

Multiplicity and Rapidity Dependent Study of (Multi)-strange Hadrons in $d+Au$ collisions using the STAR detector

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Abstract. Strangeness enhancement has long been considered as a signature of the quark-gluon plasma formation in heavy-ion collisions. Strangeness enhancement has also been observed in small systems at the Large Hadron Collider (LHC), but the underlying physics is not yet fully understood. This motivates us to study the strange hadron production in small systems at RHIC. We present new measurements of (multi-)strange hadrons (K_s^0 , Λ , Ξ and Ω) in $d+Au$ collisions at $\sqrt{s_{NN}} = 200$ GeV, collected by STAR in 2016. We report the multiplicity and rapidity dependence of strange hadron transverse momentum (p_T) spectra, p_T -integrated yields dN/dy , average transverse momentum ($\langle p_T \rangle$), yield ratios of these strange hadrons to pions, nuclear modification factors and rapidity asymmetry for these particles in $d+Au$ collisions.

1 Introduction

The production of strange hadrons in high-energy hadronic interactions provides a way to investigate the properties of quantum chromodynamics (QCD), the theory of strongly interacting matter. Unlike up (u) & down (d) quarks, strange (s) quarks are not present as valence quarks in the initial state while they are created during the collisions. Strangeness enhancement in heavy ion collisions with respect to proton-proton ($p+p$) collisions has been suggested as a signature of quark-gluon plasma (QGP) formation [1]. But creation of QGP in small systems is still under intense debate.

Asymmetric small collision systems like proton-nucleus ($p-A$) and deuteron-nucleus ($d-A$) can be considered as control experiments where the formation of an extended QGP phase is not expected. These collision systems are used for baseline measurements to study the possible effects of cold nuclear matter and disentangle them from hot dense matter effects present in collisions of heavy ions. The process of generating hadrons can be affected by various factors, including alterations in parton distribution functions within nuclei, the possibility of parton saturation, multiple scatterings, and radial flow. It is anticipated that the magnitude of these effects may vary with the rapidity of the produced particles.

The abundances of strange particles relative to pions in heavy-ion collisions for top RHIC and LHC energies do not show a significant dependence on the initial volume (collision centrality) [2].

Nuclear effects can be quantified using various variables such as nuclear modification factor and rapidity asymmetry. Nuclear modification factor is defined as the ratio of the yield

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of particle in heavy-ion collisions to its yield in $p+p$ collisions, scaled by the number of binary nucleon-nucleon inelastic collisions [3].

$$R_{AB}(p_T) = \frac{\text{Yield}_{AB}}{\langle N_{\text{bin}} \rangle \text{Yield}_{pp}} \quad (1)$$

where $\langle N_{\text{bin}} \rangle$ is the average number of binary nucleon-nucleon collisions. Comparative study of particle production in forward and backward rapidity regions is done using rapidity asymmetry (Y_{Asym}) [4]. Y_{Asym} is defined as

$$Y_{\text{Asym}}(p_T) = \frac{Y_B(p_T)}{Y_F(p_T)} \quad (2)$$

where Y_B and Y_F are particle yields in backward and forward rapidity regions, respectively. Y_{Asym} may provide unique information to help determine the relative contributions of various physics processes affecting particle production, such as multiple scattering, nuclear shadowing, recombination of thermal partons, and parton saturation.

2 Analysis Details

A successful run of $d+\text{Au}$ collisions at $\sqrt{s_{NN}} = 200$ GeV was carried out in 2016 at RHIC. A total of approximately 100 million good events have been selected in STAR for the reconstruction of K_S^0 , $\Lambda(\bar{\Lambda})$, $\Xi(\bar{\Xi})$, $\Omega(\bar{\Omega})$. K_S^0 , $\Lambda(\bar{\Lambda})$, $\Xi(\bar{\Xi})$, $\Omega(\bar{\Omega})$ are weakly decaying particles and they travel a certain distance before decaying into daughter particles ($K_S^0 \rightarrow \pi^+ + \pi^-$, $\Lambda(\bar{\Lambda}) \rightarrow p(\bar{p}) + \pi^-(\pi^+)$, $\Xi^-(\bar{\Xi}^+) \rightarrow \Lambda(\bar{\Lambda}) + \pi^-(\pi^+)$, $\Omega^-(\bar{\Omega}^+) \rightarrow \Lambda(\bar{\Lambda}) + K^-(K^+)$). Identification of daughter particles (π , K and p) is done via measuring $\langle dE/dx \rangle$ using Time Projection Chamber (TPC).

