

Polarized Proton-Proton Elastic Scattering at RHIC

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For The STAR Collaboration

International Spin Physics Symposium 2008

Charlottesville, Virginia

October 9th, 2008

Outline

- Proton-Proton Elastic Scattering
- Very-Forward Detectors at STAR Detector at RHIC
- Spin-Dependent Proton-Proton Elastic Scattering
- Simulation Results on Detector Acceptance
- Summary

Proton-Proton Elastic Scattering at RHIC

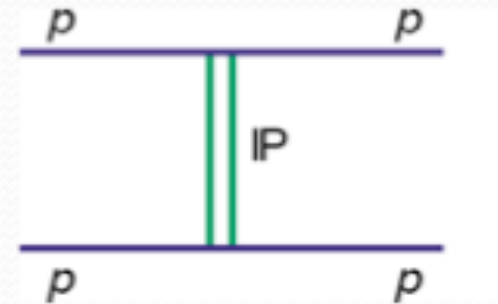
- RHIC provides polarized proton beams with a high average polarization of about **0.70** for each beam and a high luminosity reaching $L = 2 \cdot 10^{32} \text{ cm}^{-2} \text{ sec}^{-1}$ at a cms energy range of $50 \text{ GeV} \leq \sqrt{s} \leq 500 \text{ GeV}$.
- **PP₂PP Experiment at RHIC**
 - **Past**
 - Successful engineering run in **2002** (First Measurement of p - p Elastic Scattering at RHIC at $\sqrt{s} = 200 \text{ GeV}$)
 - Successful physics run in **2003** (Measurements of Spin Asymmetries at $\sqrt{s} = 200 \text{ GeV}$)
 - **Present – The PP₂PP setup has been integrated with STAR**
 - Study elastic scattering and central production by using the STAR detector.
 - During run **2008** used a dedicated beam time during last 2 hrs of the run for Roman Pot and DAQ commissioning.
 - Integration with the accelerator and STAR trigger was successful
 - **Future**
 - Ready to take data with STAR during run **2009** and beyond

Proton-Proton Elastic Scattering

- In elastic scattering protons remain intact.

$$p + p \rightarrow p + p$$

- A *Pomeron* (IP) is exchanged.
- *Pomeron* exchange in PQCD consists of a color singlet combination of gluons (two gluons).



Detect protons in very forward direction by using Roman Pots

Differential pp Elastic Cross Section

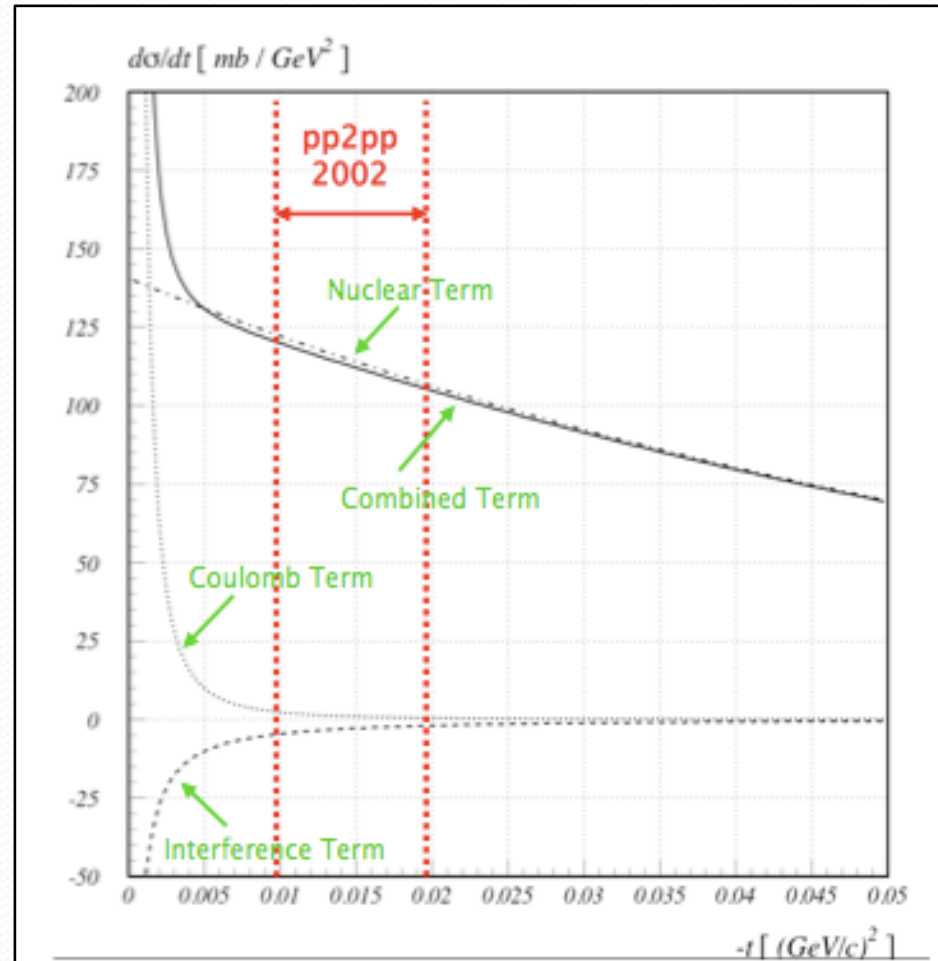
$$\frac{d\sigma_{el}}{dt} = 4\pi(\hbar c)^2 \left(\frac{\alpha G_E^2}{t} \right)^2$$

