# NCQ scaling of $f_0(980)$ elliptic flow in 200 GeV Au+Au collisions by STAR and its constituent quark content

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**Abstract.** Searching for exotic state particles and studying their properties have furthered our understanding of quantum chromodynamics (QCD). The  $f_0(980)$  resonance is an exotic state with relatively high production rate in relativistic heavy-ion collisions, decaying primarily into  $\pi\pi$ . Currently the structure and quark content of the  $f_0(980)$  are unknown with several predictions from theory being a  $q\overline{q}$  state, a  $qq\overline{q}q$  state, a  $K\overline{K}$  molecule state, or a gluonium state. We report the first  $f_0(980)$  elliptic flow  $(v_2)$  measurement from 200 GeV Au+Au collisions at STAR. The transverse momentum dependence of  $v_2$  is examined and compared to those of other hadrons (baryons and mesons). The empirical number of constituent quark (NCQ) scaling is used to investigate the constituent quark content of  $f_0(980)$ , which may potentially address an important question in QCD.

#### 1 Introduction

- <sup>2</sup> Searching for exotic state particles and studying their properties have furthered our understanding of
- quantum chromodynamics (QCD). Currently the structure and quark content of  $f_0(980)$  are unknown
- with several predictions being a  $q\overline{q}$  state, a  $qq\overline{qq}$  state, a  $K\overline{K}$  molecule state, or a gluonium state [1–6].
- 5 In contrast to the vector and tensor mesons, the identification of the scalar mesons is a long-standing
- <sub>6</sub> puzzle [7]. Previous preliminary experimental measurements [8] on the yield of  $f_0(980)$  at RHIC and
- theoretical calculation [9] suggest that it could be a  $K\overline{K}$  stat. In this analysis, the empirical number
- 8 of constituent quark (NCQ) scaling [10-12] is used to investigate the constituent quark content of
- $f_0(980)$  [13].

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## 2 Experiment setup and data analysis

The data reported here are from Au+Au collisions at a nucleon-nucleon center-of-mass energy of 200 GeV, collected by the STAR experiment [14] at Brookhaven National Laboratory in 2011, 2014 and 2016. A total of 2.4 billion minimum-bias (MB) events are selected for this analysis. The main subsystem used for the data analysis is the Time Projection Chamber (TPC) [15] with  $2\pi$  azimuthal coverage at mid-rapidity. The TPC dE/dx is used to select  $\pi^{\pm}$  candidate with  $0.2 < p_T < 5.0 \text{ GeV/}c$ .

The  $\pi^+\pi^-$  are used to reconstruct the  $f_0(980)$ . The combinatorial background subtraction is based on the mixed-event technique and the like-sign method [16]. The acceptance-corrected like-sign

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pairs [16, 17] are used to subtract the combinatorial background after being normalized to unlikesign pairs in the invariant mass  $(m_{inv})$  range beyond 1.5 GeV/ $c^2$ . Figure 1 (left) shows the background subtracted  $\pi^+\pi^-$  invariant mass distribution. The resonance peaks are parametrized with the relativistic Breit-Wigner function [18, 19]. The total fit function is given by:

$$f(m_{inv}) = \left(\sum_{X = f_0, \rho^0, f_2} \frac{A_X m_{inv} m_X \Gamma(X)}{(m_{inv}^2 - m_X^2)^2 + m_X^2 \Gamma(X)^2}\right) \times PS + bg(m_{inv})$$
(1)

where  $\Gamma(X) = \frac{\Gamma_X m_X}{m_{inv}} \left( \frac{m_{inv}^2 - 4m_\pi^2}{m_{inv}^2 - 4m_\pi^2} \right)^{J+1/2}$  [18, 19],  $PS = \frac{m_{inv}}{\sqrt{m_{inv}^2 + p_T^2}} \exp\left(-\frac{\sqrt{m_{inv}^2 + p_T^2}}{T}\right)$  is the phase space correction taking into account the  $\pi\pi$  scattering during the hadronic phase [19–22], and  $bg(m_{inv})$  is a third order polynomial function to describe the residual background.  $m_X$  and  $\Gamma_X$  are the mass and width of the corresponding resonances.  $\Gamma_{\rho^0}$  is set to 160 MeV, and  $m_{f_2}$  and  $\Gamma_{f_2}$  are set according to the PDG values [7]. T is the kinetic freeze-out temperature, set to 120 MeV [20].  $A_{f_0}$ ,  $A_{\rho^0}$ ,  $A_{f_2}$ ,  $m_{f_0}$ ,  $\Gamma_{f_0}$ , and  $m_{\rho^0}$  are free parameters.

