

# Measurements of $\gamma$ production in p+p, p+Au and Au+Au collisions with the STAR experiment

Leszek Kosarzewski, BEng, Ph.D.  
for the STAR collaboration

Czech Technical University in Prague

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YOUTH AND SPORTS



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CZ.02.2.69/0.0/0.0/16\_027/0008465

## 1 Introduction

- Upsilon as a probe of quark-gluon plasma
- Production mechanism

## 2 STAR experiment

### 3 $\Upsilon$ production in p+p

- $p_T$  and rapidity spectra
- Cross section ratios
- Event activity dependence of  $\Upsilon$  production

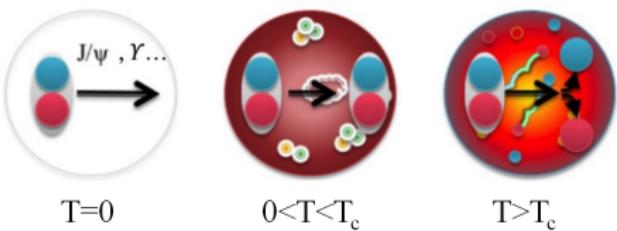
### 4 $\Upsilon$ production in p+Au

### 5 $\Upsilon$ suppression in Au+Au

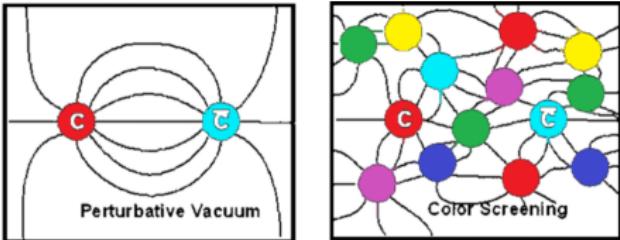
## 6 Summary

# Upsilon in quark-gluon plasma

$$\Upsilon = b \bar{b}$$



[A. Rothkopf, Hard Probes 2012]



High mass - produced early

$$m_b = 4.18_{-0.03}^{+0.04} \text{ GeV}/c^2$$

$$m_{\Upsilon(1S)} = 9460.30 \pm 0.26 \text{ MeV}/c^2$$

$$m_{\Upsilon(2S)} = 10023.26 \pm 0.31 \text{ MeV}/c^2$$

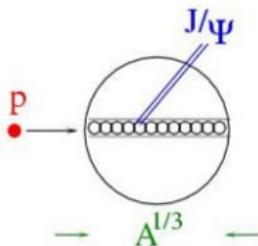
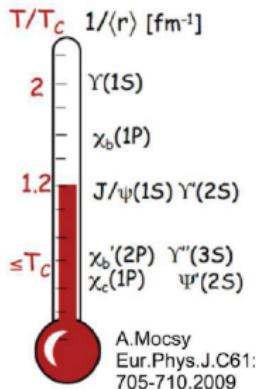
$$m_{\Upsilon(3S)} = 10355.2 \pm 0.5 \text{ MeV}/c^2$$

[Phys.Rev.D 98, 030001 (2018)]

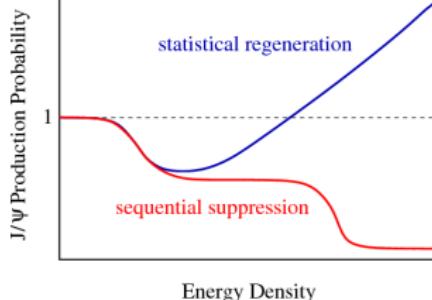
$\Upsilon$  as a probe of quark-gluon plasma

- $\Upsilon$  is sensitive to the QGP properties similarly to  $J/\psi$   
[Phys.Lett.B 178(4),416-422(1986)]
- Dissociation due to Debye-like screening when  $r_\Upsilon > r_{Debye} \propto T^{-1}$
- Suppression observed at RHIC and LHC  
[Phys.Lett.B 735,127-137(2014)], [Phys.Lett.B. 770,357-359(2017)]
- Sequential suppression, due to lower binding energy for excited  $\Upsilon$  states, expected and observed  
[Phys.Rev.D 64, 094015(2001)], [Phys.Rev.Lett 109, 222301(2012)]

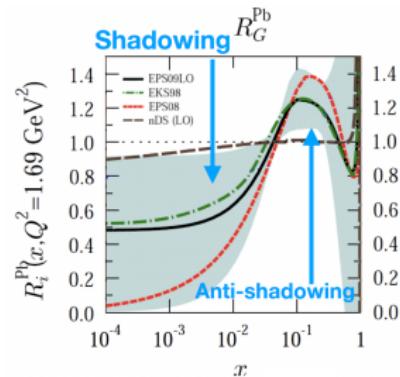
# Other effects



[L. Grandchamp, LBNL  
2005]



[Nucl.Phys.B (Proc.Suppl.) 214, 3-36(2011)]



[Phys.Rev.C 81 064911(2010)]

## Other modifications to $\Upsilon$ production

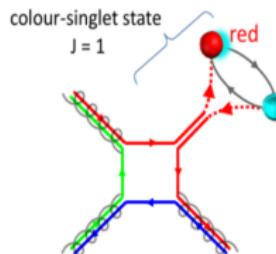
- Regeneration very small for  $\Upsilon$  at RHIC  
[Phys.Rev.C 96, 054901(2017)]
- Feed-down from excited states  $\Upsilon(nS) \rightarrow \Upsilon(1S)\pi^+\pi^-$  and  $\chi_{bn} \rightarrow \gamma\Upsilon(1S)$
- Cold Nuclear Matter effects - can be studied separately in  $p + A$  or  $d + A$  collision
  - nuclear absorption
  - comover interactions - very small for  $\Upsilon(1S)$   
[Phys.Lett.B 503, 104(2001)]
  - nuclear PDFs: shadowing, anti-shadowing

# Production mechanism

- Still not well understood: hard scattering+non-perturbative hadronization
- Quarkonium measurements provide tests of production models, and thus help to understand QCD

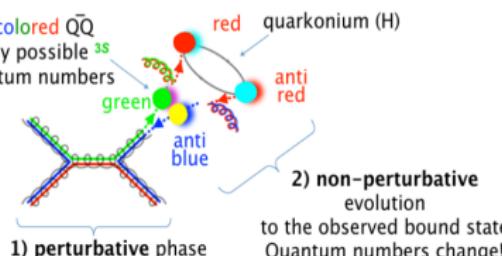
## Quarkonium production models

- Color Singlet - only  $Q\bar{Q}$  produced directly in a color neutral state can bind to form quarkonia
- Color Octet -  $Q\bar{Q}$  produced in a colored state. Gluon emissions are needed to neutralize color. This is described by long-distance matrix elements (LDMEs) which are assumed universal.
- Color Evaporation Model - color irrelevant. Fixed fractions of  $Q\bar{Q}$  pairs evolve into various quarkonium states.

