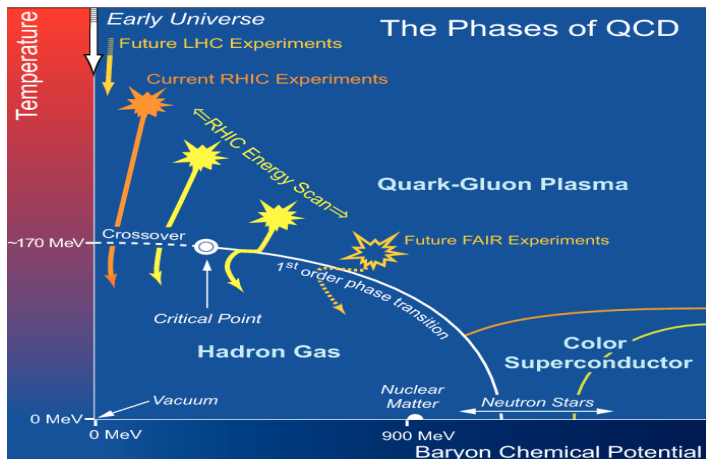




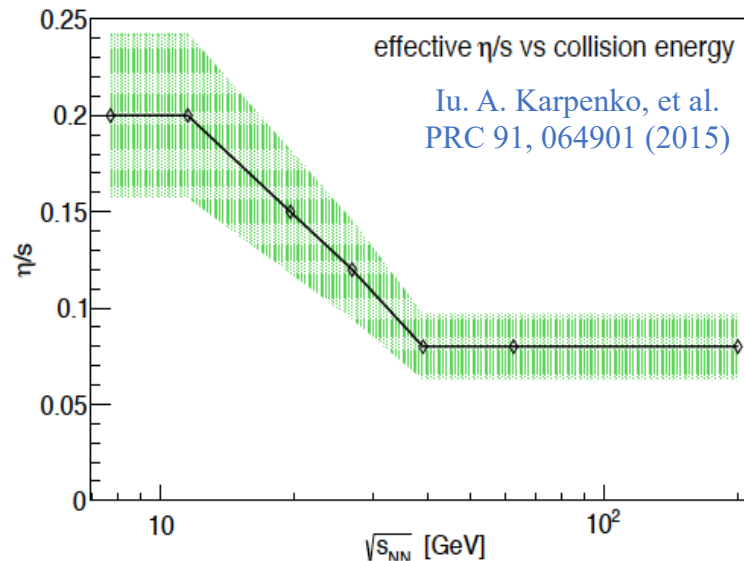
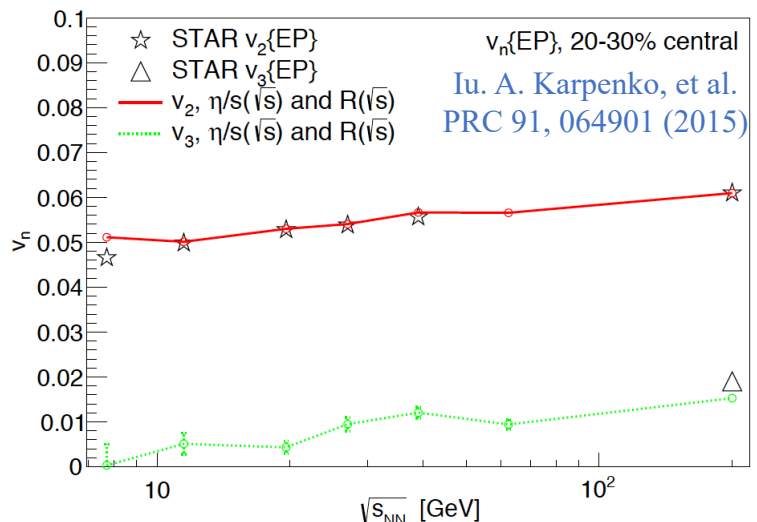
Beam-energy dependence of the viscous damping of anisotropic flow

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- Strong interest in the theoretical calculations which span a broad (T, μ_B) domain.

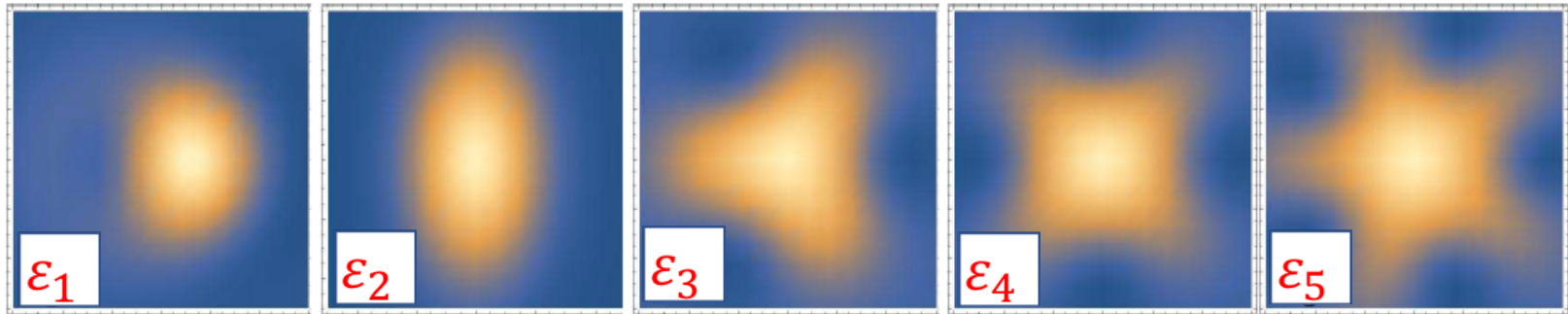
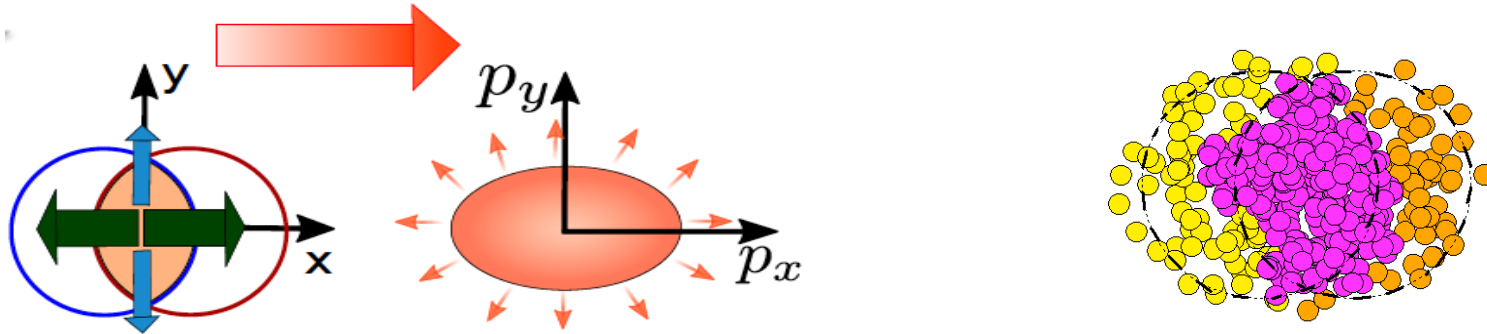


- The beam energy dependence of the anisotropic flow measurements can aid the extraction of $\frac{\eta}{s}(T, \mu_B)$

Introduction

Anisotropic flow

Asymmetry in initial geometry \rightarrow Final state momentum anisotropy (flow)



$$dN/d\varphi = 1 + 2 \sum_n^{\infty} v_n \cos(\varphi - \Psi_n)$$

- The flow harmonic coefficients v_n are influenced by eccentricities (ϵ_n), fluctuations, system size, speed of sound $c_s(\mu_B, T)$, and transport coefficient $\frac{\eta}{s}(\mu_B, T)$

- Comprehensive set of flow measurements are important to study;
 - ✓ Differentiate between initial-state models
 - ❑ Initial-state eccentricity & its fluctuations
 - ✓ Transport coefficients (η/s , etc)
 - ❑ Pin down the temperature dependence of the transport coefficients
 - ✓ Detailed flow measurements could aid ongoing efforts to search for the critical end point(CEP)

What are the respective roles of ε_n and its fluctuations, system size and transport coefficient $\frac{\eta}{s}(T, \mu_B)$?

