

Observation of ${}^3_{\Lambda}\text{H}$ and ${}^3_{\bar{\Lambda}}\bar{\text{H}}$ @ RHIC

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(for the STAR Collaboration)

- Introduction & Motivation
- Evidence for first **antihypernucleus**
 - ${}^3_{\Lambda}\text{H}$ and ${}^3_{\bar{\Lambda}}\bar{\text{H}}$ signal (for **discovery**)
 - Mass and Lifetime measurements
- Production rate and ratios
 - Yields as a measure of correlation
 - A case for RHIC energy scan
- Conclusions and Outlook

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a passion for discovery

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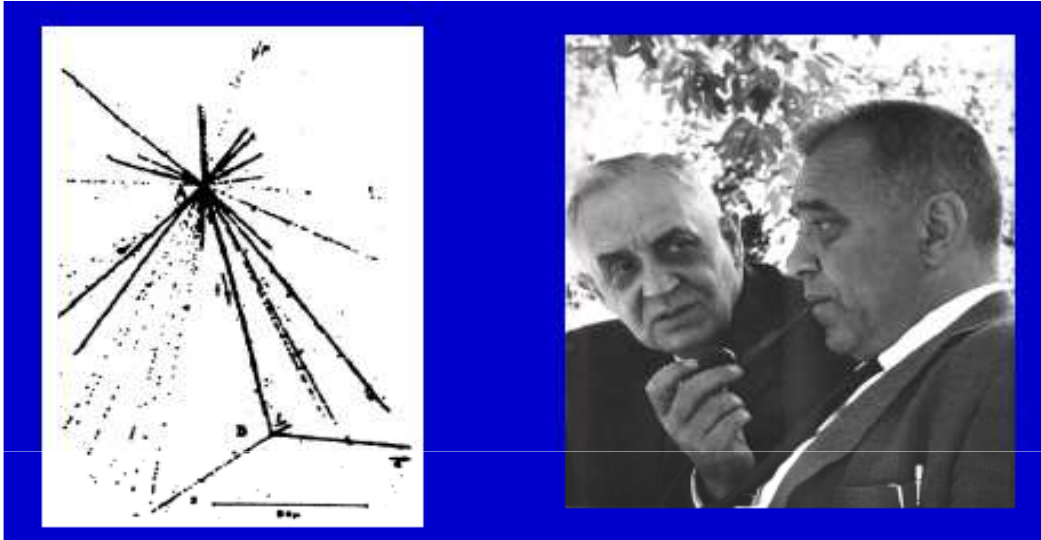
What are hypernuclei?

Nucleus which contains at least one hyperon in addition to nucleons.

Hypernuclei of lowest A

$${}^3_{\Lambda}H(n + p + \Lambda)$$

$${}^3_{\bar{\Lambda}}\bar{H}(\bar{n} + \bar{p} + \bar{\Lambda})$$



The first hypernucleus was discovered by Danysz and Pniewski in 1952. It was formed in a cosmic ray interaction in a balloon-flown emulsion plate.

M. Danysz and J. Pniewski, Phil. Mag. 44 (1953) 348

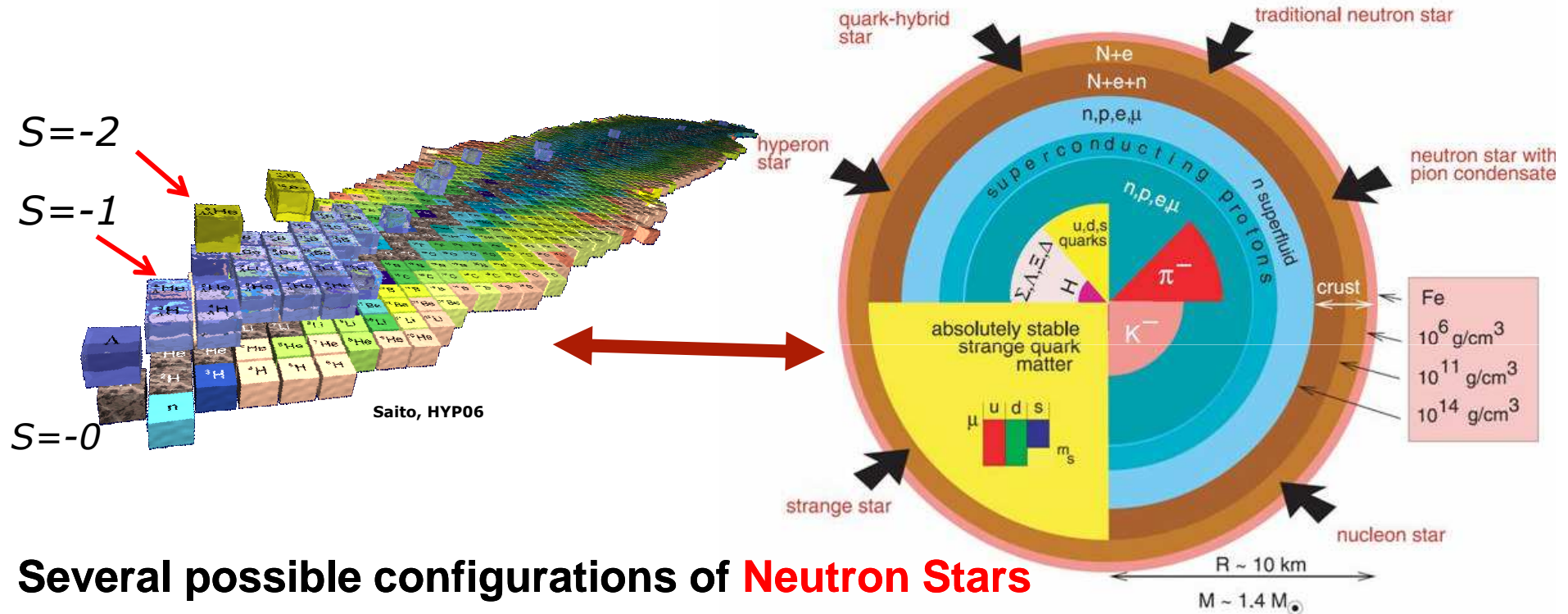
No one has ever observed
any antihypernucleus

- Y-N interaction: a good window to understand the baryon potential
- Binding energy and lifetime are very sensitive to Y-N interactions
- Hypertriton: $\Delta B = 130 \pm 50$ KeV; $r \sim 10$ fm
- Production rate via coalescence at RHIC depends on overlapping wave functions of $n+p+\Lambda$ in final state
- Important first step for searching for other exotic hypernuclei (double- Λ)



from Hypernuclei to Neutron Stars

hypernuclei $\leftarrow \Lambda$ -B Interaction \rightarrow Neutron Stars



Several possible configurations of **Neutron Stars**

- Kaon condensate, hyperons, strange quark matter

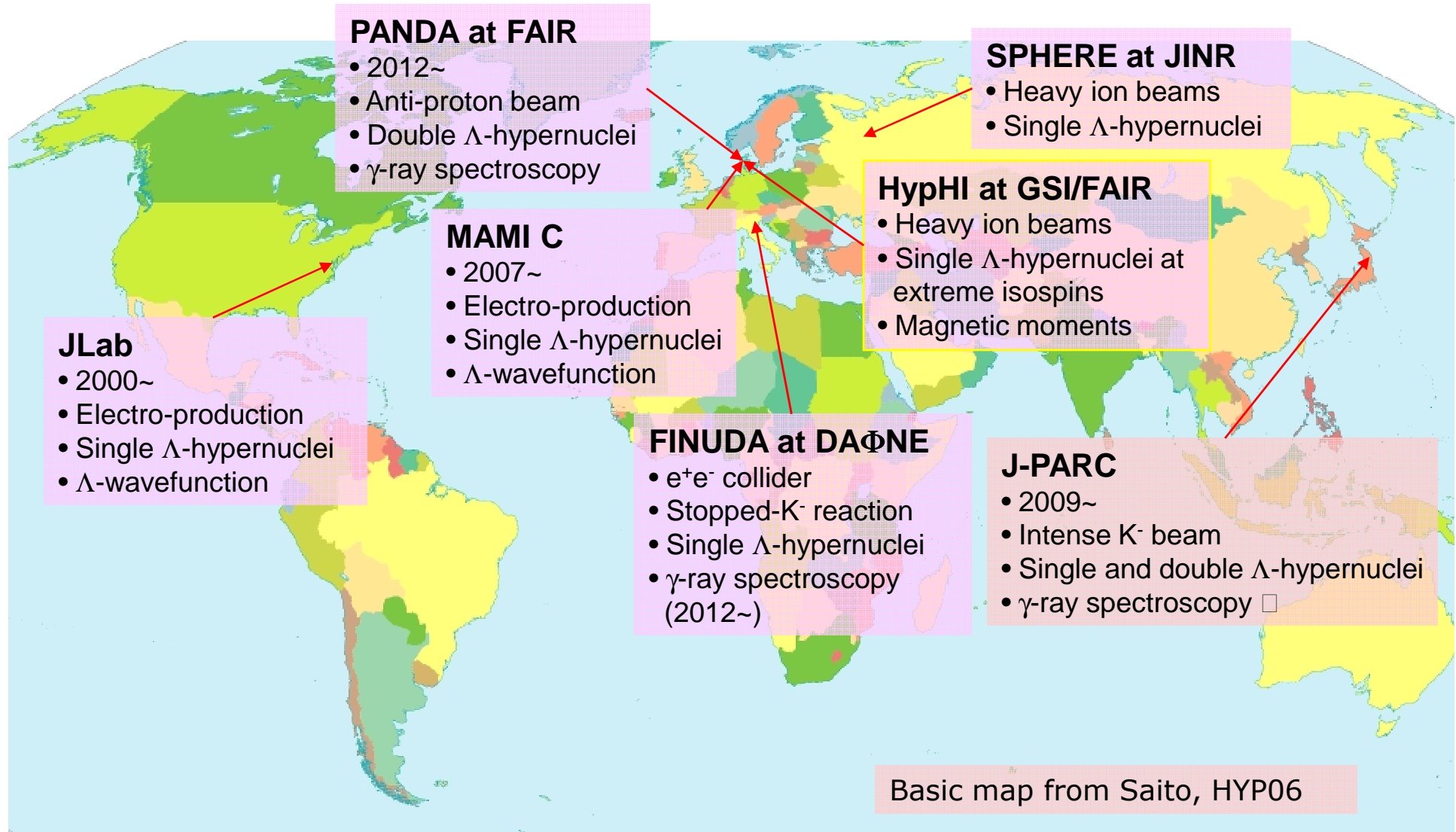
Single and **double** hypernuclei in the laboratory:

- study the **strange sector** of the baryon-baryon interaction
- provide info on EOS of neutron stars

J.M. Lattimer and M. Prakash,
"The Physics of Neutron Stars", Science 304, 536 (2004)
J. Schaffner and I. Mishustin, Phys. Rev. C 53 (1996):
Hyperon-rich matter in neutron stars



Current hypernucleus experiments





Can we observe hypernuclei at RHIC?

