

UPC2025: The second international workshop on the physics
of Ultra Peripheral Collisions
June 9-13, 2025

Investigating Spin Interference in photonuclear $\gamma A \rightarrow \pi^+ \pi^-$ and $\gamma\gamma \rightarrow \pi^+ \pi^-$ at STAR

Samuel Corey
for the STAR Collaboration



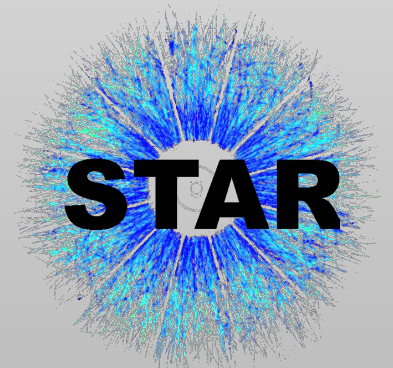
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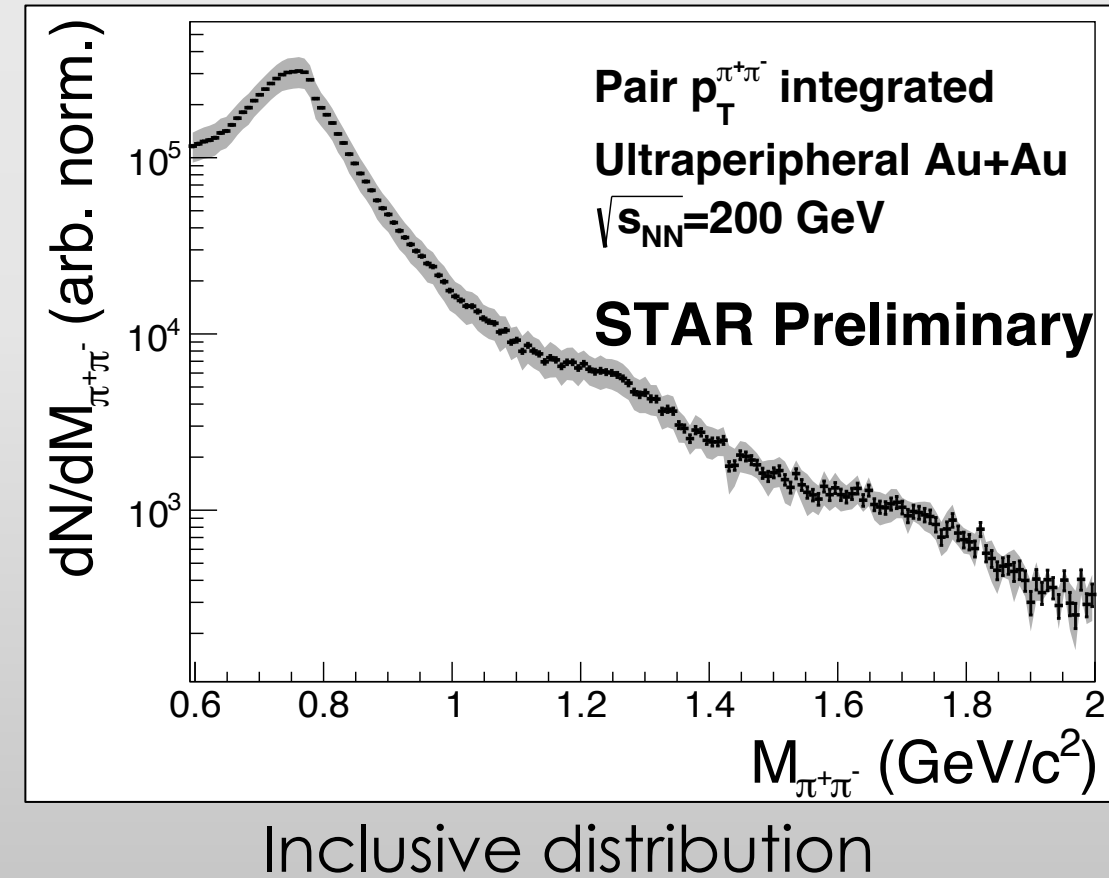
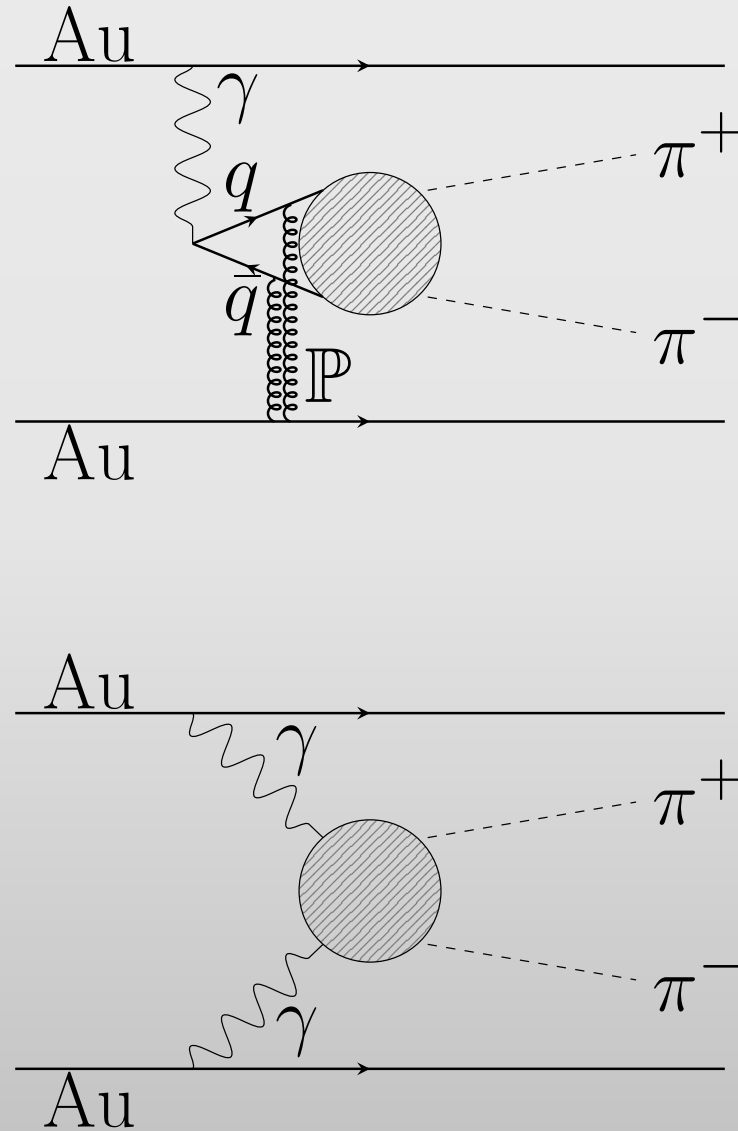


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Motivation

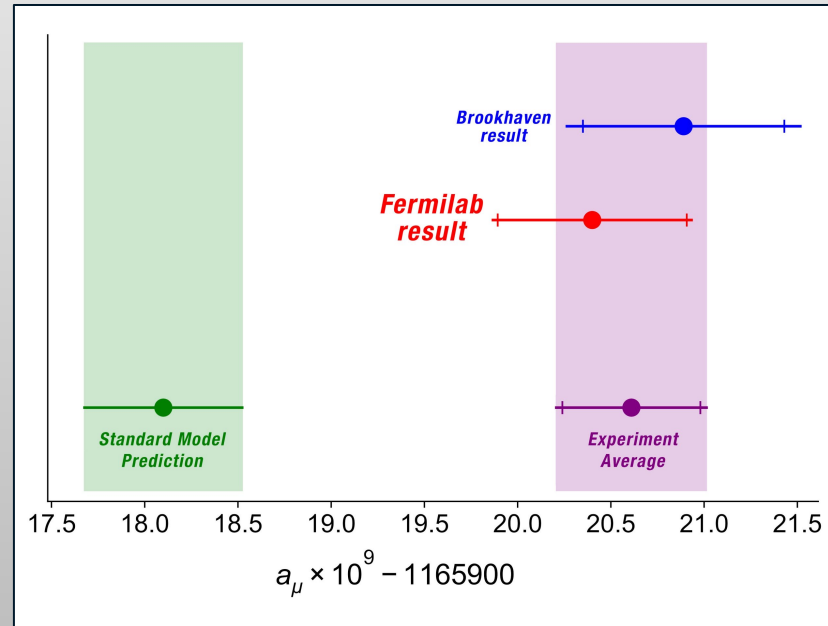
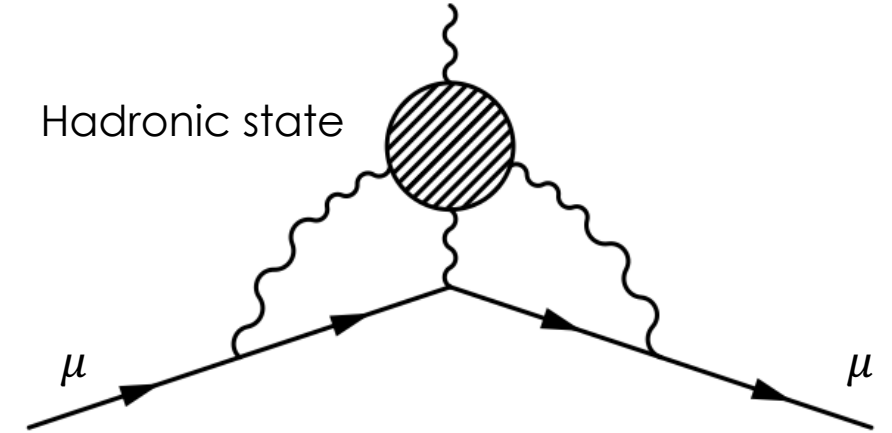
- In Ultraperipheral Collisions (UPCs), $\pi^+ \pi^-$ primarily from $\gamma A \rightarrow X \rightarrow \pi^+ \pi^-$.
- $\gamma\gamma \rightarrow \pi^+ \pi^-$ not yet measured in hadronic systems.
- $\gamma\gamma \rightarrow \pi^+ \pi^-$ makes up tiny fraction of overall $\pi^+ \pi^-$ cross section.



Hadronic Light-by-Light and a_μ

- Hadronic light-by-light (HLbyL) is one the two dominant theoretical uncertainties on a_μ .
- Related to $\gamma\gamma \rightarrow \pi^+\pi^-$ by the optical theorem.
- Previous measurements of $\gamma\gamma \rightarrow \pi\pi$ are from e^-e^+ collisions.
- In UPC, we have quasi-real photons and larger mass range

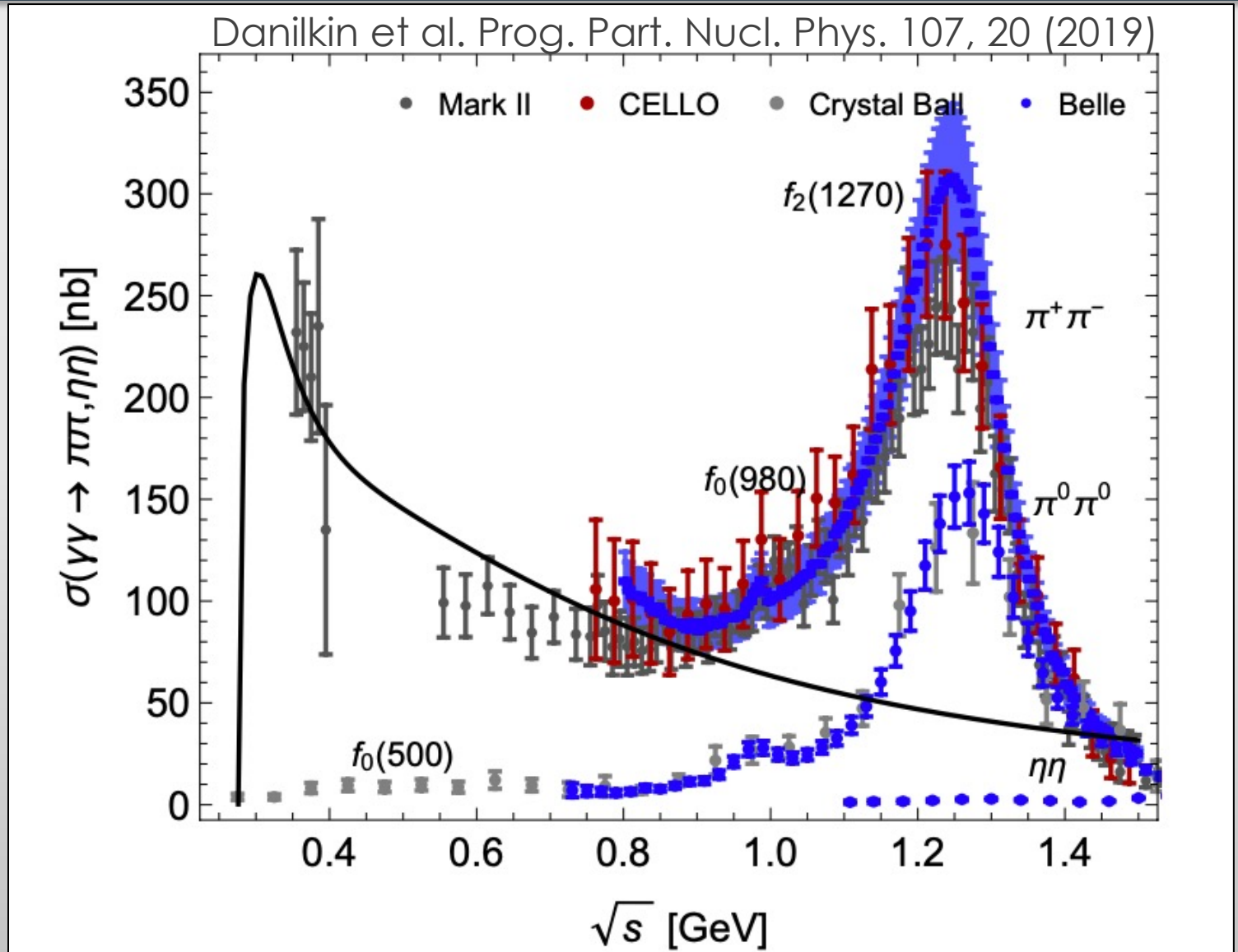
Melo-Porras et al., Particles 2024, 7(2), 327 – 381