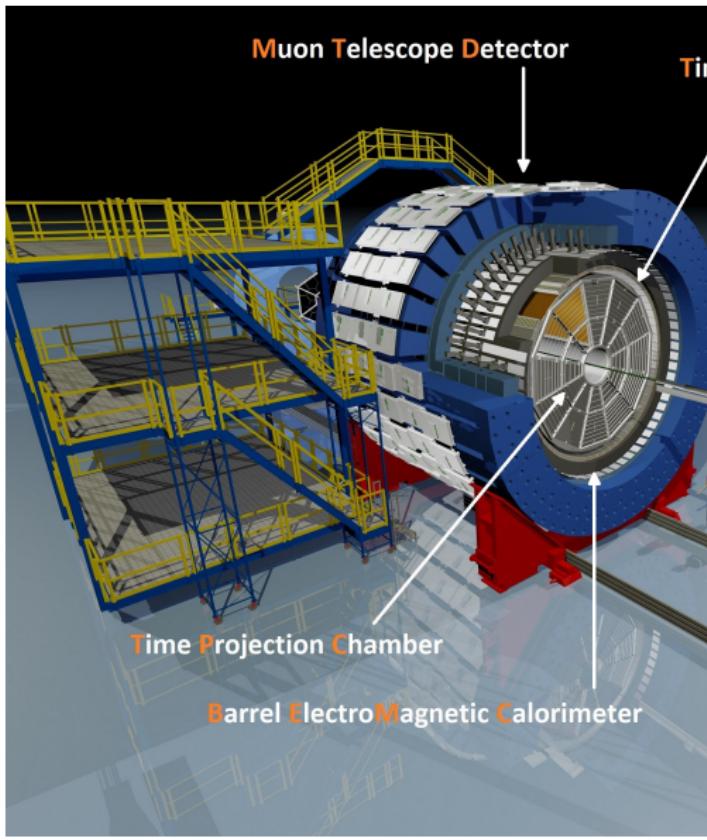


+ analogous colour  
combinations

possibly colored  $\bar{Q}Q$   
pair of any possible  $^{3S}_{+}L_J$   
quantum numbers



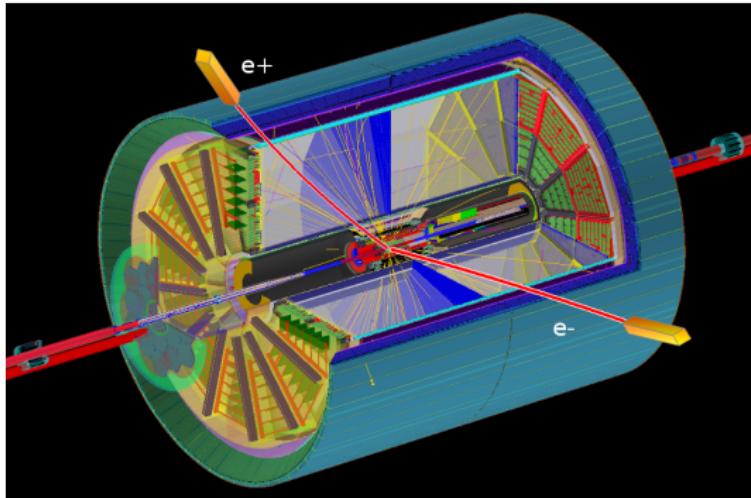
[P. Faccioli, Polarization in LHC physics, Course on Physics at the LHC 2014]



## Detectors used for $\gamma$ studies

- TPC  $|\eta| < 1, 0 < \phi < 2\pi$ 
  - Tracking - momentum measurement
  - Particle identification based on energy loss  $dE/dx$
- BEMC  $|\eta| < 1, 0 < \phi < 2\pi$ 
  - Trigger on high- $p_T$  electrons
  - Electron identification via  $E/p$  and EM shower shape
- MTD  $|\eta| < 0.5, 45\% \text{ in } 0 < \phi < 2\pi$ 
  - Dimuon trigger
  - Muon identification with time-of-flight
  - Magnet used as hadron absorber
  - Muons - less bremsstrahlung
- TOF  $|\eta| < 1, 0 < \phi < 2\pi$ 
  - Particle identification based on time-of-flight - not used for  $\gamma$
  - Fast detector used to remove pile-up for  $N_{ch}$  determination

# Upsilon measurements in STAR - dielectron channel



## $\Upsilon$ decay reconstruction in $e^+e^-$ channel

$$\Upsilon(1S) \rightarrow e^+e^-, B_{ee}^{\Upsilon(1S)} = 2.38 \pm 0.11\%$$

$$\Upsilon(2S) \rightarrow e^+e^-, B_{ee}^{\Upsilon(2S)} = 1.91 \pm 0.16\%$$

$$\Upsilon(3S) \rightarrow e^+e^-, B_{ee}^{\Upsilon(3S)} = 2.18 \pm 0.20\%$$

[*Phys.Rev.D* 98, 030001(2018)]

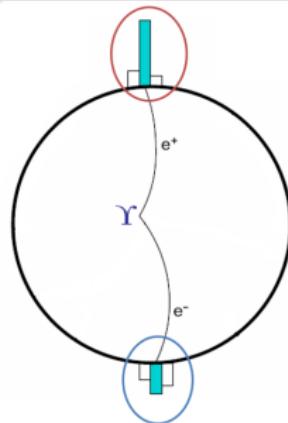
- Project TPC track to the high-energy tower in BEMC, which fired the trigger, and reconstruct a cluster
- Find a partner track and project it to BEMC cluster

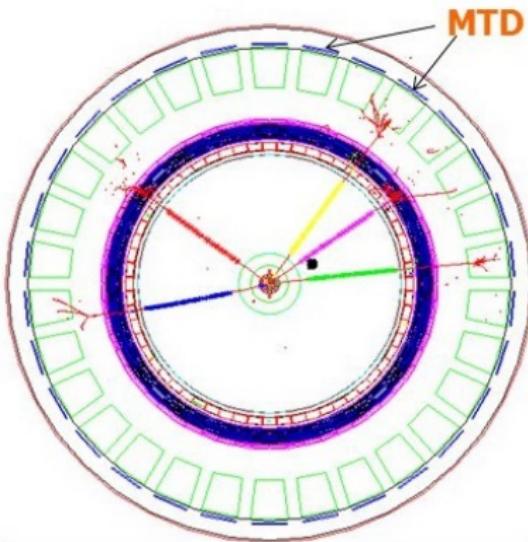
## Detectors used

- TPC+BEMC(+TOF for  $N_{ch}$ )

## STAR datasets for $\Upsilon \rightarrow e^+e^-$ NEW

- Au+Au  $\sqrt{s_{NN}} = 200$  GeV: **2011, 2010**
- p+Au  $\sqrt{s_{NN}} = 200$  GeV: **2015**
- d+Au  $\sqrt{s_{NN}} = 200$  GeV: **2008**
- p+p  $\sqrt{s} = 500$  GeV: **2011**
- p+p  $\sqrt{s} = 200$  GeV: **2015, 2009**





## $\Upsilon$ decay reconstruction in $\mu^+\mu^-$ channel

$$\Upsilon(1S) \rightarrow \mu^+ \mu^-, B_{\mu\mu}^{\Upsilon(1S)} = 2.48 \pm 0.05\%$$

$$\Upsilon(2S) \rightarrow \mu^+ \mu^-, B_{\mu\mu}^{\Upsilon(2S)} = 1.93 \pm 0.17\%$$

$$\Upsilon(3S) \rightarrow \mu^+ \mu^-, B_{\mu\mu}^{\Upsilon(3S)} = 2.18 \pm 0.21\%$$

[*Phys. Rev. D* 98, 030001 (2018)]