$$+ \frac{1 + \rho^2}{16\pi(\hbar c)^2} \sigma_{tot}^2 e^{-b|t|}$$

$$- (\rho + \Delta\Phi) \frac{\alpha G_E^2}{|t|} \sigma_{tot} e^{-\frac{1}{2}b|t|}$$

1. Coulomb Term
2. Hadronic Term
3. Coulomb Nuclear Interference Term

$d\sigma/dt$ (mb/GeV²)



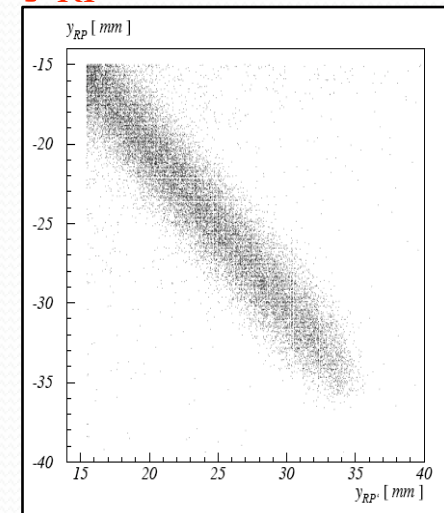
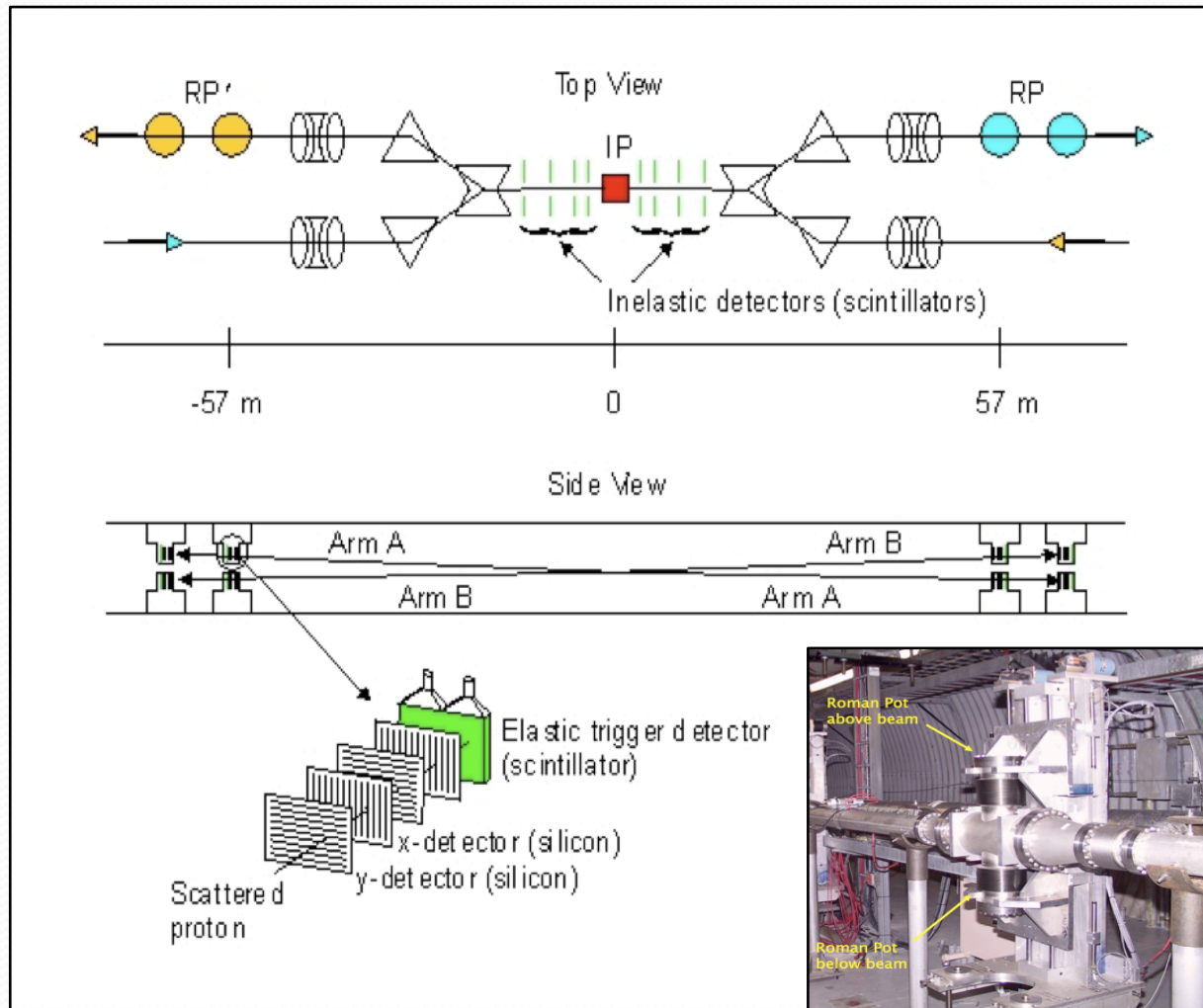
$-t$ (GeV²/c²)

Experimental Layout

RP'

RP

$y_{RP'}$ (mm)



y_{RP} (mm)

Collinearity
condition
for elastic
events

Measurement Technique

- Very forward detectors detect protons that scatter elastically at very small angles.
- Trajectories of scattered particles are determined by the beam transport equations: (* = at IP)

$$y = a_{11}y^* + L_{eff}\theta_y^*$$
$$\theta_y = a_{12}y^* + a_{22}\theta_y^*$$

- a_{11} , L_{eff} , a_{12} , a_{22} are the beam transport matrix elements
- Similar equations in x -coordinate
- Choose a_{11} to be small and L_{eff} very large, thus:

$$y \sim L_{eff}\theta_y^*$$

Spin Dependent Elastic Scattering

- ◆ The scattering amplitude with **initial** $|h_1 h_2\rangle$ and **final helicity** $|h_3 h_4\rangle$ **states** of the interacting protons (1 & 2) with momenta $p_1 = -p_2$ and scattered protons (3 & 4) with momenta $p_3 = -p_4$ is of the form:

$$\phi_n(s, t) = \langle h_1 h_2 | M | h_3 h_4 \rangle = \phi_n^{em}(s, t) + \phi_n^{had}(s, t)$$