◆ Low energy and cosmic ray experiments (wikipedia):

hypernucleus **production** via

- Λ or K capture by nuclei
- the direct strangeness exchange reaction

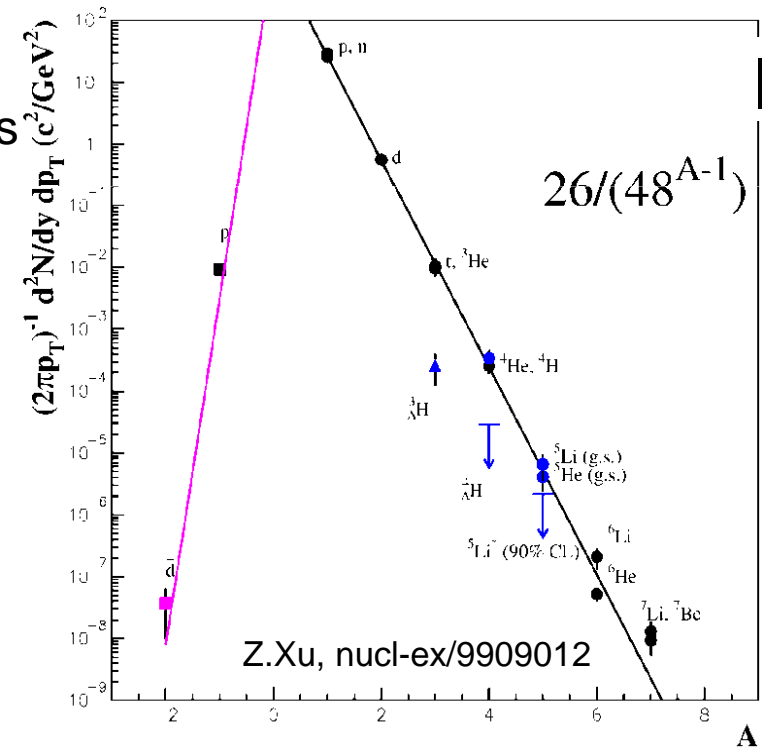
hypernuclei **observed**

- energetic but delayed decay,
- measure momentum of the K and π mesons

◆ In high energy heavy-ion collisions:

- nucleus production by coalescence, characterized by **penalty factor**.
- AGS data^[1] indicated that hypernucleus production will be further suppressed.
- What's the case at RHIC?

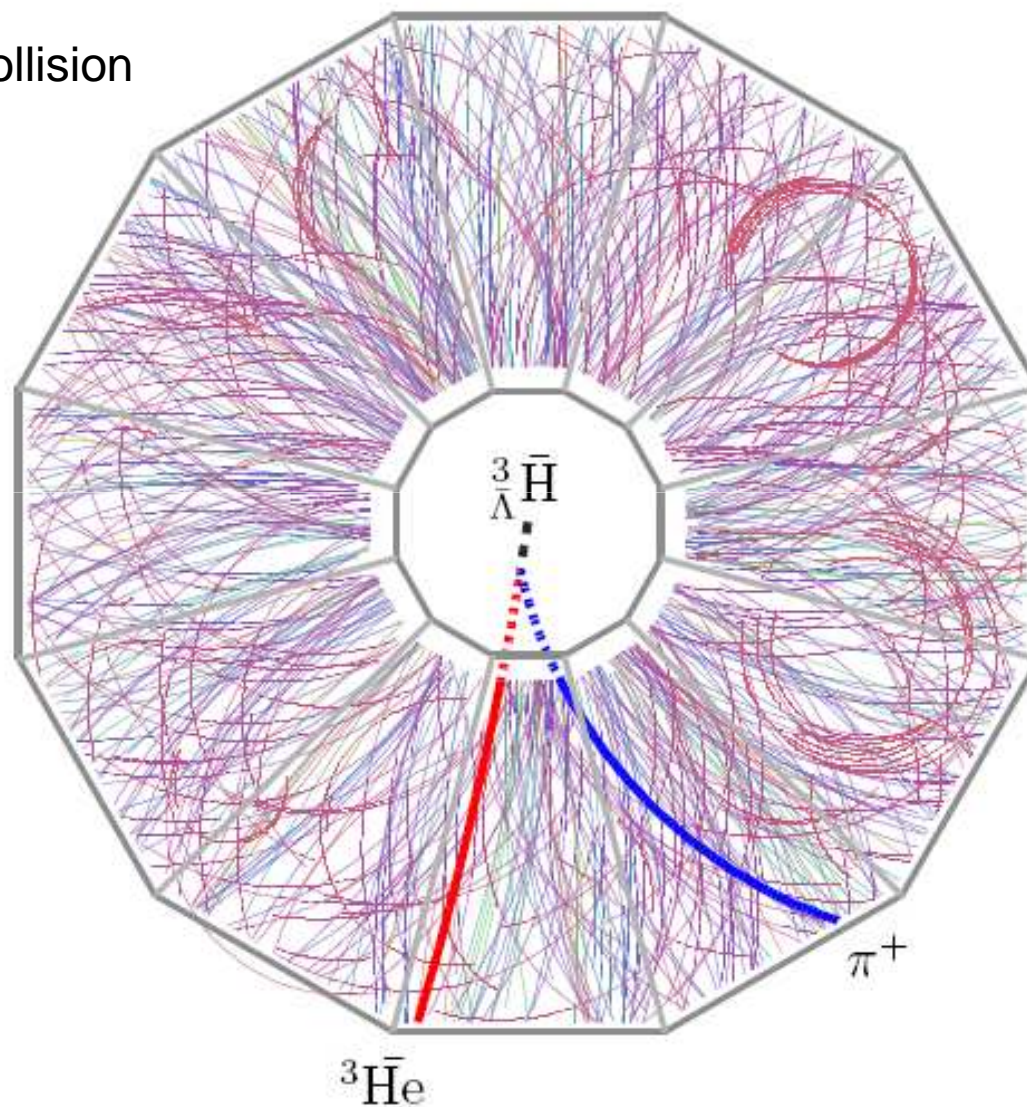
[1] AGS-E864, [Phys. Rev. C70,024902 \(2004\)](#)





A candidate event at STAR

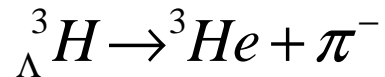
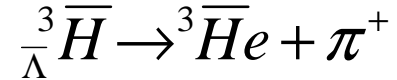
Run4 (2004)
200 GeV Au+Au collision





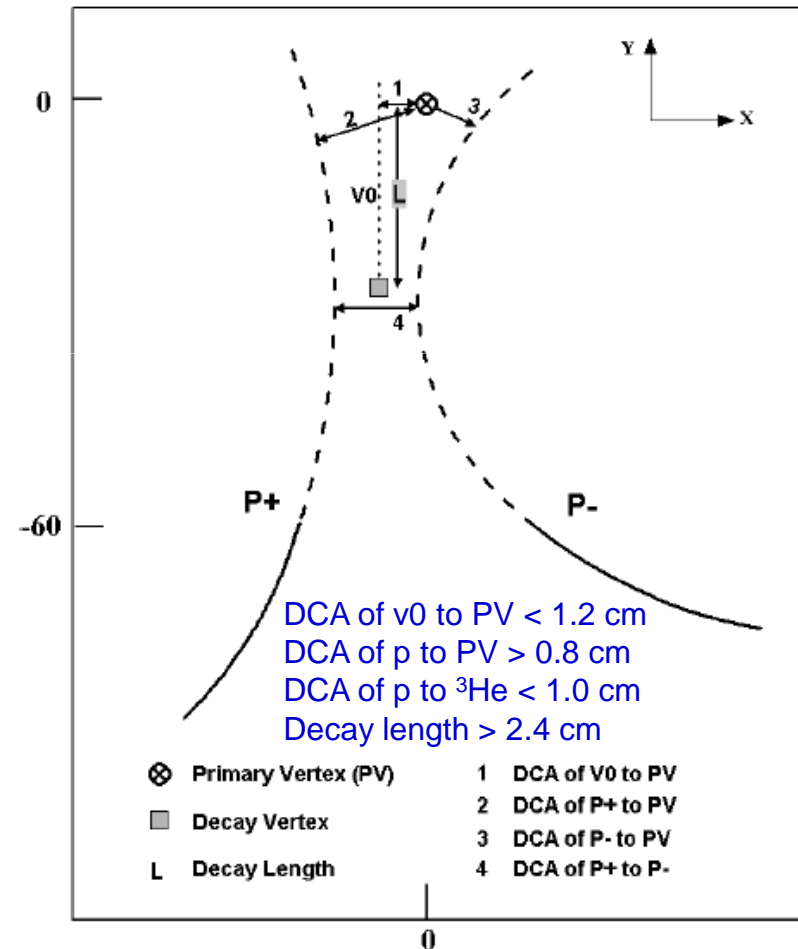
Data-set and track selection

${}^3_{\Lambda}H$ mesonic decay, $m=2.991$ GeV, B.R. 0.25;



