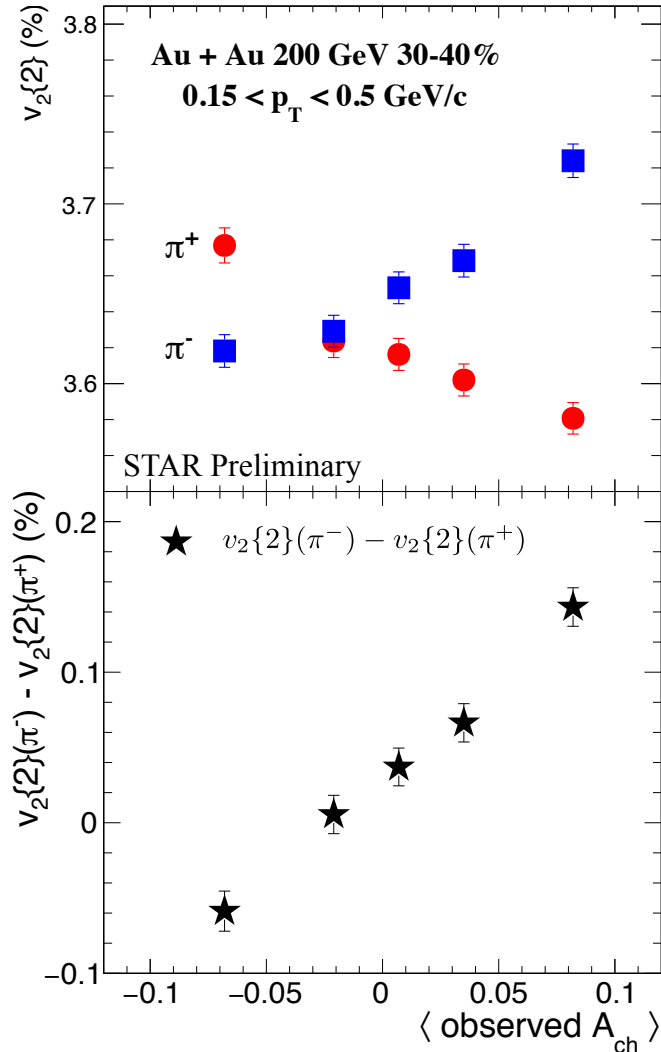
- observed  $A_{ch}$

$$A_{ch} = \frac{N_+ - N_-}{N_+ + N_-}$$

- Each bin has roughly the same number of events



# Integrated Elliptic Flow

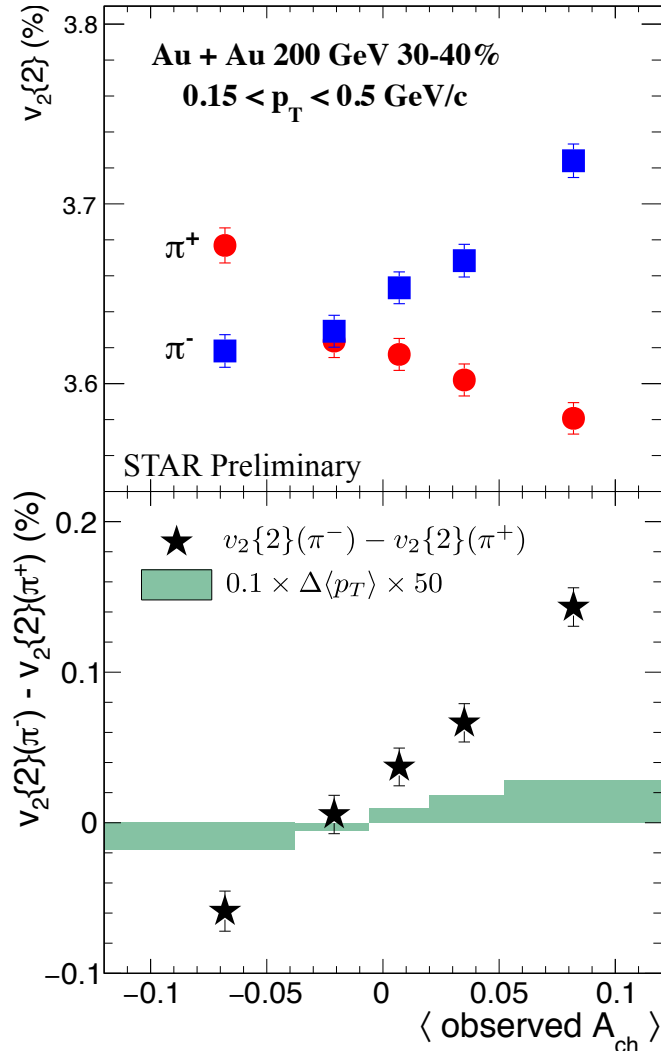


- $v_2(\pi^+)/v_2(\pi^-)$  integrated over  $0.15 < p_T < 0.5$  GeV/c
- $v_2(\pi^-) > v_2(\pi^+)$