- ◆ The **helicity amplitudes** that describe **pp elastic scattering**:

$$\phi_1(s, t) = \langle ++ | M | ++ \rangle$$

$$\phi_2(s, t) = \langle ++ | M | -- \rangle$$

$$\phi_3(s, t) = \langle +- | M | +- \rangle$$

$$\phi_4(s, t) = \langle +- | M | -+ \rangle$$

$$\phi_5(s, t) = \langle ++ | M | +- \rangle$$

N.H. Buttimore et al, Spin dependence of high energy proton scattering, Phys. Rev. D, Volume 59, 114010 (1999)

- ◆ The **differential elastic cross section** and the **total cross section** are given as:

$$\frac{d\sigma_{el}}{dt} = \frac{2\pi}{s^2} \left[|\phi_1|^2 + |\phi_2|^2 + |\phi_3|^2 + |\phi_4|^2 + 4|\phi_5|^2 \right]$$

$$\sigma_{tot} = \frac{4\pi}{s} \text{Im} \left[\phi_1(s, t) + \phi_3(s, t) \right]_{t=0}$$

Single Spin Asymmetry

- The interference between the hadronic non-flip and electromagnetic spin-flip amplitudes gives rise to the single spin asymmetry in the CNI region.

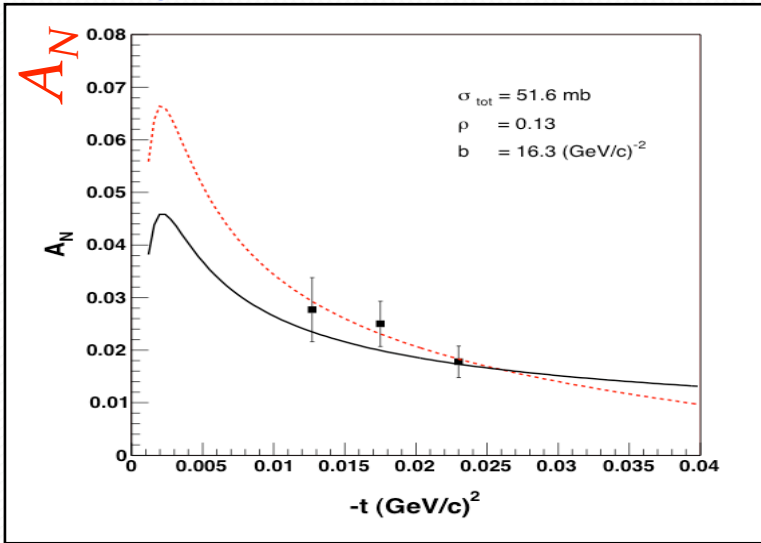
$$A_N(t) = \frac{1}{P \cos \phi_{azimuth}} \frac{N_{\uparrow\uparrow}(t) - N_{\downarrow\downarrow}(t) + N_{\uparrow\downarrow}(t) - N_{\downarrow\uparrow}(t)}{N_{\uparrow\uparrow}(t) + N_{\downarrow\downarrow}(t) + N_{\uparrow\downarrow} + N_{\downarrow\uparrow}} \quad N(t) = dN/dt$$

$$A_N(t) \approx \frac{\text{Im}\left(\phi_{flip}^{em*} \phi_{non-flip}^{had} + \phi_{flip}^{had*} \phi_{non-flip}^{em}\right)}{d\sigma/dt} \quad P: \text{ beam polarization}$$

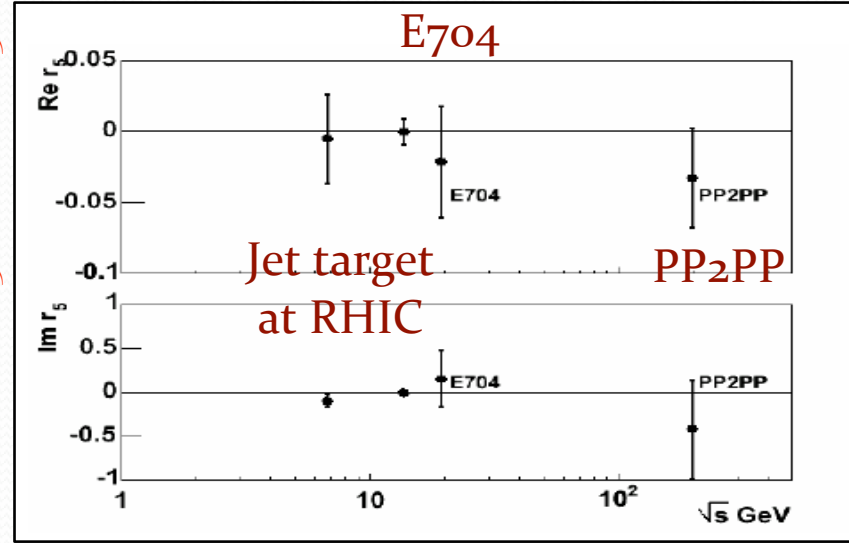
- The electromagnetic single-flip amplitude is real and calculable. Hence A_N can be calculated and compared to data.
- Study A_N and its $|t|$ -**dependence** to probe the interference term between hadronic and Coulomb interactions.

