Overview of \varUpsilon production studies performed with the STAR experiment

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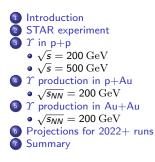
ICHEP 2020, Prague, Czech Republic 30.7.2020



EUROPEAN UNION European Structural and Investment Funds Operational Programme Research, Development and Education



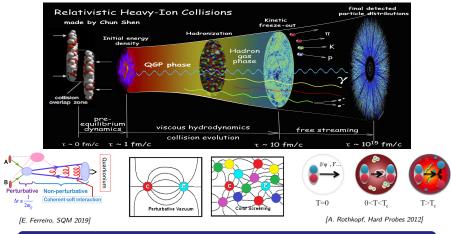
The work was also supported from European Regional Development Fund-Project "Center of Advanced Applied Science" No. CZ.02.1.01/0.0/0.0/16-019/0000778 and by the grant LTT18002 of Ministry of Education, Youth and Sports of the Czech Republic.





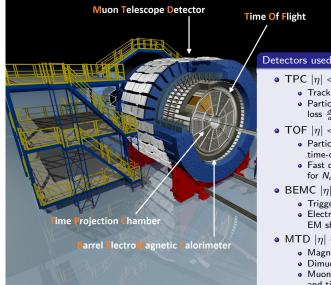


Υ - a probe of quark-gluon plasma



Quark-gluon plasma studies with \varUpsilon states

- ullet QGP can be created in heavy-ion collisions and probed using arLambda states
- $\Upsilon(nS)$ states $\Upsilon = b\bar{b} \ (m_{u,d} << m_b)$:
 - · Contain heavy quarks, created at the early stages of the collision
 - Dissociate at high T in QGP via Debye-like screening [Phys.Lett.B 178(4),416-422(1986)]



Detectors used for quarkonium studies

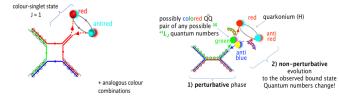
- TPC $|\eta| <$ 1, 0 $\leq \phi < 2\pi$
 - Tracking momentum measurement
 - Particle identification based on energy loss $\frac{dE}{dx}$
- TOF $|\eta| < 1$, $0 \le \phi < 2\pi$
 - Particle identification based on time-of-flight
 - Fast detector used to remove pile-up for N_{ch} determination
- BEMC $|\eta| <$ 1, 0 $\leq \phi < 2\pi$
 - Trigger on high-p_T electrons
 - Electron identification via *E/p* and EM shower shape
- MTD $|\eta| < 0.5$, 45% in ϕ
 - Magnet used as hadron absorber
 - Dimuon trigger
 - Muon identification utilizing position and time-of-flight information
 - μ less bremsstrahlung than e

L. Kosarzewski

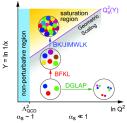
Υ production in p+p collisions



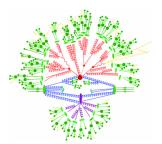
• Study of production mechanism: Color Singlet vs. Color Octet channels



- Events with high charged particle multiplicity
 - Interplay between hard and soft processes
 - Saturation effects/multiple parton interactions

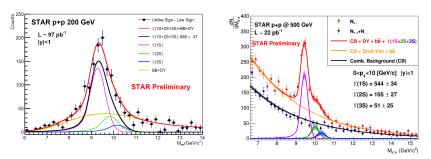






$\Upsilon ightarrow e^+e^-$ signal in p+p at 200 and 500 GeV

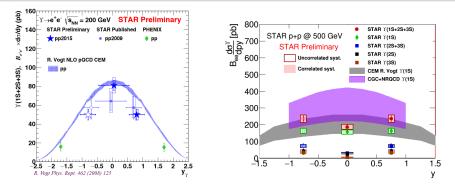
 $\Upsilon
ightarrow e^+ e^-$ in 2015 p+p $\sqrt{s} = 200 \text{ GeV}$ $\Upsilon \rightarrow e^+ e^-$ in 2011 p+p $\sqrt{s} = 500 \text{ GeV}$



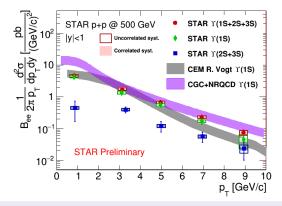
Signal extraction

- Challenging to extract individual $\Upsilon(nS)$ yields!
- Fit *m_{ee}* histograms with:
 - Signal lineshapes from STAR detector simulation
 - Backgrounds: combinatorial, bb, Drell-Yan

arLambda rapidity dependence in p+p at 200 and 500 ${ m GeV}$



- $\sqrt{s} = 200 \text{ GeV}$ STAR data:
 - Slightly narrower than Color Evaporation Model (CEM)
- $\sqrt{s} = 500 \text{ GeV}$ data:
 - Separate Υ(1S) and Υ(2S)(NEW!), Υ(3S)(NEW!) spectra.
 - Flatter rapidity spectrum compared to $\sqrt{s}=200~{
 m GeV}$
 - Dip at mid-rapidity for $\Upsilon(2S+3S) \approx 2\sigma$ level from flat, mostly due to low $\Upsilon(3S)$ yield
 - CEM model (inclusive) consistent with the measurement for $\Upsilon(1S)$ [Phys.Rev.C 92 034909(2015)]
 - CGC+NRQCD predictions for direct $\Upsilon(1S)$ are above the data for inclusive $\Upsilon(1S)$ [Phys.Rev.D 94, 014028(2016)],[Phys.Rev.Lett. 113, 192301(2014)]