Azimuthal anisotropy measurements

Correlation function

Two-particle correlation function $Cr(\Delta\varphi = \varphi_a - \varphi_b)$,

$$Cr(\Delta\varphi) = dN/d\Delta\varphi \text{ and } v_n^{ab} = \frac{\sum_{\Delta\varphi} Cr(\Delta\varphi) \cos(n \Delta\varphi)}{\sum_{\Delta\varphi} Cr(\Delta\varphi)}$$

$$n > 1$$

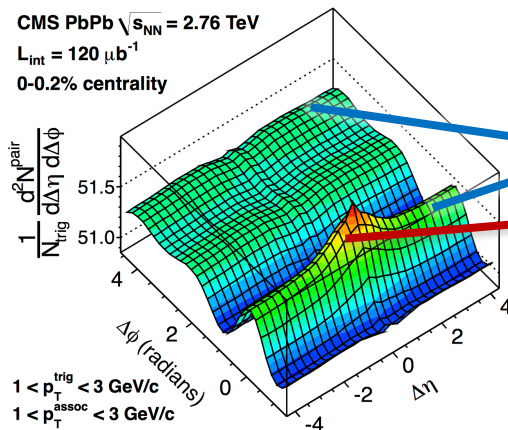
$$v_n^{ab} = v_n^a v_n^b + \delta_{short}$$