Analyzing Power Measurement 2003

S. Bueltmann et al, "First Measurement of A_N at $\sqrt{s} = 200$ GeV in Polarized Proton-Proton Elastic Scattering at RHIC", Phys. Lett. B 632, 167-172 (2006)



$Im r_5$ $Re r_5$



$-t$ (GeV^2/c^2)

- $\sqrt{s} = 200$ GeV
- 0.01 (GeV/c)² $\leq |t| \leq 0.03$ (GeV/c)²
- solid curve: calculation without hadronic spin-flip amplitude

$$A_N = 0.021 \pm 0.0023 \text{ (stat.)}$$

$$\pm 0.0018 \text{ (sys.)} \pm 0.0036 \text{ (beam)}$$

$$r_5 = \frac{m_p}{\sqrt{-t}} \frac{\phi_5^{had}}{\text{Im}(\phi_1^{had} + \phi_3^{had})/2}$$

- A possible hadronic single spin-flip amplitude would alter A_N and its effect would depend on the ratio of the single spin-flip amplitude ϕ_5 to non-flip amplitudes ϕ_1 and ϕ_3 .

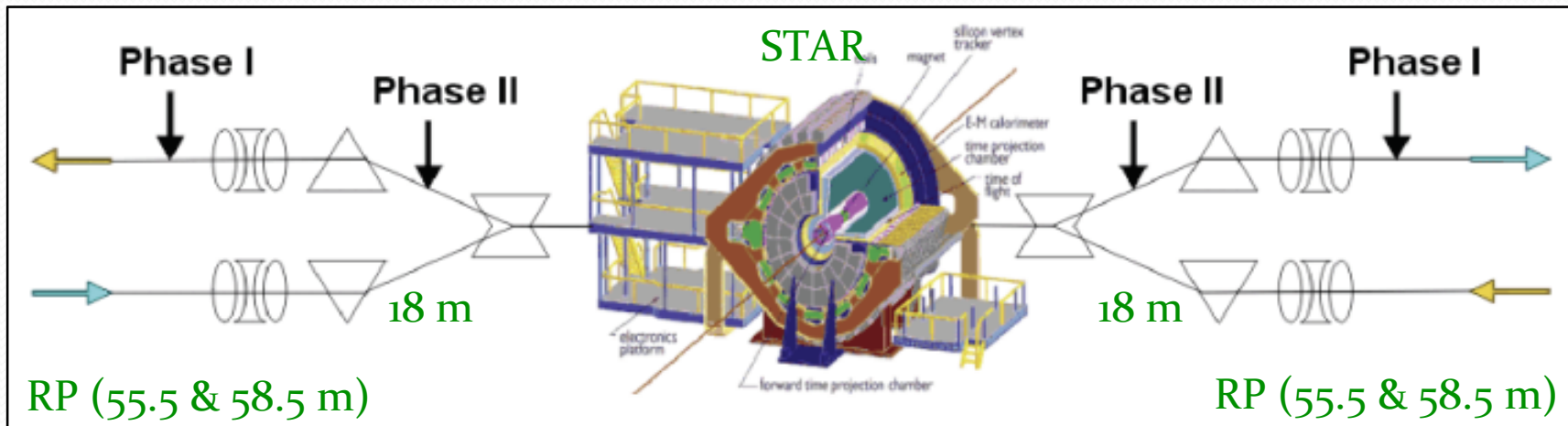
Goals for the Future

Expected uncertainties for elastic scattering (for three-four days of data taking and luminosity = $2 \cdot 10^{29} \text{ cm}^{-2} \text{ sec}^{-1}$):