## Detectors used

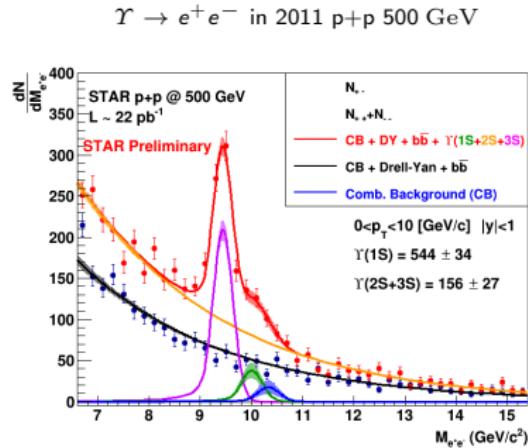
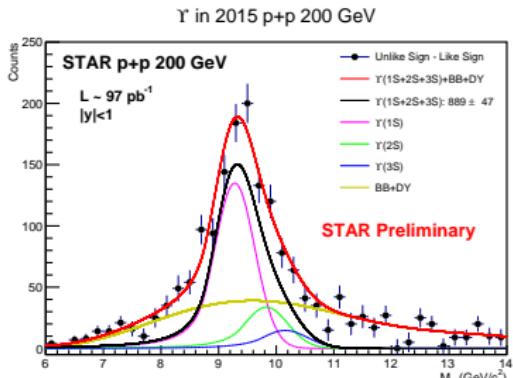
- TPC+MTD

## STAR MTD datasets for $\Upsilon \rightarrow \mu^+ \mu^-$

- 2016 Au+Au  $\sqrt{s_{NN}} = 200$  GeV
- 2014 Au+Au  $\sqrt{s_{NN}} = 200$  GeV

## Track projection to MTD hits

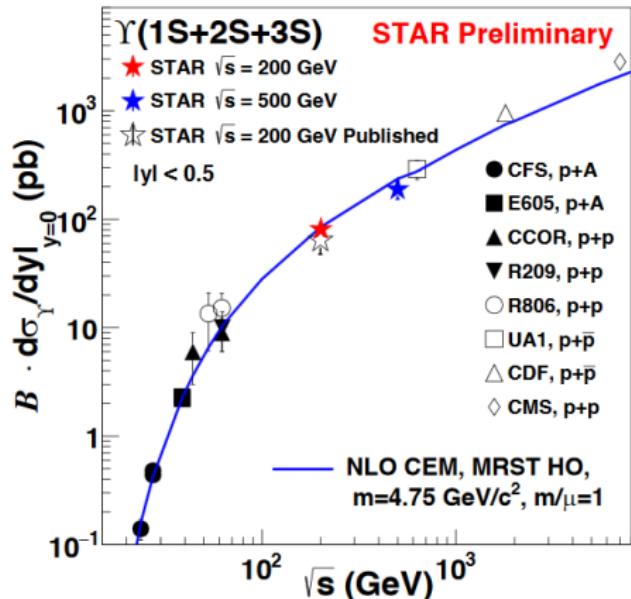
- Tracks are projected from TPC to MTD and matched to hits
- Energy loss in the magnet is included in the track projection procedure



### Signal fitting $\Upsilon \rightarrow e^+ e^-$

- $\Upsilon$  signal shapes modeled by 3 Crystal-Ball functions
- Fit to **Unlike-sign (red)** distribution consists of:
  - 3 Crystal-Ball functions (**1S, 2S, 3S** states) - fixed using MC simulation
  - $b\bar{b}$ +Drell-Yan correlated background (**orange**) determined using MC simulation

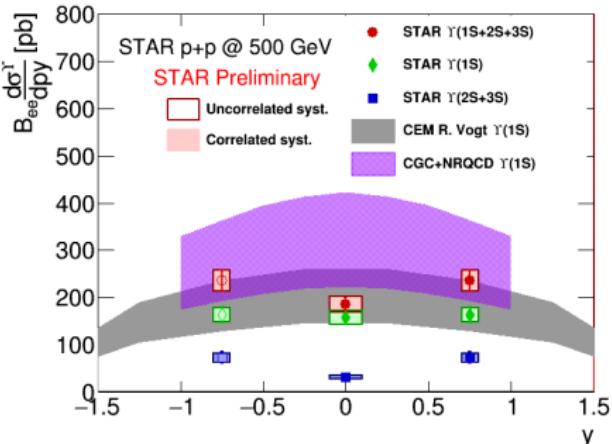
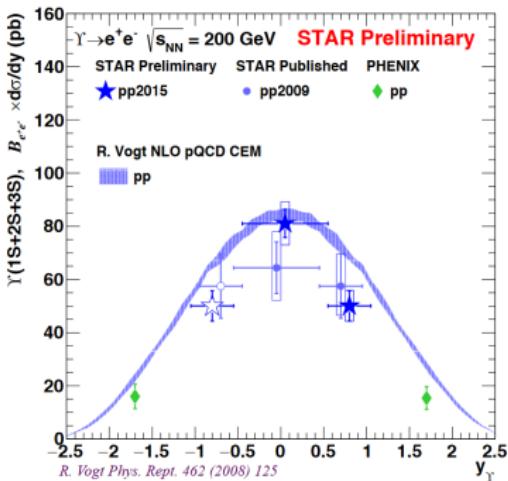
# Integrated cross section in p+p



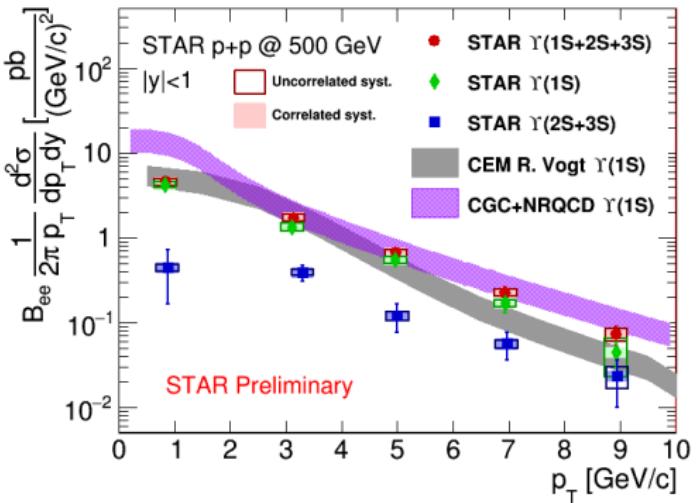
STAR  
 [Phys.Lett.B 735,127–137(2014)]  
 CDF  
 [Phys.Rev.Lett. 88,161802(2002)]  
 CMS  
 [Phys.Rev.D 83,112004(2010)]  
 CF5  
 [Phys.Rev.Lett. 39,1240–1242(1977)]  
 [Phys.Rev.Lett. 41,684–687(1978)]  
 [Phys.Rev.Lett. 42,486–489(1979)]  
 [Phys.Rev.Lett. 55,1962–1964(1985)]  
 E605  
 [Phys.Rev.D 43,2815–2835(1991)]  
 [Phys.Rev.D 39,3516(1989)]  
 CCOR  
 [Phys.Lett.B 87,398–402(1979)]  
 L. Camilleri, T.B.W. Kirk, H.D.I. Abarbanel (Eds.)  
 E866  
 [Phys.Rev.Lett. 100,062301(2008)]  
 ISR  
 [Phys.Lett.B 91,481–486(1980)]

- $B_{ee} \frac{d\sigma}{dy}|_{|y|<0.5} = 81 \pm 5(stat) \pm 8(syst)$  pb in p+p collisions at  $\sqrt{s} = 200$  GeV
- $B_{ee} \frac{d\sigma}{dy}|_{|y|<0.5} = 186 \pm 14(stat) \pm 33(syst)$  pb in p+p collisions at  $\sqrt{s} = 500$  GeV
- STAR results follow the world data trend
- Consistent with the Color Evaporation Model calculation  
 $[Phys.Rep. 462, pp.125–175(2008)]$