First Results from Fermilab's Muon g-2 Experiment Strengthen Evidence of New Physics; Fermilab, 7 April 2021.

Hadronic Light-by-Light and a_μ

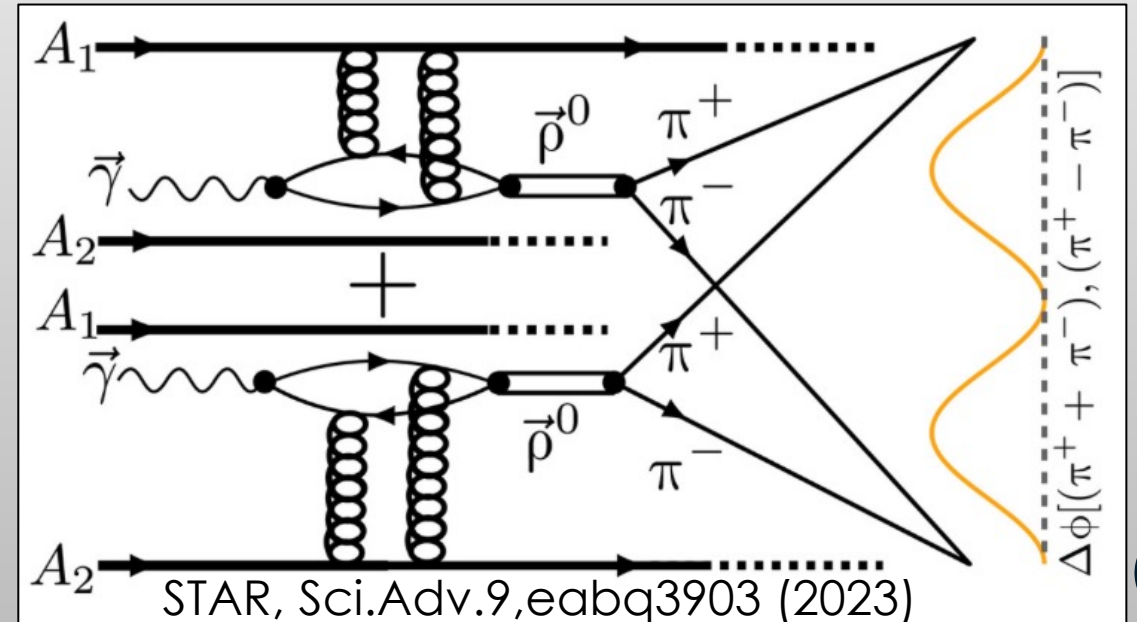
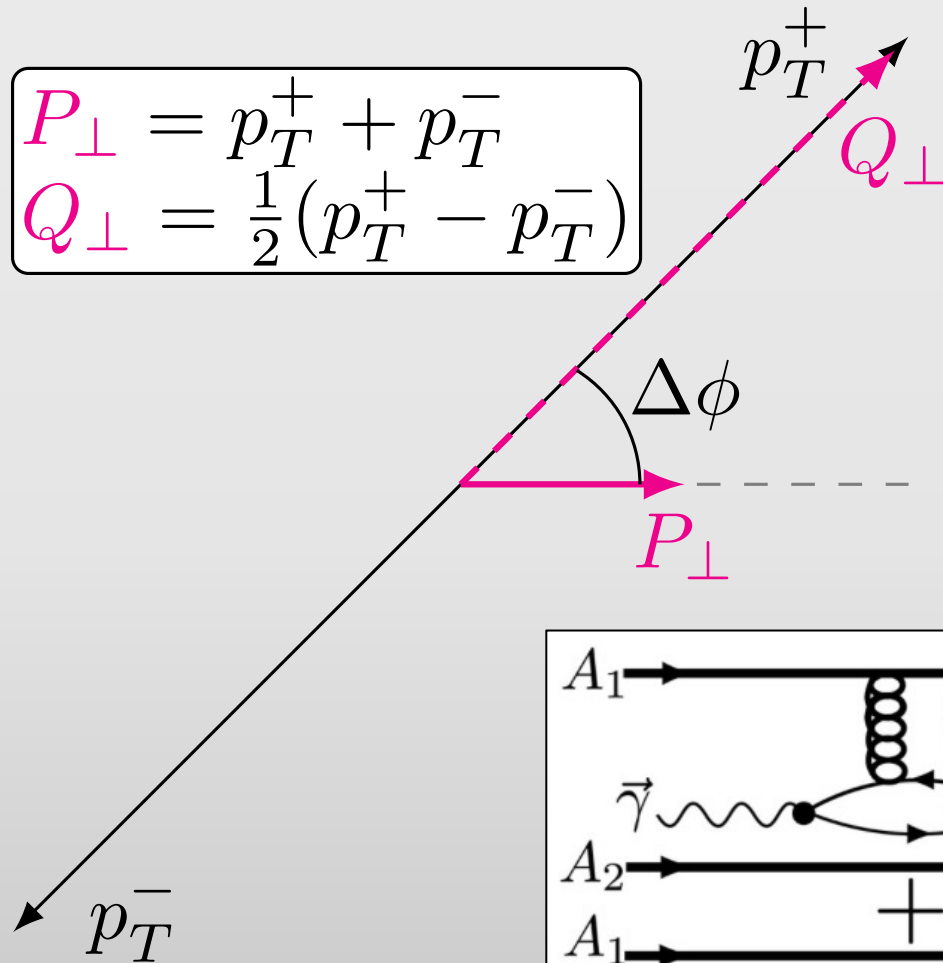
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Spin Interference in UPCs

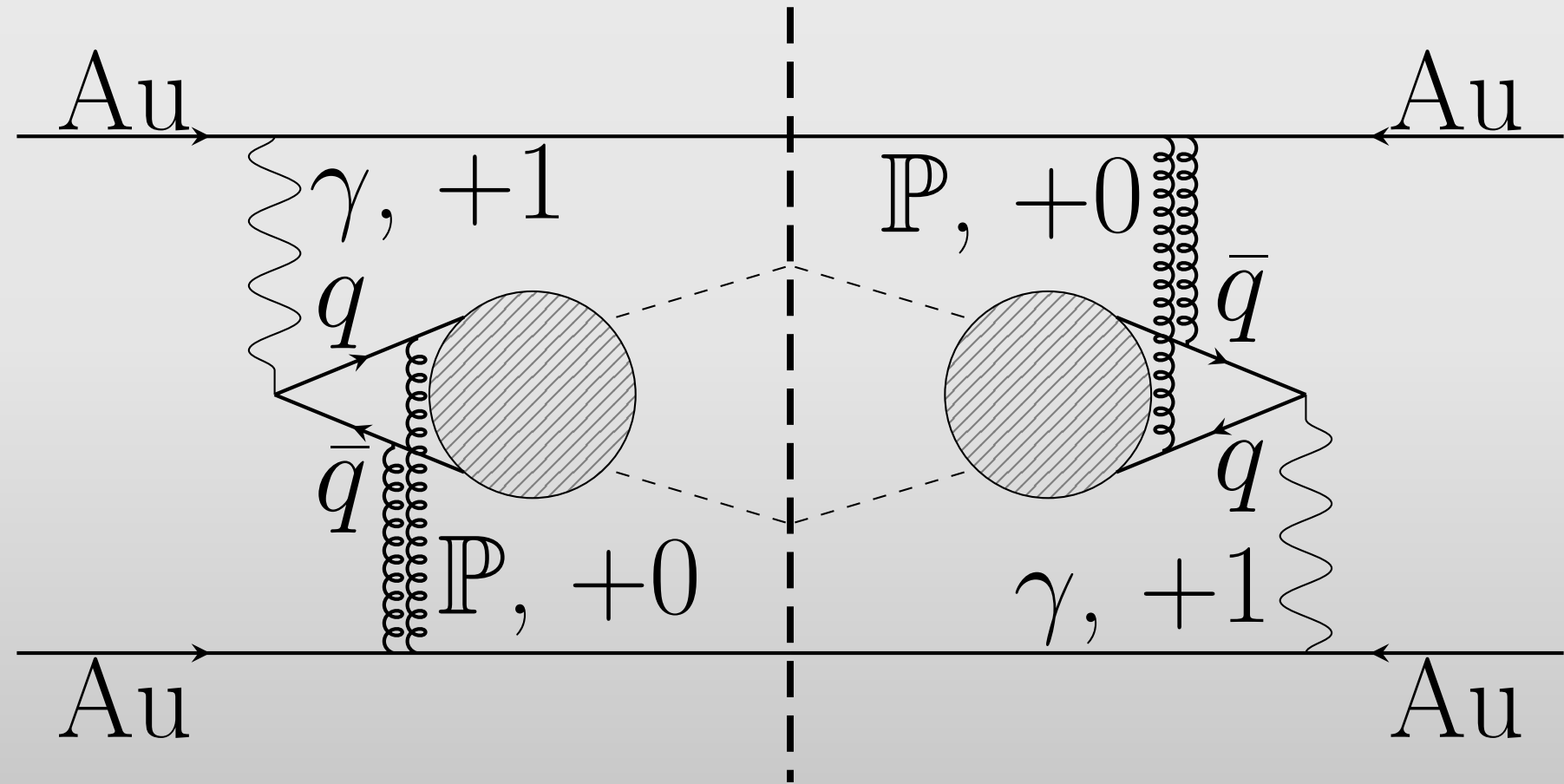
- In γA events, one nucleus acts as the emitter and the other as the target.
- Double slit-like interference leads to a $\cos(2\Delta\phi)$ angular anisotropy.