$$n = 1$$

$$v_1^{ab} = v_1^a v_1^b + \delta_{long}$$

Flow

Non-flow



Long – range

Short – range

Momentum Conservation

HBT

Di-jets

Decay

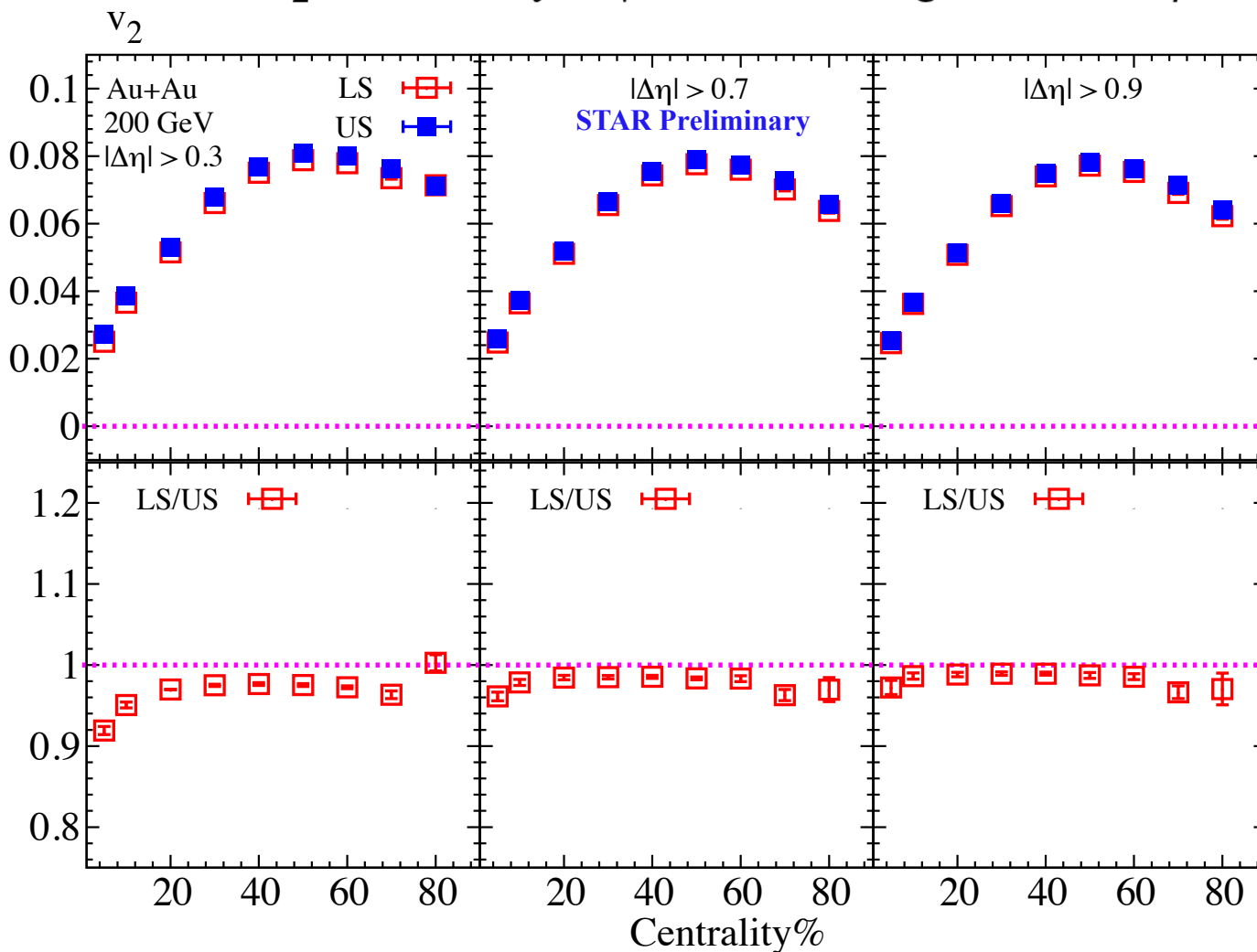
Charge

Non-flow suppression is needed

Short - range
Non-flow

Short-range non-flow suppression

The v_2 vs centrality at $\sqrt{s_{NN}} = 200$ using different $\Delta\eta$ cuts

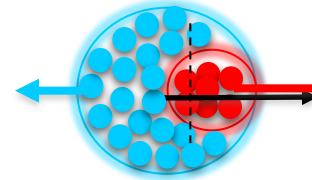


✓ Short-range non-flow effect reduced by using
 $|\Delta\eta| > 0.7$ cut

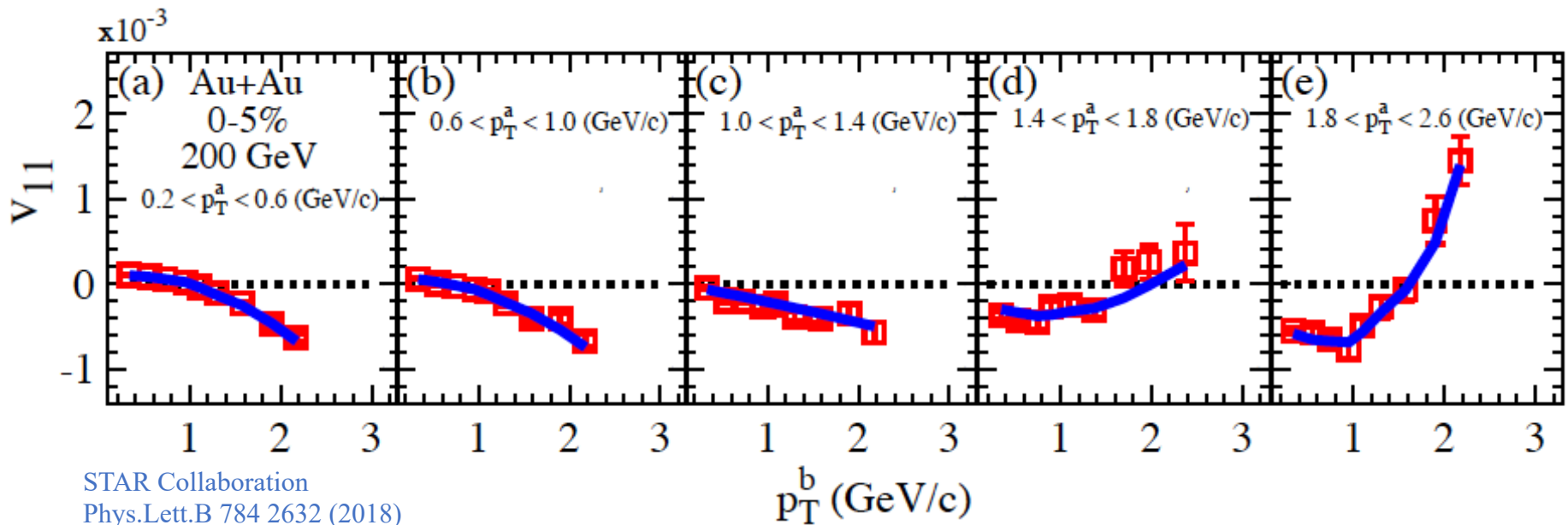
Long-range non-flow suppression

$$v_{11}^{ab} = v_1^{even}(p_T^a) v_1^{even}(p_T^b) + \delta_{long}$$

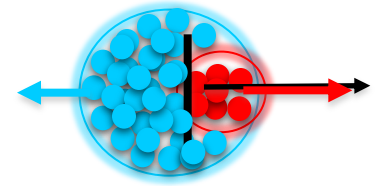
$$v_{11}(p_T^a, p_T^b) = v_1^{even}(p_T^a) v_1^{even}(p_T^b) - K p_T^a p_T^b$$



v_{11} in Eq(1) represents NxM matrix which we fit with N+1 parameters

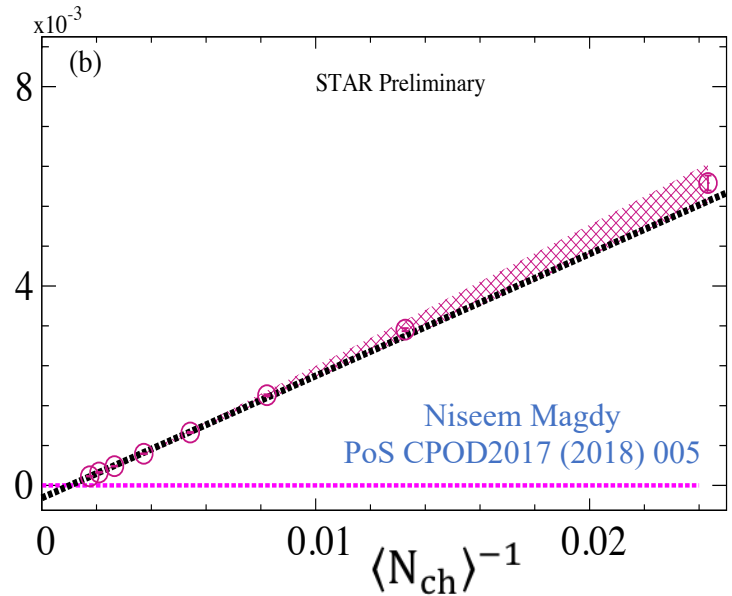
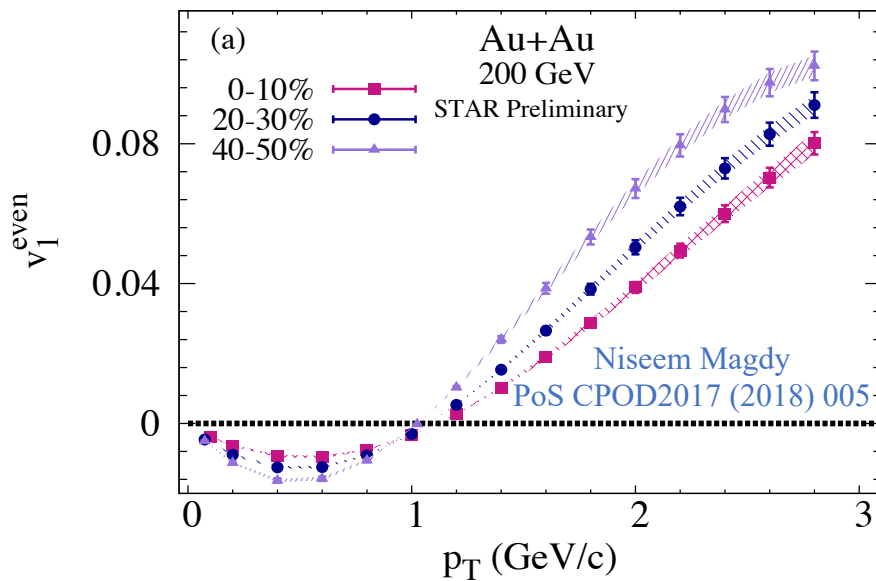


➤ v_{11} characteristic behavior gives a good constraint for $v_1^{even}(p_T)$ extraction