The K_S^0 , $\Lambda(\bar{\Lambda})$, $\Xi(\bar{\Xi})$, $\Omega(\bar{\Omega})$ signals are extracted by reconstructing the invariant mass of the decay daughter pairs. The decay topology can be used for their reconstruction and to suppress the background. We have used double Gaussian and second order polynomial function to describe the signal and background invariant mass distributions respectively. Raw yield is also estimated using functional fitting which is considered in systematic uncertainty. Raw yield is determined by using the bin counting method under the mass window of $M_0 \pm 3\sigma$, where M_0 is mass of K_S^0 (or $\Lambda/\Xi/\Omega$) and σ is the fitted width. Raw yield for each p_T interval is corrected for branching ratio, acceptance and efficiency to obtain corrected p_T spectra. Weak decay feed down correction from Ξ is applied to Λ .

3 Results and Discussions

In Fig. 1 (left), integrated yield particle-to-pion ratio for K_S^0 , Λ , Ξ and Ω as a function of multiplicity for $d+\text{Au}$ collisions at $\sqrt{s_{NN}} = 200$ GeV is shown. Smooth transition of these ratios from $p+p$ to heavy ion collisions is observed. Data from different collision systems are consistent with each other at similar multiplicities [2] and yields of particles with more strangeness content decrease faster as we move from high to low multiplicities.

In Fig. 1 (right), nuclear modification factors for K_S^0 , Λ and Ξ at mid rapidity ($|y| < 0.5$) for $d+\text{Au}$ collision is presented and Cronin-like enhancement is observed which is stronger for baryons (Ξ , Λ , p) as compared to mesons (K_S^0 , π).

In Fig. 2 (left), integrated yield and $\langle p_T \rangle$ as function of multiplicity for mid-rapidity in $p+p$, $d+\text{Au}$, $\text{Cu}+\text{Cu}$ and $\text{Au}+\text{Au}$ collisions at $\sqrt{s_{NN}} = 200$ GeV are shown. Integrated yield increases as a function of multiplicity and we have also observed that particle production is mainly driven by multiplicity and not by collision system. Hint of an increase in $\langle p_T \rangle$ is observed for heavier particles, which supports the picture of radial flow.

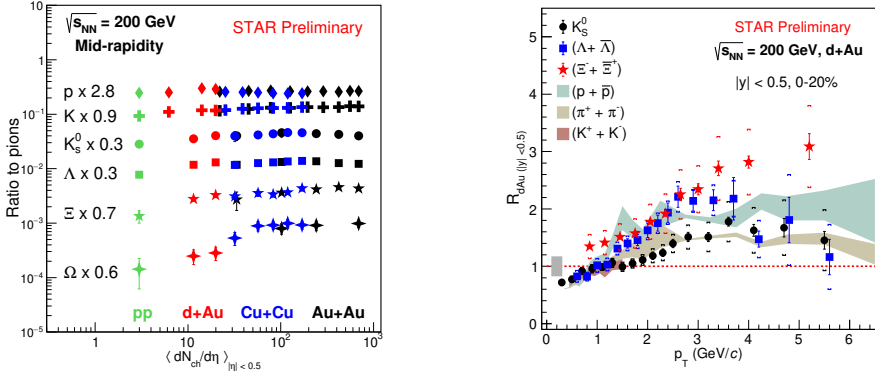


Figure 1: Integrated yield particle-to-pion ratio (left) for strange particles (K_S^0 , Λ , Ξ and Ω) as a function of multiplicity for $p+p$, $d+Au$, $Cu+Cu$ and $Au+Au$ collisions and Nuclear modification factors (R_{dAu}) (right) for strange particles (K_S^0 , Λ and Ξ) in $d+Au$ collisions at $\sqrt{s_{NN}} = 200$ GeV.

