

1 First measurement of heavy flavour femtoscopy in Au+Au 2 collisions at $\sqrt{s_{\text{NN}}} = 200$ GeV by STAR

3 *Priyanka Roy Chowdhury*^{1,*} (for the STAR Collaboration)

4 ¹Warsaw University of Technology, Faculty of Physics, Koszykowa 75, 00-662 Warsaw, Poland

5 **Abstract.** In the very beginning stages of heavy-ion collisions, hard par-
6 tonic scatterings yield heavy quarks, which go through the entire Quark-Gluon
7 Plasma medium evolution. Femtosopic correlation is characterized as a two
8 particle correlation at low relative momentum that depends on the size of the
9 region from which the correlated particles are emitted and on the final-state in-
10 teraction. Such correlations between identifiable charged hadrons and charmed
11 mesons could provide insights into their interactions with the medium and
12 charm quarks in the hadronic phase. We describe the first femtosopic cor-
13 relation measurements made by the STAR experiment between $D^0/\bar{D}^0 - \pi^\pm$,
14 $D^0/\bar{D}^0 - K^\pm$ and $D^0/\bar{D}^0 - p^\pm$ pairs at mid-rapidity in Au+Au collisions at $\sqrt{s_{\text{NN}}}$
15 = 200 GeV. The discussion of the physics consequences involves a comparison
16 to theory calculations.

17 1 Introduction

18 Heavy quarks such as charm (c) and its charge conjugate (\bar{c}) are created in the early stages
19 of ultra-relativistic collisions of heavy ions. Therefore, they probe all stages of the evolution
20 of heavy-ion collisions - Quark-Gluon Plasma (QGP), hadronization, chemical and kinetic
21 freeze-out. [1] One heavy (c/\bar{c}) and one light (\bar{u}/u) quark make up D^0/\bar{D}^0 mesons (lifetime,
22 $c\tau \approx 123 \mu\text{m}$ [2]). These charmed mesons are useful to probe how heavy quarks interact with
23 QGP medium. The STAR experiment at RHIC (Relativistic Heavy Ion Collider) observed
24 significant D^0 elliptic flow [3] and D^0 suppression at high p_T [4] in heavy-ion collisions.
25 Numerous theoretical models with varying assumptions were able to quantitatively describe
26 these findings. In order to improve our comprehension of heavy quarks (c, b) interactions
27 with the medium, new observables like the two-particle momentum correlation functions are
28 of particular interest. Phase-space evolution of emission source and final-state interactions
29 can be both studied with such measurements in heavy-ion collisions [5]. According to theory,
30 phase-space cloud of outgoing correlated pairs or so-called area of homogeneity is responsive
31 to QGP dynamics, such as collective flow [5, 6]. Size of area of homogeneity is substantially
32 lower than the size of fireball in cases of strong correlation [5].

33 According to Koonin-Pratt equation [5], femtosopic correlation function, $C(k^*) =$
34 $\int S(r)|\psi(k^*, r)|^2 d^3r$, is a combination of emission source function, $S(r)$ and pair-wave func-
35 tion, $\psi(k^*, r)$. Here k^* represents the reduced momentum difference between correlated par-
36 ticle pairs emitting from a source of size r . We computed $C(k^*)$ by taking the k^* ratio of

*e-mail: priyanka.roy_chowdhury.dokt@pw.edu.pl

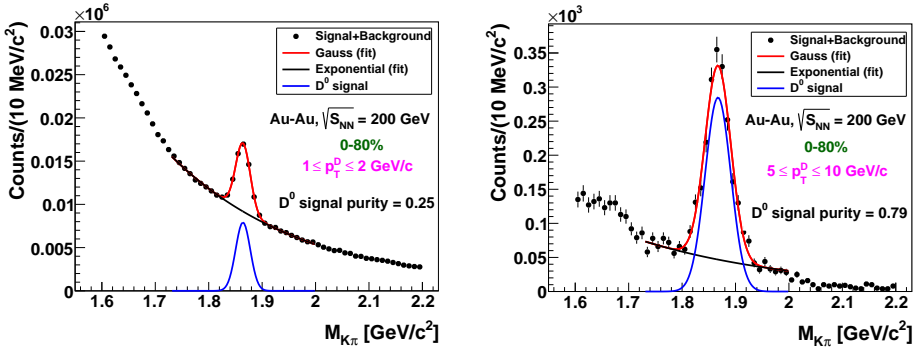
37 correlated $[A(k^*)]$ and uncorrelated $[B(k^*)]$ D^0 -hadron pairs in the rest frame of their center
 38 of mass given in Eq. 1 [5].

$$C(k^*) = N \frac{A(k^*)}{B(k^*)} \quad \text{and} \quad k^* = \frac{1}{2}(p_1 - p_2), \quad (1)$$

39 where N is normalization factor and p_1, p_2 are momenta of D^0/\bar{D}^0 and $K/\pi/p$ tracks in the
 40 pair-rest frame. Two tracks were chosen from the same event in order to compute $A(k^*)$. $B(k^*)$
 41 was calculated using event mixing technique to select a pair of tracks from different events
 42 chosen within a similar primary z-vertex position (V_z) and centrality range as the same-event
 43 pairs.