$$v_{11}(p_T^a, p_T^b) = v_1^{even}(p_T^a)v_1^{even}(p_T^b) - K p_T^a p_T^b$$

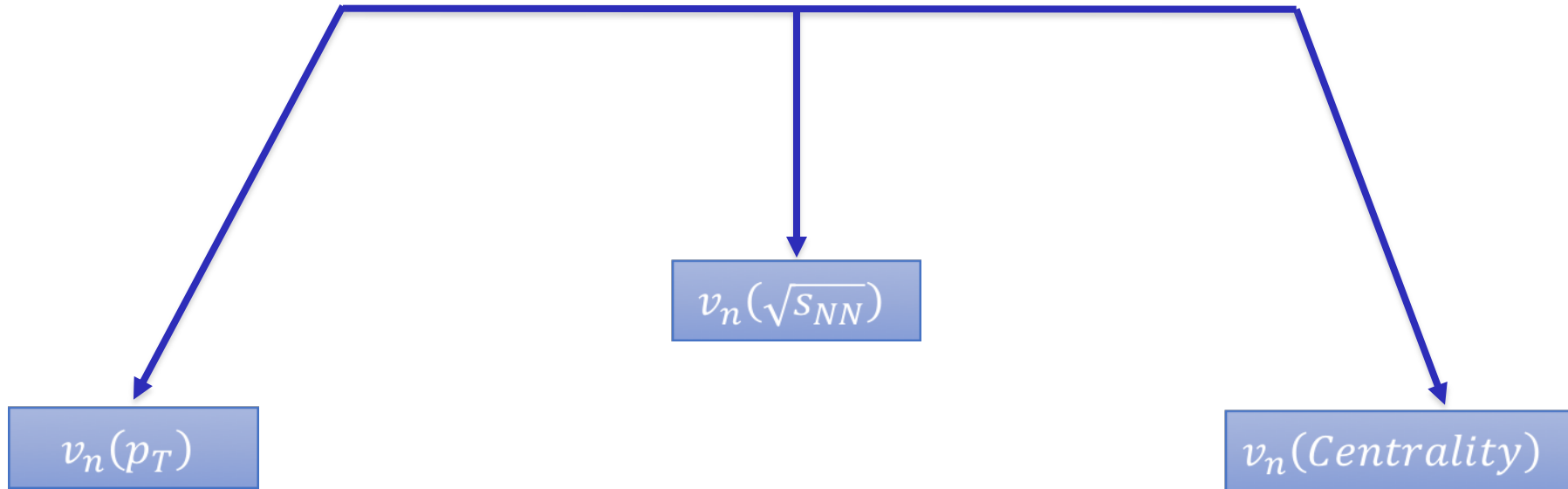
The extracted $v_1^{even}(p_T)$ and the momentum conservation parameter K at $\sqrt{s_{NN}} = 200$

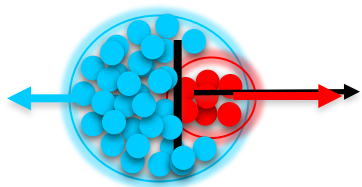


➤ The characteristic behavior of $v_1^{even}(p_T)$ shows a weak centrality dependence

➤ The momentum conservation parameter K scales as $\langle N_{ch} \rangle^{-1}$

Flow harmonics



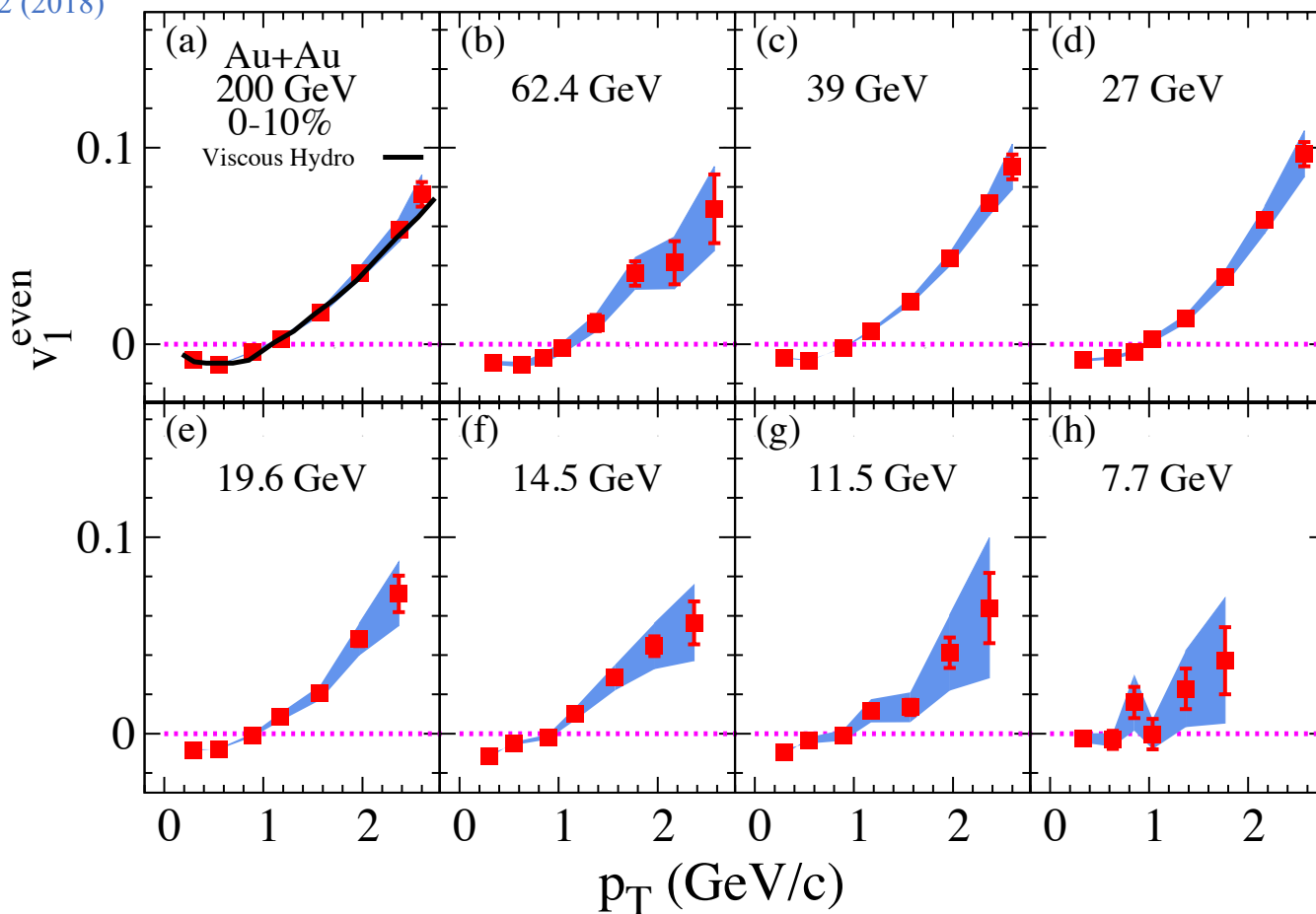


Beam Energy Dependence of v_1^{even}

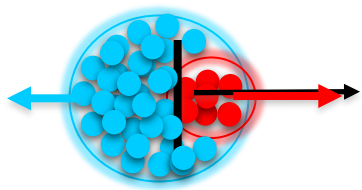
$$v_{11}(p_T^a, p_T^t) = v_1^{even}(p_T^a) v_1^{even}(p_T^t) - K p_T^a p_T^t$$

The extracted $v_1^{even}(p_T)$ at all BES energies

STAR Collaboration
PLB 784 2632 (2018)



