

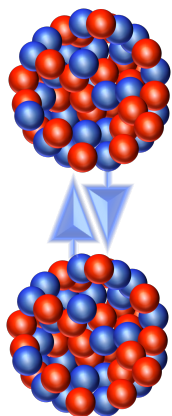
Observation of nuclear deformation in isobar collisions

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Nuclear Structure



$$\rho(r, \theta, \phi) = \frac{\rho_0}{1 + e^{(r - R_0(1 + \sum_n \beta_n Y_n^0(\theta, \phi)))/a_0}}$$

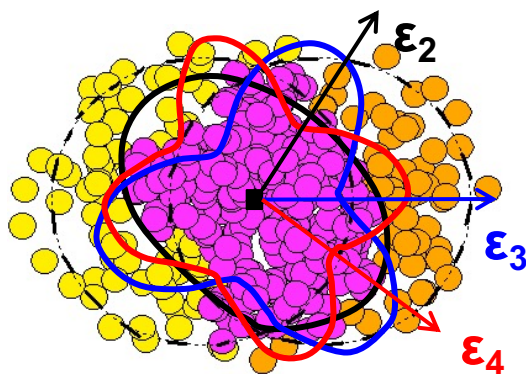
$\beta_2 \rightarrow$ Quadrupole deformation

$\beta_3 \rightarrow$ Octupole deformation

$a_0 \rightarrow$ Surface diffuseness

$R_0 \rightarrow$ Nuclear size

Initial State



Initial Size

$$R_{\perp}^2 \propto \langle r_{\perp}^2 \rangle$$

Initial Shape

$$\mathcal{E}_n \propto \langle r_{\perp}^n e^{in\phi} \rangle$$

R_0

a_0

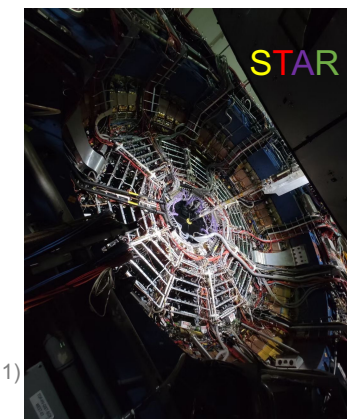
β_n

Hydrodynamic response

D. Teaney and L. Yan, PRC86, 044908(2012)

H.C. Song, S.A. Bass, U. Heinz, T. Hirano and C. Shen, PRL106, 192301(2011)

Final state particles



Radial Flow

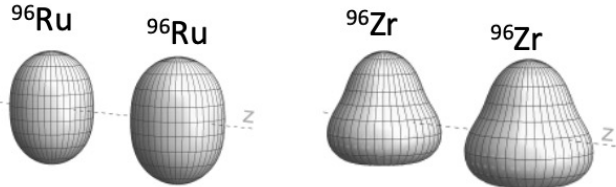
Anisotropic Flow

$$\frac{d^2 N}{d\phi dp_T} = N(p_T) \left(\sum_n V_n e^{-in\phi} \right)$$

High energy: approximate linear Response in each event

$$\frac{\delta[p_T]}{[p_T]} \propto -\frac{\delta R_{\perp}}{R_{\perp}} \quad V_n \propto \mathcal{E}_n$$