44 2 D^0/\bar{D}^0 reconstruction and signal purity estimation

45 STAR is made up of multiple detectors with distinct functions [7]. The two primary detec-
 46 tors for tracking and identifying charge particles are the Time Projection Chamber (TPC) and
 47 Time of Flight (TOF). The HFT detector with outstanding track pointing resolution was used
 48 for the reconstruction of D^0/\bar{D}^0 mesons via the $K^\mp\pi^\pm$ decay channel using the set of topo-
 49 logical criteria [4]. Fig. 1 shows the reconstructed D^0 and \bar{D}^0 invariant mass distributions
 50 using STAR data of Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV [4]. We selected D^0 candidates
 51 (and their charge conjugates) based on two criteria: $p_T > 1$ GeV/c and the ratio of signal-
 52 to-combinatorial background under D^0 peak (S/B) had to exceed 30% in the lowest p_T bin.
 53 D^0 (\bar{D}^0) signal peak lies between 1.82 to 1.91 GeV/ c^2 mass range. With rising transverse
 54 momentum, p_T , the S/B ratio increases. The D^0 signal purity ($\frac{S}{S+B}$) was calculated for each
 55 p_T bin (1-2, 2-3, 3-5, and 5-10 GeV/c). In the lowest p_T bin, the purity is 25% , gradually
 56 increasing to 80% in the highest p_T bin. D^0 purity is used to calculate the pair purity as
 discussed in detail in Sec. 3.



57 **Figure 1.** Invariant mass ($M_{K\pi}$) of D^0 and \bar{D}^0 candidates using STAR Run 2014 data. The black solid
 circles depict the invariant mass of unlike-sign (US) $K\pi$ pairs. The red curve corresponds to a Gaussian
 fit, while the black line shows an exponential fit to the US pairs. The blue curve illustrates the fit to the
 58 D^0 (\bar{D}^0) signal within the mass range of 1.73 to 2.0 GeV/ c^2 .

58 3 Correlation function calculation

59 The computed value of the correlation function can be affected by incorrectly identified cor-
 60 related pairs. We used both TPC and TOF to select the primary tracks of considered charged
 61 hadrons (π, K and p). We eliminated splitted hadron tracks and excluded self-correlation

62 between D^0 daughters in order to reduce the potential effects of the TPC detector. Combining
 63 two different tracks into one is another possible detector impact, however our investigation
 64 showed that the contribution from merged tracks was negligible. To account for the contri-
 65 bution from the background under D^0 signal peak and impurity of the selected π , K and p
 66 samples with other hadrons and electrons present in the system, we implemented a purity
 67 correction for D^0 -hadron pairs using a specific formula [8]:

$$C(k^*) = \frac{C_{\text{measured}}(k^*) - 1}{\text{Pair Purity}} + 1, \quad (2)$$

68 where $C(k^*)$ is the final correlation function after purity correction, $C_{\text{measured}}(k^*)$ represents the
 69 correlation function after correcting possible detector effects, while Pair Purity is determined
 70 by multiplying the D^0 signal purity with the average purity of the hadron sample. The purity
 71 of the hadron sample was determined using standard approach in STAR [9]. Hadrons were
 72 selected with momentum criteria, $p_\pi < 1$ GeV/c, $p_p < 1.2$ GeV/c and $p_K < 1$ GeV/c as
 73 they become difficult to distinguish from electrons and other hadrons beyond the mentioned
 74 thresholds respectively. Within selected momentum range, average purity of π and p are
 75 $(99.5 \pm 0.5)\%$ while for K sample, purity is $(97 \pm 3)\%$. Systematic uncertainty studies were
 76 made by variations in the topological cuts used for D^0 reconstruction [4], and by including
 77 the uncertainties in the purity estimation of D^0 -hadron pairs.

78 4 Results and discussions

79 Figure 2 presents the final correlation function of $D^0/\overline{D^0} - \pi^\pm$, $D^0/\overline{D^0} - p^\pm$ and $D^0/\overline{D^0} - K^\pm$
 80 pairs using STAR data from minimum bias Au+Au collision at $\sqrt{s_{\text{NN}}} = 200$ GeV within
 81 $|\eta| < 1$. The $C(k^*)$ distribution is consistent with unity, indicating an absence of significant
 82 correlation, with large statistical fluctuations. The strength of the correlation is linked to the
 size of the source [5, 6]. A weak or absent correlation suggests that the emission source