- $v_2(\pi^-)$  and  $v_2(\pi^+)$  have opposite trend as a function of  $A_{ch}$
- difference between  $v_2(\pi^-)$  and  $v_2(\pi^+)$  has a linear relationship with  $A_{ch}$

All of the above are consistent with calculation based on Chiral Magnetic Wave



# Integrated Elliptic Flow

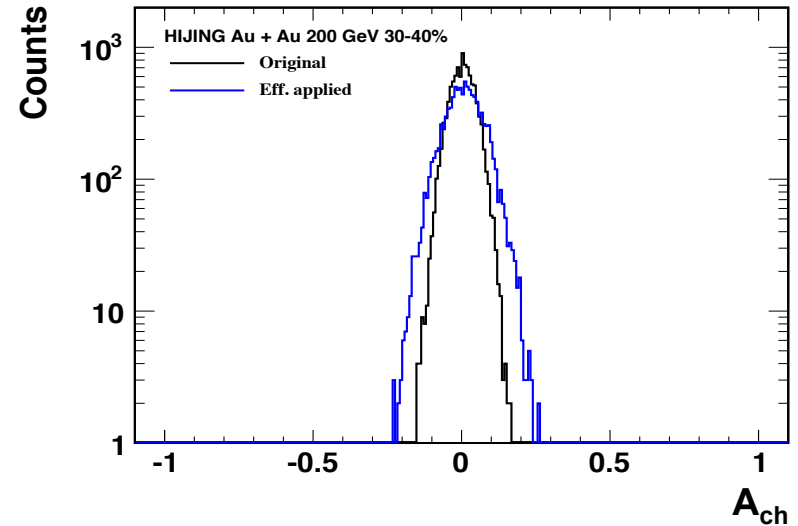
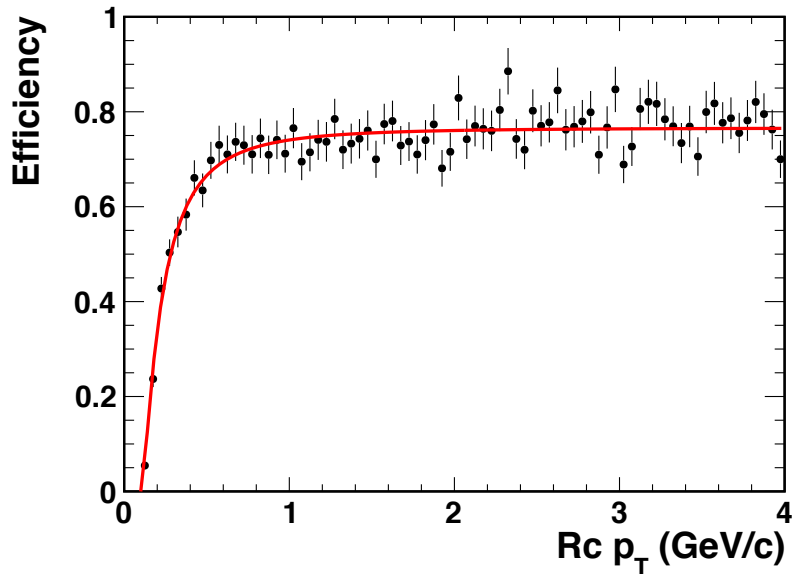


