Exploring Electromagnetic Field Effects and Constraining Transport Parameters of QGP using STAR BES-II data

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Abstract. Heavy-ion collisions undergo various stages in their evolution and 5 it is crucial to disentangle the initial- and final-stage effects. In this work, we 6 report measurements of two types of observables: (i) charge-dependent directed 7 flow (Δv_1) , which is sensitive to the initial ultra-strong electromagnetic fields, 8 and (ii) flow correlations, such as $r_n(\eta)$ which is sensitive to the initial longitu-9 dinal de-correlation, and correlations among flow harmonics. These measure-10 ments contribute to constraining the initial state of the collisions, and through 11 a comprehensive beam energy scan, we gain significant insights into the sys-12 tem evolution in the presence of initial spatial asymmetry and electromagnetic 13 fields. 14

15 1 Introduction

Ultra-strong magnetic fields on the order of 10¹⁸ Gauss, the strongest magnetic fields ob-16 served in nature, are expected to be produced in the very early stages of heavy-ion collisions 17 at RHIC energies [1, 2]. The magnetic fields are mostly created by spectators and decay very 18 fast, with time scales comparable to the passage time of the colliding nuclei [1, 2]. However, 19 the decay of the fields can be compensated by the Faraday induction effect, which depends 20 on the medium properties (such as electrical conductivity) and on the formation time of the 21 quarks. Additionally, the study of the formation and decay of the initial electromagnetic fields 22 is essential for understanding the evolution of the quark-gluon plasma (QGP) in the presence 23 of electromagnetic (EM) fields. 24

The initial state in heavy-ion collisions could have significant longitudinal de-correlation, leading to a difference between event planes reconstructed at different pseudorapidity ranges [3, 4]. Moreover, the initial-state geometry and asymmetries in energy deposition could evolve into final-state flow harmonics and event-plane angular correlations, the study of which can be used to constrain various initial-state models and understand the mechanisms of energy deposition by the colliding nuclei.

2 Experimental Results

32 2.1 EM-field Effects on Directed Flow

Quarks carry electric charges and experience various electromagnetic forces. In particular, the Lorentz force due to the Hall effect, $\mathbf{F} = q\mathbf{v} \times \mathbf{B}$, acts in the opposite direction to the

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Figure 1: $\Delta(dv_1/dy)$ as a function of centrality for pions, kaons, and protons in Au+Au collisions at $\sqrt{s_{NN}} = 200, 27, 19.6, 14.6, \text{ and } 7.7 \text{ GeV}$. At 200 GeV, model calculations for protons from UrQMD and iEBE-VISHNU are also added for comparison.



Figure 2: $\Delta v_1(p_T)$ for pions, kaons, and protons in Au+Au collisions at $\sqrt{s_{NN}} = 19.6$ GeV.

forces due to the Faraday induction effect (generated by the decaying magnetic field) and the Coulomb effect (the electric field due to spectators). As electromagnetic forces are proportional to charge, the difference in directed flow slope between positively and negatively charged particles ($\Delta dv_1/dy$) can serve as a probe to the EM-field effects in heavy-ion collisions [1, 2]. The Hall effect should impart a positive $\Delta dv_1/dy$, while the Faraday+Coulomb effect should give a negative $\Delta dv_1/dy$ to charged particles.

Furthermore, quarks transported from the colliding nuclei to midrapidities can have a different directed flow than those produced in pairs. Model studies [6–8] show that transportedquark effect could contribute a positive $\Delta dv_1/dy$ for protons and kaons, and a negative $\Delta dv_1/dy$ for pions due to their quark contents [9].

Figure 1 shows the STAR measurements of $\Delta dv_1/dy$ as a function of centrality for π^{\pm}, K^{\pm} 45 and p, \bar{p} in Au+Au collisions at $\sqrt{s_{NN}} = 200-7.7$ GeV [5, 9]. $\Delta dv_1/dy$ becomes negative in 46 peripheral collisions (except for pions at 200 GeV), which is expected from the dominance 47 of the Faraday+Coulomb effect. Other mechanisms for the centrality dependence of $\Delta dv_1/dy$ 48 are under investigation [10, 11]. Data at $\sqrt{s_{NN}} = 200$ GeV are comparable to the IEBE-49 VISHNU+EM field calculations with conductivity $\sigma = 0.023$ fm⁻¹ from lattice QCD [1, 2]. 50 The increase in the magnitude of $\Delta dv_1/dy$ in peripheral collisions with decrease in beam 51 energy, observed in data, is expected from the longer passage times $(2R/\gamma)$ and shorter life-52 times of the fireball (due to which the late-time dilution of the v_1 splitting is smaller) at lower 53 collision energies [1, 2]. 54