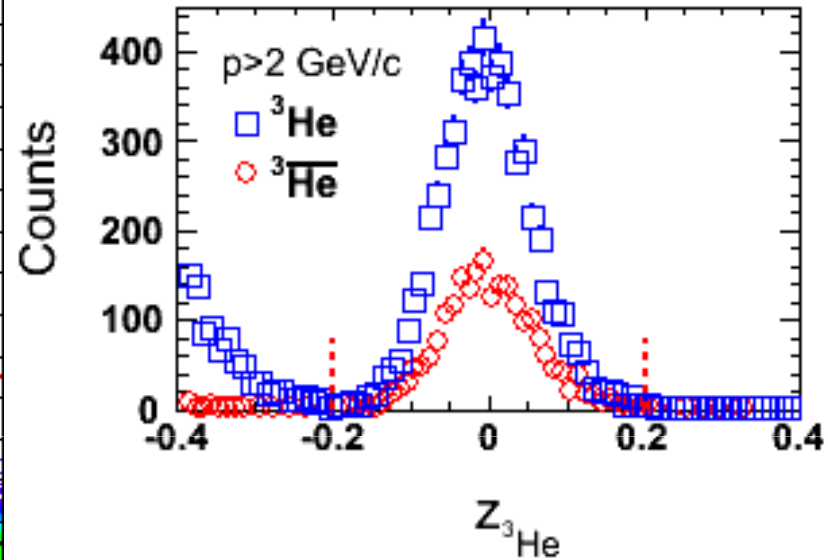
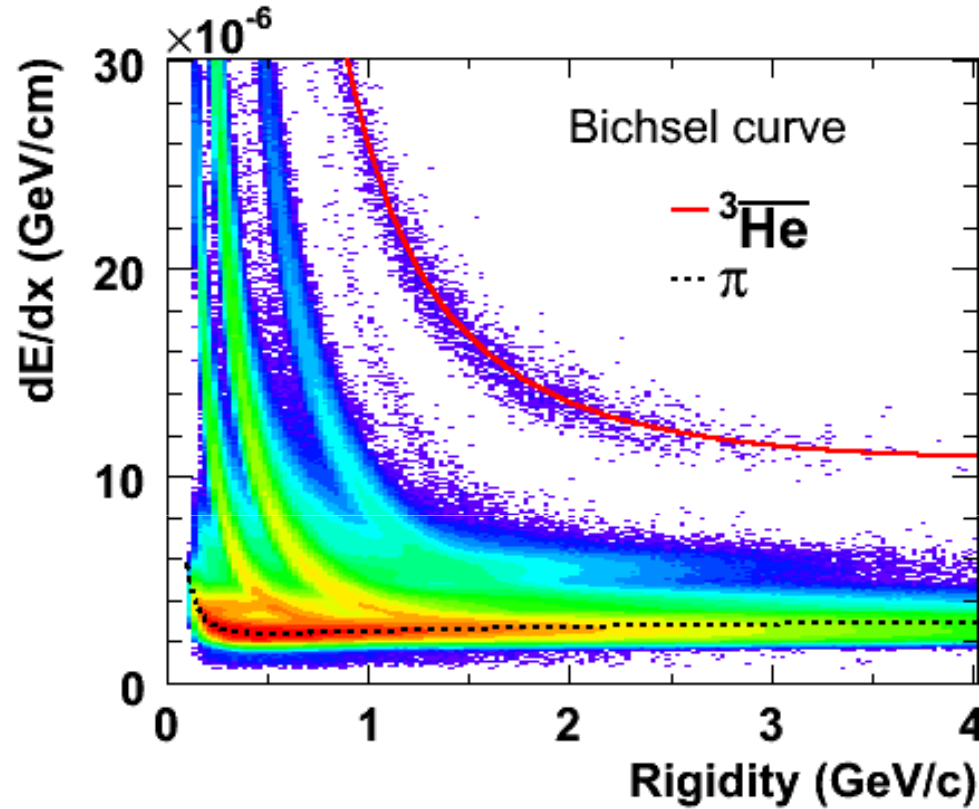
- Data-set used, Au+Au 200 GeV
 - ✓ ~67M Run7 MB,
 - ✓ ~23M Run4 central,
 - ✓ ~22M Run4 MB,
 - ✓ $|VZ| < 30\text{cm}$
- Track quality cuts, global track
 - ✓ $n\text{FitsPts} > 25$, $n\text{FitsPts}/\text{Max} > 0.52$
 - ✓ $n\text{HitsdEdx} > 15$
 - ✓ $P_t > 0.20$, $|\text{eta}| < 1.0$
 - ✓ Pion n-sigma (-2.0, 2.0)

Secondary vertex finding technique





^3He & anti- ^3He selection



$$Z = \ln\left(\frac{\langle dE/dx \rangle}{\langle dE/dx \rangle_{\text{Bichsel}}}\right)$$

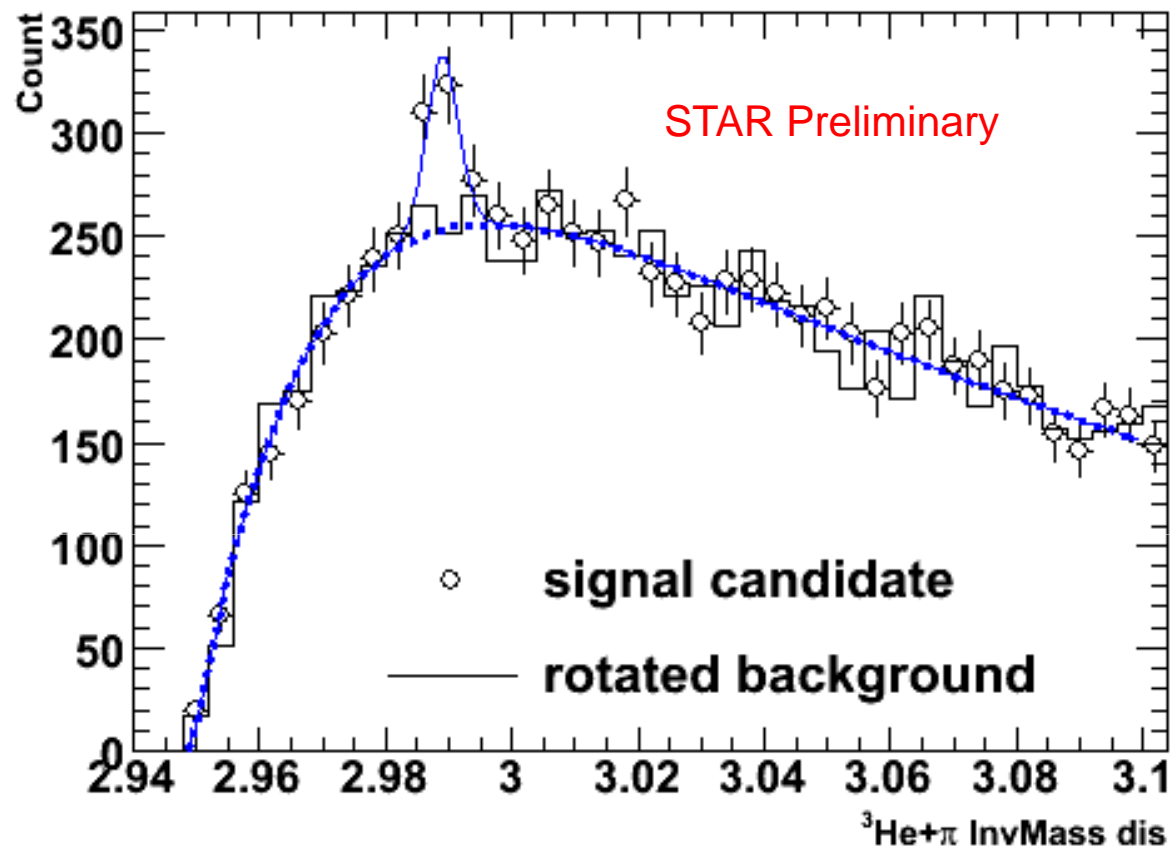
Select pure ^3He sample: $-0.2 < Z < 0.2$ & $dca < 1.0\text{cm}$ & $p > 2$ GeV