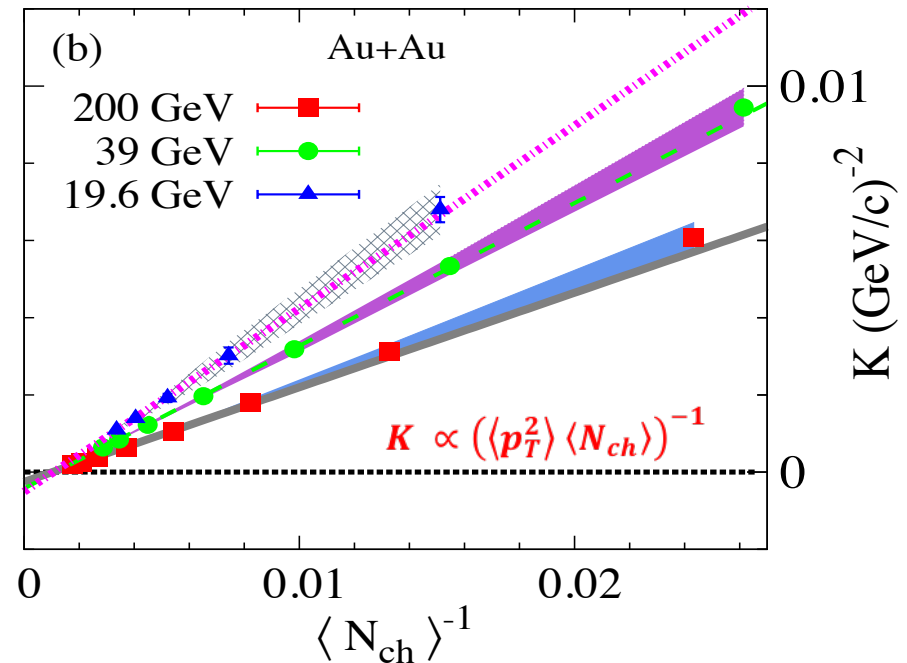
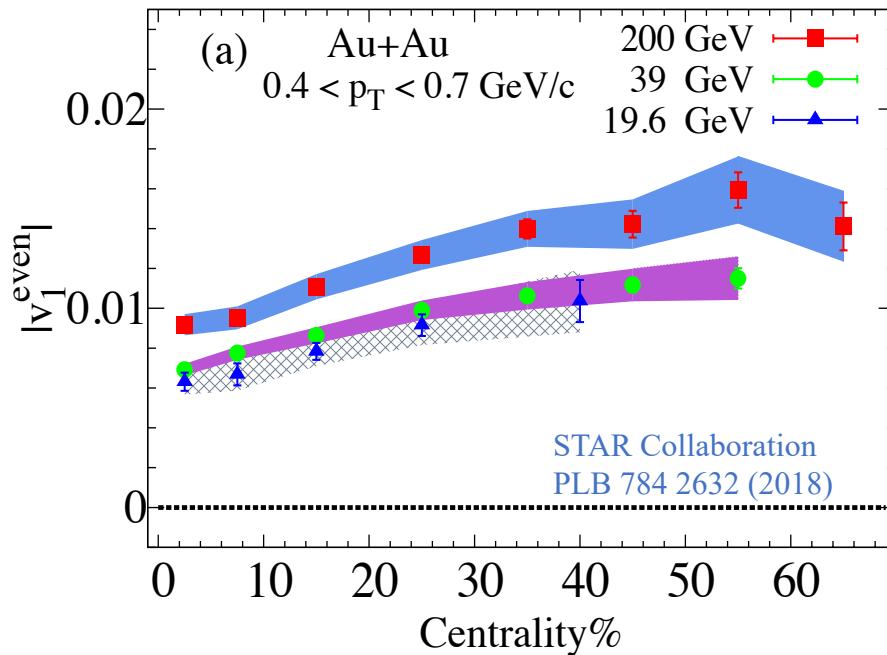
- Similar characteristic behavior of $v_1^{even}(p_T)$ at all energies
- $v_1^{even}(p_T)$ agrees with hydrodynamic calculations at 200 GeV



Beam Energy Dependence of v_1^{even}

$$v_{11}(p_T^a, p_T^t) = v_1^{even}(p_T^a)v_1^{even}(p_T^t) - K p_T^a p_T^t$$

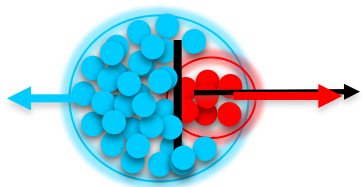
The extracted $v_1^{even}(Cent)$ and the momentum conservation parameter at different beam energies



For different beam energies;

➤ v_1^{even} increases weakly as collisions become more peripheral

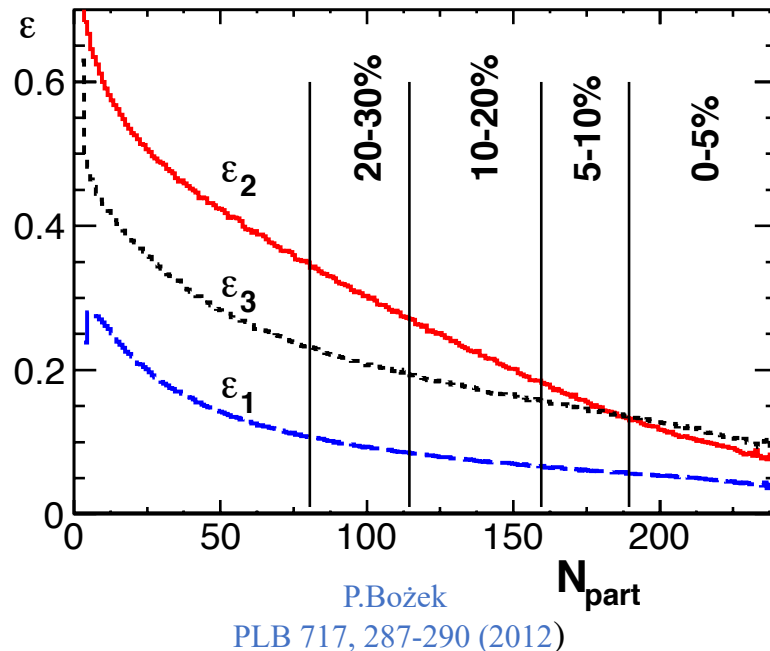
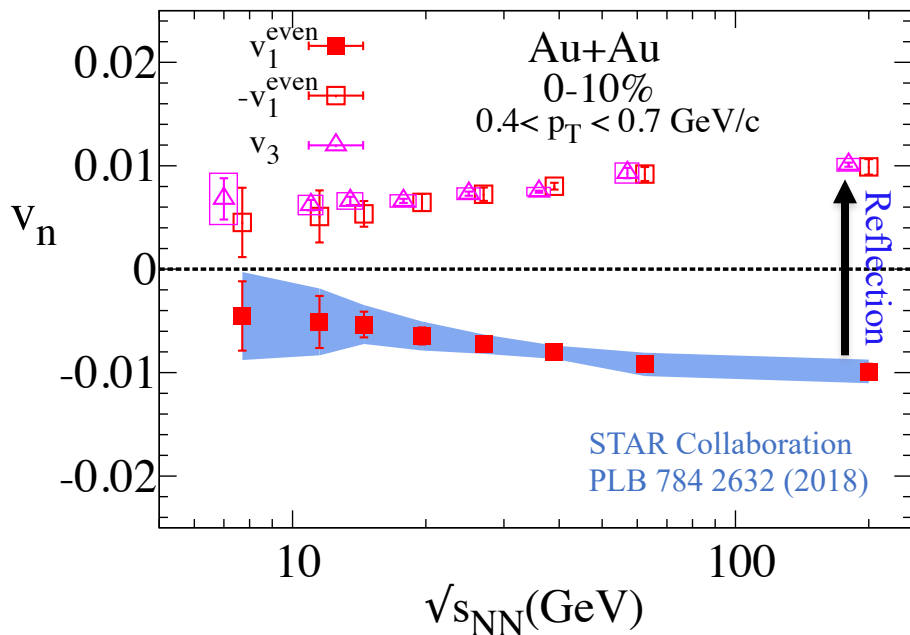
➤ Momentum conservation parameter K scales as $\langle N_{ch} \rangle^{-1}$



