

Measurement of global polarization of Λ hyperons in Au+Au $\sqrt{s_{\text{NN}}} = 7.2$ GeV fixed target collisions at RHIC-STAR experiment

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Abstract. In non-central nuclear collisions, the created matter should exhibit strong vorticity, which induces spin polarization. Global polarization is a good tool to investigate the dynamics of the system. Global polarization of Λ and $\bar{\Lambda}$ hyperons has been measured at $\sqrt{s_{\text{NN}}} = 2.4$ GeV - 5.02 TeV in various experiments. We have measured the global polarization of Λ hyperons in Au+Au collisions with fixed-target configuration at $\sqrt{s_{\text{NN}}} = 7.2$ GeV at RHIC-STAR experiment. The non-zero signal is observed and it follows the global trend of Λ global polarization that increases towards lower collision energies. The global polarization is larger in more peripheral collisions but it does not show any significant rapidity dependence as well as transverse momentum dependence within uncertainties.

1 Introduction

In non-central heavy-ion collisions, the created matter should exhibit a rotational motion because of the conservation of the initial orbital angular momentum carried by the two nuclei, which is perpendicular to the reaction plane. Particle's and anti-particle's spins are aligned due to the spin-orbit coupling [1]. The strong magnetic field would appear in the initial state and its direction is almost the same as the initial angular orbital momentum. Particle's and anti-particle's spins also could get aligned in the opposite direction by the magnetic field due to the opposite signs of their magnetic moments. Therefore, global polarization is a good tool to investigate the dynamics of the system. Global polarization can be measured with hyperons through their parity-violating weak decays. The STAR Collaboration reported the global polarization of Λ and $\bar{\Lambda}$ hyperons at $\sqrt{s_{\text{NN}}} = 7.7 - 200$ GeV and observed positive signals [2–4]. Additionally, the ALICE Collaboration measured global polarization of Λ and $\bar{\Lambda}$ hyperons at $\sqrt{s_{\text{NN}}} = 2.76, 5.02$ TeV [5]. These measurements show that global polarization becomes larger at lower collision energies. The STAR fixed-target program provides an opportunity to extend such measurements at even lower energies. We here present the measurement of global polarization of Λ hyperons in Au+Au collisions at $\sqrt{s_{\text{NN}}} = 7.2$ GeV with fixed-target configuration.

2 Analysis details

The good minimum bias data of Au+Au collisions have been collected by the STAR detector with fixed-target configuration in 2018 and 209 million events were analyzed in this study.

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37 The gold target has been installed inside the vacuum pipe at 2.0 meters away from the center
 38 of the STAR detector. The decay channel of $\Lambda \rightarrow p + \pi$ (BR: 63.9% [6]) was utilized for Λ
 39 identification. The daughter particles were identified using TPC dE/dx and TOF and then Λ
 40 hyperons were reconstructed based on the invariant mass, applying cuts on decay topology. In
 41 this parity-violating weak decay, the daughter proton preferentially decays along the Λ 's spin
 42 direction. Furthermore, since the system angular momentum is perpendicular to the reaction
 43 plane, the global polarization can be measured via the distribution of the azimuthal angle of
 44 the daughter proton in the Λ 's rest frame with respect to the reaction plane as follows [2]:

$$P_H = \frac{8}{\pi\alpha_H} \frac{\langle \sin(\Psi_1 - \phi_p^*) \rangle}{\text{Res}(\Psi_1)} \quad (1)$$

45 where α_H is Λ 's decay parameter ($\alpha_\Lambda = 0.732 \pm 0.014$ [6]), Ψ_1 is the first order event plane,
 46 and ϕ_p^* is the azimuthal angle of the daughter proton in the Λ 's rest frame. The event plane
 47 reconstructed by Event Plane Detector (EPD) and the resolution estimated by three subevent
 method [7].

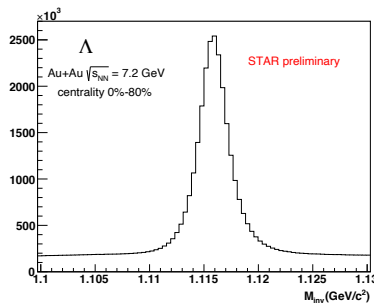


Figure 1: Invariant mass distribution of Λ hyperons in 0%-80% centrality for Au+Au collisions at $\sqrt{s_{NN}} = 7.2$ GeV.

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49 3 Result

50 Figure 2 shows global polarization of Λ and $\bar{\Lambda}$ as a function of collision energy. Solid markers
 51 indicate Λ and open markers indicate $\bar{\Lambda}$ global polarization. Because the Λ and $\bar{\Lambda}$'s decay
 52 parameter were updated recently, the previous results were rescaled by using the new values
 53 ($\alpha_{old}/\alpha_{new}$). The result at 7.2 GeV shows a non-zero positive signal. It follows the global
 54 trend of Λ global polarization that increases towards lower collision energies.

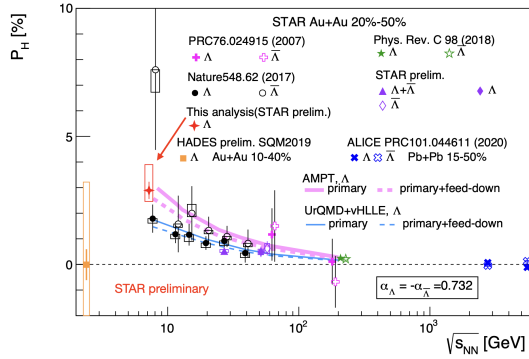


Figure 2: Global polarization of Λ (solid markers) and $\bar{\Lambda}$ (open markers) as a function of collision energy. Bold lines show the AMPT model calculations [8] and thin lines show calculations by a 3+1D cascade + viscous hydrodynamic model (UrQMD + vHELL) [9]. Primary Λ with and without the feed-down effect are indicated by dashed and solid lines, respectively.