^3He : 2931(MB07) + 2008(central04) + 871(MB04) = 5810

Anti- ^3He : 1105(MB07) + 735(central04) + 328(MB04) = **2168**



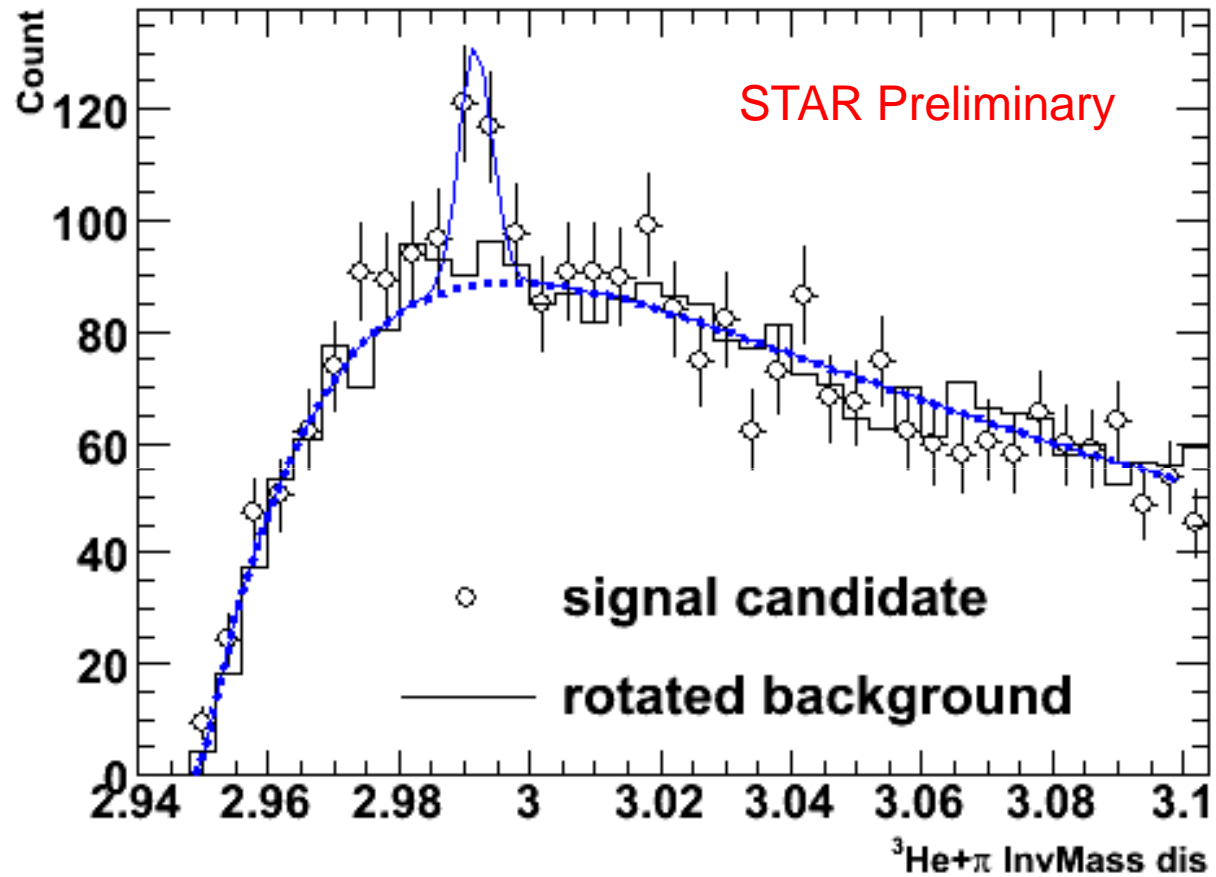
$\Lambda^3\text{H}$ signal from the data



- ◆ background shape determined from rotated background analysis;
- ◆ Signal observed from the data (bin-by-bin counting): 157 ± 30 ;
- ◆ Projection on antihypertriton yields: $\frac{^3\bar{\Lambda}\text{H}}{\Lambda^3\text{H}} = \frac{^3\bar{\text{H}} * ^3\bar{\text{He}}}{^3\text{He}} = 157 * 2168 / 5810 = 59 \pm 11$
constraint on antihypertriton yields without direct observation



$\bar{\Lambda}^3\bar{H}$ signal from the data

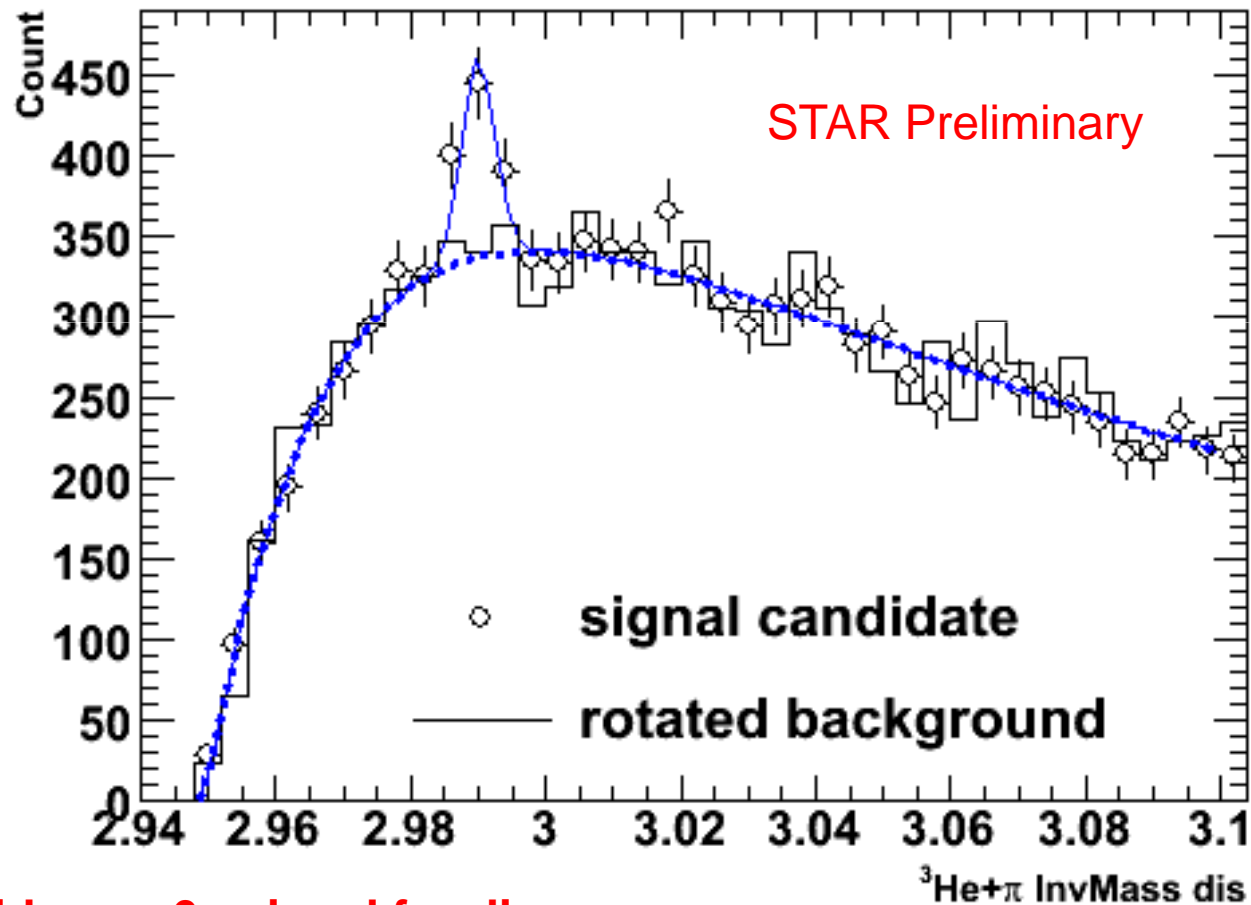


- ◆ Signal observed from the data (bin-by-bin counting): 70 ± 17 ;
Mass: 2.991 ± 0.001 GeV; Width (fixed): 0.0025 GeV;



Combined signals

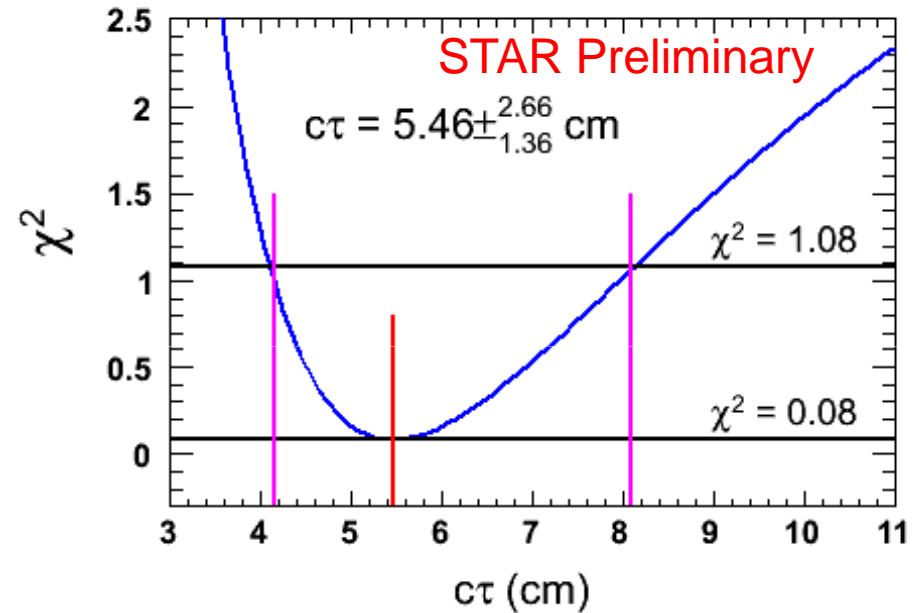
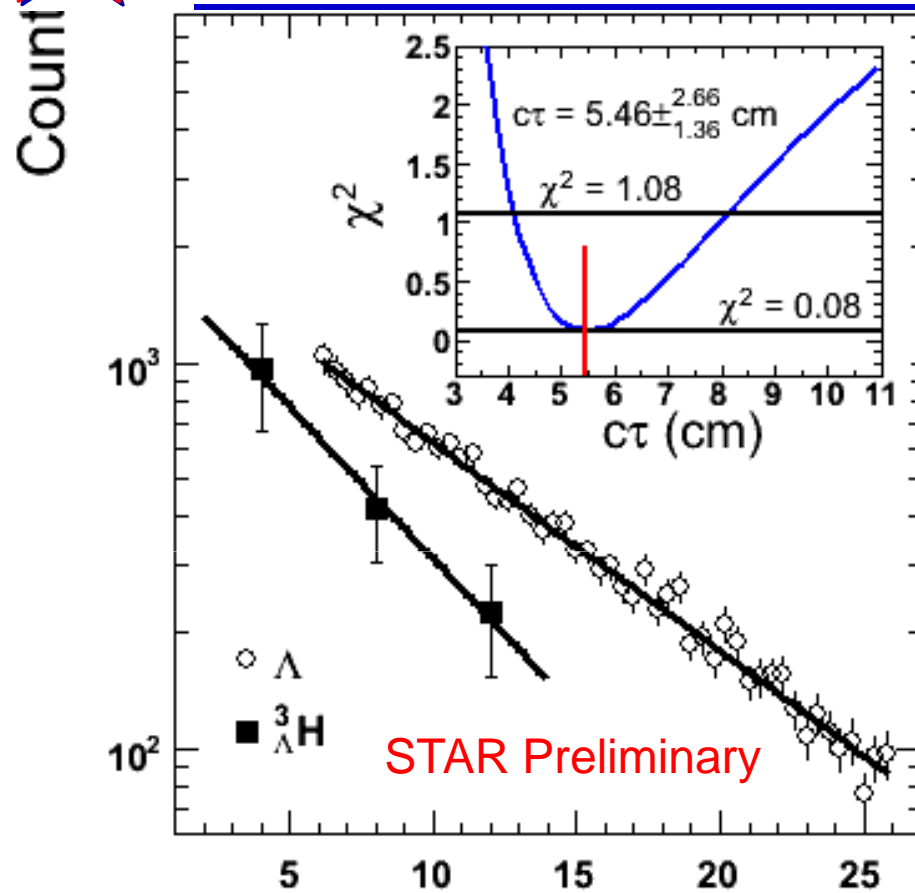
Combine hypertriton and antihypertriton signal:
 225 ± 35