1. In the small $|t|$, Coulomb-Nuclear Interference (CNI) region measure:
 - σ_{tot} and $d\sigma_{el}/dt$ and study their dependence on \sqrt{s}
 - Nuclear slope parameter b
 - Ratio of real and imaginary part of forward elastic scattering amplitude ρ
 - The analyzing power A_N
 - The transverse double spin asymmetries A_{NN} and A_{SS}
2. In the intermediate $|t|$ region, $|t| \leq 1.5 \text{ (GeV/c)}^2$
 - Study the correlation between the position of the dip and the single spin asymmetry A_N crossing zero at the same $|t|$ value

Uncertainty	Previous	Expected
$\Delta\sigma_{tot}$	-	3 mb
Δb	1.6 (GeV/c) ⁻²	0.3 (GeV/c) ⁻²
$\Delta\rho$	-	0.01
ΔA_N	0.0023	0.0017
ΔA_{NN}	0.017	0.0053
ΔA_{SS}	0.008	0.0053

Plan of the Experiment



Phase I (existing Roman Pots around 58 m)

- Requires special data taking run with **special proton beam tune** to reach small- t values.
- Run at $\sqrt{s} = 200$ GeV with $\beta^* = 20$ m and at $\sqrt{s} = 500$ GeV with $\beta^* = 20$ m

Phase II (additional Roman Pots at 18 m)

- Install additional Roman Pots between dipole magnets (18 m)
- Will extend the kinematic $|t|$ -range. For example at $\sqrt{s} = 500$ GeV to $-t < 1.5$ (GeV/c)²
- **No special running conditions** required (use standard proton beam tune $\beta^* \leq 1$ m)
- Beam pipe between dipole magnets needs to be rebuilt.

Phase I Acceptance (55.5 m and 58.5 m)

Coordinates of Scattered Protons (y versus x in mm)

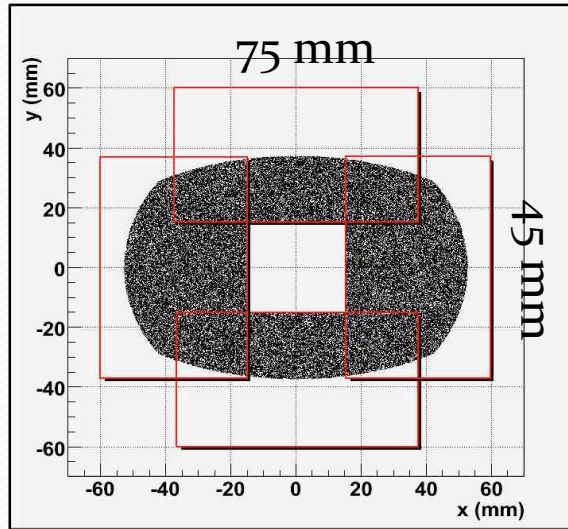
$\beta^* = 21$ m
Special running conditions