(Brandenburg et al.,
Phys. Rev. Research 7,
013131 (2025))



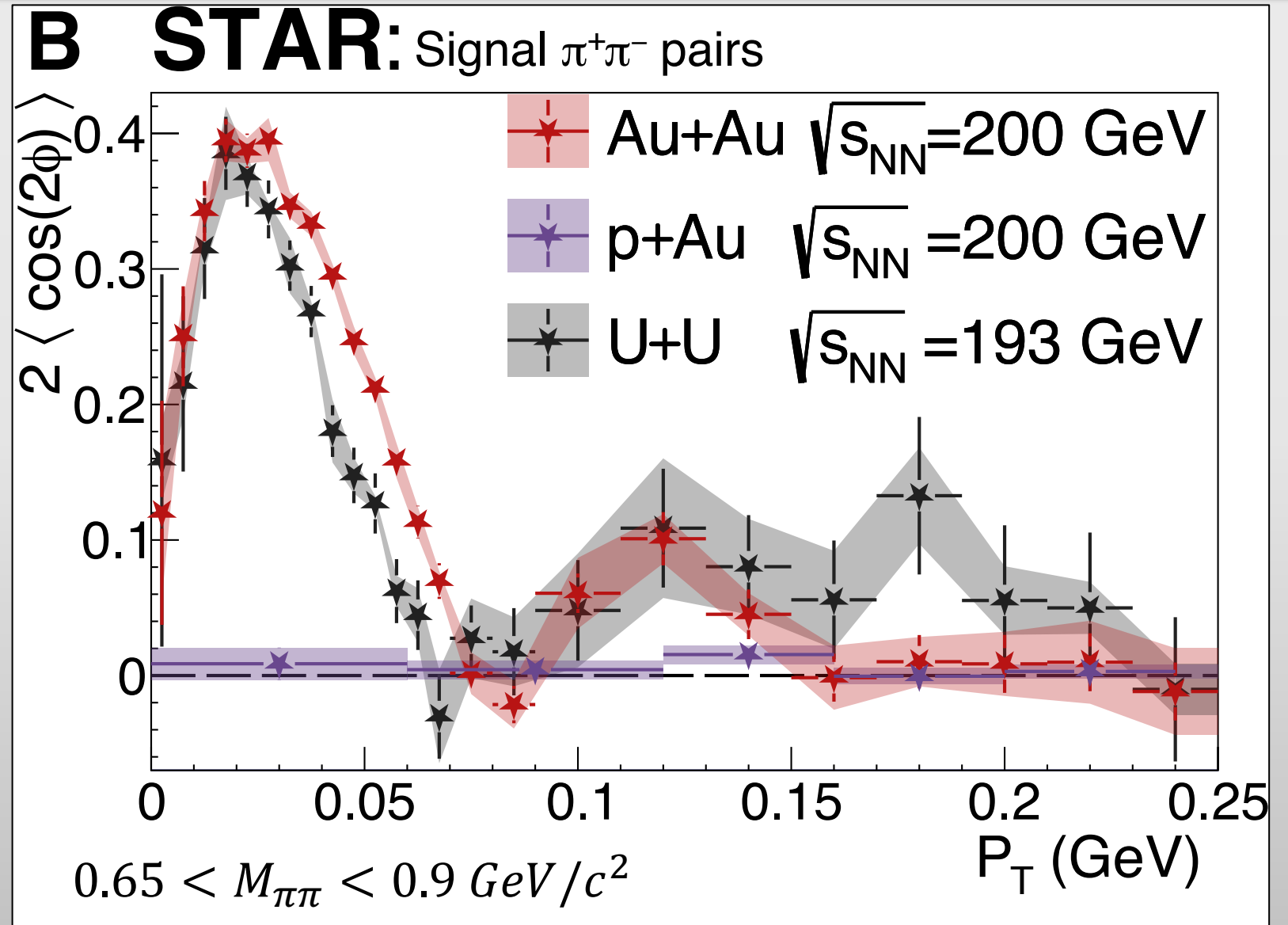
Previous $\langle 2\cos(2\Delta\phi) \rangle$ measurement

- Photon spin encodes into final state orbital angular momentum: $\cos(2\Delta\phi)$ modulation.
- Present in A+A, but not p+A--photon emission scales with Z^2



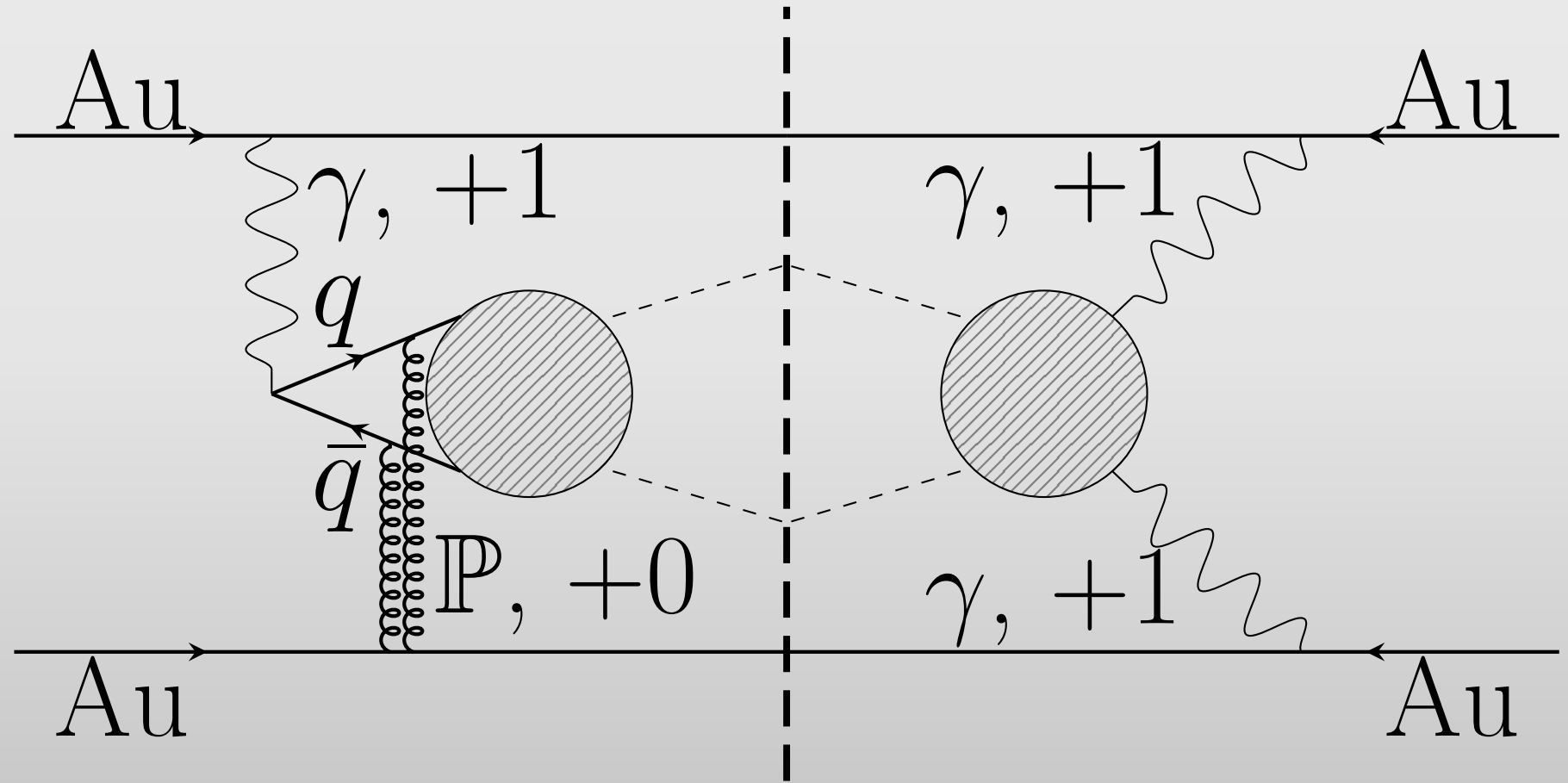
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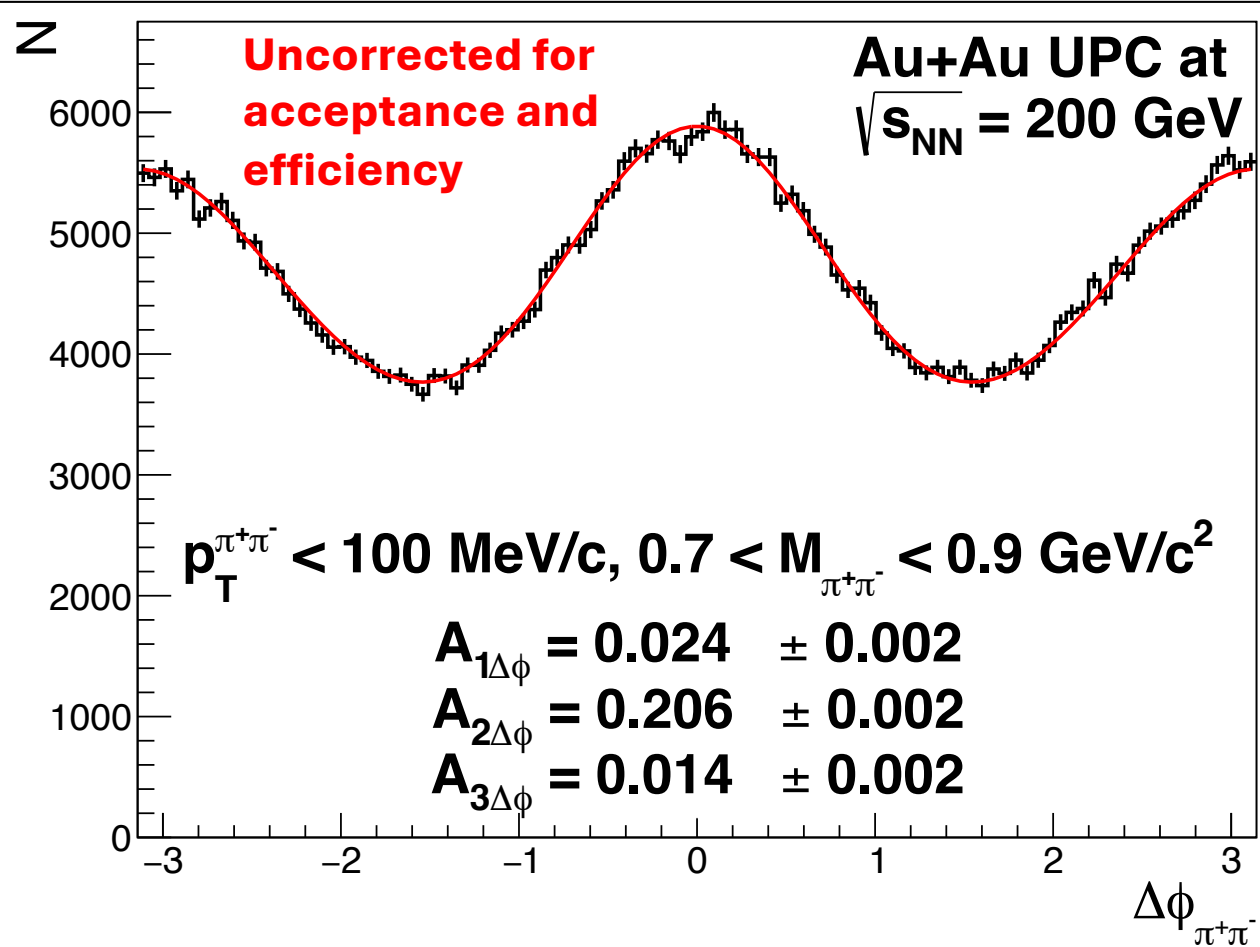


Spin Interference between γA and $\gamma\gamma$

- Interference between γA and $\gamma\gamma$ also present.
- This interference is expected to produce a $\cos(\Delta\phi)$ and $\cos(3\Delta\phi)$ anisotropy.