55 We have also performed differential measurements such as centrality, rapidity, and transverse
 56 momentum dependence of global polarization. Figure 3 shows centrality dependence of Λ
 57 global polarization. Λ global polarization increases in more peripheral collisions as expected
 from the initial collision geometry.

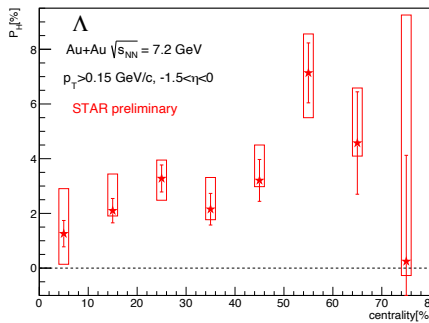


Figure 3: Global polarization of Λ hyperons as a function of centrality for $p_T > 0.15$ GeV/c, $-1.5 < \eta < 0$.

58 Theoretical models predict that the polarization depends on rapidity, although the prediction
 59 is different among the models [10–13]. Figure 4 shows rapidity dependence of Λ global polar-
 60 ization for 20%-60% centrality, shown for 4 different p_T regions. Λ global polarization is
 61 consistent with being constant within uncertainties within the acceptance of the STAR detec-
 62 tor. STAR experiment has recently installed a new detector inner Time Projection Chamber
 63 (iTPC) and will upgrade the forward detectors (2023+2025). With these upgrades, global
 64 polarization in large rapidity regions can be explored in the future. The Λ global polariza-
 65 tion is expected to decrease at low p_T due to the smearing effect caused by scattering at the
 66

67 later stage of the collision. On the other hand, it may also decrease at high p_T due to the jet
 68 fragmentation. Figure 5 shows Λ global polarization as a function of transverse momentum
 69 for 20%-60% centrality. The Λ global polarization does not show significant p_T dependence
 within uncertainties.

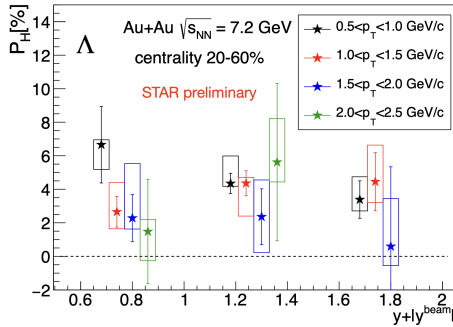


Figure 4: Global polarization of Λ hyperons as a function of rapidity for 20%-60% centrality each p_T region.

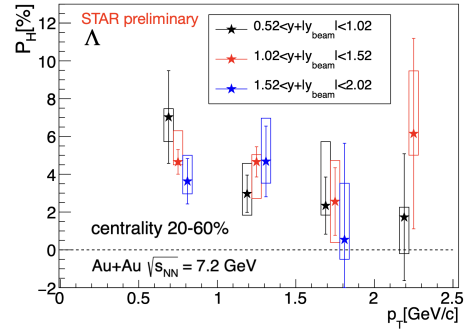


Figure 5: Global polarization of Λ hyperons as a function of transverse momentum for 20%-60% centrality each rapidity region.

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71 4 Summary

72 We present the first measurement of Λ global polarization in Au+Au collisions at $\sqrt{s_{NN}} = 7.2$
 73 GeV with STAR fixed-target configuration. The positive non-zero signal has been observed
 74 and it follows the global trend of energy dependence. We have also performed differential
 75 measurements of global polarization. The increasing trend of polarization towards peripheral
 76 collisions is observed. We have observed no significant rapidity and p_T dependence within
 77 uncertainties. With newly upgraded iTPC and future upgrade of forward detectors in 2023-
 78 2025, the global polarization will be explored in more forward regions.

79 References

- 80 [1] Z. T. Liang and X. N. Wang, Phys. Rev. Lett. 94, 102301 (2005).
 81 [2] B. I. Abelev *et al.* (STAR), Phys. Rev. C 76 024915 (2007).
 82 [3] L. Adamczyk *et al.* (STAR), Nature 548 62 (2017).
 83 [4] J. Adam *et al.* (STAR), Phys. Rev. C 98 14910 (2018).
 84 [5] S. Acharya *et al.* (ALICE), Phys. Rev. C 101 4, 044611 (2020).
 85 [6] P. A. Zyla *et al.* (Particle Data Group), Prog. Theor. Exp. Phys. 2020, 083C01 (2020).
 86 [7] A. M. Poskanzer and S. A. Voloshin, Phys. Rev. C 58, 1671 (1998).
 87 [8] H. Li, L. Pang, Q. Wang, and X. Xia, Phys. Rev. C 96, 054908 (2017).
 88 [9] I. Karpenko and F. Becattini, Eur. Phys. J. C 77, 213 (2017).
 89 [10] W. T. Deng and X. -G Huang, Phys. Rev. C 93, 064907 (2016).
 90 [11] Z. T. Liang *et al.*, Chin. Phys. C 45 1, 014102 (2021).
 91 [12] D. X. Wei, W. T. Deng and X. G. Huang, Phys. Rev. C 99.014905 (2019).
 92 [13] H. Z. Wu *et al.*, Phys. Rev. Research 1, 033058 (2019).