$p = 100$ GeV/c

$d_{\min} = 10$ mm $\approx 12 \sigma_{\text{beam}}$

100% acceptance for elastic scattering

0.003 (GeV/c) $^2 \leq |t| \leq 0.022$ (GeV/c) 2

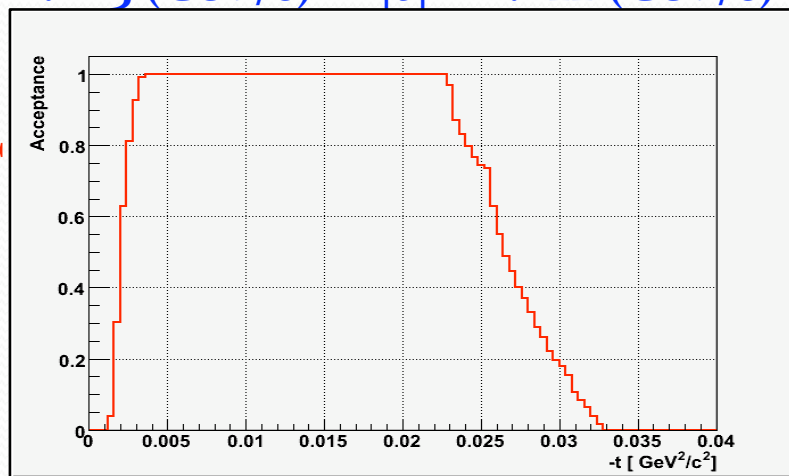


J. de Favereau, X. Rouby, K. Piotrkowski, Hector, a fast simulator for the transport of particles in beamlines, arXiv: 0707.1198 [physics.acc-ph] (2007)

$p = 250$ GeV/c

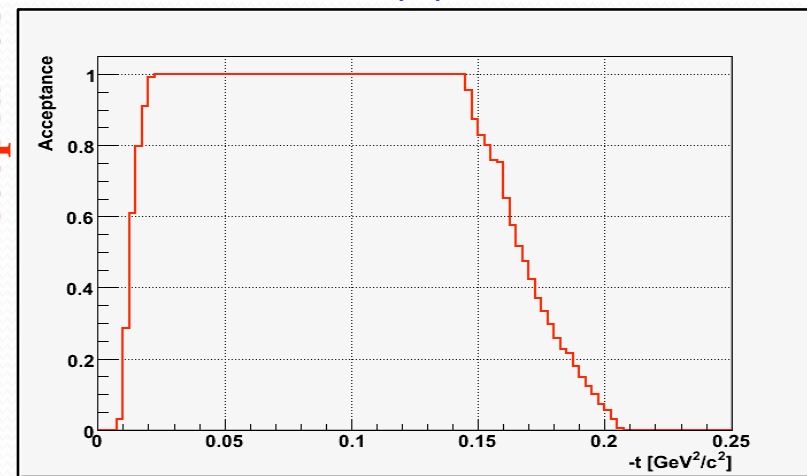
0.02 (GeV/c) $^2 \leq |t| \leq 0.145$ (GeV/c) 2

Acceptance



$|t|$ (GeV $^2/c^2$)

Acceptance



$|t|$ (GeV $^2/c^2$)

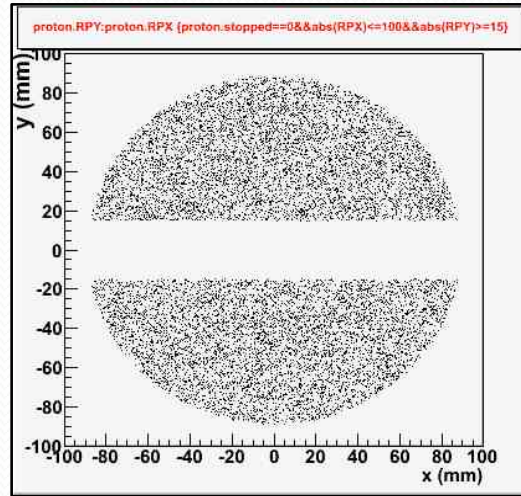
Phase II Acceptance (18 m)

Coordinates of Scattered Protons (y versus x in mm)