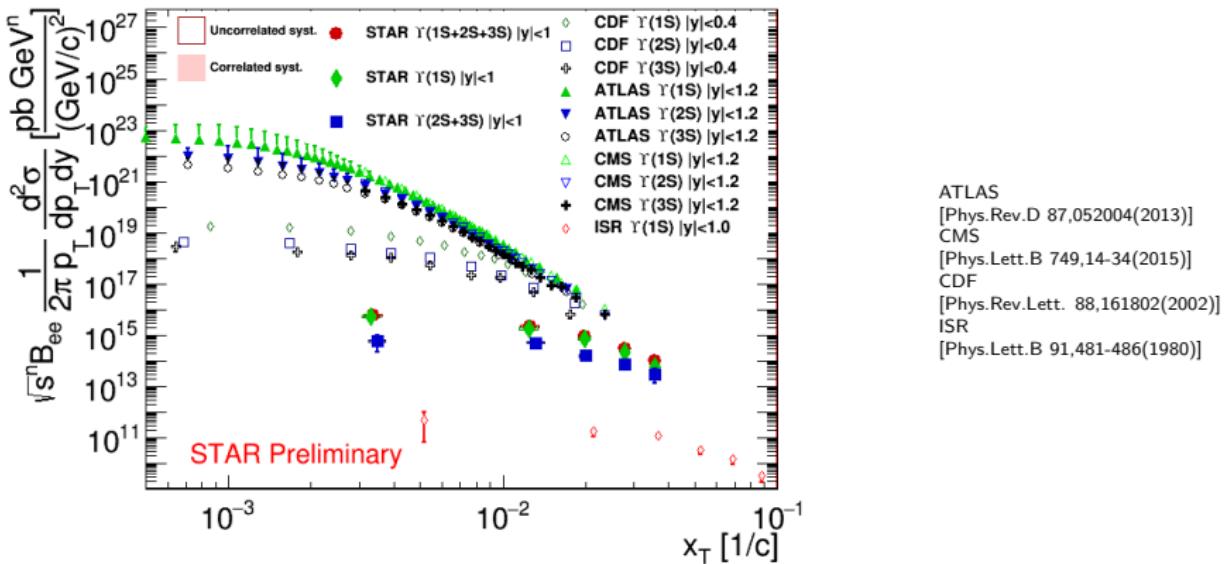
# Rapidity dependence in p+p



- STAR data slightly narrower than Color Evaporation Model (CEM) at  $\sqrt{s} = 200 \text{ GeV}$
- Flatter rapidity spectrum at  $\sqrt{s} = 500 \text{ GeV}$  compared to  $\sqrt{s} = 200 \text{ GeV}$ 
  - Dip at mid-rapidity for  $\Upsilon(2S + 3S) \approx 2\sigma$  level from flat
  - CEM model (inclusive) consistent with the measurement for  $\Upsilon(1S)$  [*Phys.Rev.C* 92 034909(2015)]
  - Non-relativistic Quantum Chromodynamics coupled with the Color-Glass Condensate formalism (CGC+NRQCD) predictions for direct  $\Upsilon(1S)$  are above the data for  $\Upsilon(1S)$  [*Phys.Rev.D* 94, 014028(2016)]. [*Phys.Rev.Lett.* 113, 192301(2014)]

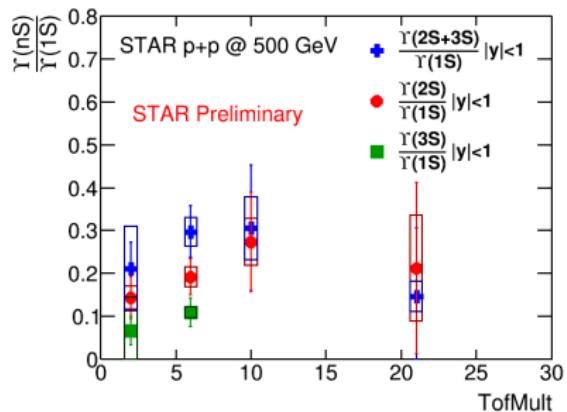
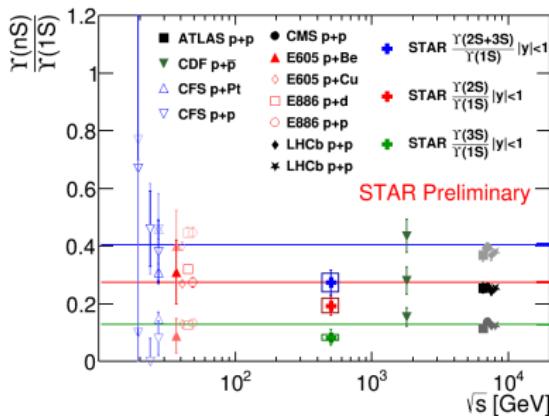


- CEM calculation for inclusive  $\Upsilon(1S)$   
[Phys.Rev.C 92 034909(2015)]
  - Agree with data reasonably well
- CGC+NRQCD for direct  $\Upsilon$   
[Phys.Rev.D 94, 014028(2016)] [Phys.Rev.Lett. 113, 192301(2014)]
  - $\Upsilon(1S)$ : model calculation is above the data points. Caveat: additional corrections are needed at low  $p_T$  according to authors.



- $x_T = \frac{2p_T}{\sqrt{s}}$ ,  $\sigma^{inv} \equiv E \frac{d^3\sigma}{d^3p} = \frac{F(x_T)}{p_T^n(x_T, \sqrt{s})} = \frac{F'(x_T)}{\sqrt{s}^n(x_T, \sqrt{s})}$   
[JHEP06,035(2010)]
- pQCD predicts that spectra of hard processes should follow  $x_T$  scaling - check with  $n = 5.6$  (number of partons taking active part in the process) obtained for  $J/\psi$   
[Phys.Rev.C 80, 041902(2009)]
- No clear scaling observed

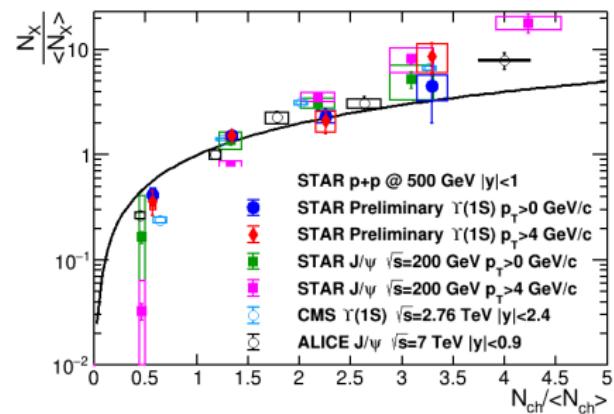
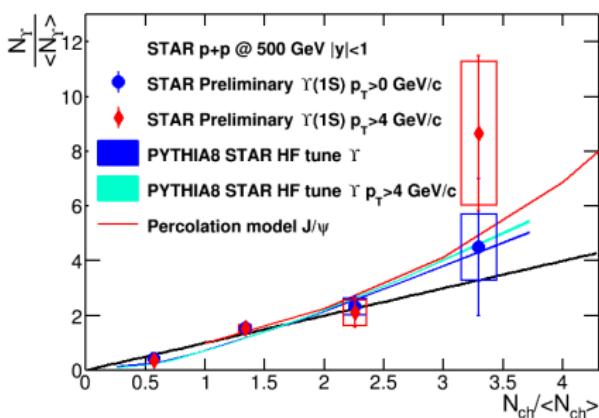
# Cross section ratios: $\Upsilon(nS)/\Upsilon(1S)$



[Phys. Rev. C 88, 067901 (2013)]

- TofMult: number of tracks matched to TOF within  $|\eta| < 1$ ,  $p_T > 0.2 \text{ GeV}/c$
- Boxes correspond to uncorrelated systematic uncertainties (correlated uncertainties largely cancel out)
- Cross section ratios measured in 500 GeV p+p collisions are slightly below (within  $2\sigma$ ) world data average, shown as solid lines in the left plot.
- Right plot: No strong multiplicity dependence observed.