$\Delta\phi$ distribution



- Uncorrected $\Delta\phi$ distribution fit with:

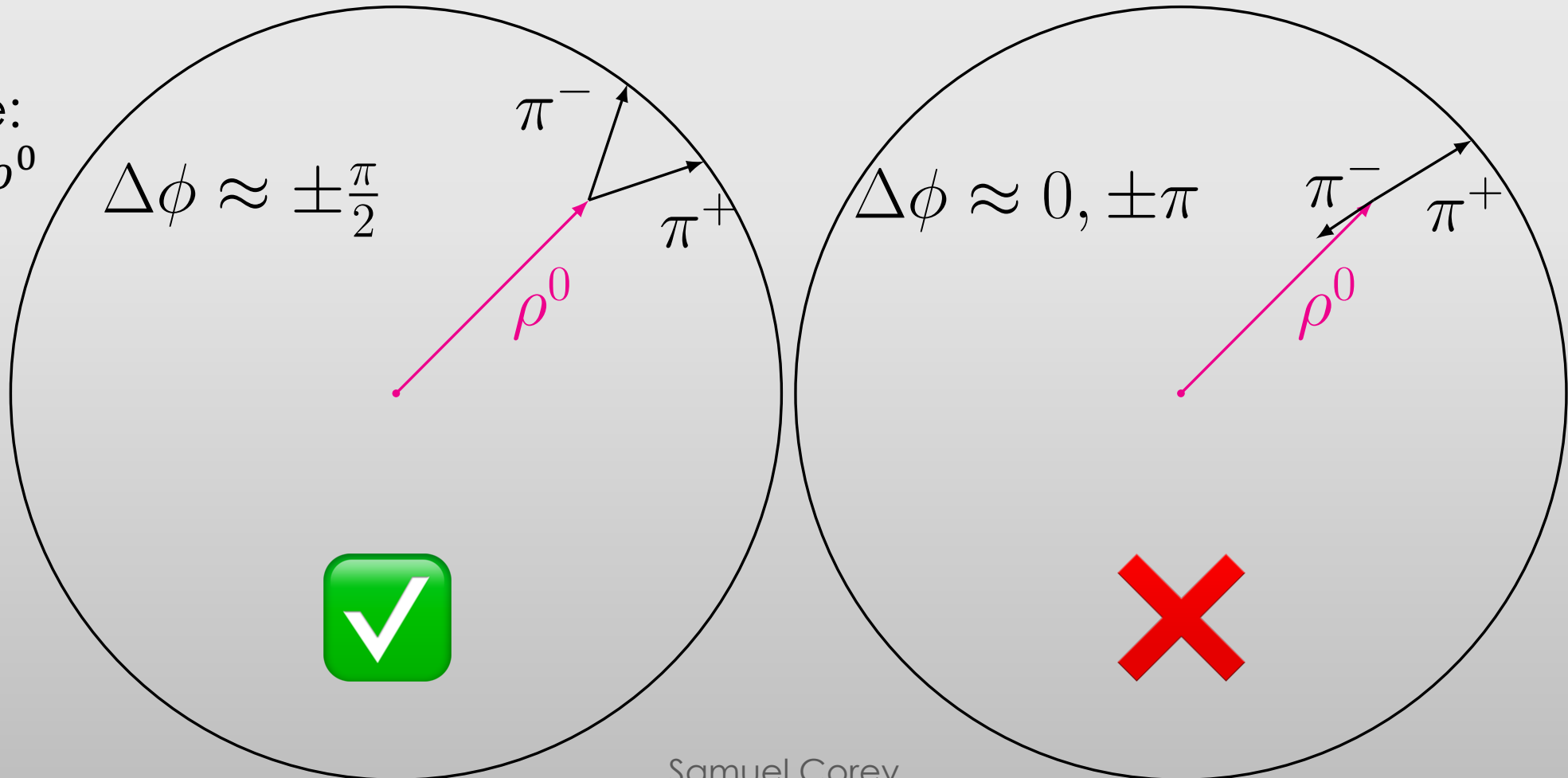
$$C[1 + A_{1\Delta\phi} \cos(\Delta\phi) + A_{2\Delta\phi} \cos(2\Delta\phi) + A_{3\Delta\phi} \cos(3\Delta\phi)]$$

- Large and highly significant $A_{2\Delta\phi}$.
- $A_{1,3\Delta\phi} \neq 0$: first hint that we have $\gamma\gamma \rightarrow \pi^+\pi^-$.

$\Delta\phi$ efficiency and acceptance correction

- Want to measure $A_{n\Delta\phi} = \langle 2\cos(n\Delta\phi) \rangle$ as a function of $p_T, M_{\pi\pi}$.
- $\Delta\phi$ has unique $p_T, M_{\pi\pi}$ -dependent acceptance effects.

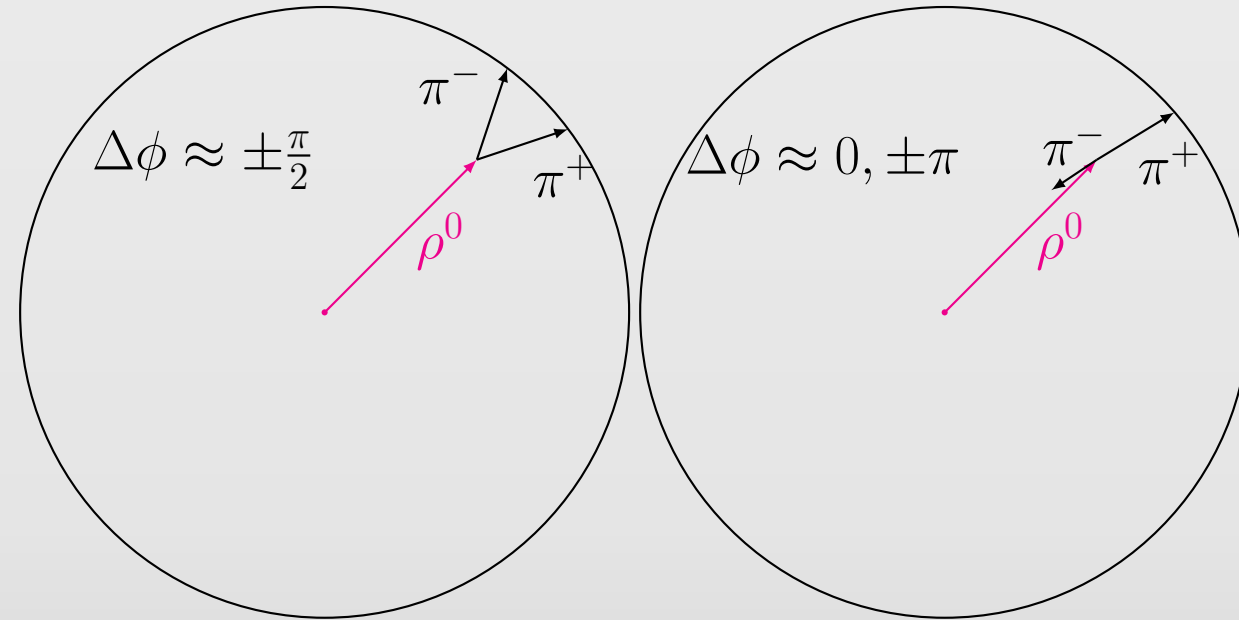
Example:
High $p_T \rho^0$



$\Delta\phi$ efficiency and acceptance correction

- Consider γ, α, ω as measured, true, and distortion $A_{n\Delta\phi}$ respectively.

$$\gamma_n = \frac{\int_{-\pi}^{\pi} \alpha(\Delta\phi) \omega(\Delta\phi) \cos(n\Delta\phi) d\Delta\phi}{\int_{-\pi}^{\pi} \alpha(\Delta\phi) \omega(\Delta\phi) d\Delta\phi}$$

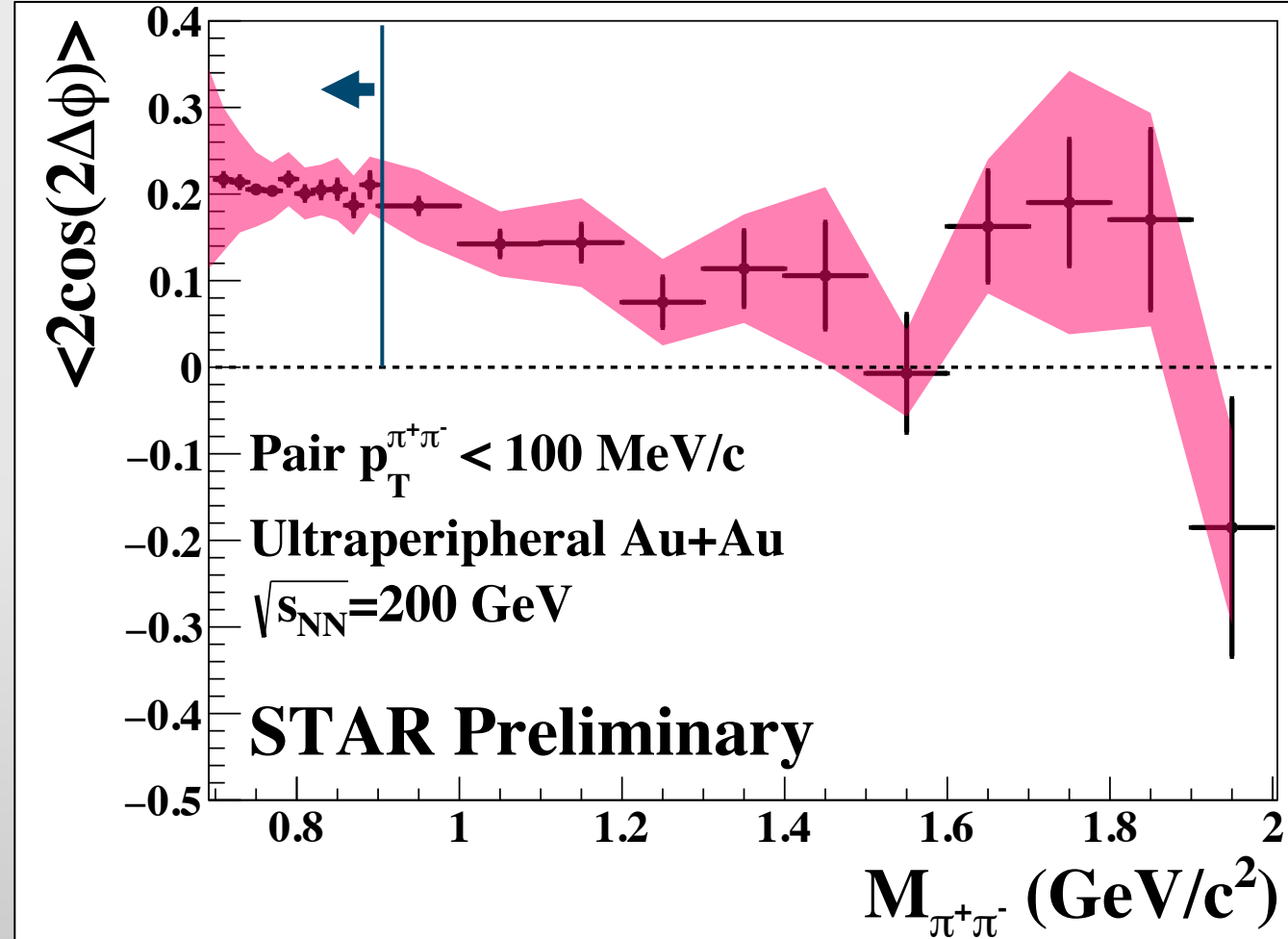
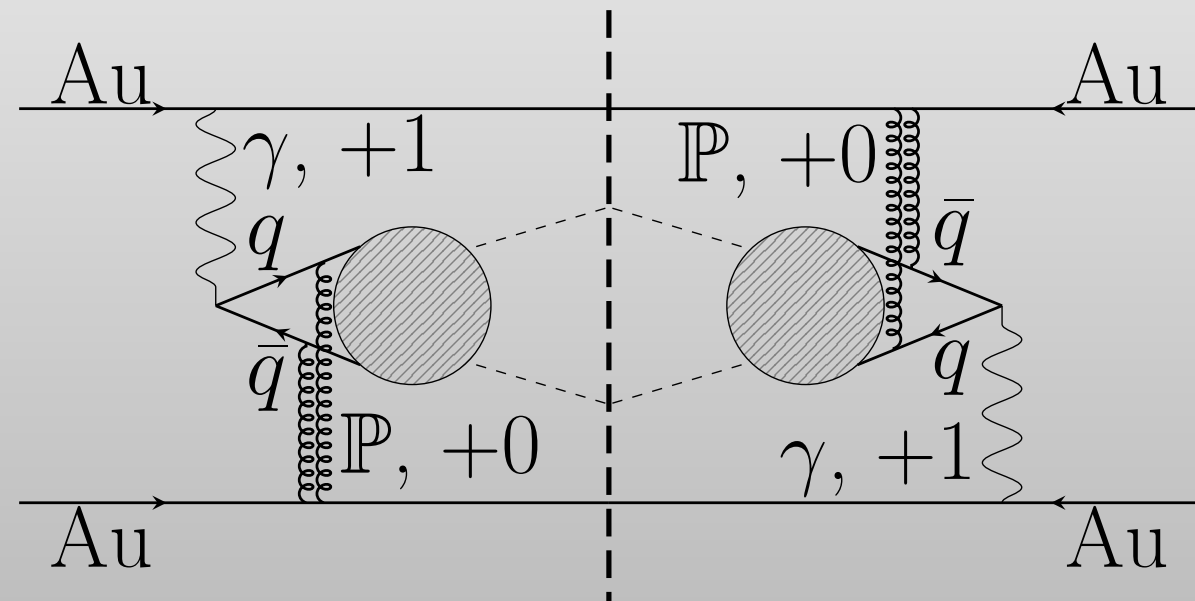