Harmonic flow and non-linear coupling coefficient

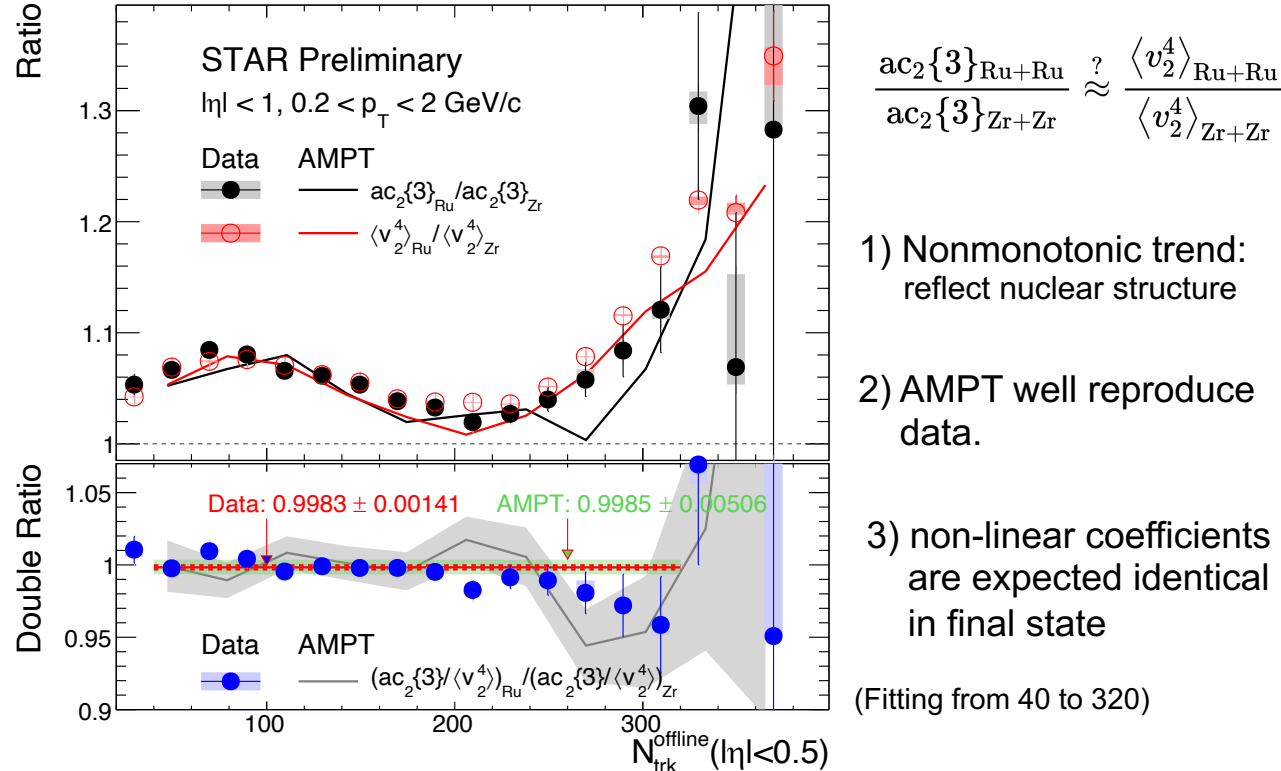
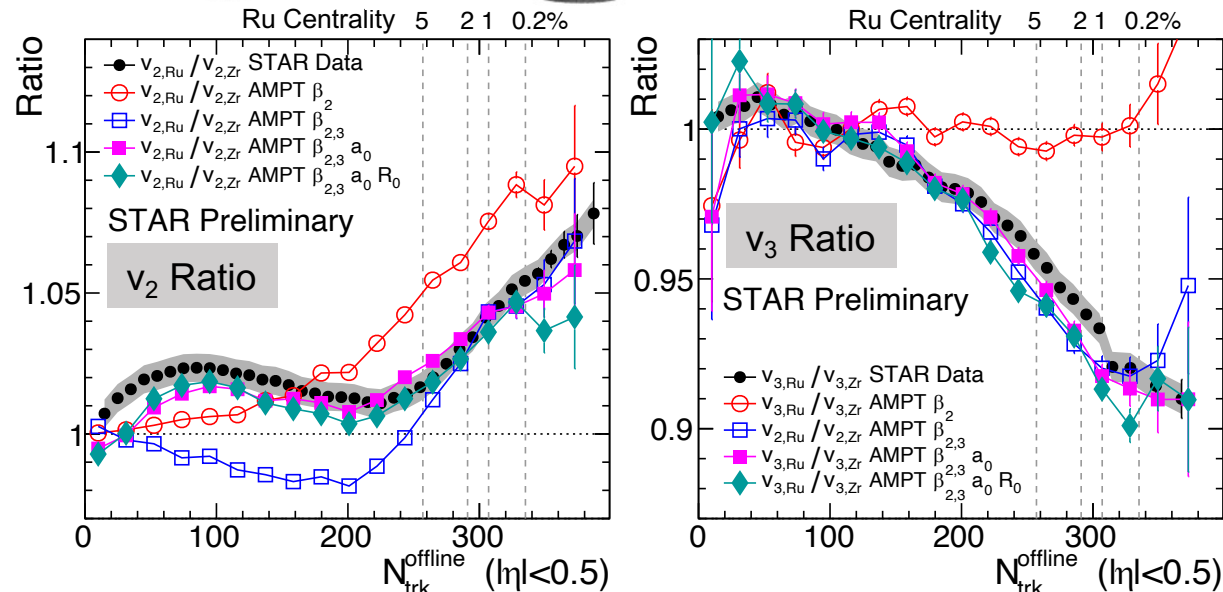