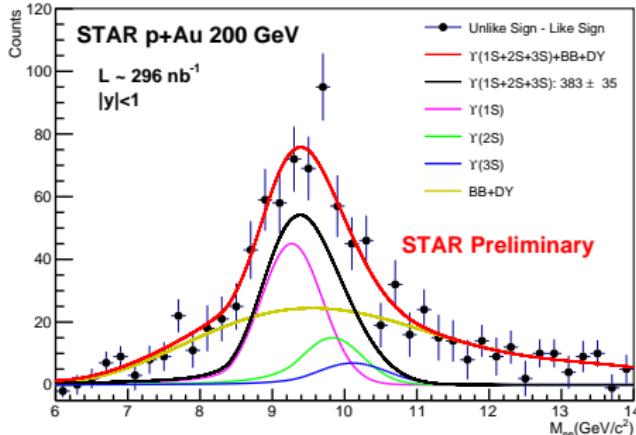
# $\Upsilon$ production vs. event activity



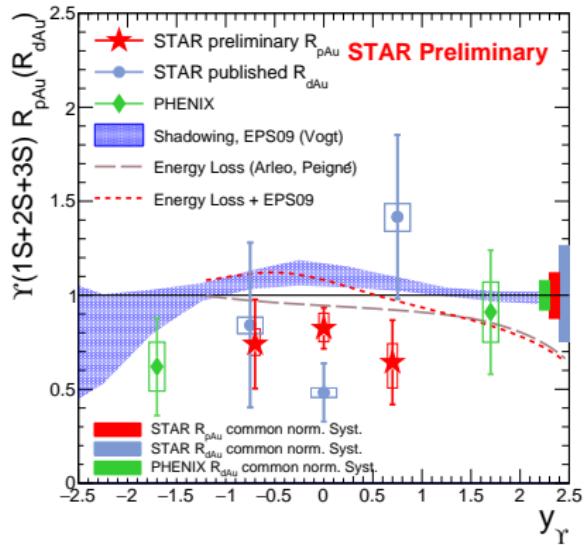
- Normalized  $\Upsilon(1S)$  yield vs. normalized multiplicity (a measure of event activity)
- Data consistent with a linear rise (black line), with a hint for stronger-than-linear rise for  $\Upsilon(1S)$  above  $p_T > 4$  GeV/c
- Similar trend at RHIC and LHC for  $\Upsilon$  and  $J/\psi$   
 $[JHEP04, 103(2014)], [Nucl. and Part. Phys. Proc., 276-278, pp. 261-264(2016)], [Phys. Lett. B 712, 165-175(2012)]$
- Indication of  $\Upsilon$  production in MPI or soft particle production being suppressed by interactions of strings of color field in high- $N_{ch}$  collisions compared to quarkonium yield  
 $[Phys. Rev. C, 86, 034903(2012)]$ 
  - Need more data to distinguish between the 2 scenarios

# $\Upsilon$ production in p+Au

$\Upsilon$  in 2015 p+Au 200 GeV



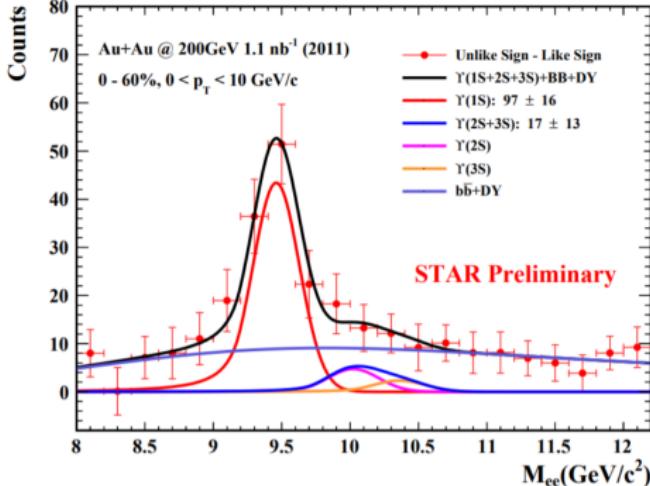
[J.Phys.Lett.B 735(2014)127],  
[Phys. Rev. C 87, 044909],  
[JHEP 03, 122(2013)]



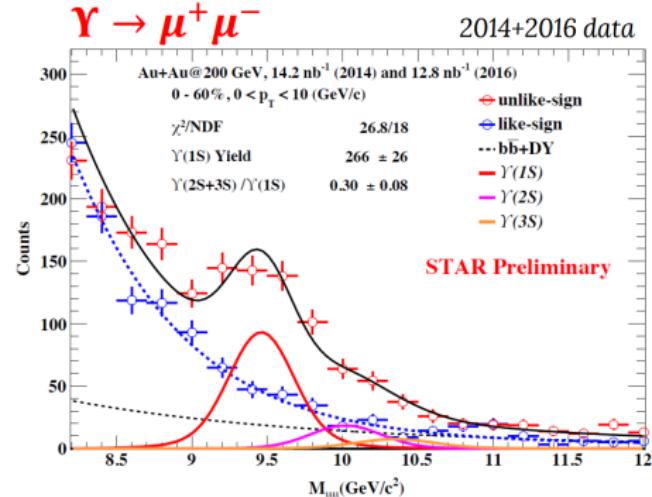
## $\Upsilon(1S + 2S + 3S)$

- Significantly improved precision over published results from  $R_{dAu}$ 
  - $R_{pAu}|_{|y|<0.5} = 0.82 \pm 0.10(\text{stat.})^{+0.08}_{-0.07}(\text{syst.}) \pm 0.10(\text{glob.})$
- Indication of  $\Upsilon(1S + 2S + 3S)$  suppression in p+Au collisions

$\Upsilon \rightarrow e^+e^-$



$\Upsilon \rightarrow \mu^+\mu^-$

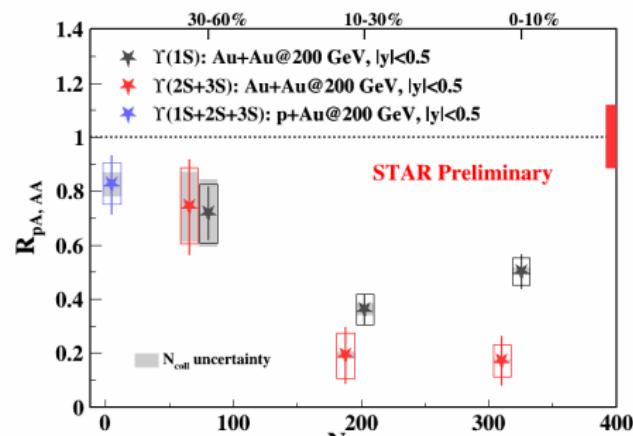
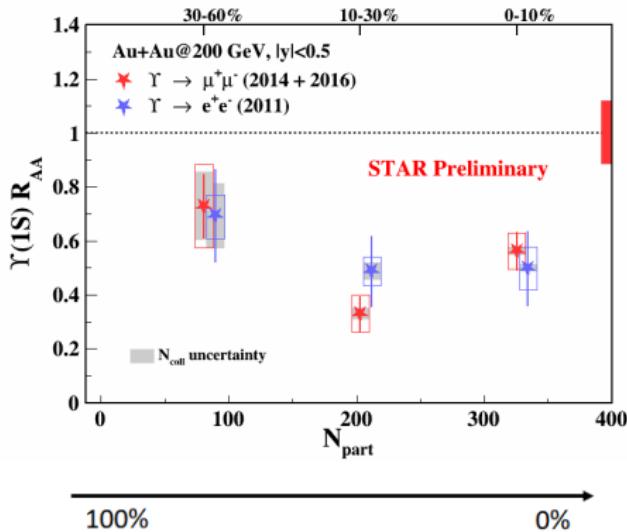