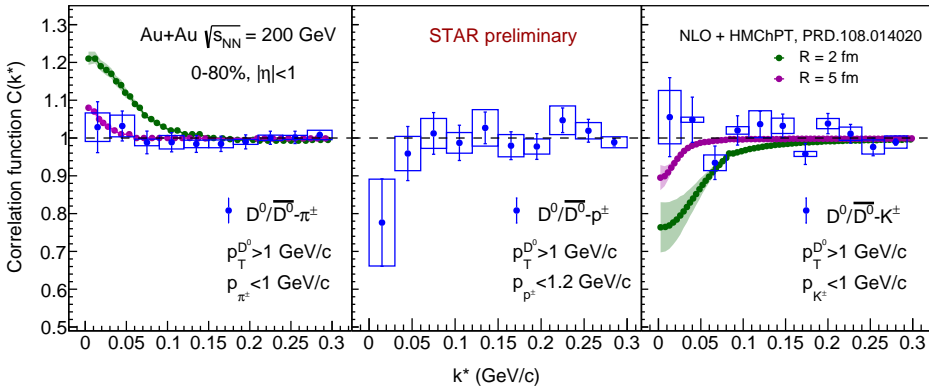


Figure 2. $C(k^*)$ for (left) $D^0/\overline{D^0} - \pi^\pm$, (mid) $D^0/\overline{D^0} - p^\pm$ and (right) $D^0/\overline{D^0} - K^\pm$ pairs. The blue solid circles, along with the boxes, indicate the STAR data and the associated systematic uncertainties. The green and pink bands illustrate the $C(k^*)$ values predicted by NLO + HMChPT model for source sizes of 2 fm and 5 fm, respectively.

83 from which the D -hadron pairs are produced is relatively large in size. In left and right
 84 panel, STAR data are compared with $C(k^*)$ calculated using NLO (next-to-leading order)-
 85 HMChPT (Heavy Meson Chiral Perturbation Theory) scheme [10–12]. In the left panel, the
 86

theory predictions are for $D^0 - \pi^+$ and $D^+ - \pi^0$ pairs while in the right panel, for $D^0 - K^+$ pairs [10]. Neither of these channels includes Coulomb interaction. The STAR data and theoretical predictions align with an emission source size of 5 fm or more. The $D^0 K^+$ channel exhibits a threshold near 0.083 GeV where a cusp effect is observed. The prediction shows a marked depletion around k^* between 0 to 0.05 GeV/c, which is attributed to the $D_S^*(2317)^\pm$ bound state and becomes more pronounced as the source radius decreases [10]. However, the resonance effect of this state is not evident in the STAR results due to either the large source size or significant experimental uncertainties. As a next step, we plan to update these results by combining data from Runs 2014 and 2016. This approach is expected to improve the precision of the correlation function measurements and offer more definitive conclusions about the source size. Additionally, theoretical inputs are necessary for a better interpretation of these data.

5 Acknowledgement

Priyanka Roy Chowdhury expresses gratitude for the financial support provided by the National Science Centre, Poland (NCN) under grant no. 2018/30/E/ST2/0008, as well as partial support from the U.S. Department of Energy (DOE).

References

- [1] X. Dong *et al.*, Open Heavy-Flavor Production in Heavy-Ion Collisions, *Ann. Rev. Nucl. Part. Sci.* **69**, 417-445 (2019).
- [2] R. L. Workman *et al.* [Particle Data Group], Review of Particle Physics, *PTEP* **2022**, 083C01 (2022).
- [3] L. Adamczyk *et al.* [STAR], Measurement of D^0 Azimuthal Anisotropy at Midrapidity in Au+Au Collisions at $\sqrt{s_{NN}}=200$ GeV, *Phys. Rev. Lett.* **118**, no.21, 212301 (2017).
- [4] J. Adam *et al.* [STAR], Centrality and transverse momentum dependence of D^0 -meson production at mid-rapidity in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV, *Phys. Rev. C* **99**, no.3, 034908 (2019).
- [5] M. A. Lisa *et al.*, Femtoscopy in relativistic heavy ion collisions, *Ann. Rev. Nucl. Part. Sci.* **55**, 357-402 (2005).
- [6] S. V. Akkelin and Y. M. Sinyukov, The HBT interferometry of expanding sources, *Phys. Lett. B* **356**, 525-530 (1995).
- [7] Special Issue on RHIC and Its Detectors, edited by M. Harrison, T. Ludlam and S. Ozaki, *Nucl. Instrum. Methods Phys. Res., Sect. A* **499**, no.2-3 (2003).
- [8] J. Adams *et al.* [STAR], Proton - lambda correlations in central Au+Au collisions at $\sqrt{s_{NN}} = 200$ -GeV, *Phys. Rev. C* **74**, 064906 (2006).
- [9] L. Adamczyk *et al.* [STAR], *Phys. Rev. C* **95**, no. 3, 034907 (2017).
- [10] M. Albaladejo, J. Nieves and E. Ruiz-Arriola, Femtoscopic signatures of the lightest S-wave scalar open-charm mesons, *Phys. Rev. D* **108**, no.1, 014020 (2023).
- [11] F. K. Guo, C. Hanhart and U. G. Meissner, Interactions between heavy mesons and Goldstone bosons from chiral dynamics, *Eur. Phys. J. A* **40**, 171-179 (2009).
- [12] L. S. Geng, N. Kaiser, J. Martin-Camalich and W. Weise, Low-energy interactions of Nambu-Goldstone bosons with D mesons in covariant chiral perturbation theory, *Phys. Rev. D* **82**, 054022 (2010).