- Expand $\alpha(\Delta\phi), \omega(\Delta\phi)$ as Fourier series, and assume only one term nonzero.
- Computing integral and inverting for α_n leaves:

$$\alpha_n = \frac{-2(\gamma_n - \omega_n)}{\gamma_n \times \omega_n - 2}$$

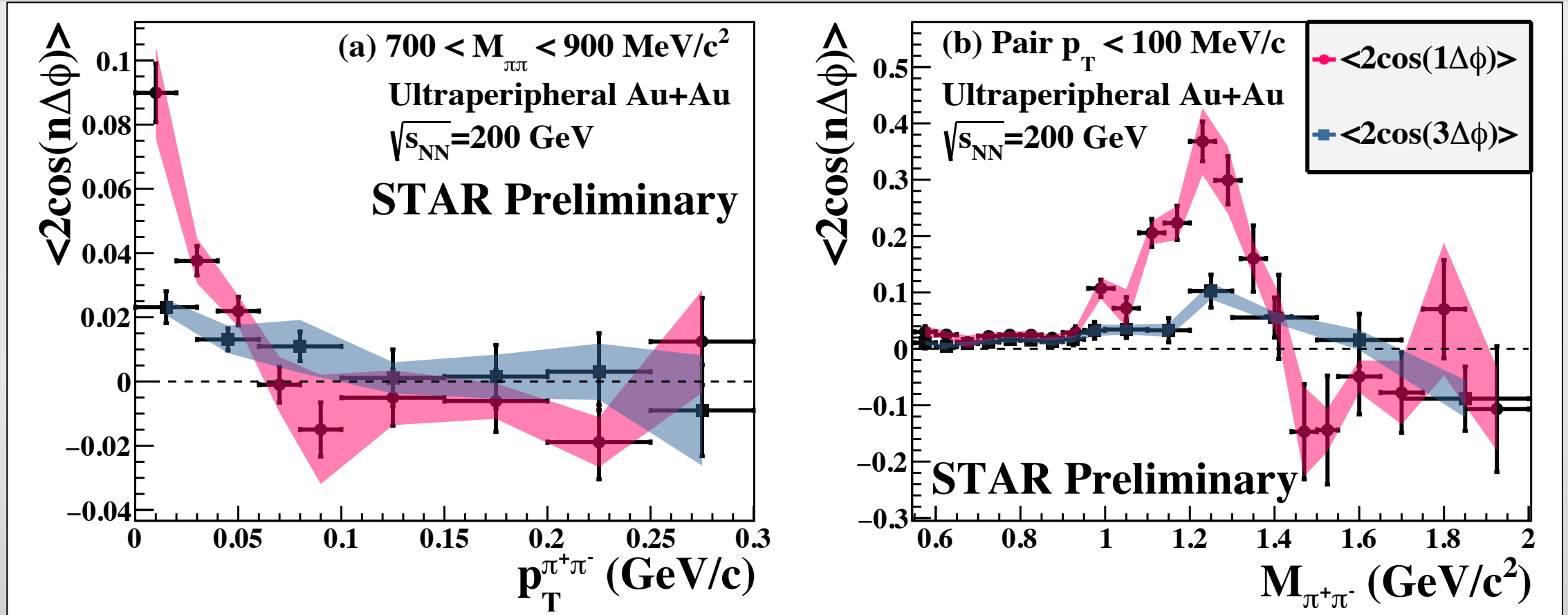
$\langle 2\cos(2\Delta\phi) \rangle$ vs. $M_{\pi\pi}$

- ω_n calculated with toy model using STAR acceptance.
- Correction applied in each bin in $p_T, M_{\pi\pi}$, then projected.



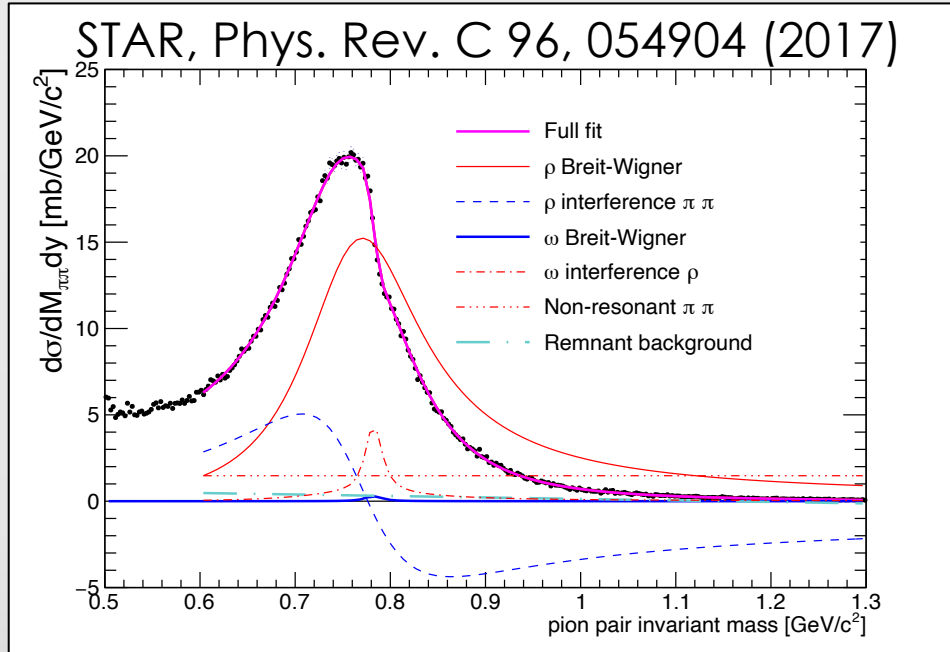
Left of solid line: previous STAR result

$\langle 2\cos(1,3\Delta\phi) \rangle$ measurement

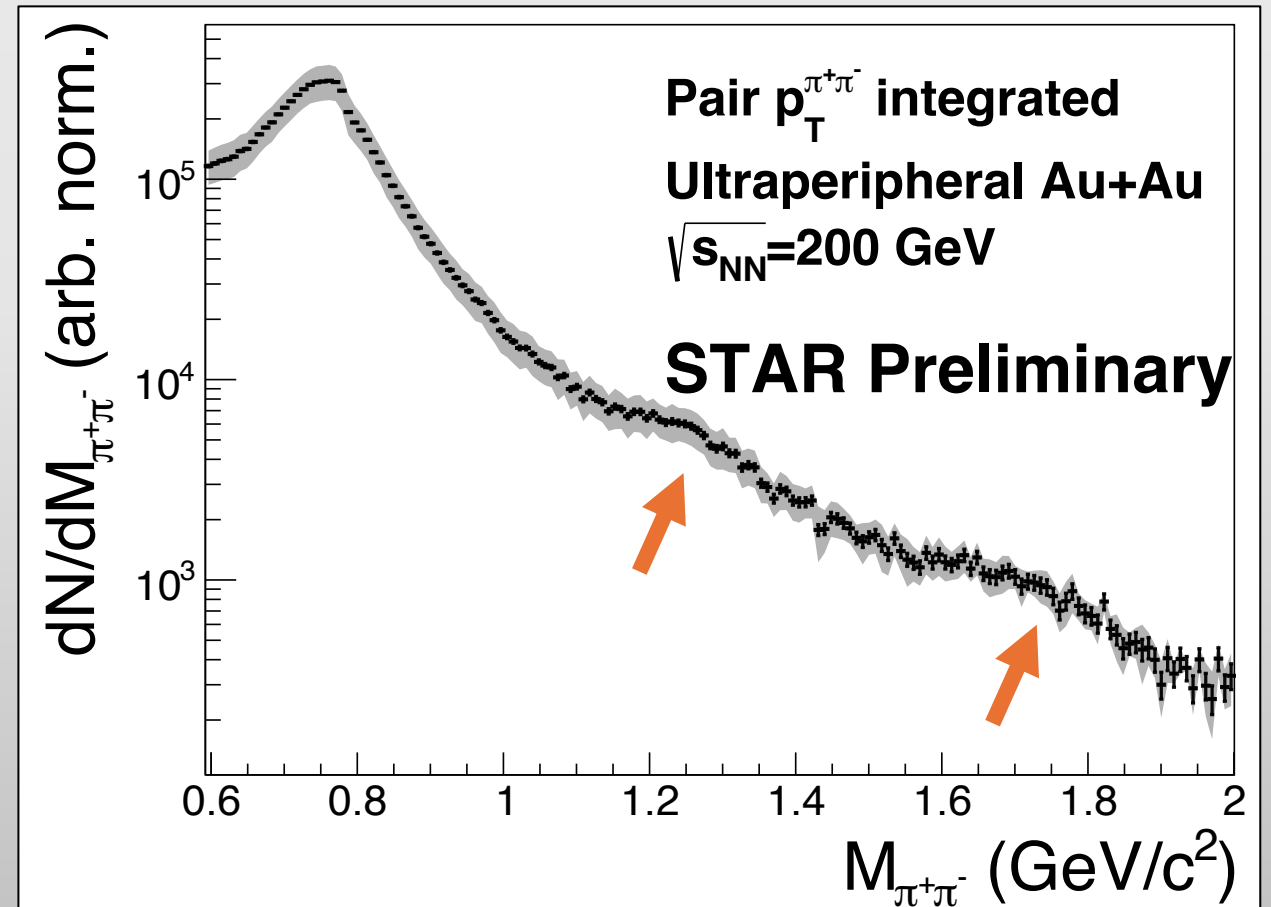


- Interference between $\gamma A \rightarrow \rho^0 \rightarrow \pi^+ \pi^-$ and non-resonant $\gamma\gamma \rightarrow \pi^+ \pi^-$ near ρ^0 mass.
- Large feature near $1270 \text{ MeV}/c^2$: **$f_2(1270)$ resonance**

UPC 2-pion invariant mass distribution



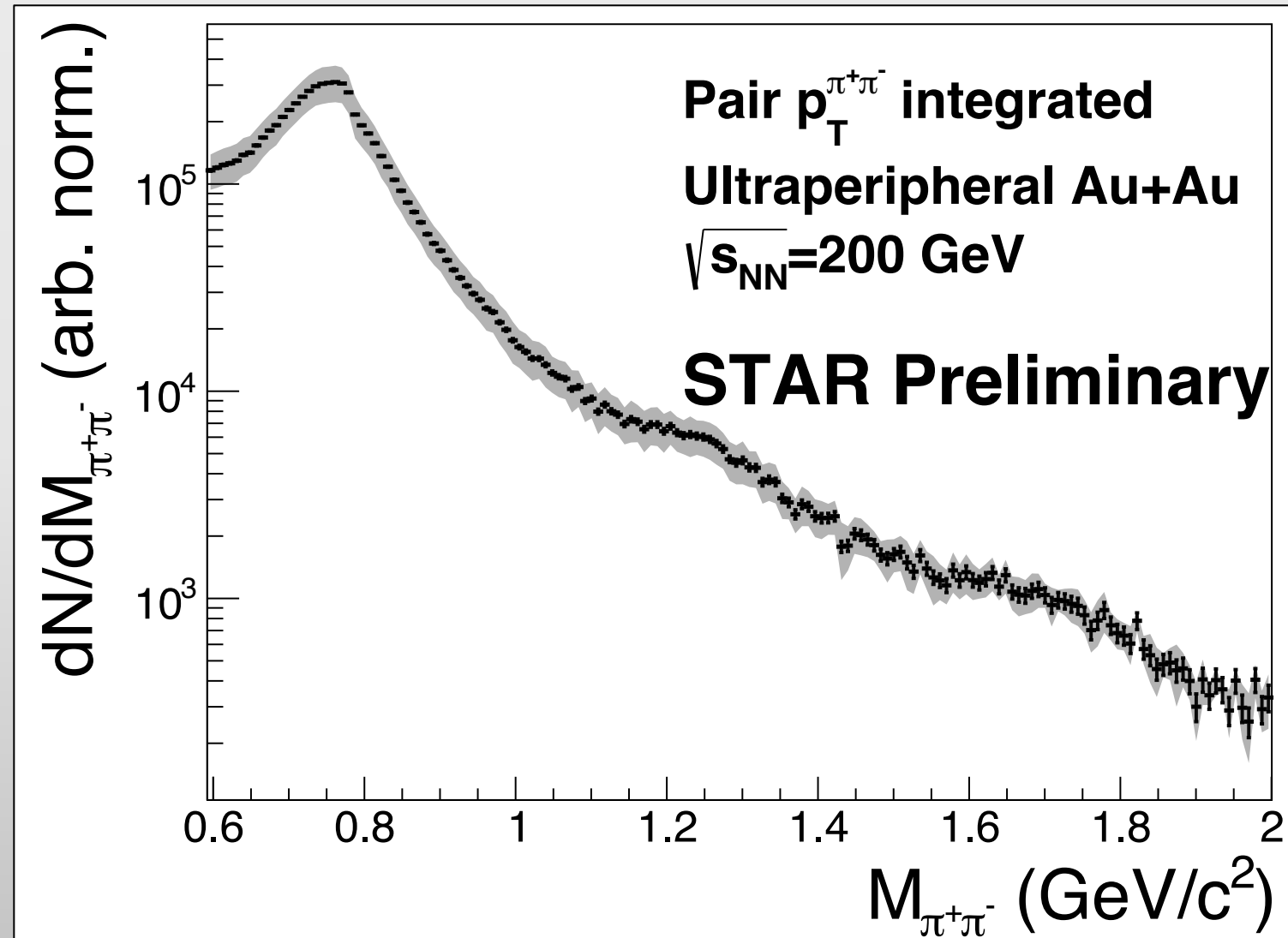
- To apply this to this analysis' mass range, we need to add **two more resonances**



- Previous STAR fit includes:
 $|\rho^0 + \omega + B_{\pi\pi}|^2$
- Where $B_{\pi\pi}$ is the Drell-Söding process, modeled by a constant.
- The interference terms are very important to determine the shape.