Nuclear parameters used in AMPT:

Species	β_2	β_3	a_0 (fm)	R_0 (fm)
Ru	0.162	0	0.46	5.09
Zr	0.06	0.20	0.52	5.02

Asymmetric cumulant: $ac_2\{3\} = \langle V_2^2 V_4^* \rangle = \langle v_2^2 v_4 \cos 4(\Phi_2 - \Phi_4) \rangle$

Non-linear coefficient: $\chi_{4,22} = \frac{\langle V_2^2 V_4^* \rangle}{\langle v_2^4 \rangle}$ $V_4 = U_4 + \chi_{4,22} V_2^2$



Heavy-ion expectation: $v_2^2 = a_2 + b_2 \beta_2^2 + b_{2,3} \beta_3^2$, $v_3^2 = a_3 + b_3 \beta_3^2$

$$\frac{v_{2,Ru}^2}{v_{2,Zr}^2} \approx 1 + \frac{b_2}{a_2} (\beta_{2,Ru}^2 - \beta_{2,Zr}^2) - \frac{b_{2,3}}{a_2} \beta_{3,Zr}^2$$

$$\frac{v_{3,Ru}^2}{v_{3,Zr}^2} \approx 1 - \frac{b_3}{a_3} \beta_{3,Zr}^2 < 1$$

Cancellation expected in non-central collisions

AMPT extractions:

$\beta_2^{Ru} = 0.16 \pm 0.02$ $\beta_3^{Zr} = 0.20 \pm 0.02$ $\Delta a_{0,Ru-Zr} = -0.06$ fm
 Direct indication and well constrain the nuclear deformation

Data: $\frac{\chi_{4,22}^{Ru+Ru}}{\chi_{4,22}^{Zr+Zr}} = 0.9983 \pm 0.00141$ AMPT: $\frac{\chi_{4,22}^{Ru+Ru}}{\chi_{4,22}^{Zr+Zr}} = 0.9985 \pm 0.00506$
 non-linear coupling are comparable between Ru and Zr

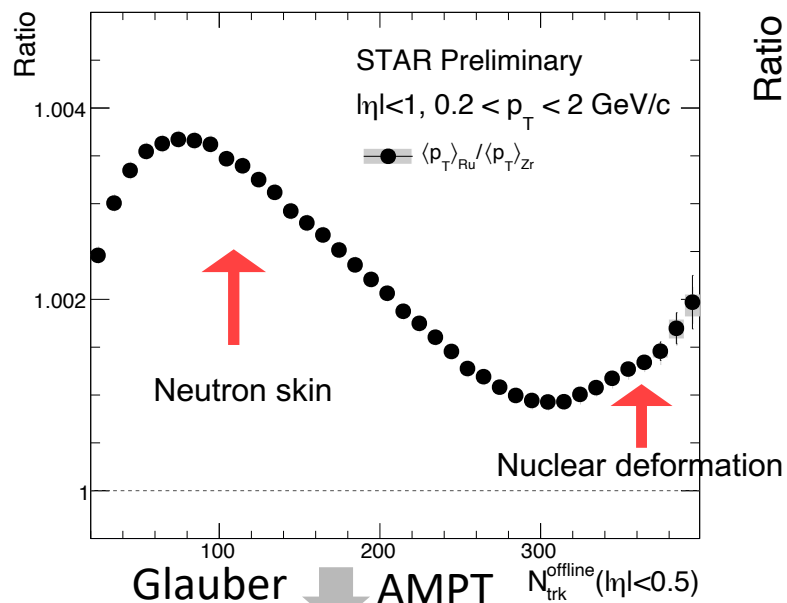
- 1) Nonmonotonic trend: reflect nuclear structure
 - 2) AMPT well reproduce data.
 - 3) non-linear coefficients are expected identical in final state
- (Fitting from 40 to 320)