[Pengfei Wang, QM2018]

### Signal fits

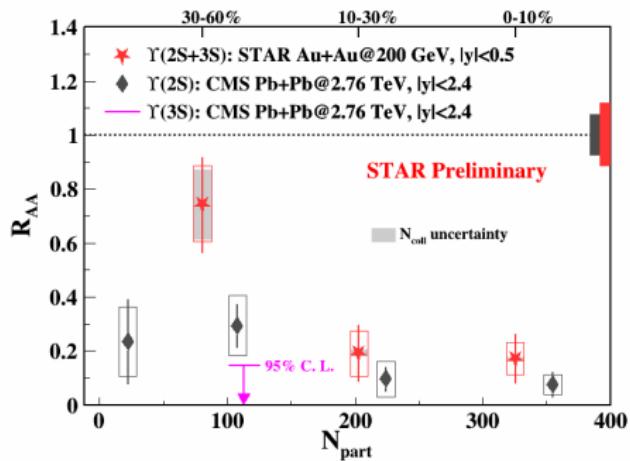
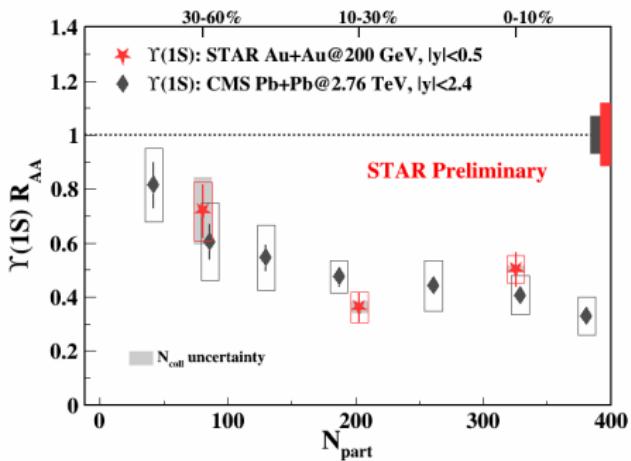
- 3 Crystal Ball fits for  $\Upsilon \rightarrow e^+e^-$
- 3 Gaussian fits for  $\Upsilon \rightarrow \mu^+\mu^-$ , because of less bremsstrahlung

# $R_{AuAu}$ vs. $N_{part}$



## $R_{AuAu}$ measured by STAR

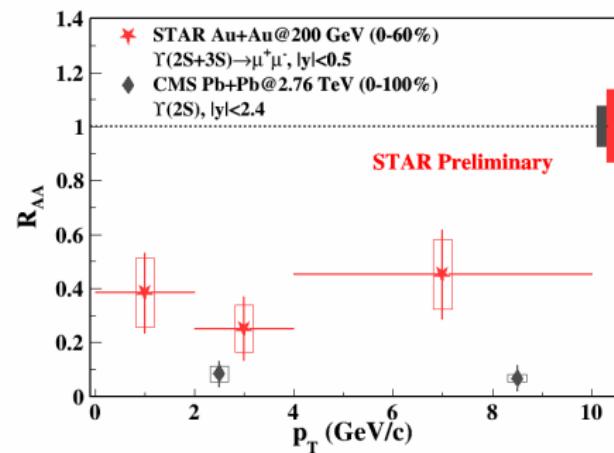
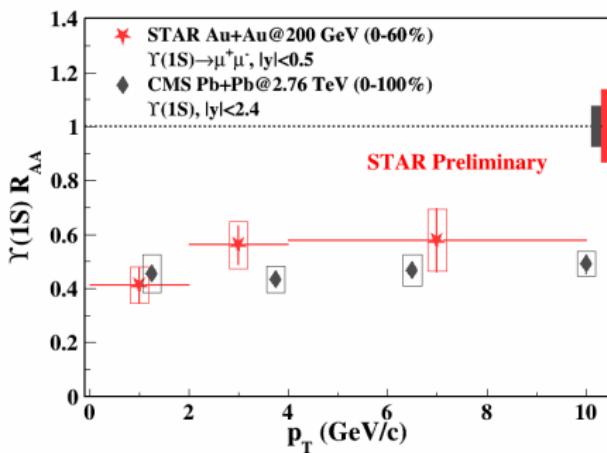
- Consistent results from dielectron and dimuon channels
- Both results combined in order to achieve better precision
- Similar level of suppression in peripheral collisions as in  $p + Au$
- Stronger suppression of  $\Upsilon(2S + 3S)$  than  $\Upsilon(1S)$  in central collisions



[Phys.Lett.B 770, 357-379(2017)]

### STAR vs. CMS

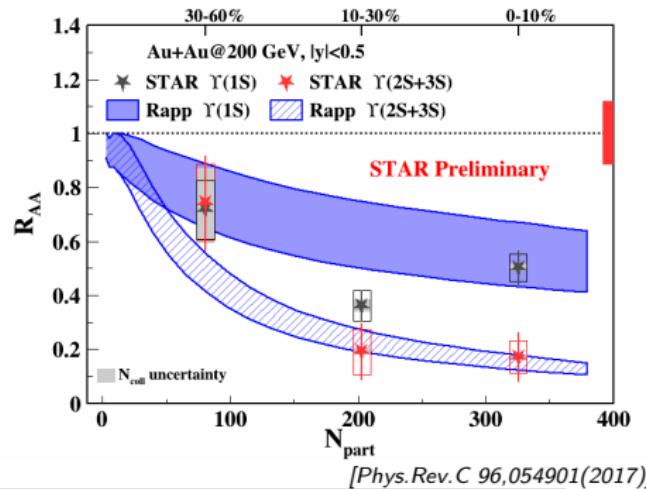
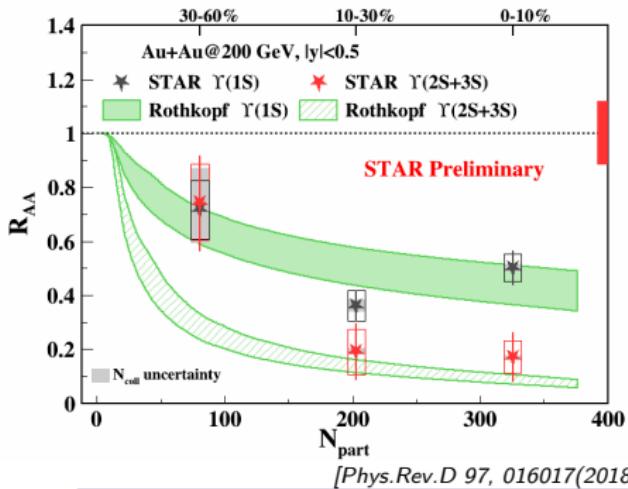
- Similar suppression for  $\Upsilon(1S)$ , despite higher medium temperature at the LHC
  - Regeneration? Larger at LHC than at RHIC
  - CNM effects
- Indication of smaller suppression for  $\Upsilon(2S + 3S)$  at RHIC than at LHC



[Phys.Lett.B 770, 357-379(2017)]

### Transverse momentum dependence

- Similar suppression for  $\Upsilon(1S)$  at RHIC and LHC
- Indications of stronger suppression of high- $p_T$   $\Upsilon(2S + 3S)$  at LHC than at RHIC
- Both consistent with flat dependence vs.  $p_T$