Beam Energy Dependence of v_1^{even}

$$v_{11}(p_T^a, p_T^t) = v_1^{even}(p_T^a)v_1^{even}(p_T^t) - K p_T^a p_T^t$$

The extracted v_1^{even} vs. $\sqrt{s_{NN}}$ at 0%-10% centrality



➤ $|v_1^{even}|$ shows similar values to v_3 at $0.4 < p_T < 0.7(GeV/c)$

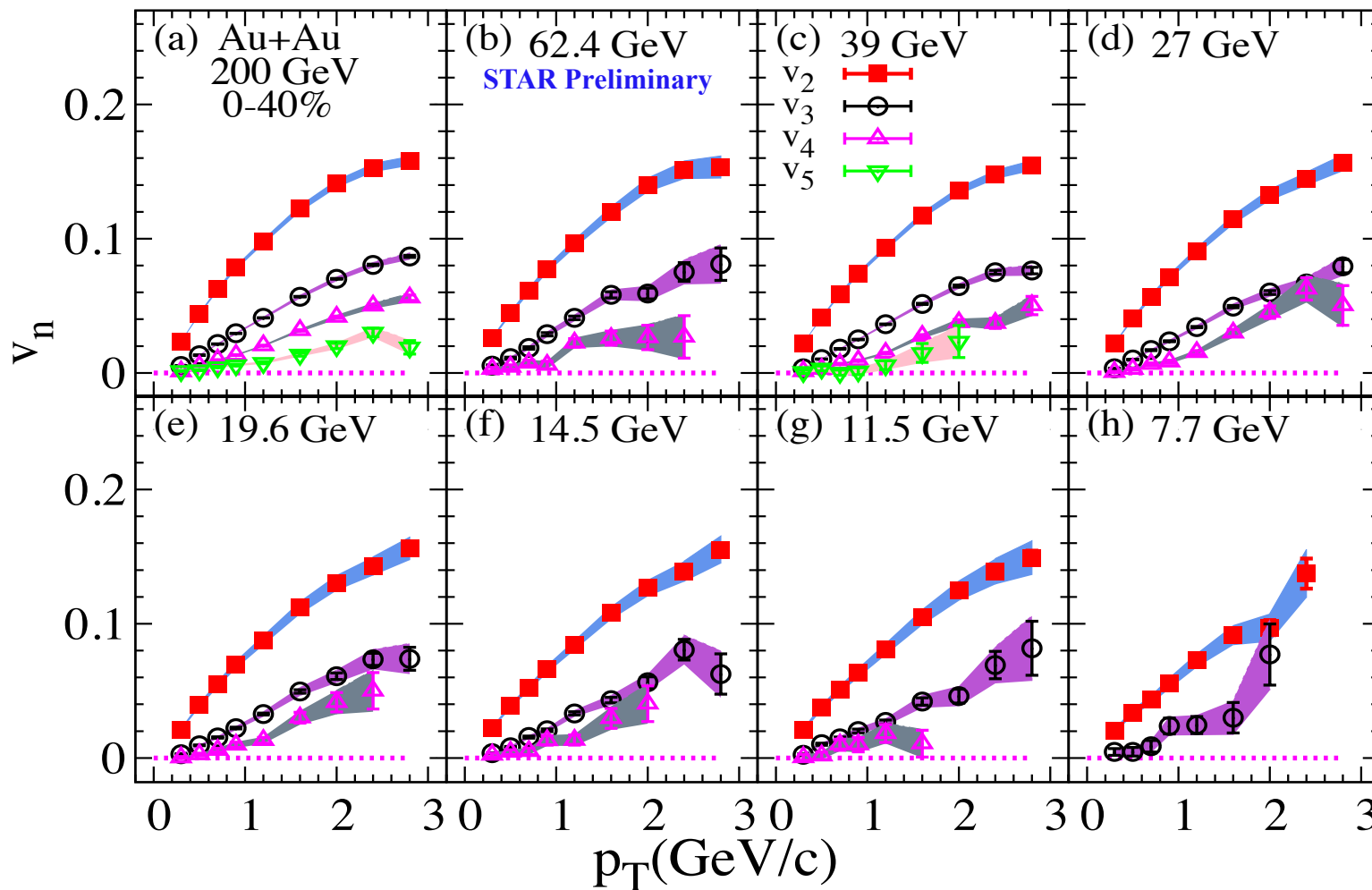
➤ $\epsilon_3 > \epsilon_1$

✓ v_3 has larger viscous damping effect than v_1^{even}

Beam Energy Dependence of v_n

$|\eta| < 1$ and $|\Delta\eta| > 0.7$

The extracted $v_{n>1}(p_T)$ at all BES energies



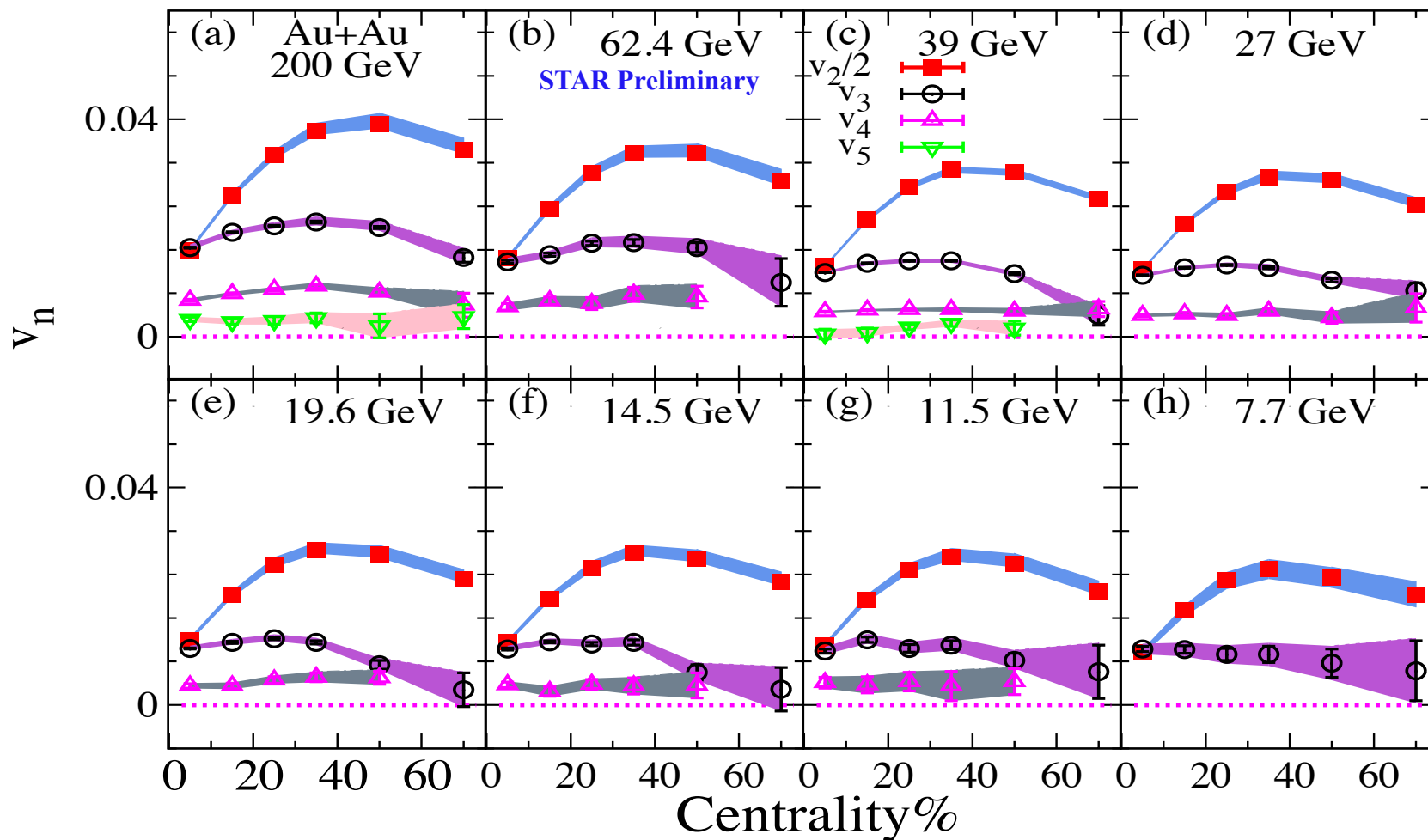
➤ $v_n(p_T)$ has similar trends for different beam energies.