The event-plane method [23] is used to study the elliptic flow  $(v_2)$  of  $f_0(980)$ . The event-plane is reconstructed by all charged particles in the TPC with pseudorapidity  $|\eta| < 1$  and transverse momentum  $0.2 < p_T < 5.0 \text{ GeV/}c$ . For each  $\pi\pi$  pair, the two  $\pi$  candidates are removed from the event-plane reconstruction to avoid auto-correlation. The event-plane resolution is calculated by the correlation between two randomly divided sub-events from the full TPC [23]. Wide centrality bin effect is corrected by weighting the event-plane resolution with the  $f_0(980)$  yield in each narrow centrality bin of 10% size [24]. Figure 1(right) shows the  $f_0(980)$  yield as function of the azimuthal angle difference between the  $\pi\pi$  pair  $(\phi)$  and the event-plane direction  $(\Psi)$  in an example  $p_T$  bin.

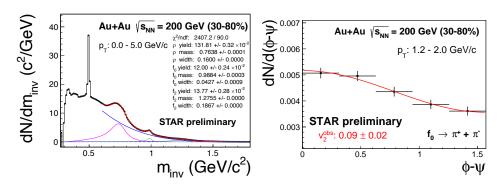
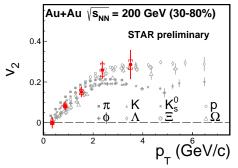


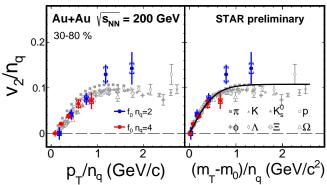
Figure 1. (Color online) (Left) The background subtracted  $\pi^+\pi^-$  invariant mass distribution over the  $p_T$  range of  $0 < p_T < 5.0$  GeV/c in 30–80% Au+Au collisions at  $\sqrt{s_{NN}}$ =200 GeV. The red line is the result of fit. The pink, green, violet lines represent the resonance peaks of the relativistic Breit-Wigner function. The solid blue line represents the residual background using a third order polynomial function. (Right)  $f_0(980)$  yield as function of  $\phi - \Psi$  in a given  $p_T$  bin. Errors are statistical. The red line represents a fit ( $\propto (1 + 2v_2^{obs}\cos(2\phi - 2\Psi))$ ) to the data.

Figure 2 shows  $f_0(980)$   $v_2$  as a function of  $p_T$  in 30-80% centrality Au+Au collisions at  $\sqrt{s_{NN}} = 200$  GeV. Results are compared with other identified particles:  $\pi$ , K, p,  $K_s^0$ ,  $\Lambda$ ,  $\Xi$ ,  $\Omega$ ,  $\phi$  [24]. In the low  $p_T$  region, the  $f_0(980)$   $v_2$  seems to follow the mass ordering. In the higher  $p_T$  region, the  $f_0(980)$   $v_2$  seems closer to the baryon band.

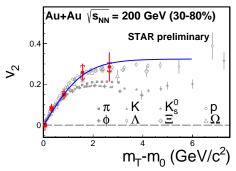
Figure 3 shows the number of constituent quark  $(n_q)$  scaled  $v_2$  as function of the  $n_q$  scaled  $p_T$  (left) and  $m_T - m_0$  (right). Here the  $f_0(980)$  is assumed to have either 2 quarks or 4 quarks. The data are



**Figure 2.** (Color online)  $f_0(980) v_2$  as a function of  $p_T$  in 30-80% centrality Au+Au collisions at  $\sqrt{s_{NN}}$  = 200 GeV. Statistical uncertainties are shown by the vertical bars and systematic uncertainties are shown by the caps. Results of other particles are taken from Ref. [24].



**Figure 3.** (Color online)  $f_0(980)$   $v_2$  divided by  $n_q$  as a function of  $p_T/n_q$  (left) and  $(m_T - m_0)/n_q$  (right) in 30-80% centrality Au+Au collisions at  $\sqrt{s_{NN}} = 200$  GeV. Results of other particles are taken from Ref. [24]. Black line in the right panel represents a fit to results of other particles using a NCQ scaling inspired function (Eq. 2).



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**Figure 4.** (Color online)  $f_0(980)$   $v_2$  as a function of  $m_T - m_0$  in 30-80% centrality Au+Au collisions at  $\sqrt{s_{NN}} = 200$  GeV. The blue curve represents the NCQ inspired fit (Eq. 2), where the only free parameter is the  $n_q$  of  $f_0(980)$  and all other parameters are fixed according to the fit in the right panel of the Fig. 3.

compared to the fit of other particles [24] using a NCQ scaling inspired function [25]:

$$f_{v_2}(n_q) = \frac{an_q}{1 + \exp(-((m_T - m_0)/n_q - b)/c)} - dn_q.$$
 (2)

The 2-quarks (4-quarks) scaled  $f_0(980)$   $v_2$  seems to deviate from the fit, above (below) the fit by  $\sim 1\sigma$  for the last one or two points at high  $(m_T - m_0)/n_q$ .