## Models

- Kroupaa, Rothkopf, Strickland
  - Lattice QCD-vetted potential for heavy quarks in hydrodynamic-modeled medium
  - No regeneration, no CNM effects
- De, He, Rapp
  - Quarkonium in-medium binding energy described by thermodynamic T-matrix calculations with internal energy potential (strongly bound scenario)
  - Includes both regeneration and CNM effects
- Both models agree with STAR  $\Upsilon(1S)$  data
- Rothkopf's model underestimates the STAR  $\Upsilon(2S + 3S)$  results for 30 – 60% centrality

# Summary

## p+p collisions at $\sqrt{s} = 200 \text{ GeV}$ and $\sqrt{s} = 500 \text{ GeV}$

- The  $\Upsilon(1S)$  spectra can be reasonably described by CEM calculations.
- Flatter rapidity distribution for  $\Upsilon$  at  $\sqrt{s} = 500 \text{ GeV}$  than at  $\sqrt{s} = 200 \text{ GeV}$ .
- Measured  $\frac{\Upsilon(nS)}{\Upsilon(1S)}$  vs. multiplicity at 500 GeV - no strong dependence.
  - Ratios slightly lower than world data.
- Dependence of  $\Upsilon$  production on event activity.
  - Similar trends observed for  $J/\psi$  and  $\Upsilon$  at RHIC and LHC.
  - Predictions from PYTHIA8 and Percolation model can qualitatively describe the trend observed in data.

## p+Au collisions at $\sqrt{s} = 200 \text{ GeV}$

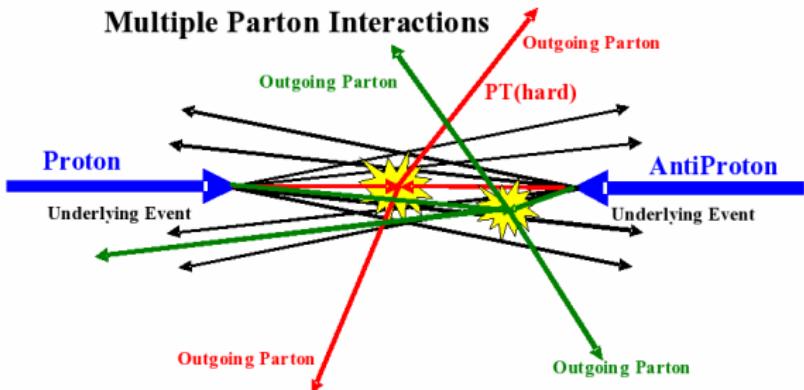
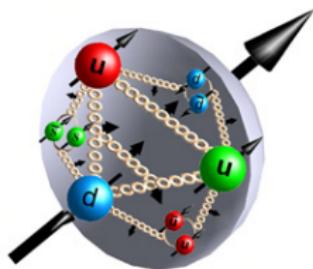
- Indication of  $\Upsilon$  suppression

## A+A collisions at $\sqrt{s} = 200 \text{ GeV}$

- Consistent results from dielectron and dimuon channels - combined for better precision
- Similar suppression of  $\Upsilon(1S)$  at RHIC and LHC
- Stronger suppression of  $\Upsilon(2S + 3S)$  than  $\Upsilon(1S)$  in central collisions
  - Sequential suppression
  - Hint of smaller suppression at RHIC than at LHC
- Data consistent with model calculations

**Thank you for attention!**

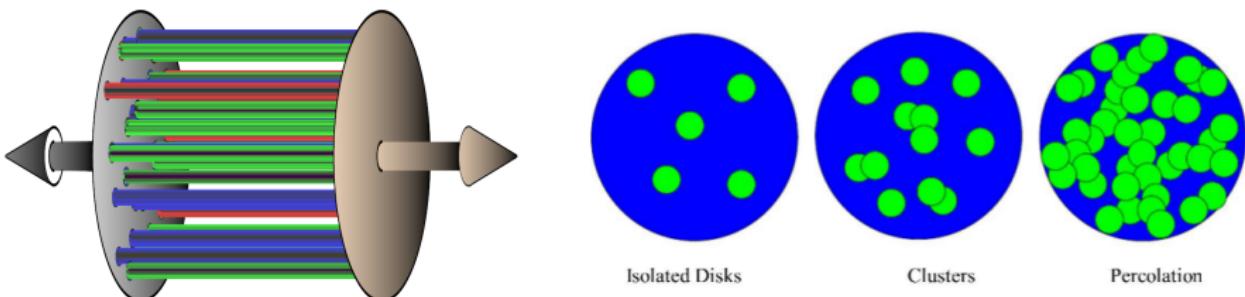
# Multiple parton interactions (MPI)



<https://www.bnl.gov/rhic/images/proton-with-gluons-300px.jpg>

<http://www.desy.de/~jung/multiple-interactions/may06/mi-rick.gif>

- Protons are complex objects consisting of constituent quarks, sea quarks and gluons.
- Multiple parton interactions (MPI) may happen in  $p + p$  collision - implemented in PYTHIA.
  - Besides the main hard process, there may be additional hard and soft processes in MPI.
- As implemented in PYTHIA8, heavy quarks can also be produced during MPI.
- MPI together with initial- (ISR), final-state radiation (FSR) and beam remnants define the event activity, which can be characterized experimentally using the charged particle multiplicity.



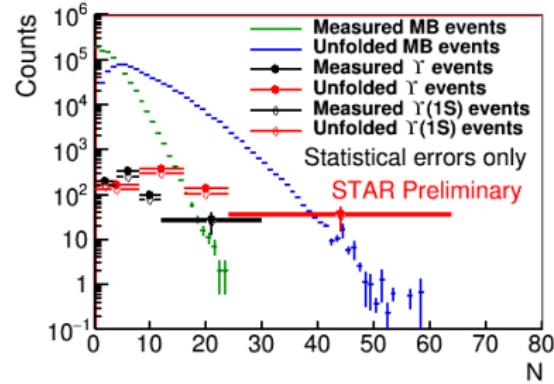
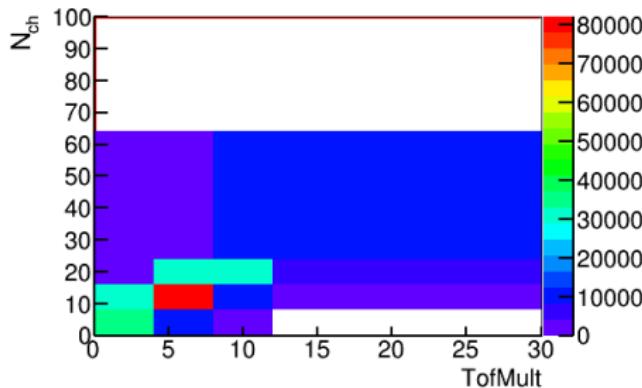
[Ann.Rev.Nucl.Part.Sci.60, 463-489(2010)] [Proc.of SPIE, 100313U(2016)]

- Models particle production originating from strings of color field formed in  $p + p$  collisions.
- Soft particle production damped by interaction of overlapping strings.
- Predicts quadratic dependence of normalized yield for particles from hard processes vs. normalized charged particle multiplicity in high multiplicity events.

$$\frac{N_{hard}}{\langle N_{hard} \rangle} = \langle \rho \rangle \left( \frac{\frac{dN_{ch}}{d\eta}}{\left\langle \frac{dN_{ch}}{d\eta} \right\rangle} \right)^2 \quad [\text{Phys.Rev. C, 86, 034903 (2012)}]$$