➤ $v_n(p_T)$ decreases with harmonic order n .

Beam Energy Dependence of v_n

$|\eta| < 1$ and $|\Delta\eta| > 0.7$

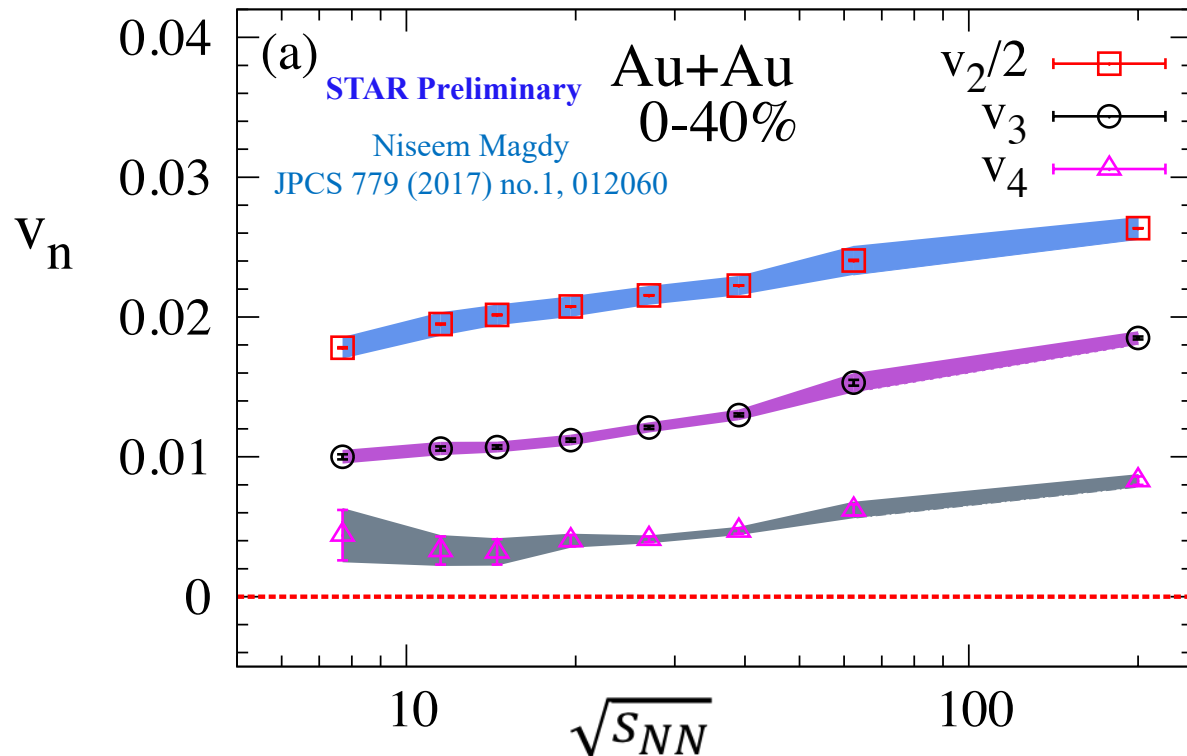
The extracted $v_{n>1}$ (Centrality) at all BES energies



- v_n (Centrality) has similar trends for different beam energies.
- v_n (Centrality) decreases with harmonic order n .

Beam Energy Dependence of v_n

The extracted $v_{n>1}$ vs. $\sqrt{s_{NN}}$ at 0-40% centrality



- $v_n(\sqrt{s_{NN}})$ shows a monotonic increase with beam-energy.
- $v_n(\sqrt{s_{NN}})$ decreases with harmonic order n (**viscous effects**).

Beam Energy Dependence of Flow Fluctuations

$$\langle\langle 2m \rangle\rangle_n = \left\langle \left\langle e^{in \sum_{j=1}^m (\phi_{2j-1} - \phi_{2j})} \right\rangle \right\rangle$$

$$c_n\{2\} = \langle\langle 2 \rangle\rangle_n$$

$$c_n\{4\} = \langle\langle 4 \rangle\rangle_n - 2 \langle\langle 2 \rangle\rangle_n \langle\langle 2 \rangle\rangle_n$$

$$c_n\{6\} = \langle\langle 6 \rangle\rangle_n - 9 \langle\langle 2 \rangle\rangle_n \langle\langle 4 \rangle\rangle_n + \langle\langle 2 \rangle\rangle_n^3$$

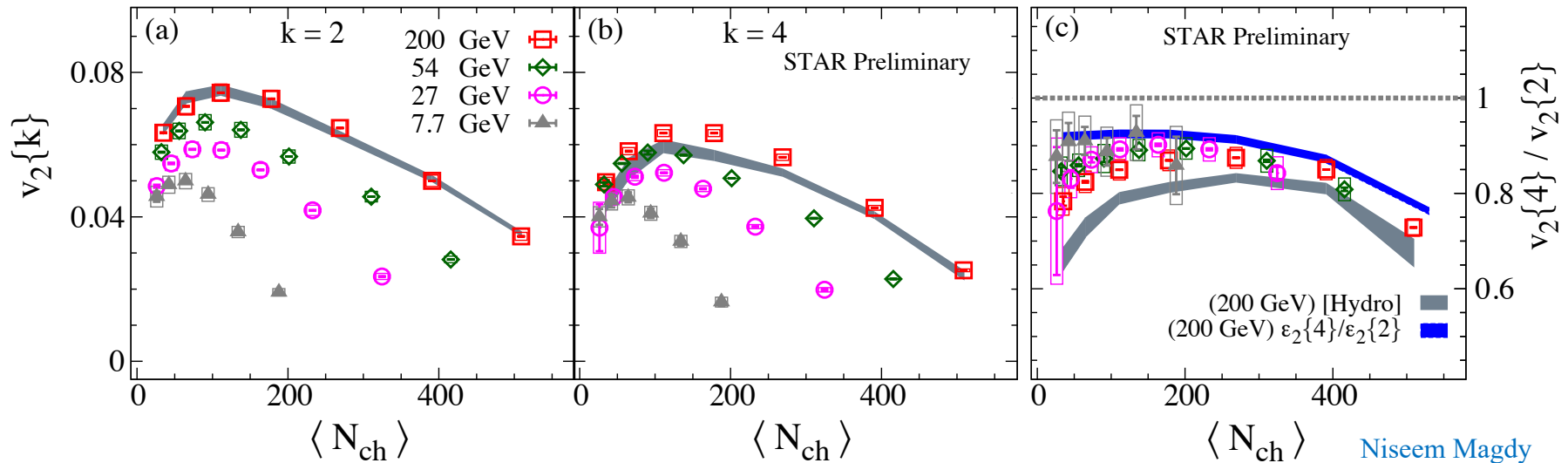
$$v_n^2\{2\} = c_n\{2\}$$

$$v_n^4\{4\} = -c_n\{4\}$$

$$v_n^6\{6\} = \frac{1}{6} c_n\{6\}$$

Beam Energy Dependence of Flow Fluctuations

Hydro
P. Alba, et al.
Phys. Rev. C 98, 034909 (2018)



Niseem Magdy
NPA 982 (2019) 255-258

➤ The elliptic flow fluctuations, $\frac{v_2\{4\}}{v_2\{2\}}$ show:

- ✓ Modest dependence on centrality
- ✓ Weak dependence on beam energy

➤ The model calculations for $(v_2\{4\}/v_2\{2\})$ and $(\epsilon_2\{4\}/\epsilon_2\{2\})$ bracket the data at 200 GeV

Viscous Attenuation

- Acoustic ansatz
 - ✓ Sound attenuation in the viscous matter reduces the magnitude of $v_{n=2,3}$.