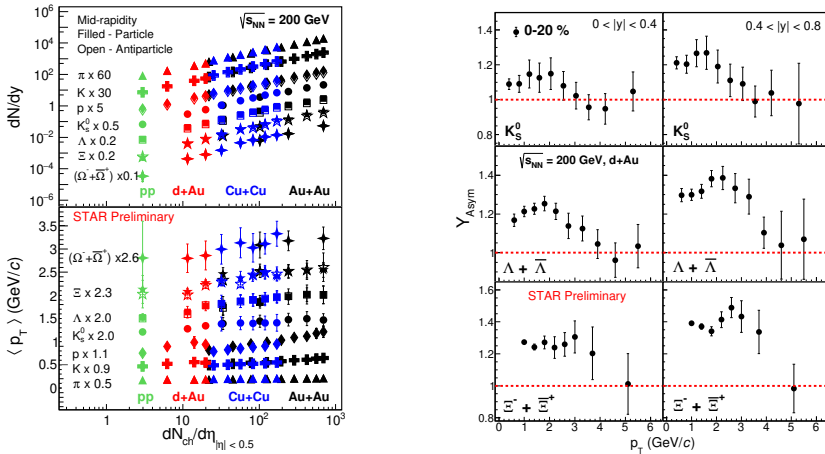


Figure 2: Integrated yield (dN/dy) and mean transverse momentum ($\langle p_T \rangle$) (left) for strange particles (K_S^0 , Λ , Ξ and Ω) as a function of multiplicity for $p+p$, $d+Au$, $Cu+Cu$ and $Au+Au$ collisions and rapidity asymmetry (Y_{Asym}) as a function of p_T (right) for strange particles (K_S^0 , Λ and Ξ) in $d+Au$ collisions at $\sqrt{s_{NN}} = 200$ GeV.

78 We have also measured the transverse momentum dependence of Y_{Asym} for K_S^0 , Λ and Ξ
79 in two rapidity intervals ($0 < |y| < 0.4$, $0.4 < |y| < 0.8$) in $d+Au$ collision at $\sqrt{s_{NN}} = 200$ GeV.
80 Y_{Asym} values deviate from unity at low p_T (< 3 GeV/c), suggesting the presence of a rapidity
81 dependence in the nuclear effects. Y_{Asym} is consistent with unity at high p_T , providing a hint
82 that nuclear effects become weaker at high p_T as shown in Fig. 2 (right). Y_{Asym} is more
83 prominent for larger rapidity interval and for heavier particles.

84 In Fig. 3, integrated yield (dN/dy) and mean transverse momentum ($\langle p_T \rangle$) are calculated
85 for different rapidity intervals ($0 < |y| < 0.4$, $0.4 < |y| < 0.8$). For strange particles K_S^0 ,

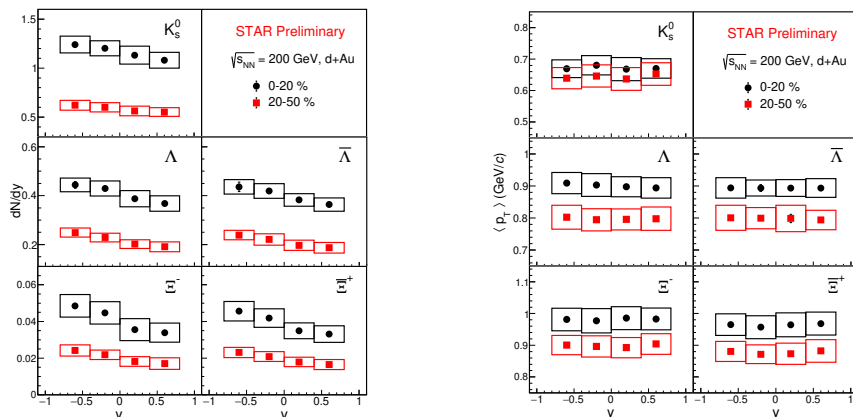


Figure 3: dN/dy (left) and $\langle p_T \rangle$ (Right) of strange particles (K_S^0 , $\Lambda(\bar{\Lambda})$ and $\Xi(\bar{\Xi})$) as a function of rapidity in $d+Au$ collisions at $\sqrt{s_{NN}} = 200$ GeV for multiplicity classes 0-20% and 20-50%.

86 $\Lambda(\bar{\Lambda})$ and $\Xi(\bar{\Xi})$, integrated yield (dN/dy) decreases slightly from negative (Au going side)
 87 to positive (d going side) rapidity but mean transverse momentum ($\langle p_T \rangle$) is constant in the
 88 measured rapidity range which suggests a similar radial flow for rapidity ranging from $-0.8 <$
 89 $y < 0.8$.

90 4 Summary

91 We present new measurements of multiplicity and rapidity dependent study of (multi-)strange
 92 hadrons (K_S^0 , $\Lambda(\bar{\Lambda})$, $\Xi(\bar{\Xi})$ and $\Omega(\bar{\Omega})$) production in $d+Au$ collisions at $\sqrt{s_{NN}} = 200$ GeV with
 93 the STAR experiment. p_T -integrated yields of strange and multi-strange particles relative to
 94 pions increase significantly with multiplicity. Results suggest that strange particle produc-
 95 tion is independent of the collision system and mainly driven by multiplicity. The observed
 96 enhancement increases with strangeness content rather than with mass or baryon number of
 97 the hadron. Hint of Cronin-like enhancement is observed for 0-20% centrality for K_S^0 , Λ and
 98 Ξ . $Y_{Asym} > 1$ at low p_T , indicating the presence of the nuclear effects. The Y_{Asym} is more
 99 prominent for heavier mass and larger rapidity interval.

100 References

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