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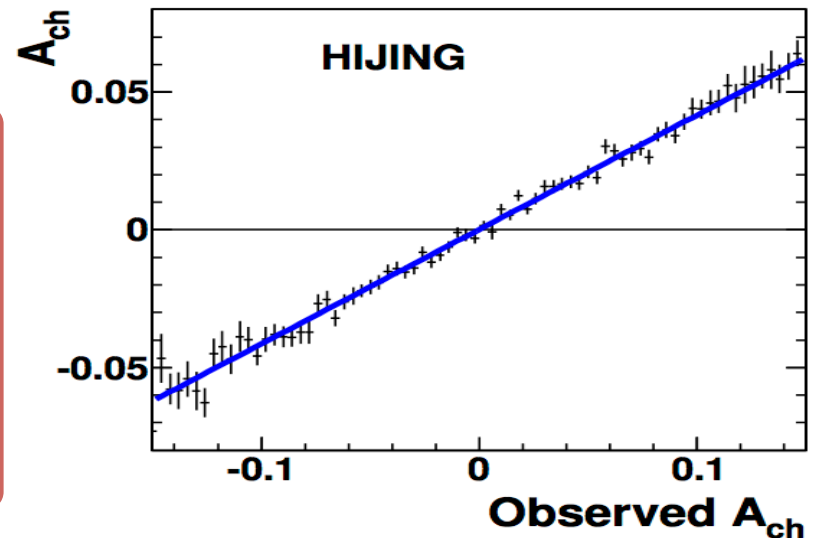
All of the above are consistent with calculation based on Chiral Magnetic Wave

- $v_2 \approx 0.1 \times p_T$  with  $p_T < 1$  GeV/c
- Mean  $p_T$  difference between  $\pi^-$  and  $\pi^+$  gives about 1% of the observed  $v_2$  difference

