



Femtocopic measurements of strange hadrons in Au+Au collisions at the STAR experiment

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Abstract

Relativistic heavy-ion collisions can study properties of nuclear matter in high-energy experiments like the STAR experiment. Femtoscopy, which relies on information carried by the particles produced in the collisions, is one of methods to learn about the bulk matter. By studying the quantum statistical effects and final state interactions between two particles, one can study spatial and temporal extents of particle emitting source. For the case of kaons, the correlation functions are sensitive to the early stage of the collision evolution and provide different information about particle-emitting sources compared to pions. Information on the final state interactions amongst the particles under study can also be extracted from the measurement. Especially, in the case of strange particle correlations, one could investigate hyperon-nucleon interactions which is little known.





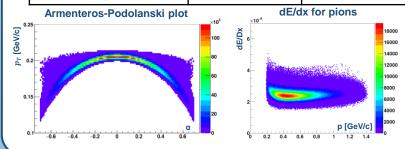


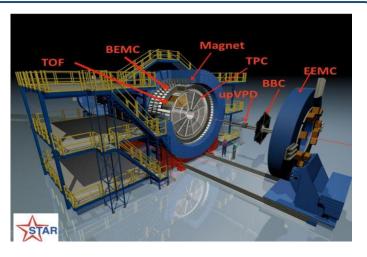
The STAR experiment



Neutral kaon selection criteria

Cuts	200 GeV	39 GeV	
$p_T [{\sf GeV/c}]$	0.4 - 2.0		
lηl	< 0.5		
DCA V ⁰ to PV [cm]	0.0 – 0.3		
DCA of daughters [cm]	0.6	0.8	
decay lenght [cm]	3	2	
Armenteros q_T [GeV/ c]	0.12 – 0.22		
Armenteros α	< 0.7		
mass range [GeV/c ²]	0.488 – 0.51	0.475 - 0.525	





Numbers of events

Energy/centrality	0-10%	10-70%	0-70%
200 GeV	~107M	~154M	~261M
39 GeV	~12M	~71M	~83M





Basics of femtoscopy



Femtoscopy – method to examine **the particle emitting source** sizes (of the order of 10^{-15} fm and lifetime 10^{-23} s) by measurements of relative momentum characteristics.

The correlation function (CF) – the ratio of probability of observing two particles with specific momenta p_1 and p_2 at the same place and time to the product of probabilities to find them separately:

$$CF(\overrightarrow{p_1},\overrightarrow{p_2}) = \frac{P_2(\overrightarrow{p_1},\overrightarrow{p_2})}{P_1(\overrightarrow{p_1})P_1(\overrightarrow{p_2})}$$

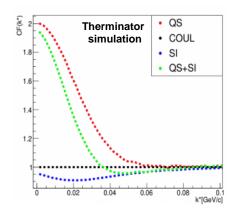
The experimental correlation function:

$$CF(q_{inv}) = \frac{A(q_{inv})}{B(q_{inv})}$$

$$q_{inv} = \sqrt{(p_1 - p_2)^2 - (E_1 - E_2)^2}$$

 $A(q_{inv})$ – the signal distribution, $B(q_{inv})$ - the background distribution.

Neutral kaons correlations



The correlation function depends on:

- Quantum statistics (QS)
- Final State Interactions (FSI)
 - Coulomb Interaction (COUL)
 - Strong Interaction (SI)





Parametrizations



Gaussian density distribution (includes only QS effects): $CF(q_{inv}) = 1 + \lambda e^{[-R_{inv}^2 q_{inv}^2]}$

 λ – the correlation strenght, R_{inv} - the size of the particle-emitting source.

Lednicky and Lyuboshitz model includes strong FSI: Sov. J. Nucl. Phys. 35, 770 (1982)

$$CF(q_{inv}) = 1 + \lambda \left(e^{[-R_{inv}^2 q_{inv}^2]} + \frac{1}{2} \left[\left| \frac{f(k^*)}{R_{inv}} \right|^2 + \frac{4\Re f(k^*)}{\sqrt{\pi} R_{inv}} F_1(q_{inv} R_{inv}) - \frac{2\Im f(k^*)}{\sqrt{\pi} R_{inv}} F_2(q_{inv} R_{inv}) \right] \right)$$

QS effects

strong FSI through the $f_0(980)$ and $a_0(980)$ resonances

$$\begin{split} F_1(z) &= \int_0^z dx \, \frac{e^{x^2}-z^2}{z}, \quad F_2(z) = \frac{1-e^{-z^2}}{z} \\ f(k^*) &= \frac{1}{2} [f_0(k^*) + f_1(k^*)], \quad f_I(k^*) = \frac{\gamma_r}{m_r - s - i \gamma_r k^* - i \gamma_r' k_r'}, \quad s = 4 \left(m_K^2 + k^{*2} \right) \end{split}$$

	$m_{f_0} \left[\frac{\text{GeV}}{\text{c}^2} \right]$	$\gamma_{f_0 K \overline{K}}$	$\gamma_{f_0\pi\pi}$	$m_{a_0}[rac{{ m GeV}}{{ m c}^2}]$	$\gamma_{a_0 K \overline{K}}$	$\gamma_{a_0\pi\pi}$
Antonelli [1]	0.973	2.763	0.5283	0.985	0.4038	0.3711
Achasov2001 [2]	0.996	1.305	0.2684	0.992	0.5555	0.4401
Achasov2003 [3]	0.996	1.305	0.2684	1.003	0.8365	0.4580
Martin [4]	0.978	0.792	0.1990	0.974	0.3330	0.2220

[1] eConf C020620, THAT06 (2002), [2] Phys. Rev. D 63, 094007 (2001)

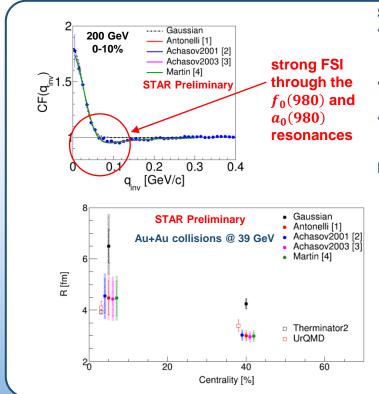
[3] Phys. Rev. D 68, 014006 (2003), [4] Nucl. Phys. B 121, 514–530 (1977)





Results





Summary:

- The strong final-state interaction has a significant effect on the $K_s^0K_s^0$ correlation due to the near-threshold $f_0(980)$ and $a_0(980)$ resonanses
- The radii of the source depend on centrality and increase with increasing collision energy
- Comparison with model calculations shows better compability for UrQMD model

Future plans: $K_s^0 K^{ch}$ correlation functions