Mean transverse momentum fluctuations

Mean

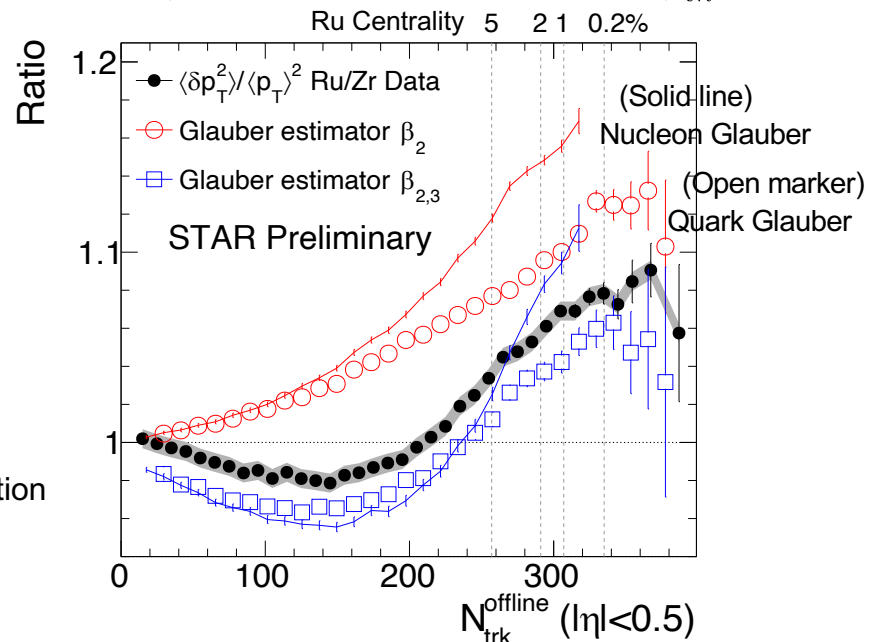
$$[p_T] \equiv \frac{\sum_i w_i p_{T,i}}{\sum_i w_i}, \langle\langle p_T \rangle\rangle \equiv \langle\langle p_T \rangle\rangle_{\text{evt}}$$

w_i is track weight



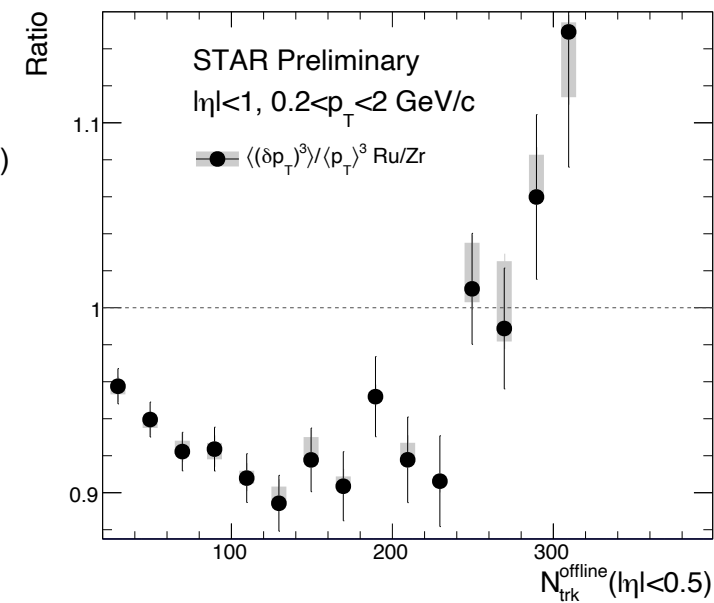
Variance

$$\langle\langle (\delta p_T)^2 \rangle\rangle = \left\langle \frac{\sum_{i \neq j} w_i w_j (p_{T,i} - \langle p_T \rangle)(p_{T,j} - \langle p_T \rangle)}{\sum_{i \neq j} w_i w_j} \right\rangle_{\text{evt}}$$



Skewness

$$\langle\langle (\delta p_T)^3 \rangle\rangle = \left\langle \frac{\sum_{i \neq j \neq k} w_i w_j w_k (p_{T,i} - \langle p_T \rangle)(p_{T,j} - \langle p_T \rangle)(p_{T,k} - \langle p_T \rangle)}{\sum_{i \neq j \neq k} w_i w_j w_k} \right\rangle_{\text{evt}}$$