- Anisotropic flow attenuation:

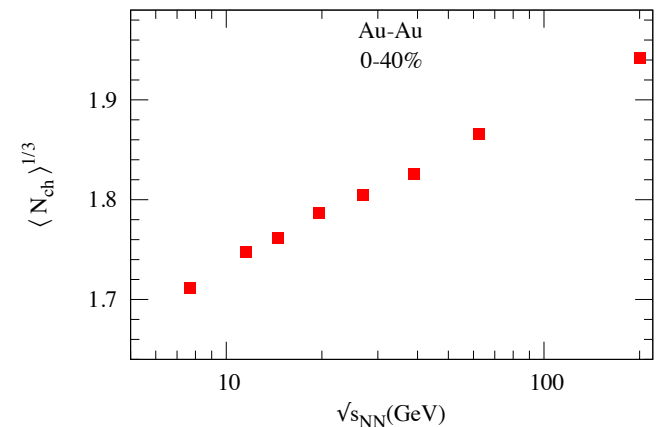
$$v_n \propto k \varepsilon_n, \quad k = e^{-\beta n^2}$$

$$\frac{v_n}{\varepsilon_n} \propto e^{-\beta n^2}, \quad \beta \propto \frac{\eta}{s} \frac{1}{RT}$$

- From macroscopic entropy considerations:

$$S \sim (RT)^3 \sim \langle N_{Ch} \rangle \text{ then } RT \sim \langle N_{Ch} \rangle^{1/3}$$

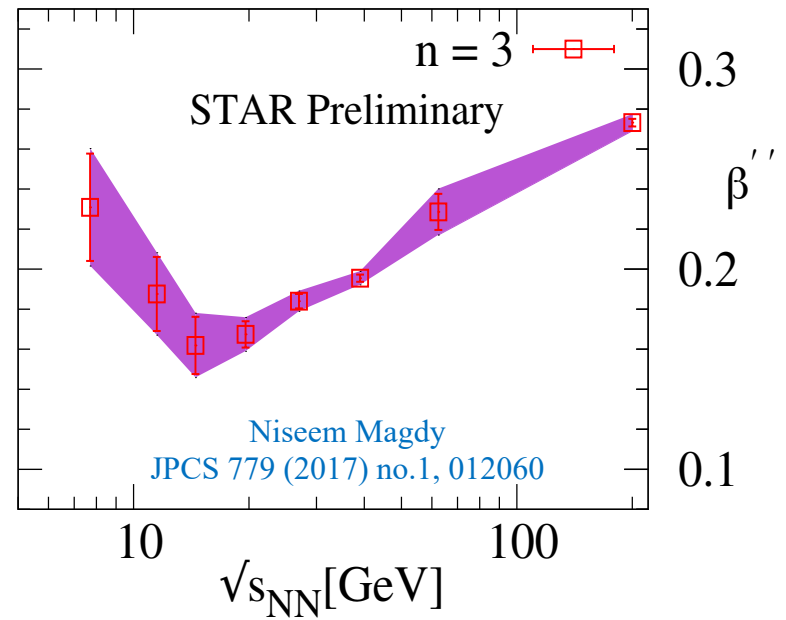
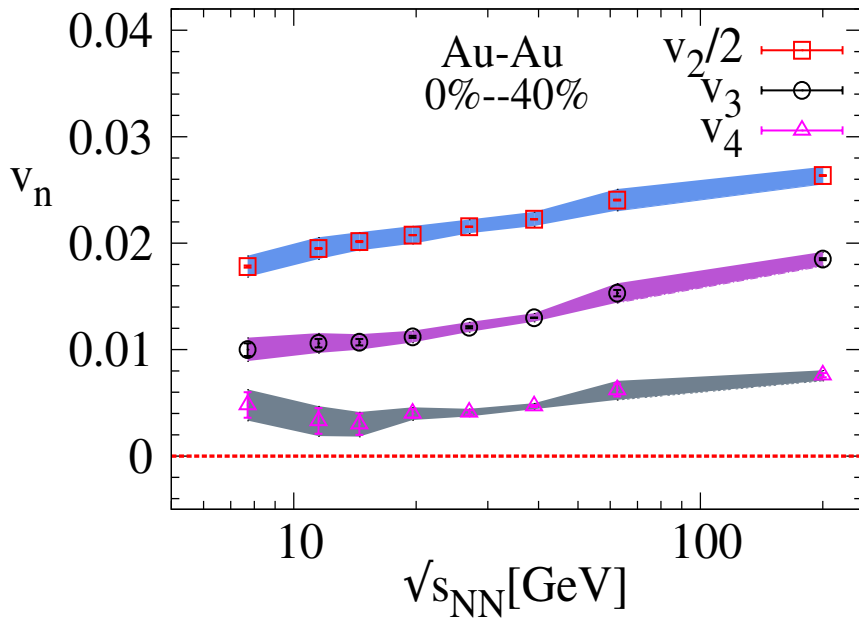
$$\ln \left(\frac{v_n}{\varepsilon_n} \right) \propto - \left(\frac{\eta}{s} \right) \langle N_{Ch} \rangle^{-1/3}$$



What are the respective roles of ε_n and its fluctuations, system size ($\langle N_{Ch} \rangle^{-1/3}$) and transport coefficient $\frac{\eta}{s}(T, \mu_B)$?

Viscous coefficient

$$\beta'' = \ln\left(\frac{v_n^{1/n}}{v_2^{1/2}}\right) \langle N_{Ch} \rangle^{1/3} (n-2)^{-1} = A \frac{\eta}{s}$$



➤ The viscous coefficient shows a non-monotonic behavior with beam-energy

Summary

Comprehensive set of flow measurements were studied for Au+Au collision system at all BES energies with one set of cuts.

➤ For v_n :

- ✓ v_n vs centrality indicates a similar trend for different beam energies.
- ✓ Momentum conservation parameter K scales as $\langle N_{\text{ch}} \rangle^{-1}$
- ✓ $v_n(\sqrt{s_{NN}})$ shows a monotonic increase with beam-energy.

➤ The elliptic flow fluctuations, $\frac{v_2\{4\}}{v_2\{2\}}$ show:

- ✓ Modest dependence on centrality
- ✓ Weak dependence on beam energy

➤ The viscous coefficient shows a non-monotonic behavior with beam-energy

For different beam energies, these comprehensive measurements provide additional constraints for theoretical models, as well as $\frac{\eta}{s}$ extraction.

THANK YOU