# Multiplicity distribution via unfolding

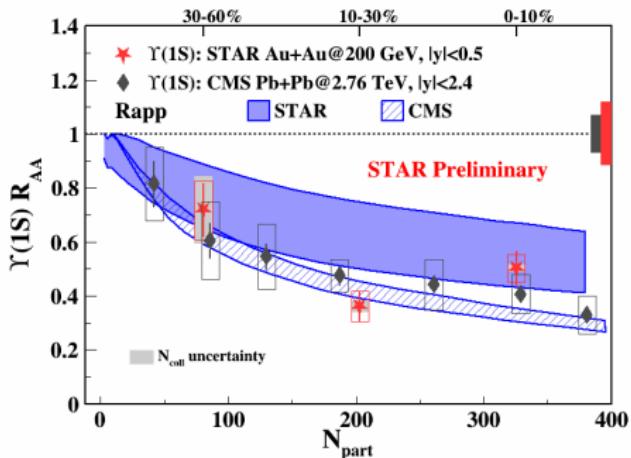
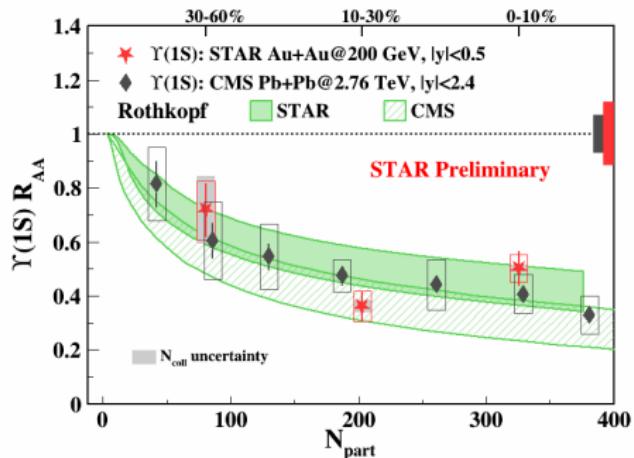
Response matrix for  $\gamma$  events



## Unfolding method used for multiplicity dependent studies

- ① A response matrix is obtained using the PYTHIA8 event generator for both min-bias and  $\gamma$  events taking into account reconstruction efficiency
- ② The measured distributions are unfolded with their respective response matrices
- ③ This procedure yields the unfolded (true) distribution

# STAR and CMS $\Upsilon(1S)$ vs. models

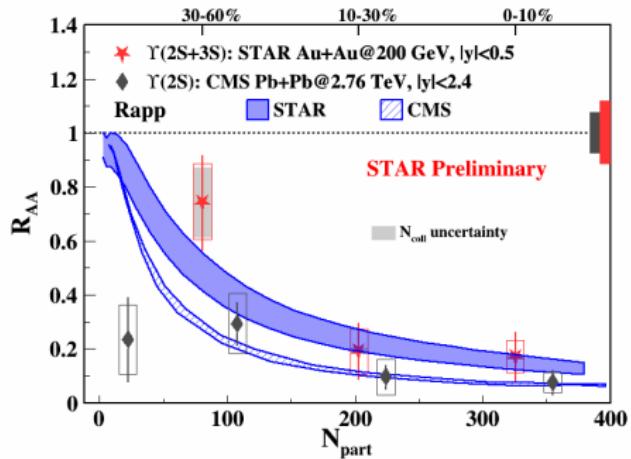
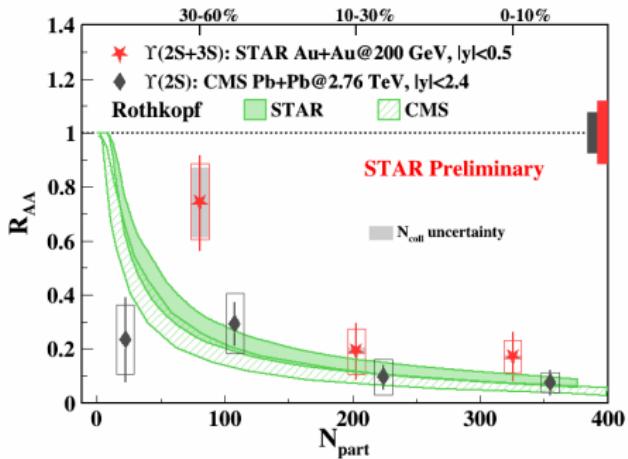


[Phys. Rev.D 97,(2018)016017], [Phys. Rev.C 96,(2017)054901]

## $\Upsilon(1S)$ vs. models

- Both models consistent with the data

# STAR and CMS $\Upsilon(2S + 3S)$ vs. models

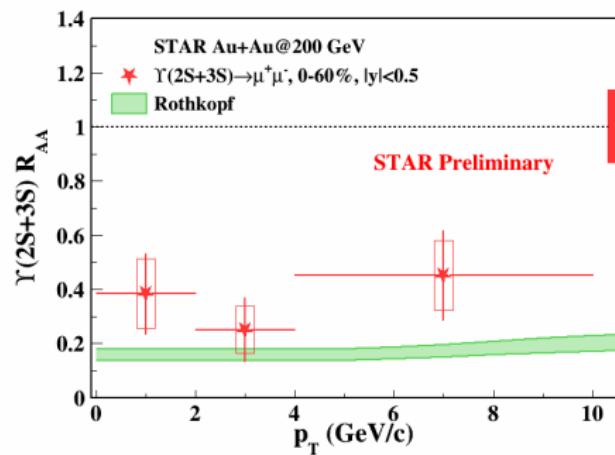
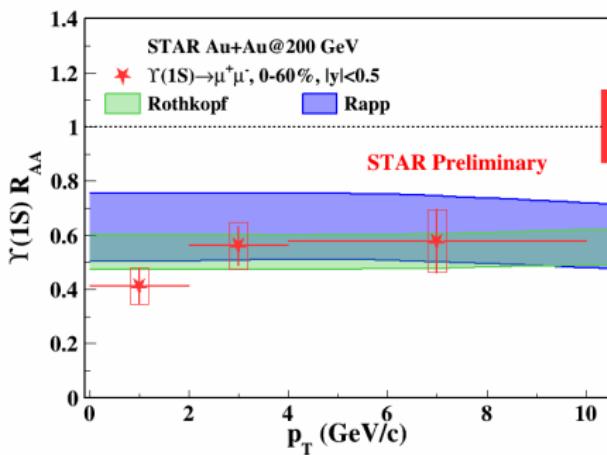


[Phys. Rev.D 97,016017(2018)], [Phys. Rev.C 96,054901(2017)]

## $\Upsilon(2S + 3S)$ vs. models

- Both models consistent with the data in central and semi-central collisions

# $R_{AA}$ vs. $p_T$ vs. models

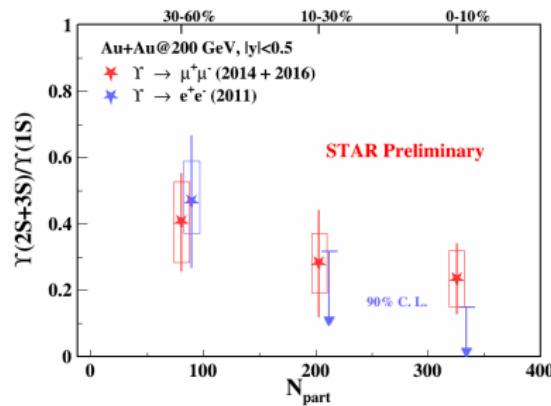


[Phys.Rev.D 97,016017(2018)], [Phys.Rev.C 96,054901(2017)]

## $R_{AA}$ vs. $p_T$ vs. models

- Both models consistent with the data
- Rothkopf's model slightly lower than  $\Upsilon(2S + 3S)$
- Flat vs.  $p_T$

# Ratio vs. $N_{part}$



$\frac{Y(2S+3S)}{Y(1S)}$  vs.  $N_{part}$

- Both channels consistent