This provides a $>6\sigma$ signal for discovery



Lifetime



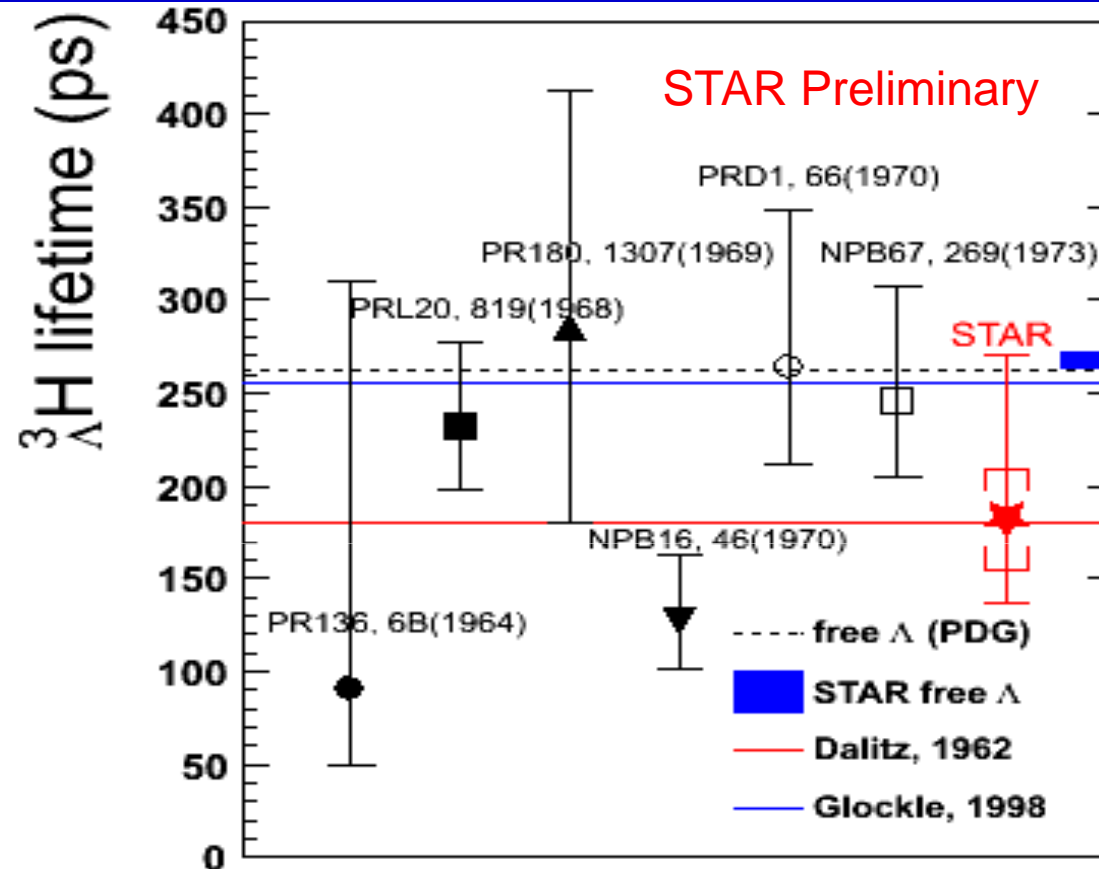
decay-length/ $(\beta\gamma)$ (cm)

◆ Our data: $\tau = 182 \pm_{45}^{89} \pm 27 ps$

◆ Consistency check on Λ lifetime yields $\tau(\Lambda) = 267 \pm 5 ps$ [PDG: 263 ps].



Comparison to world data



- ◆ Lifetime related to binding energy
- ◆ Theory input: the Λ is lightly bound in the hypertriton

[1] R. H. Dalitz, *Nuclear Interactions of the Hyperons* (Oxford Uni. Press, London, 1965).

[2] R.H. Dalitz and G. Rajasekharan, *Phys. Letts.* 1, 58 (1962).

[3] H. Kamada, W. Glockle at al., *Phys. Rev. C* 57, 1595(1998).



Measured invariant yields and ratios

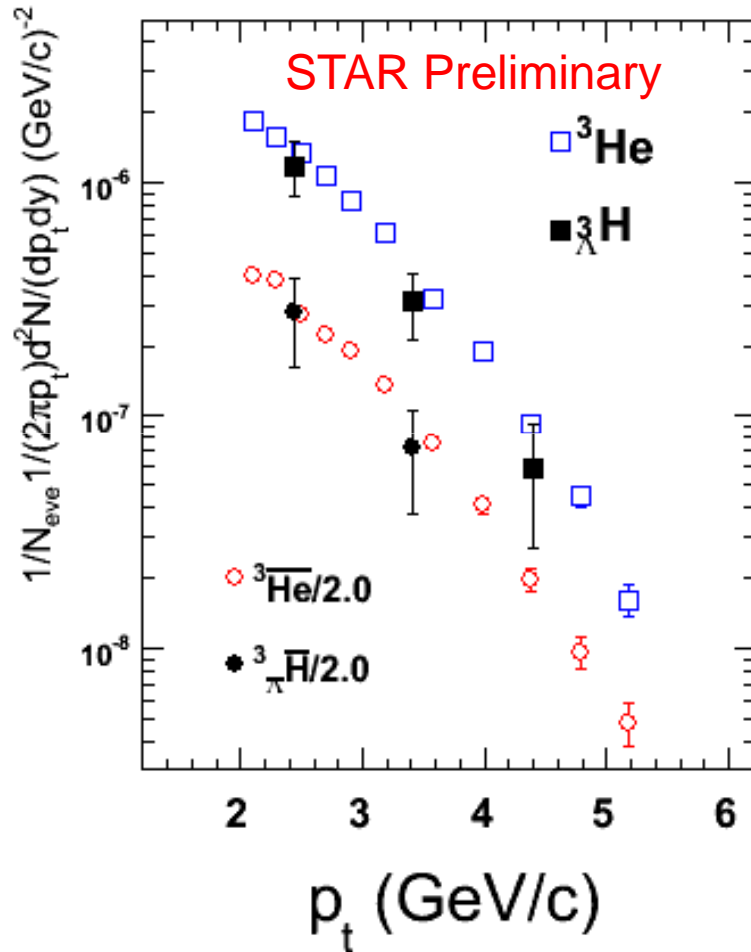


TABLE I: Particle ratios from Au+Au collisions at $\sqrt{s_{NN}} = 200 \text{ GeV}/c$. The ${}^3\text{He}$ (${}^3\bar{\text{He}}$) yield have been corrected for ${}^3_{\Lambda}\text{H}$ (${}^3_{\Lambda}\bar{\text{H}}$) feed-down contribution.

Particle type	Ratio
${}^3_{\Lambda}\bar{\text{H}}/{}^3_{\Lambda}\text{H}$	$0.49 \pm 0.18 \pm 0.07$
${}^3\bar{\text{He}}/{}^3\text{He}$	$0.45 \pm 0.02 \pm 0.04$
${}^3_{\Lambda}\bar{\text{H}}/{}^3\bar{\text{He}}$	$0.89 \pm 0.28 \pm 0.13$
${}^3_{\Lambda}\text{H}/{}^3\text{He}$	$0.82 \pm 0.16 \pm 0.12$

In a coalescence picture:

$${}^3_{\Lambda}\bar{\text{H}}/{}^3_{\Lambda}\text{H} \propto (\bar{p}/p)(\bar{n}/n)(\bar{\Lambda}/\Lambda)$$

$${}^3\bar{\text{He}}/{}^3\text{He} \propto (\bar{p}/p)^2(\bar{n}/n)$$

$$0.45 \sim (0.77)^3$$



Antinuclei in nature (new physics)

To appreciate just how rare nature produces antimatter (strange antimatter)

RHIC: an antimatter machine

Antimatter Galaxies

Where is all the antimatter in the universe? The Universe appears to have been created 13.7 billion years ago, in an instant. In the early Universe, particles of matter and antimatter were produced in equal numbers. As the Universe cooled, particles and antiparticles annihilated each other. Only a tiny fraction of matter survived. This is why we see only matter in the universe today. High-energy particle physics studies how particles behave when they collide. Particle-physicists can't make a proton without an antiproton, or an electron without an antielectron.

electron-positron pair, though, or electrons. There is no way to somehow "violate" baryon number, antimatter, or b) the Baryon number matter zone. While physics

Brookhaven National Laboratory News

Media Advisory: The Science of "Angels & Demons" Revealed

Could antimatter really destroy the Vatican? Brookhaven scientist, local educators discuss science myths and facts in movie

2009

During a free and public event, a physicist from the U.S. Department of Energy's Brookhaven National Laboratory and local educators will discuss the science facts from the science fiction movie "Angels & Demons," a major motion picture on Dan Brown's best-selling novel. The film, which opens nationally in theaters today, focuses on a plot to destroy the Vatican using antimatter from the Large Hadron Collider (LHC) at the European particle physics laboratory CERN. Brookhaven speakers will explain the real science of the LHC, including antimatter – oppositely charged cousins of ordinary matter with intriguing properties.

WHEN: Wednesday, May 27, 2009, 5:30 p.m.

WHERE: Berkner Hall Auditorium, Brookhaven National Laboratory. Brookhaven Lab is on William Floyd Parkway, one-and-a-half miles north of Exit 68 on the Long Island Expressway.