Can run under
normal running
conditions

$$\theta^* \leq 1 \text{ m}$$

$$p = 100 \text{ GeV}/c$$
$$-t \leq 0.246 \text{ (GeV}/c^2)^2$$



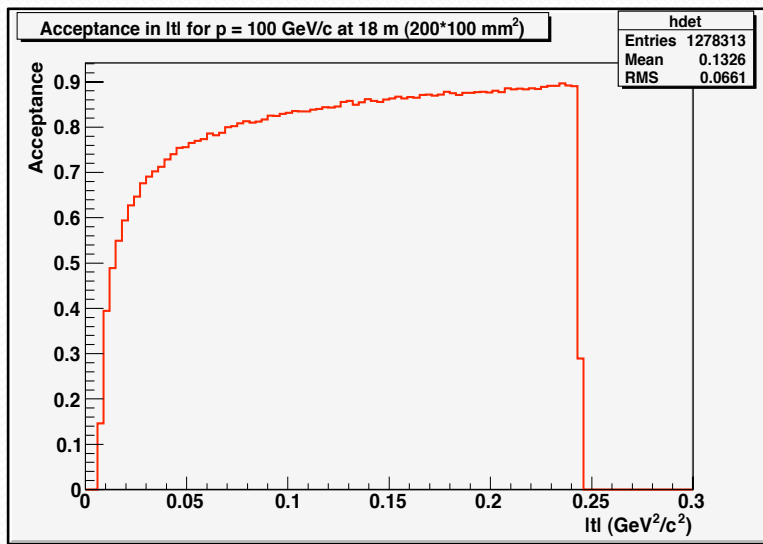
Area of Detectors
(one RP station)

$$200^*100 \text{ mm}^2$$

$$d_{\min} \approx 12 \sigma_{\text{beam}}$$

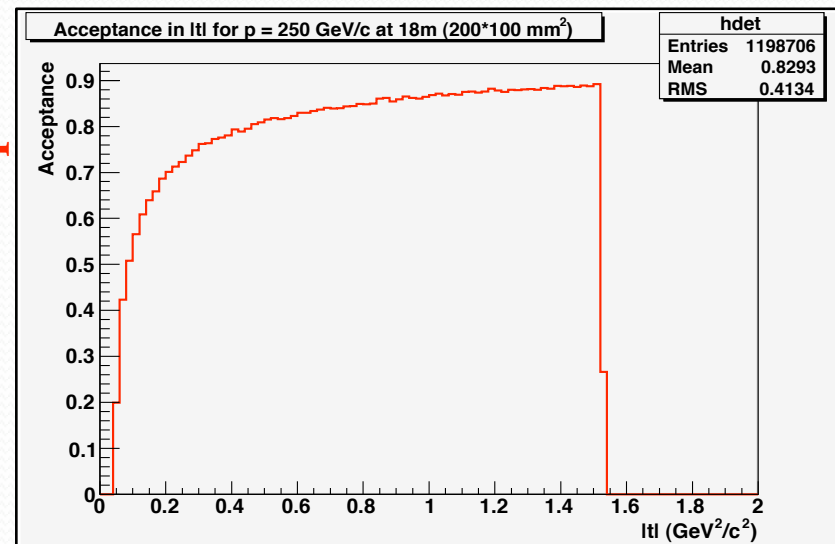
$$p = 250 \text{ GeV}/c$$
$$-t \leq 1.5 \text{ (GeV}/c^2)^2$$

Acceptance



$|t| \text{ (GeV}^2/c^2)$

Acceptance



$|t| \text{ (GeV}^2/c^2)$

Summary

- **Phase I (RP region)**

- Study **elastic** scattering by using only the existing Roman Pots on each side of STAR.
- $0.003 \text{ (GeV/c)}^2 \leq |t| \leq 0.022 \text{ (GeV/c)}^2$ for $\sqrt{s} = 200 \text{ GeV}$
- $0.02 \text{ (GeV/c)}^2 \leq |t| \leq 0.145 \text{ (GeV/c)}^2$ for $\sqrt{s} = 500 \text{ GeV}$
- Requires special data taking run with **special proton beam tune**.
- Study **inelastic** scattering with STAR as a central detector. STAR has a good acceptance and particle ID to measure the central system.

- **Phase II (18 m)**

- Install additional Roman Pots in the region between dipole magnets (18 m)
- **Will extend the kinematic $|t|$ range at:**
 - $\sqrt{s} = 200 \text{ GeV}$ to $-t \leq 0.246 \text{ (GeV/c)}^2$
 - $\sqrt{s} = 500 \text{ GeV}$ to $-t \leq 1.5 \text{ (GeV/c)}^2$
- **No special running conditions** required.
- Beam pipe between dipole magnets needs to be rebuilt.



Thank You

RHIC at BNL