Tracking efficiency for  $\pi^+$  and  $\pi^-$

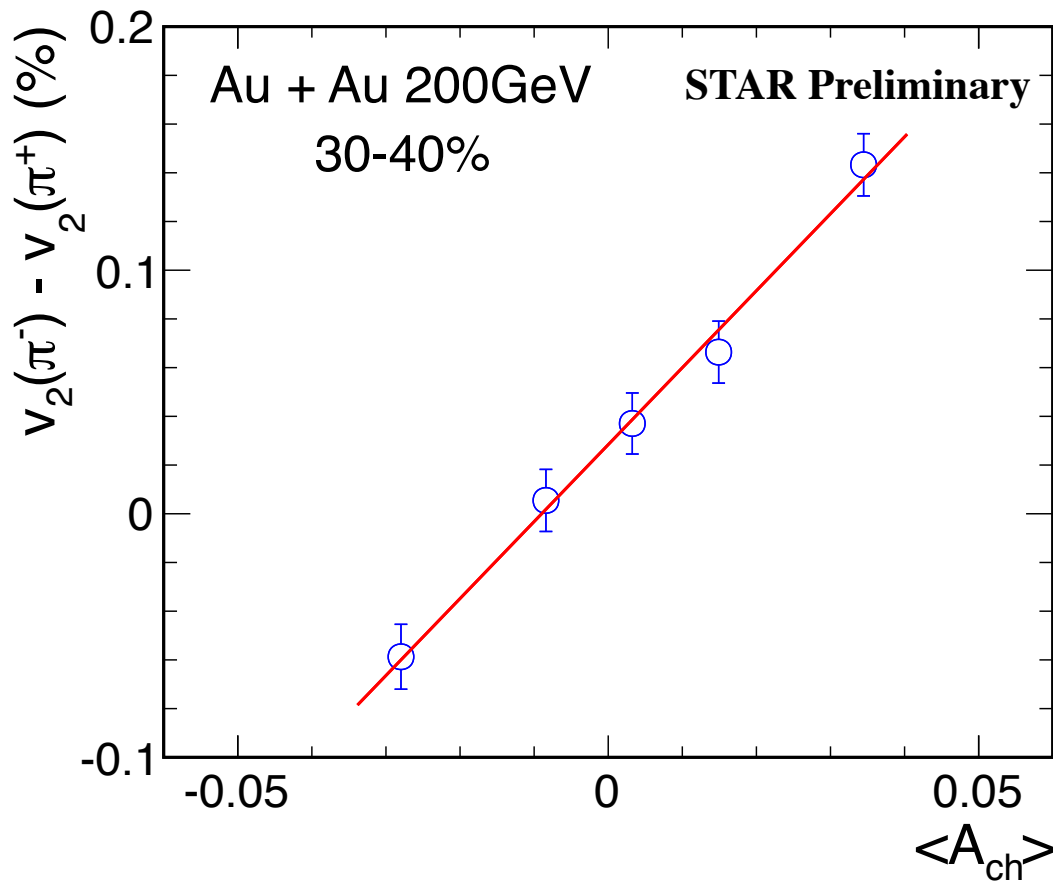


- use pion tracking efficiency as charged particle tracking efficiency
- apply same cuts to calculate  $A_{ch}$  for Monte-Carlo and real data