Including spin information

- Without spin information, nature of higher resonances is ambiguous.
- Examples:
 - $f_2(1270)$ vs. $\rho(1450)$
 - $\rho(1700)$ vs. $\rho_3(1690)$
- Including $\Delta\phi$ information allows us to differentiate these.

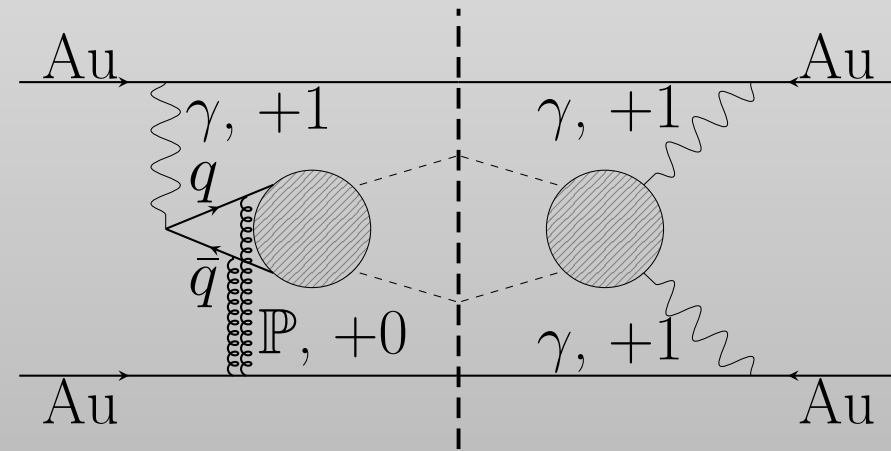
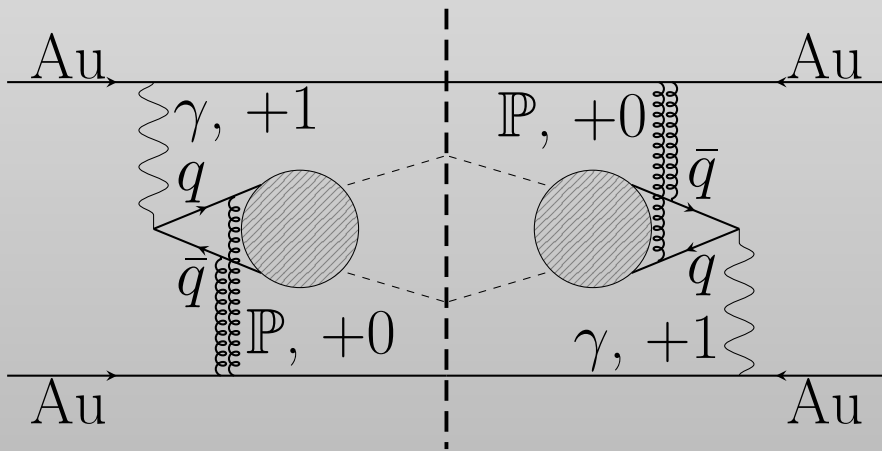


Describing $\langle 2\cos(n\Delta\phi) \rangle$ vs. $M_{\pi\pi}$

- Consider the UPC $AA \rightarrow AA\pi^+\pi^-$ cross section as:

$$\propto |\rho^0 + \omega + f_2(1270) + \rho(1700) + B_{\pi\pi}|^2$$

- Each cross term contributes to $A_{n\Delta\phi} \equiv \langle 2\cos(n\Delta\phi) \rangle$ according to the spin of the interfering states.
- Assume that **each cross term contributes a constant value, and the observed $A_{n\Delta\phi}$ is the average of these constants** weighted by the contribution of that term to the total cross section.



Simultaneous Fit Procedure

- In this formulation, fit the invariant mass spectrum, $A_{1\Delta\phi}$, $A_{2\Delta\phi}$, and $A_{3\Delta\phi}$.

$$\bullet \frac{d\sigma}{dM_{\pi\pi}} = \left| \frac{A_{\rho}\sqrt{M_{\pi\pi}M_{\rho}\Gamma_{\rho}}}{M_{\pi\pi}^2 - M_{\rho}^2 + iM_{\rho}\Gamma_{\rho}} + \frac{(A_{\omega}^R + iA_{\omega}^I)\sqrt{M_{\pi\pi}M_{\omega}\Gamma_{\omega\rightarrow\pi^+\pi^-}}}{M_{\pi\pi}^2 - M_{\omega}^2 + iM_{\omega}\Gamma_{\omega}} + \frac{(A_{f_2}^R + iA_{f_2}^I)\sqrt{M_{\pi\pi}M_{f_2}\Gamma_{f_2\rightarrow\pi^+\pi^-}}}{M_{\pi\pi}^2 - M_{f_2}^2 + iM_{f_2}\Gamma_{f_2}} + \frac{(A_{\rho(1700)}^R + iA_{\rho(1700)}^I)\sqrt{M_{\pi\pi}M_{\rho(1700)}\Gamma_{\rho(1700)\rightarrow\pi^+\pi^-}}}{M_{\pi\pi}^2 - M_{\rho(1700)}^2 + iM_{\rho(1700)}\Gamma_{\rho(1700)}} + B_{\pi\pi} \right|^2$$

Where $\Gamma_X = \Gamma_{0,X} \frac{M_X}{M_{\pi\pi}} \left(\frac{M_{\pi\pi}^2 - 4m_{\pi}^2}{M_X^2 - 4m_{\pi}^2} \right)^{2j+1/2}$ and $\Gamma_{X\rightarrow\pi\pi} = Br(X \rightarrow \pi\pi)\Gamma_X$

Orange: Relativistic Breit-Wigner distributions.

Blue: Drell-Söding process modeled as a constant.

Previous results use a function like this, but struggle to differentiate the heavier resonances.

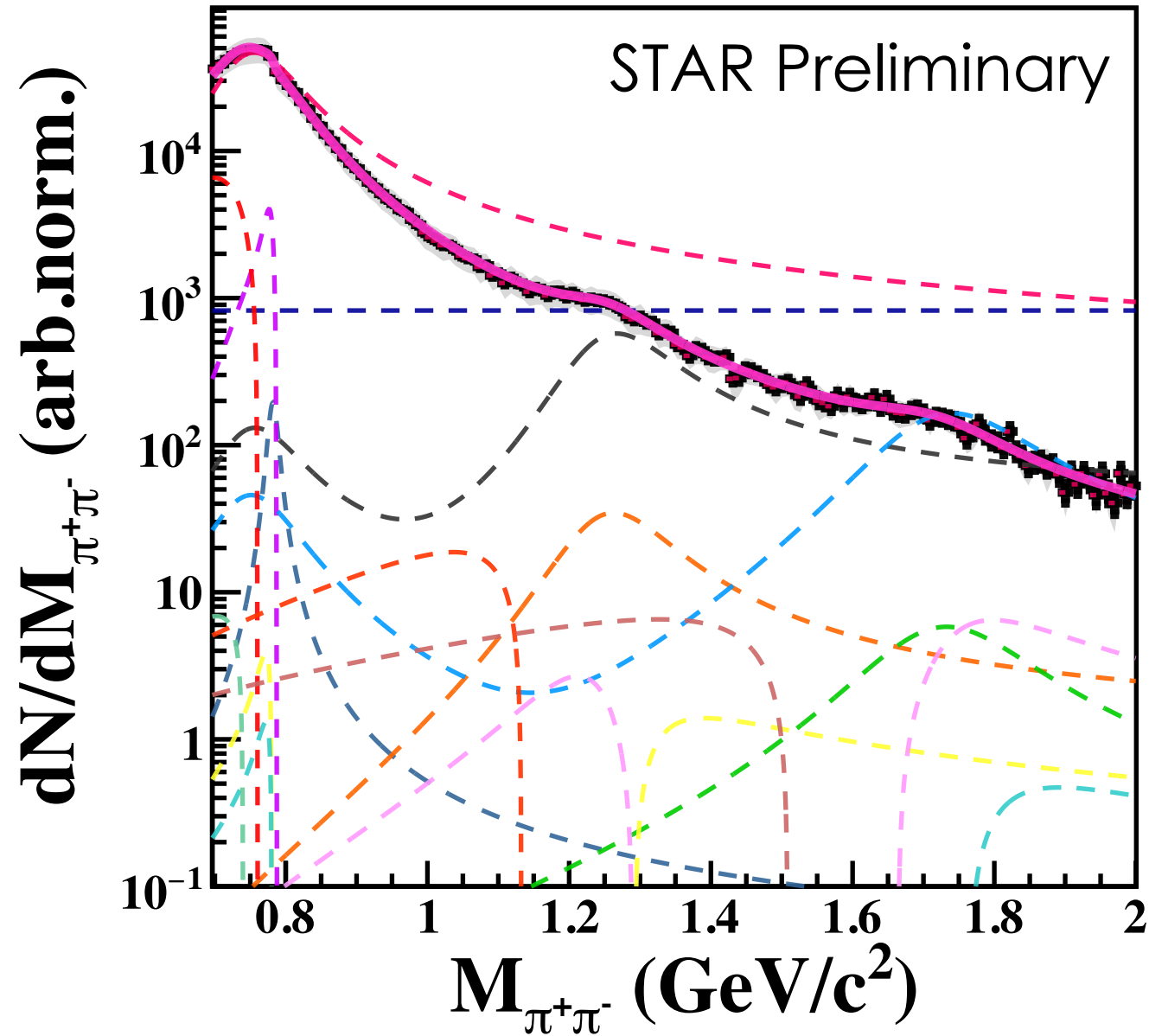
Simultaneous Fit Procedure

- In this formulation, fit the invariant mass spectrum, $A_{1\Delta\phi}$, $A_{2\Delta\phi}$, and $A_{3\Delta\phi}$.