- Separate $\Upsilon(1S)$ and $\Upsilon(2S+3S)$ spectra.
- CEM calculation for inclusive $\Upsilon(1S)$

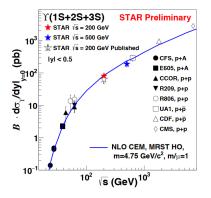
[Phys.Rev.C 92 034909(2015)]

- Agree with data reasonably well
- CGC+NRQCD for direct \varUpsilon

[Phys.Rev.D 94, 014028(2016)] [Phys.Rev.Lett. 113, 192301(2014)]

• $\Upsilon(1S)$: model calculation is above the data points. Caveat: additional corrections are needed at low p_T according to authors.

$\Upsilon \rightarrow e^+e^-$ integrated cross section in p+p

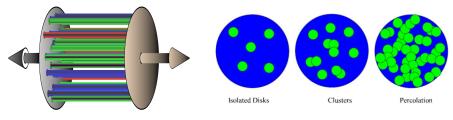


STAR [Phys.Lett.B 735,127–137(2014)] CDF [Phys.Rev.Lett. 88,161802(2002)] CMS [Phys.Rev.L 88,161802(2002)] CFS [Phys.Rev.Lett. 39,1240-1242(1977)] CFS [Phys.Rev.Lett. 41,646–687(1978)] CFS [Phys.Rev.Lett. 42,486–489(1979)] CFS [Phys.Rev.Lett. 45,1962–1964(1985)] E605 [Phys.Rev.D 43,2815–2835(1991)]] E605 [Phys.Rev.D 43,2815–2835(1991)]] E605 [Phys.Rev.D 43,2815–2835(1991)]] ECOR [Phys.Lett.B 73,394–402(1979)] L. Camilleri, T.B.W. Kirk, H.D.I. Abarbanel (Eds.) E666 [Phys.Rev.Lett. 91,481–486(1980)]

Integrated cross section

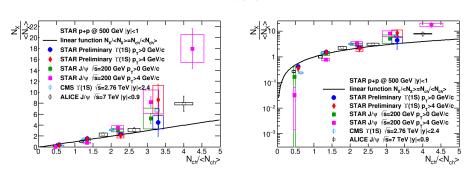
- $B_{ee} \frac{d\sigma}{dy}|_{|y|<0.5} = 81 \pm 5(stat) \pm 8(syst)$ pb in p+p collisions at $\sqrt{s} = 200$ GeV
- $B_{ee} \frac{d\sigma}{dy}|_{|y|<0.5} = 186 \pm 14(stat) \pm 33(syst)$ pb in p+p collisions at $\sqrt{s} = 500$ GeV
- STAR results follow the world data trend
- Consistent with the Color Evaporation Model calculation [Phys.Rep. 462, pp.125–175(2008)]

Event activity in Percolation model



[Ann.Rev.Nucl.Part.Sci.60, 463-489(2010)] [Proc.of SPIE, 100313U(2016)]

- Pairs of partons form elongated strings of color field
- In high charged particle multiplicity (N_{ch}) events, many strings overlap and interact (percolation)
 - Interactions lower N_{ch} yield from soft processes
 - Hard processes mostly unaffected increase faster with N_{ch}
- Click to see backup slide 26 for more info!
- Similar to CGC



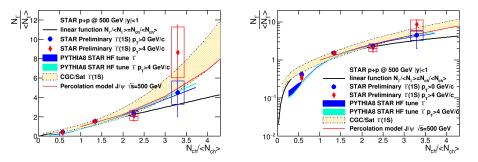
 $\gamma \rightarrow e^+ e^-$

[JHEP04,103(2014)],[Nucl.and Part.Phys. Proc., 276-278, pp.261–264(2016)],[Phys.Lett.B 712,165–175(2012)],[Phys.Lett.B 786,87-93(2018)]

- Distributions of N_{ch} fully corrected using unfolding procedure. See backup: 27
- ullet Similar trends at RHIC and LHC for \varUpsilon and J/ψ

Υ production vs. event activity - models

$$\Upsilon
ightarrow e^+e^-$$



- PYTHIA8 and String Percolation models reproduce the trend in the data [E. G. Ferreiro, C. Pajares, Phys. Rev. C, 86, 034903(2012)]
- CGC/Saturation model describes the data within large uncertainties [E. Levin, M. Siddikov, EPJC, 97(5), 376(2019)], [M. Siddikov, et al, arXiv:1910.13579 [hep-ph]]
- Suggest Υ production in MPI or saturation effects