# Fit the Slope Parameter



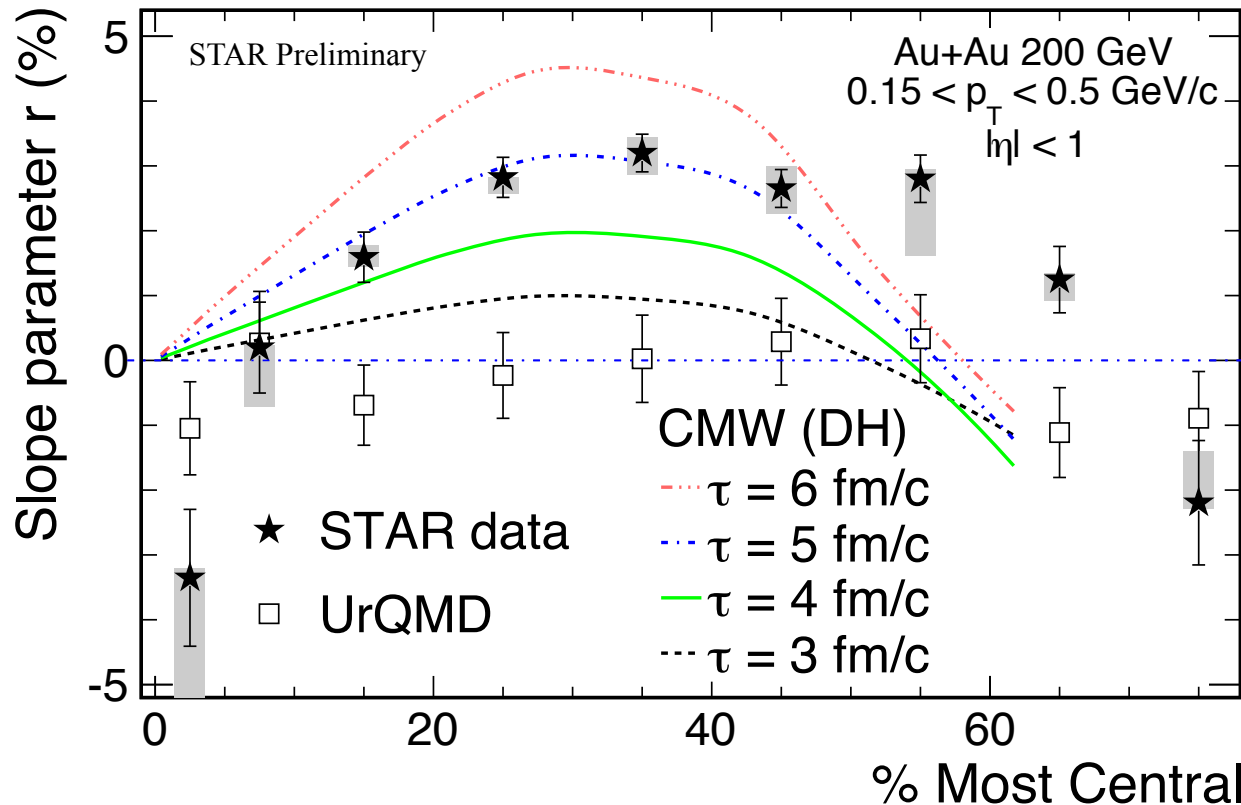
- Prediction based on CMW

$$v_2^- - v_2^+ = 2 \left( \frac{q_e}{\bar{\rho}_e} \right) A_{ch} + C_0$$

- Fit  $v_2(\pi^-) - v_2(\pi^+)$  vs.  $\langle A_{ch} \rangle$  to a straight line
- The slope parameter will be extracted for all centralities



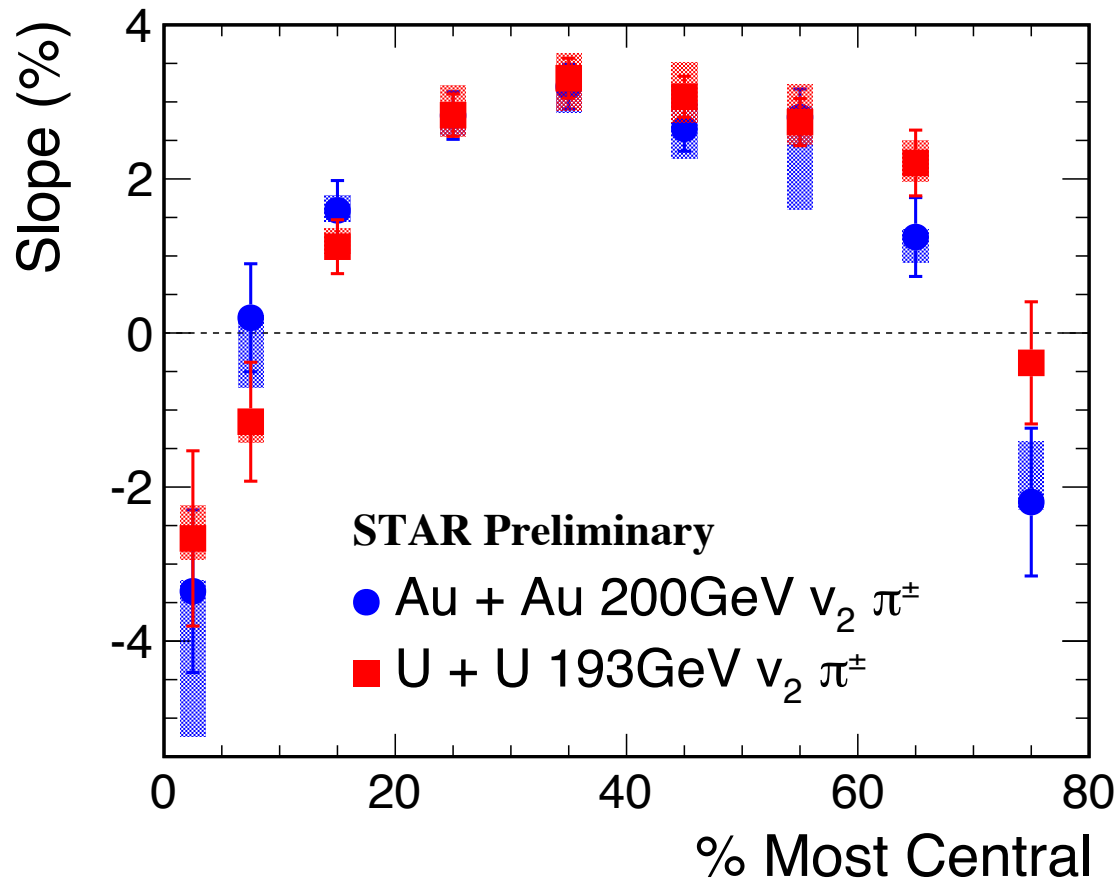
# Slope vs. Centrality



- Similar trends between data and theoretical calculations with CMW, which use different magnetic field calculations, i.e. DH and BS(not shown here), in arXiv: 1208.2537.
- There is no specific beam energy input for the theoretical calculation.
- UrQMD cannot reproduce the slopes at  $\sqrt{s_{NN}} = 200$  GeV.

