Charge dependent azimuthal correlations of K-π pairs at STAR **Charles Kent Riley, for the STAR Collaboration** Department of Physics, Yale University, New Haven CT

Local Parity Violation in Heavy-ion Collisions



STAR

Charge dependent azimuthal correlations have been used to probe for local parity violation (LPV) in the quark-gluon plasma (QGP) believed to be formed during Au+Au collisions at $\sqrt{s_{NN}}$ = 200 GeV at the BNL Relativistic Heavy Ion Collider (RHIC) [1]. It was suggested [2] that P- and CP-odd domains form in the hot dense QCD matter created **during the initial** stages of the collisions, where the chiral magnetic effect (CME) would act as the driving force behind a separation of charges across the reaction plane (RP).

- 2011 Data of $\sqrt{s_{NN}}$ = 200 GeV Au+Au Collisions 5.
- > Three-particle correlator (= $v_2\gamma$) uses a third particle, c (of any type and charge), to take the role of the RP
- > The difference between the opposite-sign and same-sign can be taken to be a charge separation signal



- > Comparing γ using all charged particles (i.e. mostly pions) from 2007 [5] and pairing K-π (2011)
- > The shaded region represents systematic errors of data from 2011 compared with 2007, the error bars are statistical

This electric dipole moment will be either parallel or anti-parallel to the strong magnetic field produced in the collision, meaning its expectation value over many events would vanish, however, the variance should reflect the fact that parity is broken event-by-event.

2. Measuring the Charge Separation

> Charge dependent azimuthal correlations were proposed [3] as a P-even observable that can detect these dipole fluctuations:

> $\gamma = \langle \cos(\Delta \varphi_a + \Delta \varphi_b) \rangle = \langle \cos \Delta \varphi_a \cos \Delta \varphi_b \rangle - \langle \sin \Delta \varphi_a \sin \Delta \varphi_b \rangle,$ where $\Delta \phi = \phi - \Psi_{RP}$ and a, b = + or - particles. $\langle \cos(\varphi_a - \varphi_b) \rangle < 0$

- \succ γ looks at whether charges are emitted in **Reaction Plane** same or opposite directions with respect to one another while taking the RP into $\langle \cos(\Delta \varphi_a + \Delta \varphi_b) \rangle > 0$ account
- > Another useful correlator:

 $\delta = \langle \cos(\varphi_a - \varphi_b) \rangle$

 \succ δ represents the tendency for two particles to be emitted in the same direction (positive) or in opposite directions (negative)

 $\langle \cos(\varphi_a - \varphi_b) \rangle > 0$ $\langle \cos(\varphi_a - \varphi_b) \rangle > 0$ **Reaction Plane** $\langle \cos(\Delta \varphi_a + \Delta \varphi_b) \rangle < 0$ $\langle \cos(\varphi_a - \varphi_b) \rangle < 0$





- > Comparing γ (of K- π pairs) before and after applying an elliptic flow correction
- Opposite-sign correlations are pulled more negative while same-sign correlations increase slightly – a slight reduction in the signal

- Charge separation signal is larger in K- π correlations for more peripheral bins – one cause is that selected kaons have higher transverse momentum
- \succ Comparing γ (of all charge pairs, 2011) before and after applying an elliptic flow correction
- Opposite-sign correlations are pulled more negative while same-sign correlations are largely unaffected – a slight reduction in the signal
- \succ Changes in signal reflects δ for all charge correlations





- > There may be a tendency for $\pi^+ \pi^$ pairs to be created during the later stages of hadronization, after elliptic flow is established [4]
 - Implies many pion pairs will be emitted at small angles with respect to one another, obscuring the LPV signal
- > It is expected that K- π pairs should suppress this effect



 $(1 + v_2)$

- Modifying according to flow
- We apply a first-order correction to compensate for the elliptic flow $\gamma = \langle \cos(\Delta \varphi_a + \Delta \varphi_b) \rangle$ $= \langle \cos \Delta \varphi_a \cos \Delta \varphi_b \rangle$
 - $-\langle \sin \Delta \varphi_a \sin \Delta \varphi_b \rangle$
- Applying a correction: $\langle \cos \Delta \varphi_a \cos \Delta \varphi_b \rangle (1 - v_2)$ $-\langle \sin \Delta \varphi_a \sin \Delta \varphi_b \rangle (1 + v_2)$ $=\gamma-\delta v_2,$ where v_2 is the elliptic flow

Changes in signal reflects δ for K- π correlations



- Comparing δ using all charged particles (i.e. mostly pions) from 2011 with K- π pairs
- Opposite-sign correlations are pulled more negative in K- π correlations since these will be removing significant contributions from late stage creation of pions
- The difference in same-sign correlations is intriguing

Conclusion **6**.

- > Charge separation effect similar to, if not greater than, previous methods using all charged pairings.
- \succ Looking at K- π pairs and applying flow corrections help suppress possible P-even sources, further refining previously published results[1].

3. Experimental Details

- \succ The Time Projection Chamber (TPC) is used to track charged particles having a 2π azimuthal coverage
- \triangleright Particles were analyzed within a pseudo-rapidity range $|\eta| < 1.0$ and transverse momentum range $0.2 < p_t < 2.0 \text{ GeV/c}$
- Particles were identified using the Tim-Of-Flight (TOF) detector using a mass squared cut: $0.001 < m^2 < 0.1 (\text{GeV/c}^2)^2$ for pions, $0.2 < m^2 < 0.3 (\text{GeV/c}^2)^2$ for kaons > TOF identified particles then had to pass a TPC cut: $n_{\sigma\pi} < 2.0$ for pions, $n_{\sigma K}$ < 2.0 for kaons
- Data indicates signal may be more due to particle emission preferentially occurring in the vicinity of in-plane more than out.
 - > In fact, one has to consider the possibilities of jets and other strong-interaction backgrounds which may dominate the CME signal.
- NB: K* resonances may contribute to opposite-sign K- π correlations, however, preliminary simulation results show they have little effect on the signal.
- The use of balance functions as a means to explain these results as a consequence of a combination of the blastwave model and charge conservation are also being studied - see Hui Wang's poster



[1] B. I. Abelev et al. (STAR Collaboration), Phys. Rev. Lett. 103, 251601 (2009) [2] D. E. Kharzeev, L. D. McLerran, and H. J. Warringa, Nucl. Phys. A803, 227 (2008) [3] S. A. Voloshin, Phys. Rev. C 70, 057901 (2004) [4] S. Pratt (2012) arXiv:1109.3647v3 [nucl-th] [5] D. Gangadharan (STAR Collaboration), J. Phys. G: Nucl. Part. Phys. 38, 124166 (2011)

The STAR Collaboration: <u>http://drupal.star.bnl.gov/STAR/presentations</u>

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