Seeing a mere antiproton or antielectron does not mean much— after all, these particles are byproducts of high-energy particle collisions. However, complex nuclei like **anti-helium** or anti-carbon are almost never created in collisions.

1. [\(introduction\)](#)
2. [antimatter galaxies](#)
3. [cold dark matter](#)
4. [strangelets](#)

examine these particles' tracks and see how they'd behave. A single really clear helium will bend to the right. A single really clear [\(back to top\)](#) - [\(back to AMS Tour\)](#)

Dark Matter, Black Hole → antinuclei production via coalescence



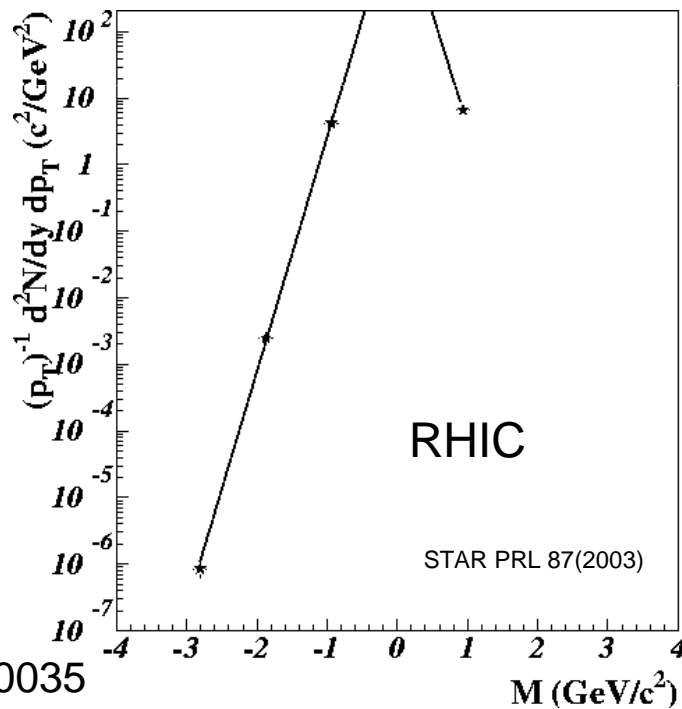
Matter and antimatter are not created equal

But we are getting there !

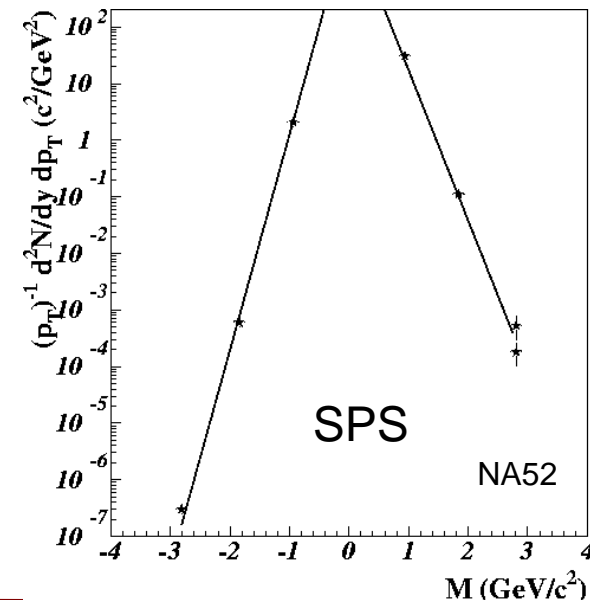
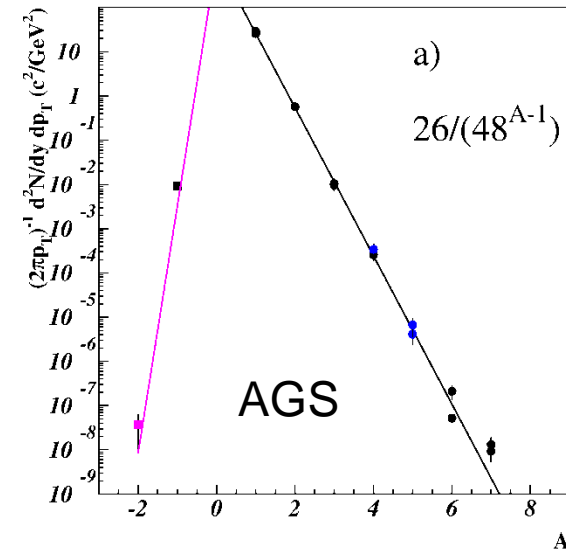
$${}^3\bar{He}/{}^3He \approx 10^{-11} \text{ (AGS, Cosmic)}$$

$${}^3\bar{He}/{}^3He \approx 10^{-3} \text{ (SPS)}$$

$${}^3\bar{He}/{}^3He \approx 0.5 \text{ (RHIC)}$$



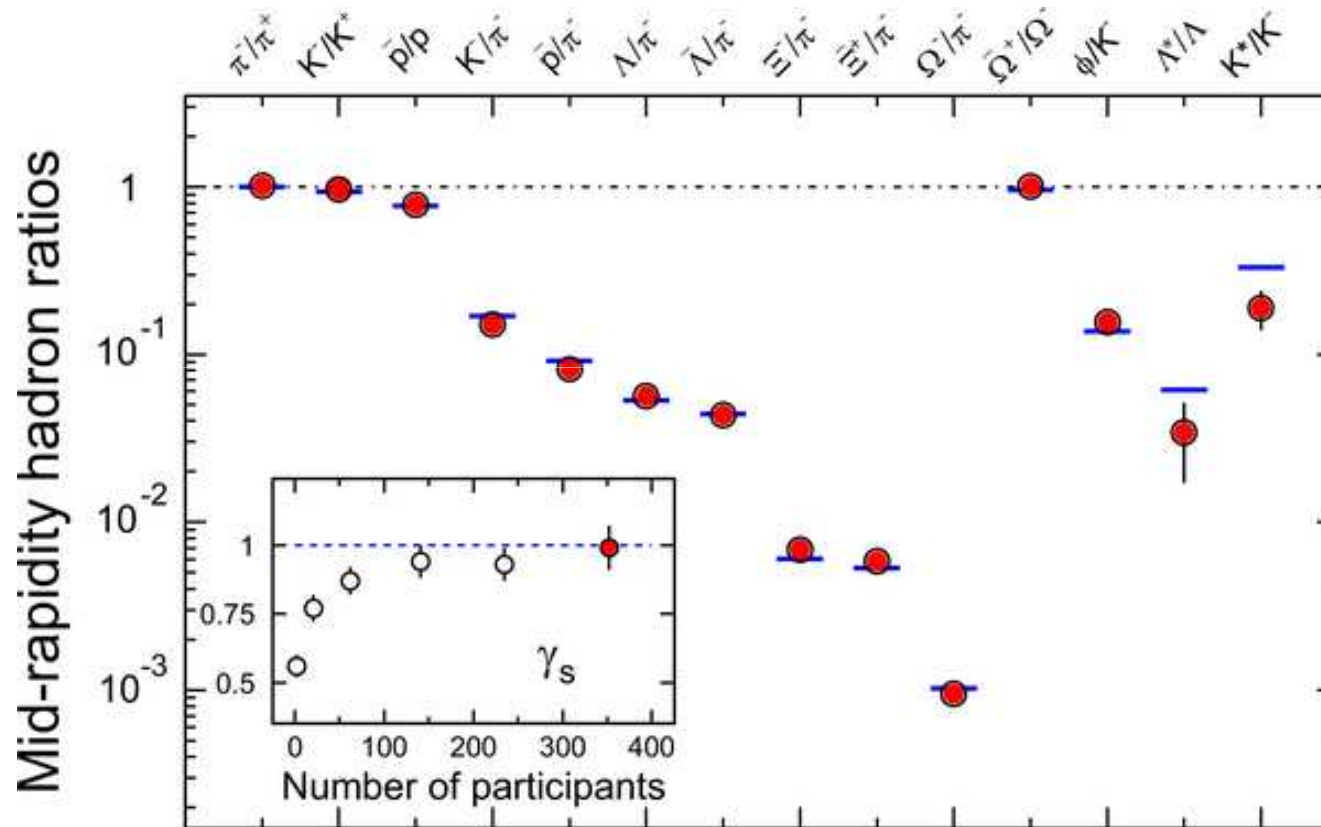
Nucl-ex/0610035





Flavors (u,d, s) are not created equal except in possible QGP

J. Rafelski and B. Muller, Phys.Rev.Lett.48:1066,1982



STAR whitepaper, NPA757(2005)

200 GeV $^{197}\text{Au} + ^{197}\text{Au}$ central collision



Yields as a measure of correlation

Volume 98B, number 3

PHYSICS LETTERS

8 January 1981

probability $F_d(\mathbf{p})$, which is determined by the deuteron internal wave function ψ_d and the spatial distribution functions D_p and D_n through eq. (9). If the k -dependences of $P_p(k)$ and $P_n(k)$ are weak compared with the p -dependence of $F_d(\mathbf{p})$, eq. (8) becomes equivalent to eq. (1) for $Z = N = 1$ with

$$\frac{4}{3} \pi p_0^3 = \int d\mathbf{p} F_d(\mathbf{p}). \quad (10)$$

Using eq. (9), one can express the coalescence volume in terms of ψ_d , D_p and D_n as

$$\frac{4}{3} \pi p_0^3 = 2^3 \cdot \frac{3}{4} \cdot (2\pi)^3 \int d\mathbf{r} |\psi_d(\mathbf{r})|^2 D_2(\mathbf{r}), \quad D_2(\mathbf{r}) \equiv \int d\mathbf{r}' D_p(\mathbf{r} - \mathbf{r}') D_n(\mathbf{r}'). \quad (11,12)$$

$D_2(\mathbf{r})$ gives the distribution of the p-n relative coordinate in the HX and is closely related to the interaction volume introduced by Mekjian [6]. In fact, if the spatial size of the internal wave function ψ_d is much smaller than that of the HX, then eq. (11) gives

$$\frac{4}{3} \pi p_0^3 \approx 2^3 \cdot \frac{3}{4} \cdot (2\pi)^3 D_2(0). \quad (13)$$

A=2 → Baryon density $\langle \rho_B \rangle$

$D_2(0)$ thus corresponds to the inverse of the interaction volume. In the actual situation, however, the size of the deuteron is comparable to that of the HX and therefore one has to use eq. (11) to relate the coalescence volume with the spatial size of the HX.