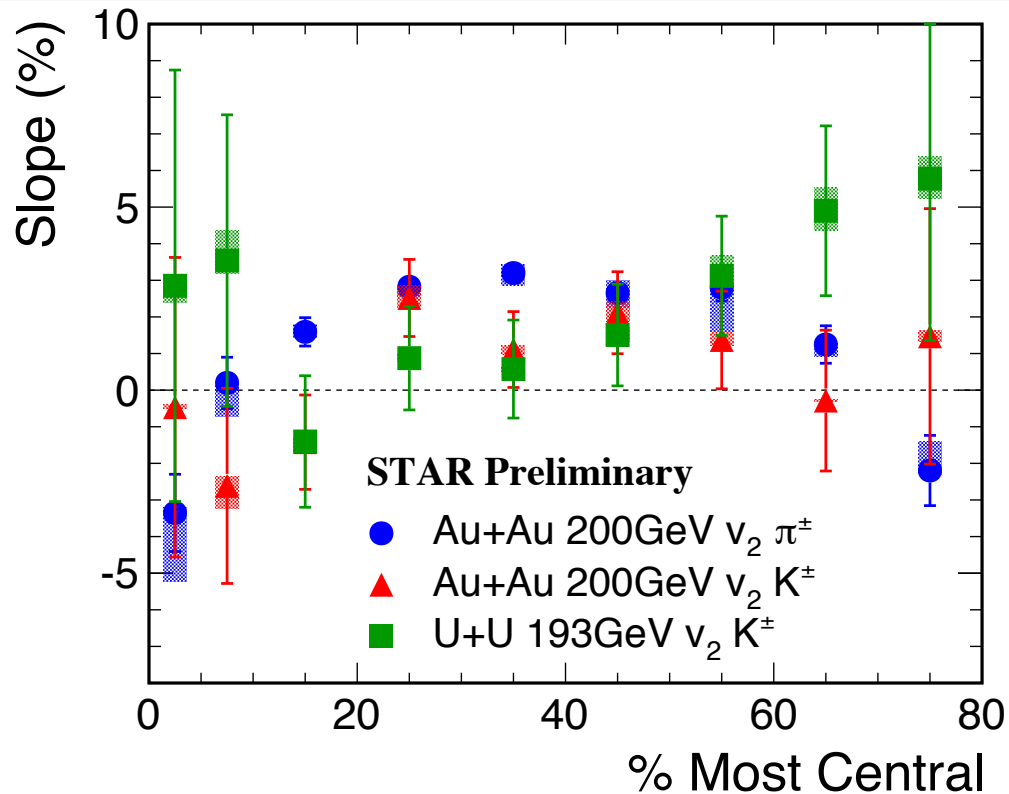


# Slope vs. Centrality in U+U Collisions



- A similar centrality dependence of the slope parameter has been seen in U + U collisions at  $\sqrt{s_{NN}} = 193$  GeV