Figure 4 shows  $f_0(980)$   $v_2$  as a function of  $m_T - m_0$  with a fit according to the function shown in Eq. 2. In the fit, only the  $n_q$  of  $f_0(980)$  is treated as a free parameter and all other parameters are fixed according to the fit in the right panel of the Fig. 3. This NCQ scaling fit of the  $f_0(980)$   $v_2$  yields  $n_q = 3.0 \pm 0.7$  (stat)  $\pm 0.5$  (syst).

With the current uncertainty, our result is not able to determine whether  $f_0(980)$  is a  $q\bar{q}$ ,  $qq\bar{q}q$ ,  $K\bar{K}$  molecule, gluonium state, or produced through  $\pi\pi$  coalescence. It could also be given by some combined states as well. Future measurements, e.g. the  $f_0(980)$  yields, could also provide different aspect to understand it.

### 53 3 Summary

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Preliminary results on the  $f_0(980)$   $v_2$  in 30-80% centrality Au+Au collisions at  $\sqrt{s_{\mathrm{NN}}}=200$  GeV are presented. In the low  $p_{\mathrm{T}}$  region ( $p_{\mathrm{T}}<2$  GeV/c), the  $f_0(980)$   $v_2$  seems to follow the mass ordering. In the higher  $p_{\mathrm{T}}$  region ( $p_{\mathrm{T}}>2$  GeV/c), the  $f_0(980)$   $v_2$  seems closer to the baryon band. A NCQ scaling inspired function was used to fit the  $f_0(980)$   $v_2$ . The extracted quark content of  $f_0(980)$  is  $n_q=3.0\pm0.7$  (stat)  $\pm0.5$  (syst). More data are needed to understand whether  $f_0(980)$  is a  $q\overline{q}$ ,  $qq\overline{q}q$ ,  $K\overline{K}$  molecule, gluonium state, or produced through  $\pi\pi$  coalescence. Our study indicates that heavyion collisions can be a useful place to examine the quark content of scalar mesons. The isobar data taken in 2018 at RHIC and the 8-fold increase in Au+Au data expected in 2023-2025 would provide more insights.

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#### References

- [1] J.D. Weinstein, N. Isgur, Phys. Rev. D 27, 588 (1983)
- [2] J.D. Weinstein, N. Isgur, Phys. Rev. D 41, 2236 (1990)
- [3] F. Kleefeld, E. van Beveren, G. Rupp, M.D. Scadron, Phys. Rev. D 66, 034007 (2002)
- [4] V. Baru et al., Phys. Lett. B **586**, 53 (2004)
- [5] M. Ablikim et al. (BESIII), Phys. Rev. D **92**, 052003 (2015)
- [6] N.N. Achasov et al., Phys. Rev. D 103, 014010 (2021)
- [7] P.A. Zyla et al. (Particle Data Group), Prog. Theor. Exp. Phys. **083C01** (2020)
- [8] P. Fachini, J. Phys. G **30**, S735 (2004)
- [9] S. Cho et al. (ExHIC), Phys. Rev. Lett. **106**, 212001 (2011)
- [10] D. Molnar, S.A. Voloshin, Phys. Rev. Lett. **91**, 092301 (2003)
- [11] S.S. Adler et al. (PHENIX), Phys. Rev. Lett. **91**, 182301 (2003)
- [12] J. Adams et al. (STAR), Phys. Rev. Lett. 92, 052302 (2004)
- [13] A. Gu, T. Edmonds, J. Zhao, F. Wang, Phys. Rev. C 101, 024908 (2020)
- [14] K.H. Ackermann et al. (STAR), Nucl. Instrum. Meth. **A499**, 624 (2003)
- [15] M. Anderson et al., Nucl. Instrum. Meth. A499, 659 (2003)
- [16] L. Adamczyk et al. (STAR), Phys. Rev. C 92, 024912 (2015)
- [17] L. Adamczyk et al. (STAR), Phys. Rev. Lett. 113, 022301 (2014)
- [18] C. Adler et al. (STAR), Phys. Rev. Lett. 89, 272302 (2002)
- [19] J. Adams et al. (STAR), Phys. Rev. Lett. 92, 092301 (2004)
- [20] E.V. Shuryak, G.E. Brown, Nucl. Phys. A 717, 322 (2003)
- [21] P.F. Kolb, M. Prakash, Phys. Rev. C 67, 044902 (2003)
- [22] R. Rapp, Nucl. Phys. A **725**, 254 (2003)
- [23] A.M. Poskanzer, S.A. Voloshin, Phys. Rev. C 58, 1671 (1998)
- [24] L. Adamczyk et al. (STAR), Phys. Rev. Lett. **116**, 062301 (2016)
- [25] X. Dong, S. Esumi, P. Sorensen, N. Xu, Z. Xu, Phys. Lett. B **597**, 328 (2004)