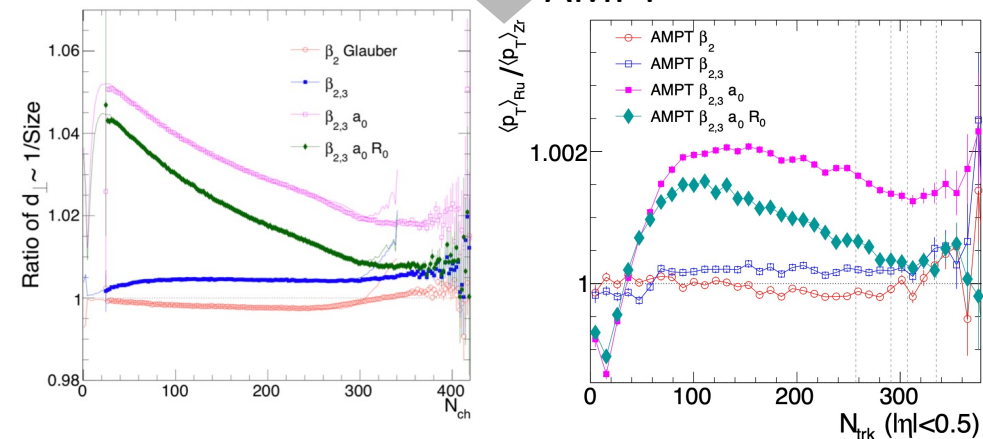


1) Mean: nonmonotonic trend reflects neutron skin and nuclear deformation.

2) Variance and skewness:

Enhancement in mid-central due to $\beta_{2,\text{Ru}}$ and large suppression due to $\beta_{3,\text{Zr}}$.

A complementary probe to decipher the Ru and Zr structure.



Pearson correlation coefficient

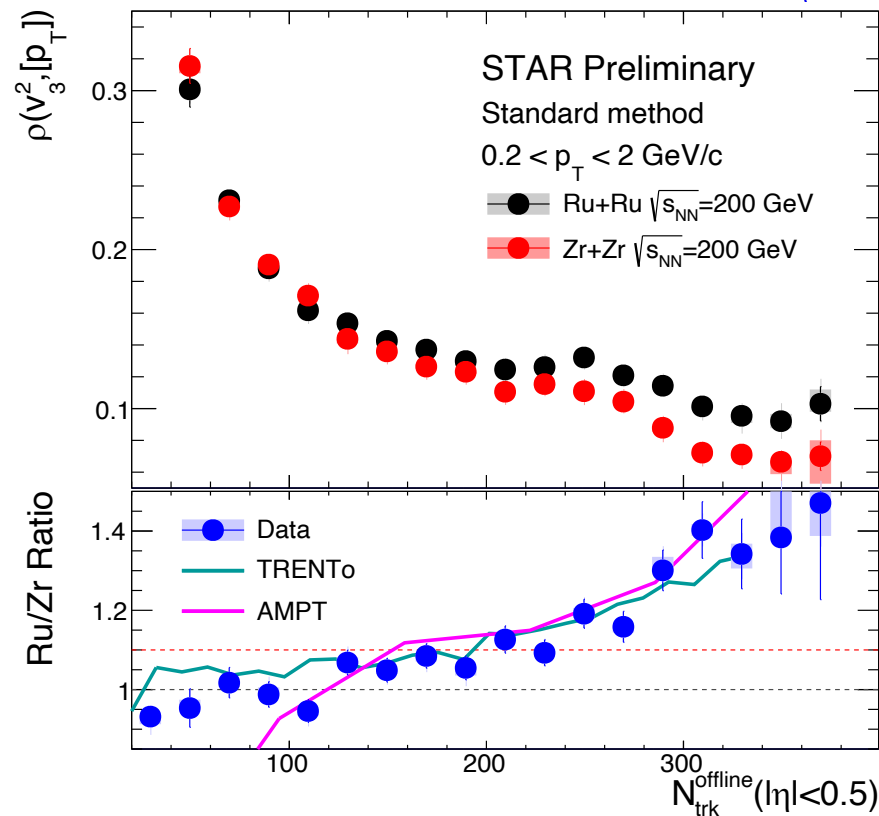
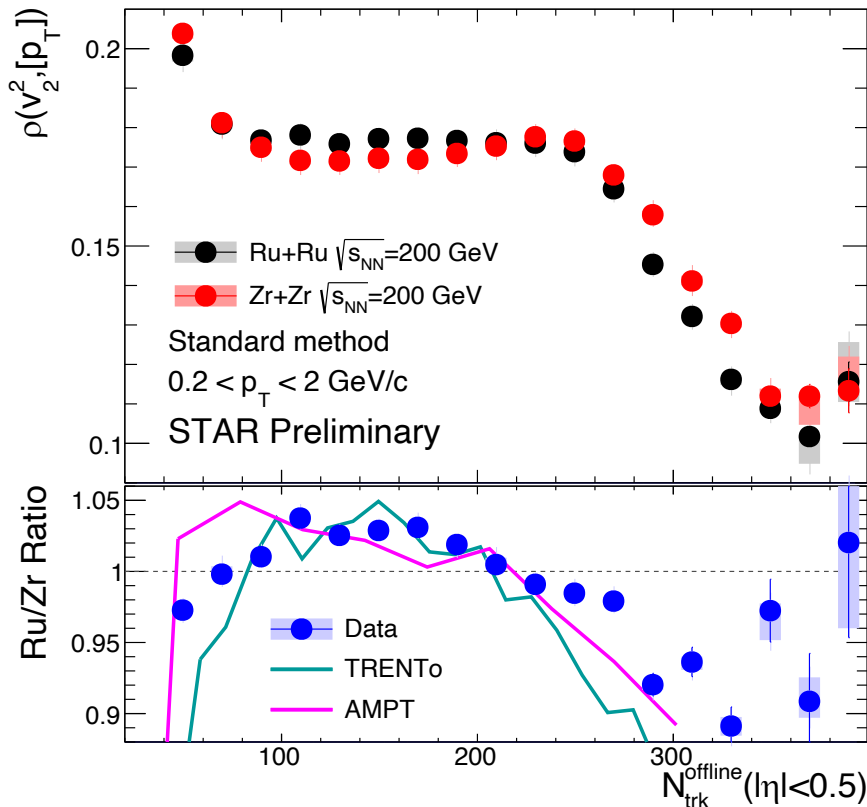
v_n -[p_T] three particle correlator:

$$\rho(v_n^2, [p_T]) = \frac{\text{cov}(v_n^2, [p_T])}{\sqrt{\text{Var}(v_n^2)_{\text{dyn}} \langle \delta p_T \delta p_T \rangle}}$$

$$\text{cov}(v_n^2, [p_T]) \equiv \left\langle \frac{\sum_{i \neq j \neq k} w_i w_j w_k e^{in\phi_i} e^{-in\phi_j} (p_{T,k} - \langle p_T \rangle)}{\sum_{i \neq j \neq k} w_i w_j w_k} \right\rangle_{\text{evt}}$$

$$\text{Var}(v_n^2)_{\text{dyn}} = v_n\{2\}^4 - v_n\{4\}^4$$

$$\langle \delta p_T \delta p_T \rangle = \left\langle \frac{\sum_{i \neq j} w_i w_j (p_{T,i} - \langle p_T \rangle) (p_{T,j} - \langle p_T \rangle)}{\sum_{i \neq j} w_i w_j} \right\rangle_{\text{evt}}$$



1) Ru/Zr ratio also reflects the possible nuclear structure.

2) Mostly dominated by the harmonic flow contributions.

3) TRENTo and AMPT can reproduce data

TRENTo assumption: $v_n \propto \epsilon_n$ $[p_T] \propto \frac{E}{S}$
(TRENTo calc. from Giacalone)

Conclusions and outlooks

- v_n ratios as a new probe to constrain nuclear structure parameters:

$$\beta_2^{\text{Ru}} = 0.16 \pm 0.02 \quad \beta_3^{\text{Zr}} = 0.20 \pm 0.02 \quad \Delta a_{0,\text{Ru-Zr}} = -0.06 \text{ fm} \quad \text{from AMPT estimation}$$



- Experimental test on the non-linear coupling coefficient: identical for Ru+Ru and Zr+Zr in final state as expected

$$\text{Data : } \frac{\chi_{4,22}^{\text{Ru+Ru}}}{\chi_{4,22}^{\text{Zr+Zr}}} = 0.9983 \pm 0.00141 \quad \text{AMPT : } \frac{\chi_{4,22}^{\text{Ru+Ru}}}{\chi_{4,22}^{\text{Zr+Zr}}} = 0.9985 \pm 0.00506$$

- Mean p_T fluctuations also as a complementary probe to decipher nuclear structure:

Nonmonotonic trend in mean, variance and skewness ratios

- Pearson correlation coefficient also reflects possible nuclear structure dominated by flow in isobar.

TRENTo and AMPT reproduce data

Thank you @