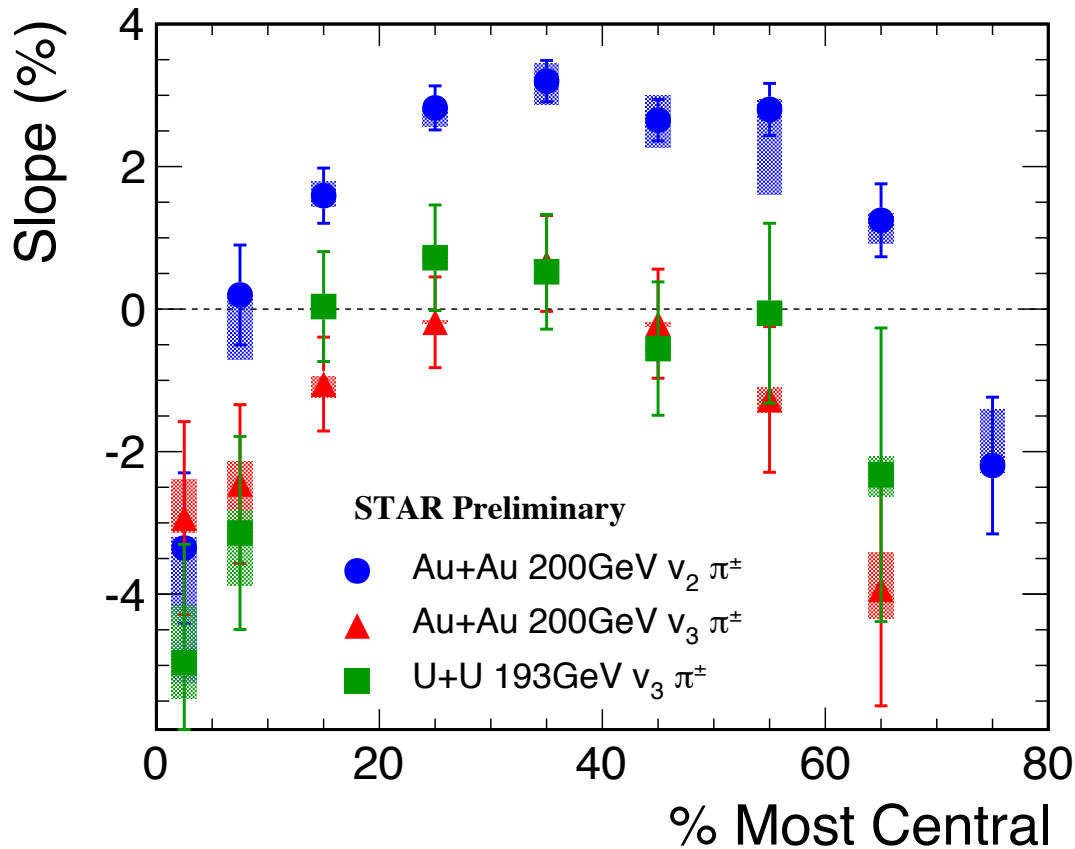
# Slope vs. Centrality for $v_2$ of $K^\pm$



- Theoretical calculations suggest the slope of  $v_2(A_{ch})$  for  $K$  can be different from that of  $\pi$
- The slopes for  $v_2(A_{ch})$  for  $K$  have the same magnitude as that from  $\pi$  in the middle centralities with large error bars. More experimental data are needed to make the measurements of  $K$  conclusive.