Expressions analogous to eqs. (8–10) can be obtained for the other composite particles such as ^3H , ^3He and ^4He . In the case of triton (^3H) one gets

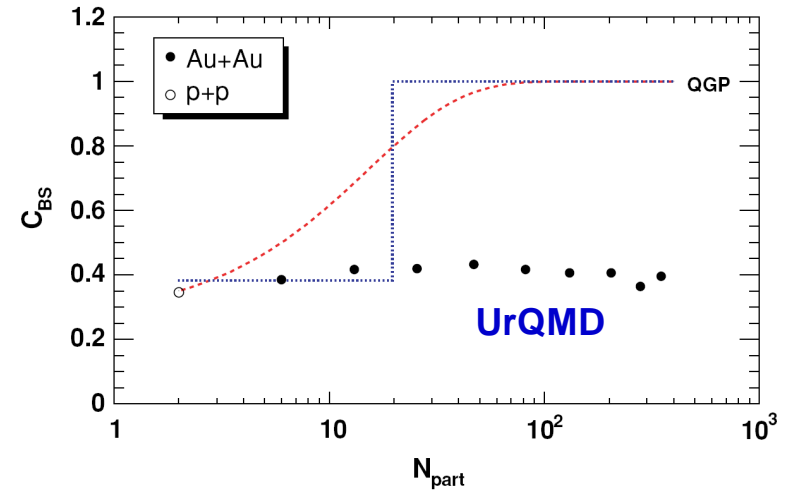
$$P(1, 2; \mathbf{k}) = \int d\mathbf{p}_1 d\mathbf{p}_2 F_t(\mathbf{p}_1, \mathbf{p}_2) P_p(\mathbf{k} + \mathbf{p}_1) P_n(\mathbf{k} - \frac{1}{2}\mathbf{p}_1 + \mathbf{p}_2) P_n(\mathbf{k} - \frac{1}{2}\mathbf{p}_1 - \mathbf{p}_2), \quad (14)$$

$$F_t(\mathbf{p}_1, \mathbf{p}_2) = 3^3 \cdot \frac{1}{4} \int \frac{d\mathbf{q}_1 d\mathbf{q}_2}{(2\pi)^6} \tilde{\psi}_t^*(\mathbf{p}_1 + \frac{1}{2}\mathbf{q}_1, \mathbf{p}_2 + \frac{1}{2}\mathbf{q}_2) \tilde{\psi}_t(\mathbf{p}_1 - \frac{1}{2}\mathbf{q}_1, \mathbf{p}_2 - \frac{1}{2}\mathbf{q}_2) \times \tilde{D}_p(\mathbf{q}_1) \tilde{D}_n(-\frac{1}{2}\mathbf{q}_1 + \mathbf{q}_2) \tilde{D}_n(-\frac{1}{2}\mathbf{q}_1 - \mathbf{q}_2), \quad (15)$$

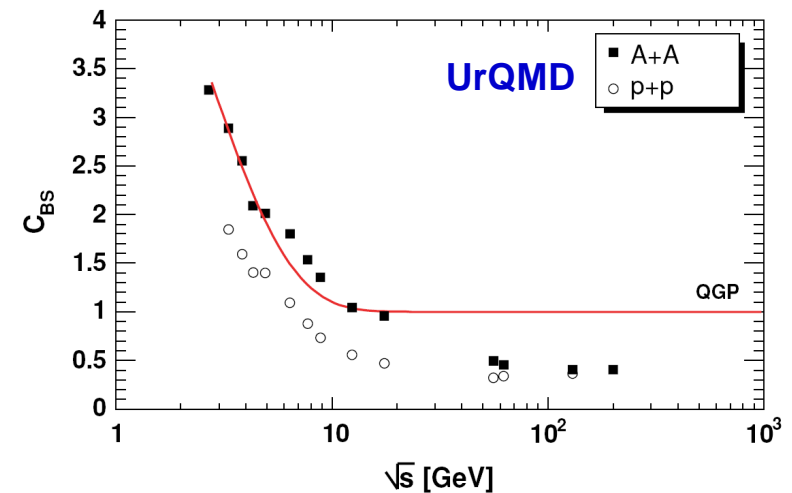
where $\tilde{\psi}_t$ is the Fourier transform of the triton internal wave function ψ_t . The coalescence volume is related to F_t as

$$\frac{1}{2} \left(\frac{4}{3} \pi p_0^3 \right)^2 = \int d\mathbf{p}_1 d\mathbf{p}_2 F_t(\mathbf{p}_1, \mathbf{p}_2). \quad (16)$$

A=3 → $\langle \rho_B^2 \rangle$; $\langle \rho_A \rho_B \rangle$



S. Haussler, H. Stoecker, M. Bleicher, PRC73



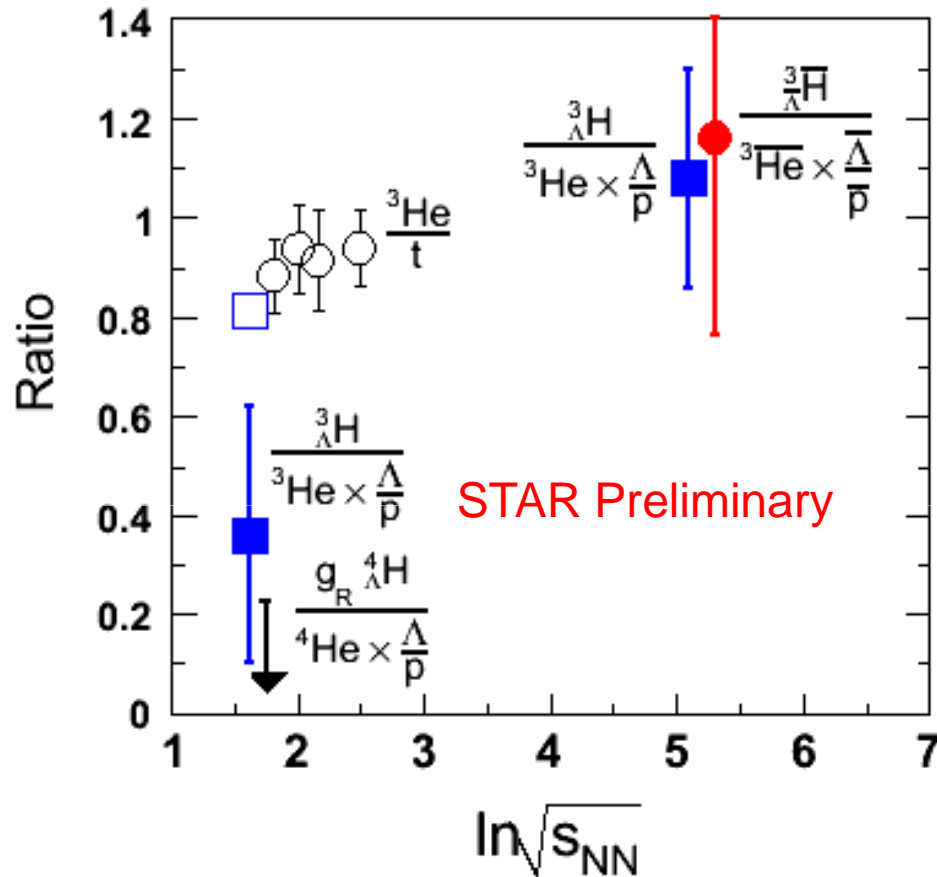
Caution:

measurements related to **local** (strangeness baryon)-baryon correlation

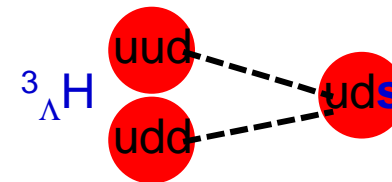
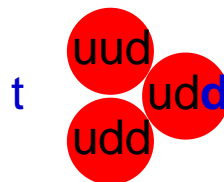
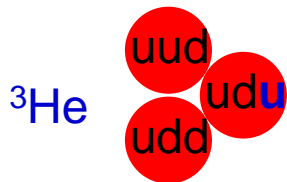
Simulations of (all strangeness)—(all baryon) correlation