Figure 3: $r_2(\eta)$ and $r_3(\eta)$ in 10–40% centrality Au+Au collisions at $\sqrt{s_{NN}} = 19.6-54.4$ GeV. 2.5< $\eta_{ref} < 4.0 \ (2.1 < |\eta_{ref}| < 5.1)$ is selected for 54.4 GeV (27 and 19.6 GeV).



Figure 4: Six-particle (normalized) cumulants as a function of centrality for Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV.

Figure 2 shows Δv_1 (p_T) for pions, kaons, and protons in Au+Au collisions at $\sqrt{s_{NN}} =$

⁵⁶ 19.6 GeV. The negative Δv_1 and the increase in $|\Delta v_1|$ with p_T are expected from theory calcu-

⁵⁷ lations [1, 2] and could arise from a decrease in the Hall effect at higher p_T (or smaller p_z).

Similar observations hold at $\sqrt{s_{NN}} = 14.6$ and 7.7 GeV.

59 2.2 Longitudinal De-correlation

The $r_n(\eta)$ observable [3, 4] is sensitive to the de-correlation between event planes reconstructed from the pseudorapidity ranges η and $-\eta$, respectively. $r_n(\eta)$ is expected to be 1 when there are no de-correlation or nonflow effects. Figure 3 shows the measurements of $r_2(\eta)$ and $r_3(\eta)$ in the 10–40% centrality range, significantly deviating from unity at RHIC energies. AMPT [12] shows stronger deviation for $r_2(\eta)$ and comparable deviation for $r_3(\eta)$ compared to experimental data.

66 2.3 Constraining Initial State Using Correlations

⁶⁷ Symmetric Cumulants (SC) are expected to be sensitive to the interplay between initial- and

final-state effects. In contrast, Normalized Symmetric Cumulants (NSC) are predicted to be

dominated by initial-state effects. Thus, simultaneous measurements of SC and NSC can constrain the initial- and final-state effects [14]. The six-particle SC and NSC are given as

$$SC(n.m)\{6\} = \langle 6 \rangle_{nnm} - \langle 4 \rangle_{nn} \langle 2 \rangle_m - 2\langle 4 \rangle_{nm} \langle 2 \rangle_n + 2\langle 2 \rangle_n^2 \langle 2 \rangle_m, \tag{1}$$

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$$NSC(n.m)\{6\} = \frac{SC(n,m)\{6\}}{\langle 2\rangle_n^{sub} \langle 2\rangle_n^{sub} \langle 2\rangle_m^{sub}}.$$
(2)

Figure 4 shows the six-particle (normalized) symmetric cumulants as a function of centrality in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV [13, 14]. Anti-correlation between v_2 and v_3 is observed which is expected from the anti-correlation between ϵ_2 and ϵ_3 . A positive correlation between v_2 and v_4 is consistent with mode coupling between v_2 and v_4 . These measurements can help constrain various initial-state models.

77 3 Summary

We have presented measurements of Δv_1 , $r_n(\eta)$, and flow harmonic correlations in Au+Au 78 collisions at $\sqrt{s_{NN}} = 7.7 - 200$ GeV in the Beam Energy Scan phase-II program of the STAR 79 experiment. The Δv_1 results are compatible with models using strong electromagnetic fields 80 and conductivity from lattice QCD. Proton $\Delta dv_1/dy$ is negative in peripheral collisions, sup-81 porting the dominance of the Faraday+Coulomb effect, and becomes more negative at lower 82 collision energies, which is expected from the corresponding longer lifetime of the electro-83 magnetic field and shorter lifetime of the fireball. $r_2(\eta)$ and $r_3(\eta)$ significantly deviate from 84 unity at RHIC energies and show beam energy dependence. From the six-particle (normal-85 ized) symmetric cumulant measurements, anti-correlation is observed between v_2 and v_3 , and 86 positive correlation is observed between v_2 and v_4 . 87

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