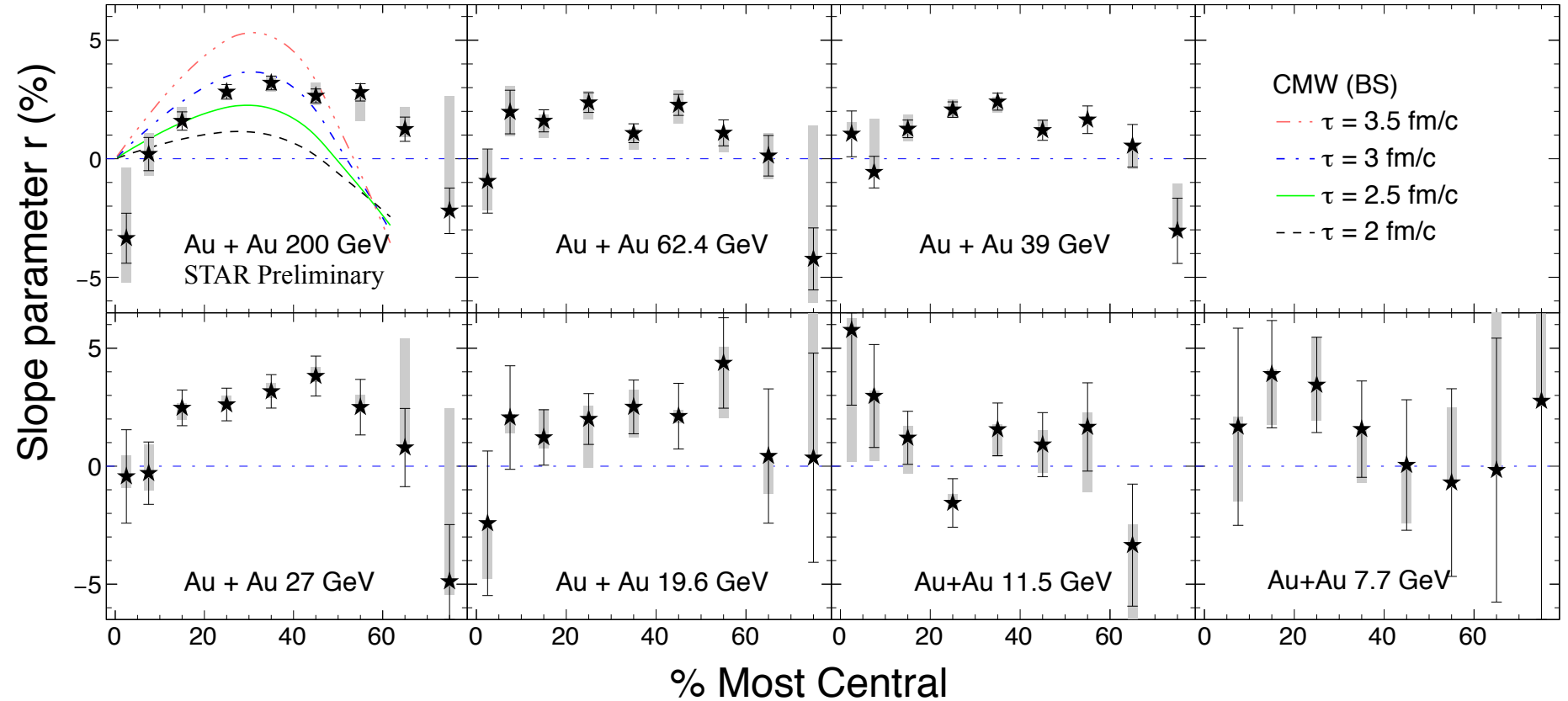


# Slope vs. Centrality from $v_3$ of $\pi^\pm$



- The trend of  $\Delta v_3(\pi)$  slope is similar to that of  $\Delta v_2(\pi)$ , which needs more study to understand
- The  $\Delta v_3(\pi)$  slope is consistent with zero in mid-central collisions in contrast to finite values

A. Bzdak and P. Bozek, Physics Letters B **726** (2013) 239



- The slope parameter  $r$  shows a rise and fall feature from central to peripheral collisions
- Slope as a function of centrality shows weak energy dependence down to 27 GeV
- Similar trend to the theoretical calculations with CMW

- The difference between  $v_2(\pi^-)$  and  $v_2(\pi^+)$  shows a linear dependence on charge asymmetry in Au + Au collisions at  $\sqrt{s_{NN}} = 200, 62.4, 39$  and  $27$  GeV and in U + U collisions at  $\sqrt{s_{NN}} = 193$  GeV. The UrQMD model calculations cannot reproduce this feature in Au + Au collisions at  $\sqrt{s_{NN}} = 200$  GeV.
- As a function of collision centrality, the slope parameter  $r$  for  $\pi^\pm$  shows a rise and fall feature from central to peripheral collisions and the energy dependence seems weak from 200 to 27 GeV.
- The centrality dependence of the  $\Delta v_2(\pi)$  slope shows a similar trend of the calculations based on Chiral Magnetic Wave at 200 GeV. However, model calculations do not reproduce the centrality dependence exactly.
- The  $v_3(\pi^\pm)$  as a function of  $A_{ch}$  has also been studied in Au + Au and U + U collisions. The centrality dependence of  $\Delta v_3(\pi)$  slope is similar to that of  $\Delta v_2(\pi)$  slope. However, the  $\Delta v_3(\pi)$  slope is consistent with zero in mid-centrality collisions.
- At low energies, i.e.  $\sqrt{s_{NN}} = 19.6, 11.5$  and  $7.7$  GeV, the slopes are consistent with zero with large bar errors.