$(^3\text{He}, t, ^3_\Lambda\text{H}) \rightarrow (u, d, s)$

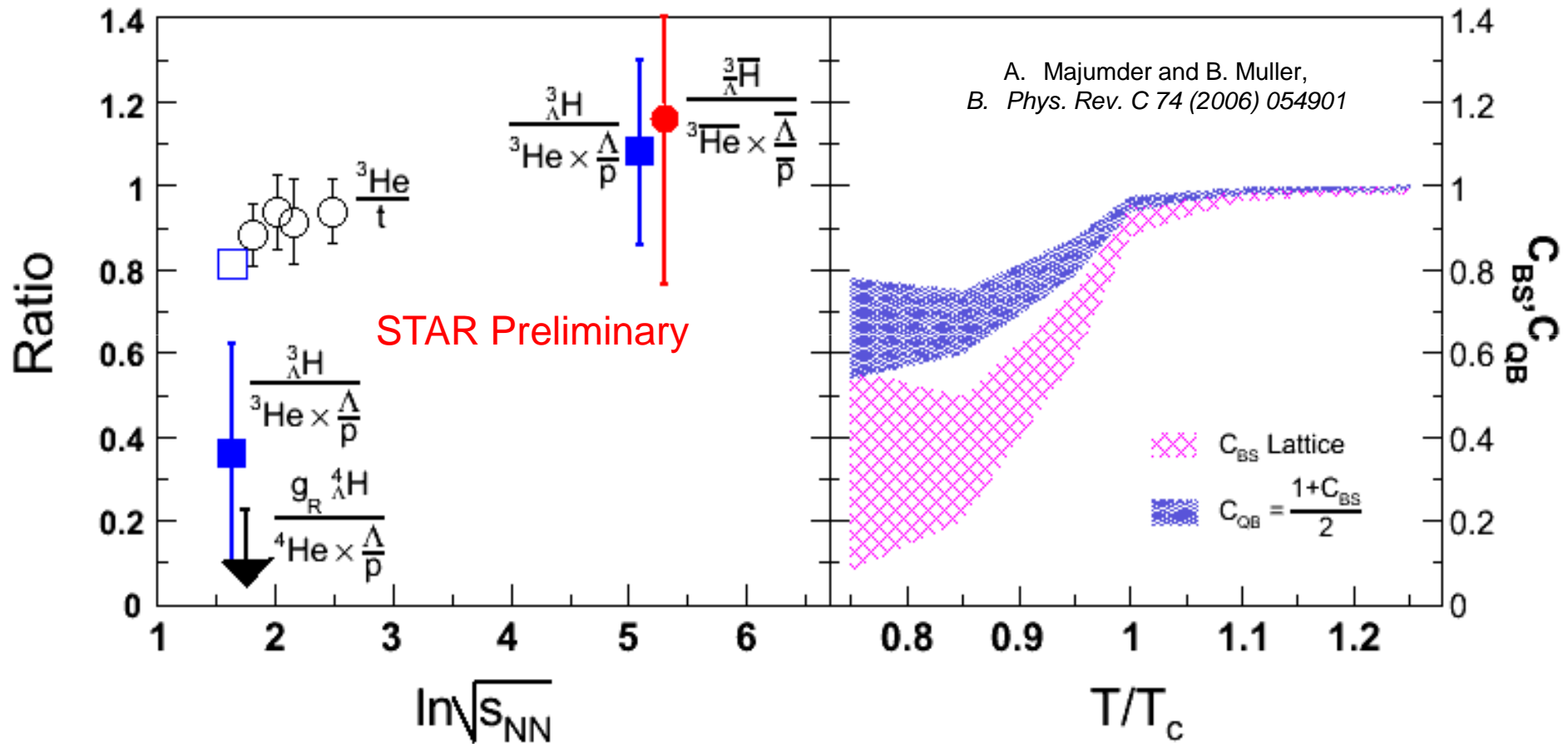


- $A=3$, a **simple** and **perfect** system
9 valence quarks,
 $(^3\text{He}, t, ^3_\Lambda\text{H}) \rightarrow (u, d, s) + 4u + 4d$
- Ratio measures **Lambda-nucleon** correlation
- RHIC: Lambda-nucleon similar phase space
- AGS: systematically lower than RHIC
- ➔ Strangeness phase-space equilibrium
- $^3\text{He}/t$ measures **charge-baryon** correlation





${}^3\Lambda\text{H}/{}^3\text{He}$: Primordial Λ -B correlation



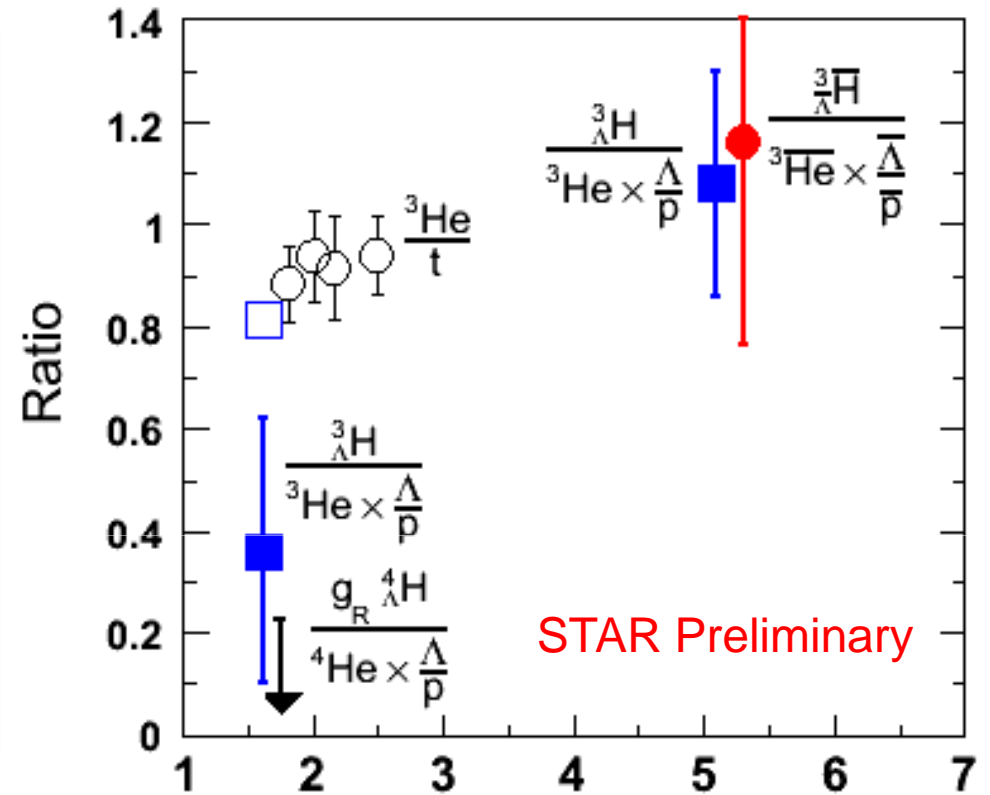
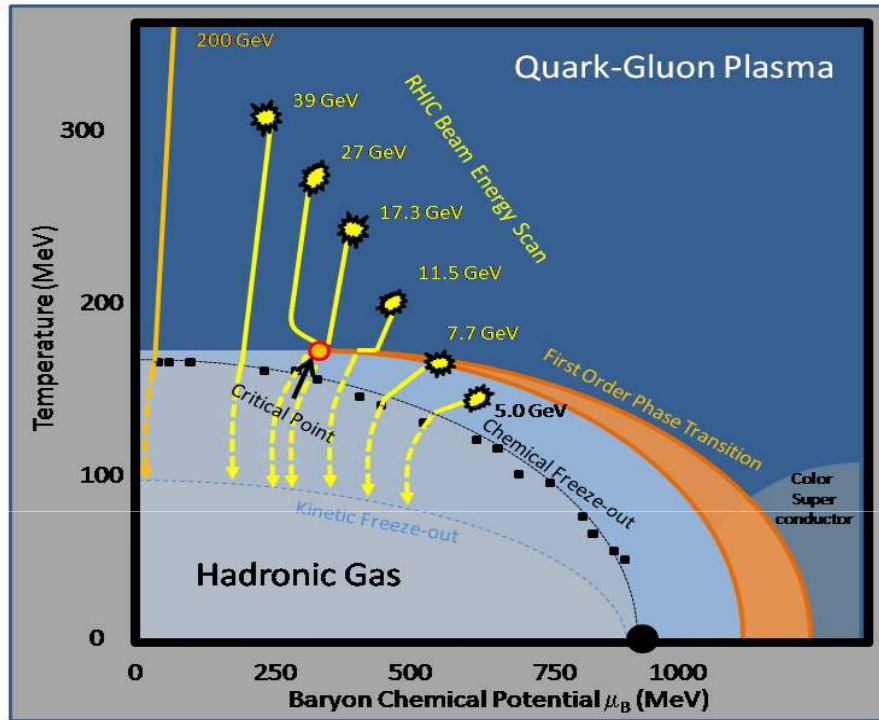
Caution:

measurements related to local (strangeness baryon)-baryon correlation

Lattice Simulations of (all strangeness)—(all baryon) correlation at zero chemical potential



Energy scan to establish the trend



Beam energy	200(30—200) GeV	~17 (10—30)GeV	~5 (5-10) GeV
Minbias events# (5σ)	300M	~10—100M	~1—10M
Penalty factor	1448	368	48
${}^3\text{He}$ invariant yields	1.6×10^{-6}	2×10^{-4}	0.01
${}^3_{\Lambda}\text{H}/{}^3\text{He}$ assumed	1.0	0.3	0.05

$\ln\sqrt{s_{NN}}$

Hypertriton only
STAR: DAQ1000+TOF



Conclusions

- ◆ $\frac{3}{\Lambda}\bar{\text{H}}$ has been **observed** for 1st time; significance $\sim 4\sigma$.
- ◆ Consistency **check** has been done on $\frac{3}{\Lambda}\text{H}$ analysis; significance is $\sim 5\sigma$
- ◆ The **lifetime** is measured to be $\tau = 182 \pm_{45}^{89} \pm 27 \text{ ps}$
- ◆ The $\frac{3}{\Lambda}\bar{\text{H}} / \frac{3}{\Lambda}\text{H}$ ratio is measured as 0.49 ± 0.18 , and $\frac{3}{\Lambda}\bar{\text{H}} / \frac{3}{\Lambda}\text{He}$ is 0.45 ± 0.02 , favoring the **coalescence** picture.
- ◆ The $\frac{3}{\Lambda}\bar{\text{H}} / \frac{3}{\Lambda}\text{He}$ ratio is determined to be 0.89 ± 0.28 , and $\frac{3}{\Lambda}\text{H} / \frac{3}{\Lambda}\text{He}$ is 0.82 ± 0.16 . No extra penalty factor observed for hypertritons at RHIC.
Strangeness phase space equilibrium



Outlook

◆ **Lifetime:**

- data samples with larger statistics

◆ **Production rate:**

- **Strangeness and baryon correlation**
Need specific model calculation for this quantity
- **Establish trend from AGS—SPS—RHIC—LHC**

◆ $\Lambda^3\text{H} \rightarrow d+p+\pi$ channel measurement: d and $d\bar{b}$ via ToF.

◆ Search for other hypernucleus: ${}^4_{\Lambda}\text{H}$, double Λ -hypernucleus.

◆ Search for anti- α

◆ **RHIC: best antimatter machine ever built**



International Hyper-nuclear network