- $$A_{1,3\Delta\phi} = \left(\frac{\rho \times f_2}{total} \times A_{1,3\Delta\phi}^{\rho \times f_2} \right) + \left(\frac{B_{\gamma A \rightarrow \pi\pi} \times f_2}{total} \times A_{1,3\Delta\phi}^{B_{\gamma A \rightarrow \pi\pi} \times f_2} \right)$$

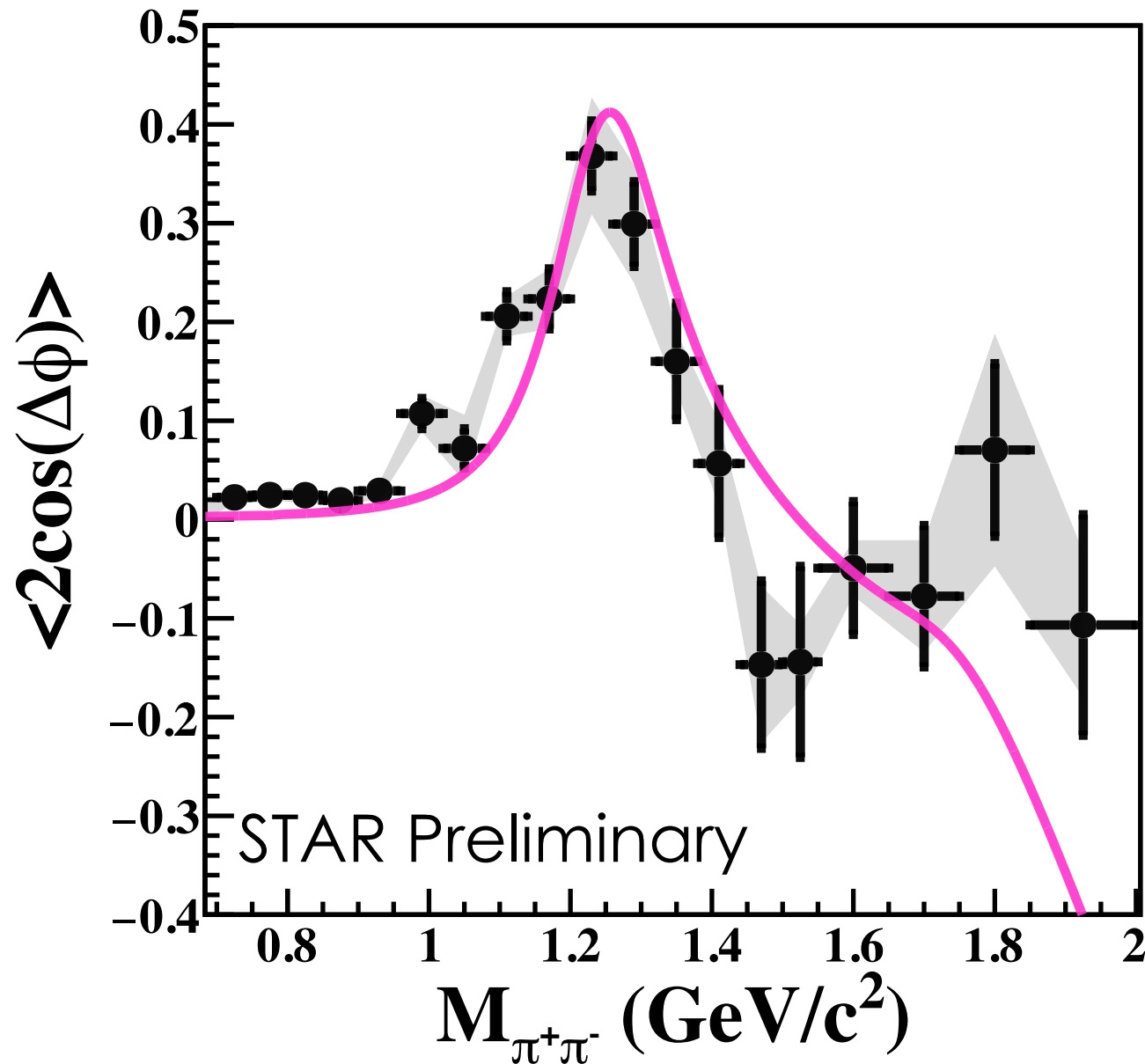
- $$A_{2\Delta\phi} = \left(\frac{\rho^2}{total} \times A_{2\Delta\phi}^{\rho^2} \right) + \left(\frac{B_{\gamma A \rightarrow \pi\pi} \times \rho}{total} \times A_{2\Delta\phi}^{B_{\gamma A \rightarrow \pi\pi} \times \rho} \right) + \left(\frac{B_{\gamma A \rightarrow \pi\pi}^2}{total} \times A_{2\Delta\phi}^{B_{\gamma A \rightarrow \pi\pi}^2} \right) +$$
$$\left(\frac{\rho(1700) \times \rho}{total} \times A_{2\Delta\phi}^{\rho(1700) \times \rho} \right) + \left(\frac{B_{\gamma A \rightarrow \pi\pi} \times \rho(1700)}{total} \times A_{2\Delta\phi}^{B_{\gamma A \rightarrow \pi\pi} \times \rho(1700)} \right)$$

Simultaneous Fit: $dN/dM_{\pi\pi}$



Invariant mass fit
constrained by $A_{n\Delta\phi}$

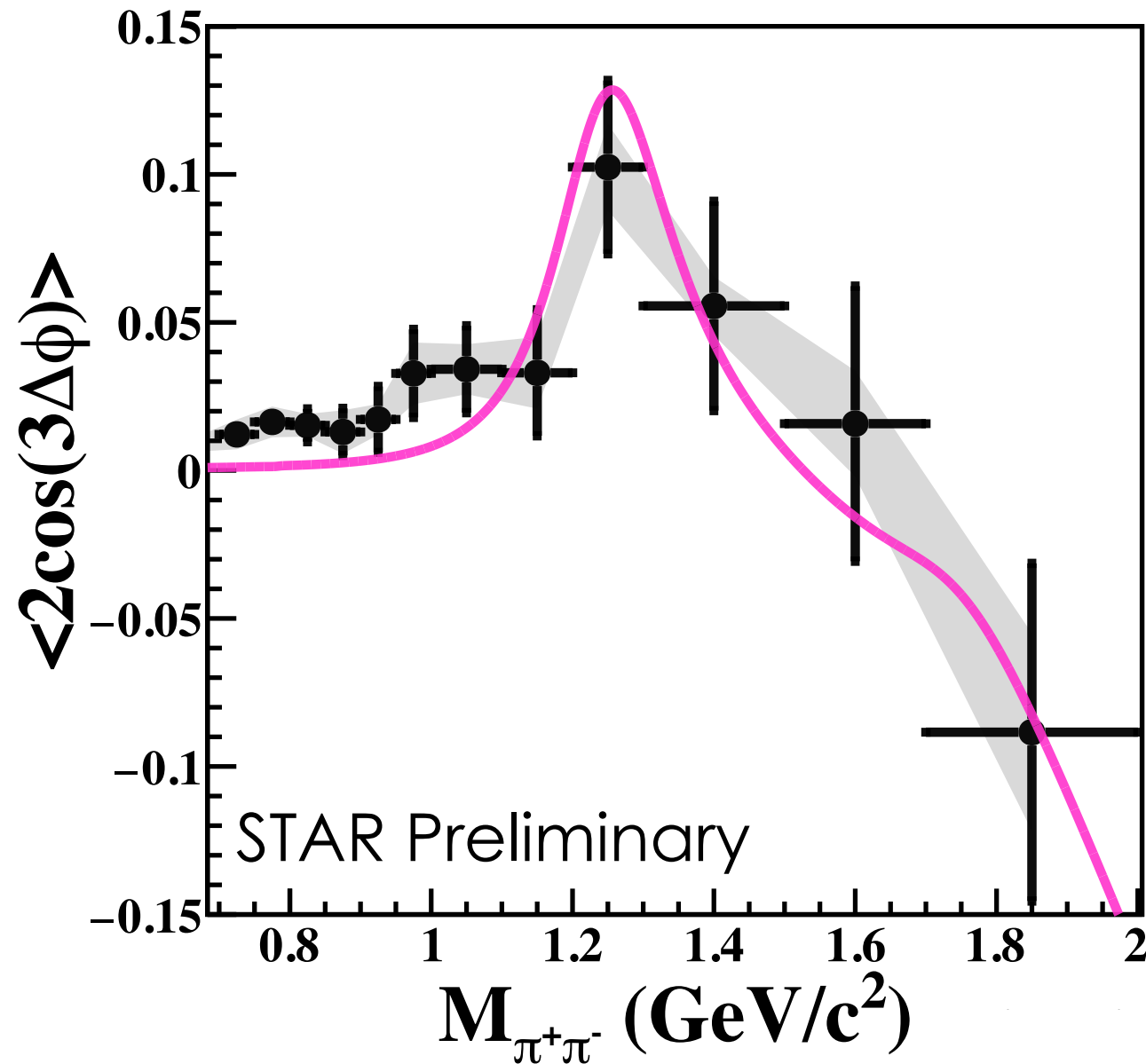
Simultaneous Fit: $A_{1\Delta\phi}$



Contributions from:

- $\rho^0 \times f_2(1270)$
- Drell-Söding $\times f_2(1270)$

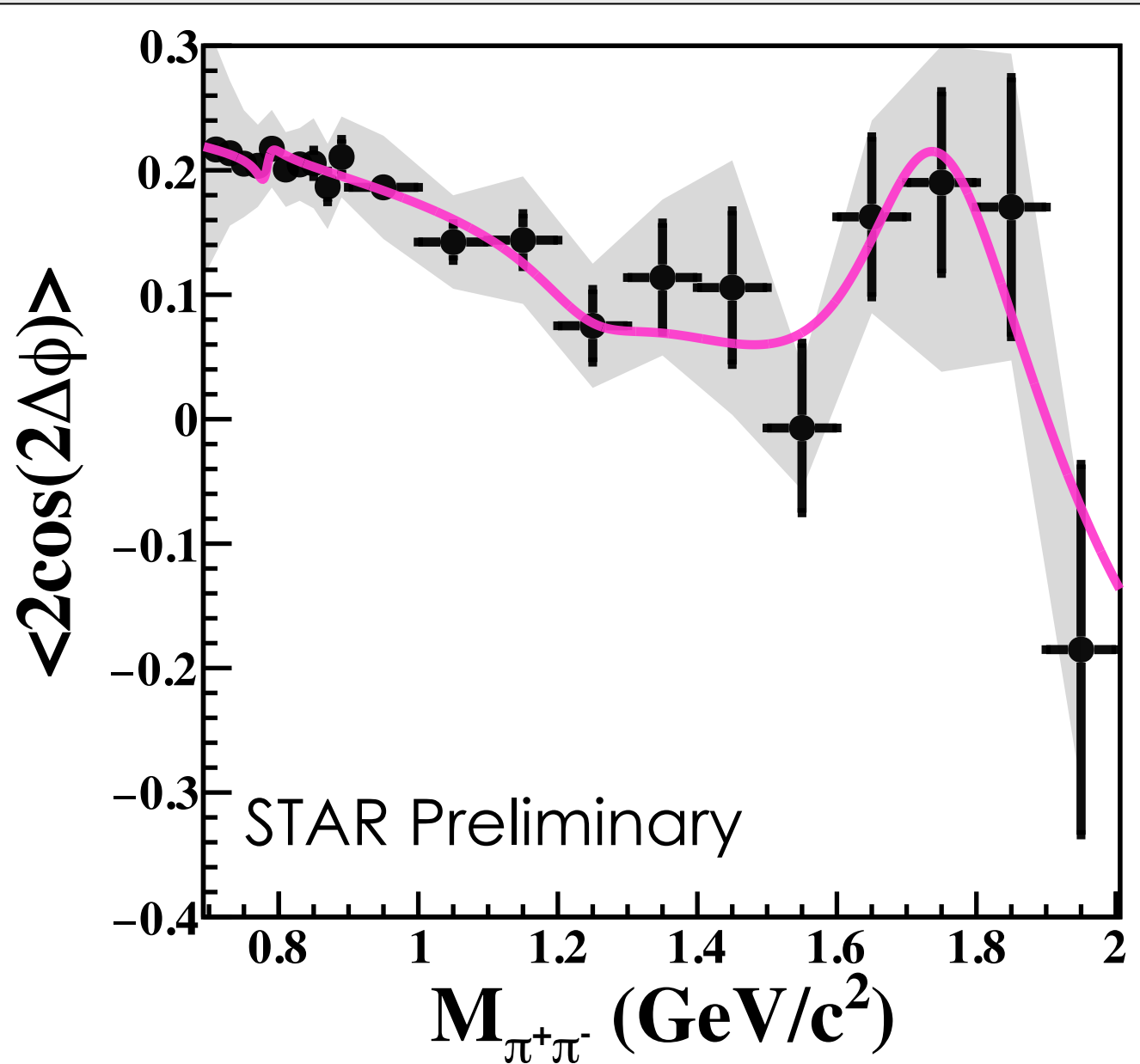
Simultaneous Fit: $A_{3\Delta\phi}$



Contributions from:

- $\rho^0 \times f_2(1270)$
- Drell-Söding $\times f_2(1270)$

Simultaneous Fit

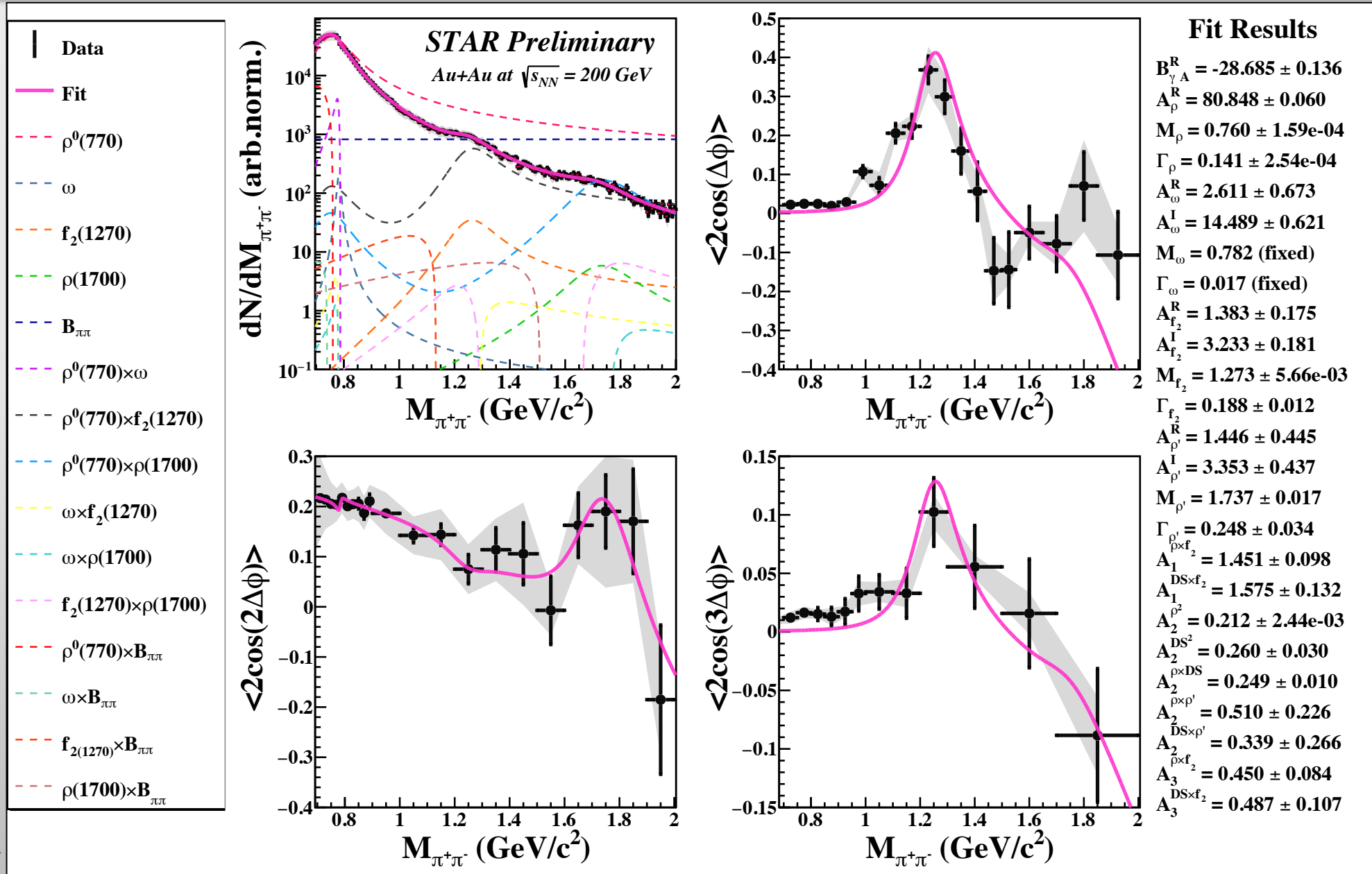


Contributions from:

- $\rho^0 \times \rho^0$
- Drell-Söding $\times$ Drell-Söding
- $\rho^0 \times$ Drell-Söding
- $\rho^0 \times \rho(1700)$
- Drell-Söding $\times \rho(1700)$

Simultaneous Fit

$$\frac{\chi^2}{ndf} = 2.9$$



Tabulated Results

This fit (STAR Preliminary)

$$M_\rho = 760.4 \pm 0.16 \pm 11 \frac{\text{MeV}}{c^2}$$

$$\Gamma_\rho = 142.0 \pm 0.25 \pm 14 \text{ MeV}$$

$$M_\omega = 782 \frac{\text{MeV}}{c^2} \text{ (fixed)}$$

$$\Gamma_\omega = 17 \text{ MeV (fixed)}$$

$$M_{f_2(1270)} = 1272.4 \pm 5.6 \pm 44 \frac{\text{MeV}}{c^2}$$

$$\Gamma_{f_2(1270)} = 188.5 \pm 11.6 \pm 115 \text{ MeV}$$

$$M_{\rho(1700)} = 1737.3 \pm 17.3 \pm 71 \frac{\text{MeV}}{c^2}$$

$$\Gamma_{\rho(1700)} = 248.0 \pm 34 \pm 13 \text{ MeV}$$

PDG mass, width:

$$M_{f_2(1270)} = 1275.5 \pm 0.8 \frac{\text{MeV}}{c^2}$$

$$\Gamma_{f_2(1270)} = 185.9 \pm 2.5 \text{ MeV}$$

$$M_{\rho(1700)} = 1720 \pm 20 \frac{\text{MeV}}{c^2}$$

$$\Gamma_{\rho(1700)} = 250 \pm 100 \text{ MeV}$$

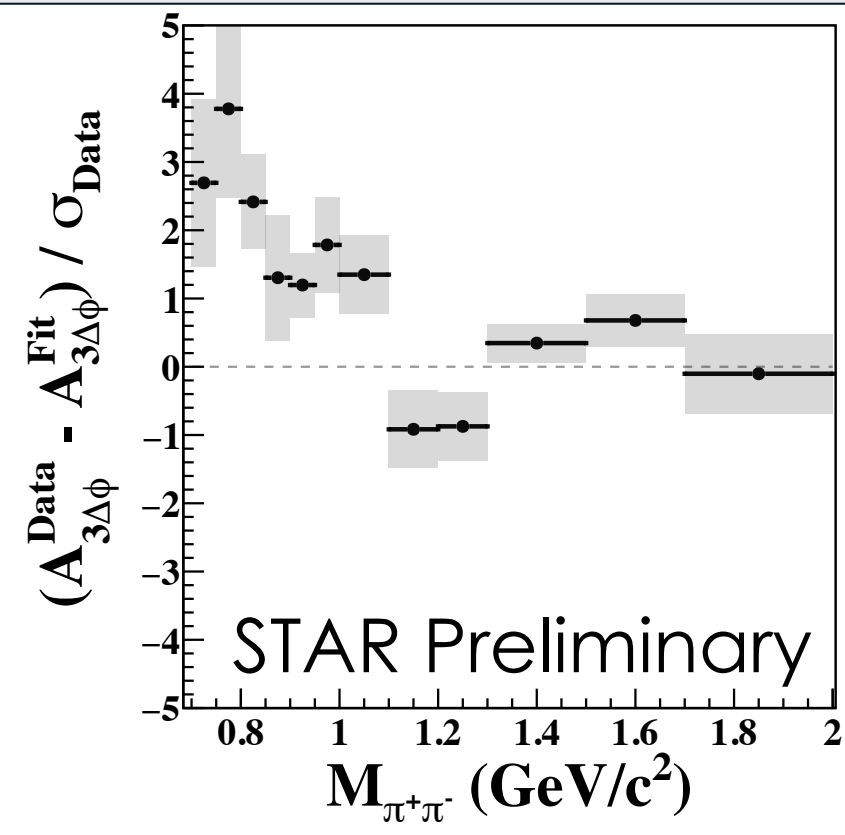
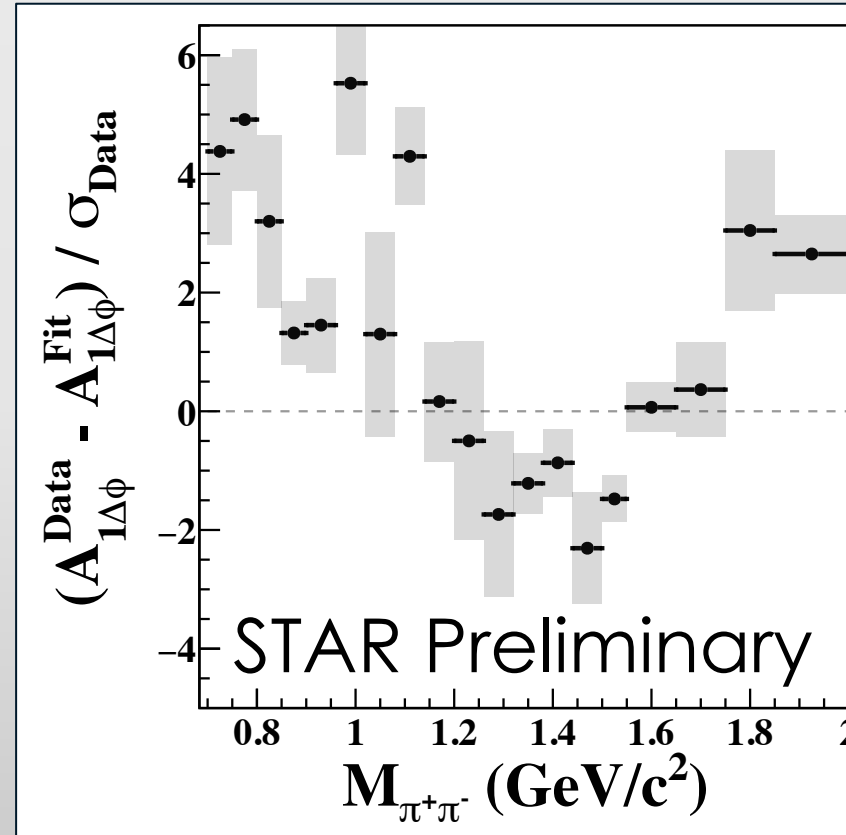
ω fixed to results from STAR, Phys. Rev. C 96, 054904 (2017)

Results reported as: *value* $\pm$ *statistical uncertainty* $\pm$ *systematic uncertainty*

- Spin interference parameters constrain on the mass and width of $f_2(1270)$ and $\rho(1700)$.
 - Close agreement with PDG values for $f_2(1270)$.
 - $\rho(1700)$ width from this fit has smaller uncertainty than PDG (2022).

Non-resonant $\gamma\gamma \rightarrow \pi^+\pi^-$

- Non-resonant $\gamma\gamma \rightarrow \pi^+\pi^-$ not included in fit.
- Largest contribution to $A_{1\Delta\phi}$ and $A_{3\Delta\phi}$ around ρ^0 mass.
- **Excess in data compared to fit in this region as a result.**



Conclusions

- $A_{1\Delta\phi}$ and $A_{3\Delta\phi}$ were measured for the first time in STAR Au+Au UPCs at 200 GeV.
 - **Spin interference between $\gamma A \rightarrow \pi^+\pi^-$ and $\gamma\gamma \rightarrow \pi^+\pi^-$.**
 - $A_{1\Delta\phi}$ and $A_{3\Delta\phi}$ **constrain $\gamma\gamma \rightarrow \pi^+\pi^-$ cross section, phase** (Hagiwara et al., Phys. Rev. D 103, 074013 (2021)).
 - $A_{1\Delta\phi}$ and $A_{3\Delta\phi}$ feature at $\sim 1270 \text{ MeV}/c^2$ **indicates $\gamma\gamma \rightarrow f_2(1270) \rightarrow \pi^+\pi^-$.**
- $A_{2\Delta\phi}$ as a function of $M_{\pi\pi}$ was also measured in the same system for the first time.
- The $A_{n\Delta\phi}$ and $\frac{dN}{dM_{\pi\pi}}$ were fit simultaneously with a **new technique that distinguishes resonances on $\frac{dN}{dM_{\pi\pi}}$ according to their spin.**
 - Extracted values of the mass and width of $f_2(1270)$ and $\rho(1700)$.
 - Values of $f_2(1270)$ consistent with known results **confirms $\gamma\gamma \rightarrow \pi^+\pi^-$ in UPC.**
 - **Next step: extract $d\sigma_{\gamma\gamma \rightarrow \pi^+\pi^-} / dM_{\pi\pi}$**

References

STAR, Phys. Rev. C 96, 054904 (2017)

First Results from Fermilab's Muon g-2 Experiment Strengthen Evidence of New Physics; Fermilab, 7 April 2021.

Danilkin, C. F. Redmer, and M. Vanderhaeghen, Prog. Part. Nucl. Phys. 107, 20 (2019).

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STAR, Sci.Adv.9,eabq3903 (2023)

Hagiwara, Zhang, Zhou, and Zhou, Phys. Rev. D 103, 074013 (2021)