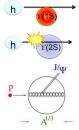


p+A collisions (Cold Nuclear Matter (CNM) effects):

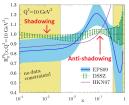
 Comover interactions - very small for Υ(1S) [Phys.Lett.B 503, 104(2001)]

• Nuclear absorption: σ_{abs}

- Nuclear PDFs: shadowing, anti-shadowing
- Studied by measuring Nuclear Modification Factor: $R_{pA} = \frac{\sigma_{inel}}{\langle N_{coll} \rangle} \frac{d^2 N_{p+Au}/dp_T dy}{d^2 \sigma_{pp}/dp_T dy}$

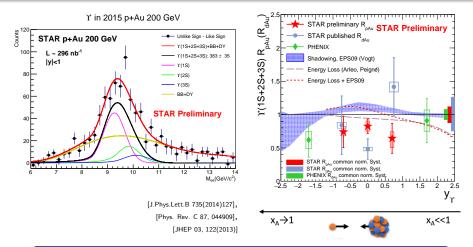


[L. Grandchamp, LBNL 2005]



[Nucl.Phys.A 926 24-33(2014)]

Υ production in p+Au at $\sqrt{s_{NN}} = 200 \text{ GeV}$



$\Upsilon(1S+2S+3S)$

- Improved precision over published results from R_{dAu}

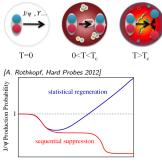
 - ~ 50% smaller statistical uncertainty vs. y $R_{PAu}|_{|y|<0.5} = 0.82 \pm 0.10(stat.)^{+0.08}_{-0.07}(syst.) \pm 0.10(glob.)$
- Indication of $\Upsilon(1S+2S+3S)$ suppression in p+Au collisions



Heavy ion (A+A) collisions (QGP+CNM effects):

 Quarkonium states dissociate at high temperature in QGP via Debye-like screening [Phys.Lett.B 178(4),416-422(1986)]

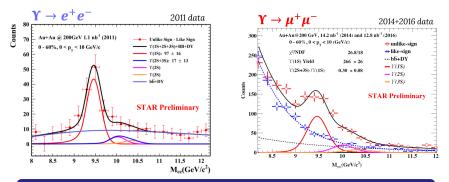
• Sequential suppression due to each $\Upsilon(nS)$ state dissociating at different $T \rightarrow$ estimate of T [Phys.Rev.D 64, 094015(2001)]



Energy Density

[Nucl.Phys.B (Proc.Suppl.) 214, 3-36(2011)]

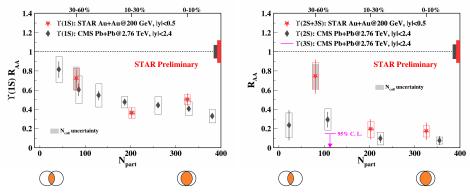
- Modified feed-down pattern
- Regeneration negligible at RHIC!



Υ signal

- Υ measured in both e^+e^- and $\mu^+\mu^-$
- Combined R_{AA} for better precision

Υ suppression - R_{AA}



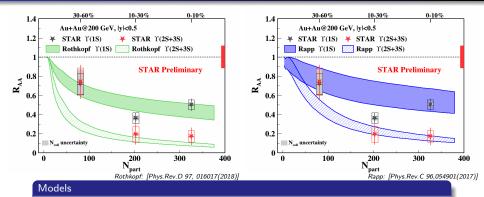
CMS: [Phys.Lett.B 770, 357-379(2017)]

STAR vs. CMS

- Similar suppression for $\Upsilon(1S)$, despite higher medium temperature at the LHC
 - · Suppression of excited states contribution
 - Regeneration? Larger at LHC than at RHIC
 - CNM effects need better constraints
- Indication of smaller suppression for $\Upsilon(2S+3S)$ at RHIC than at LHC

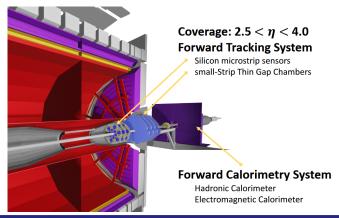
L. Kosarzewski

Υ : STAR vs. models



- Kroupaa, Rothkopf, Strickland
 - Lattice QCD-vetted potential for heavy quarks in hydrodynamic medium
 - No regeneration, no CNM effects
- De, He, Rapp
 - Quarkonium in-medium binding energy described by thermodynamic T-matrix calculations with internal energy potential (strongly bound scenario)
 - Includes both regeneration and CNM effects
- Both models agree with STAR $\Upsilon(1S)$ data
- Indication that Rothkopf's model underestimates the STAR $\varUpsilon(2S+3S)$ results for 30-60% centrality

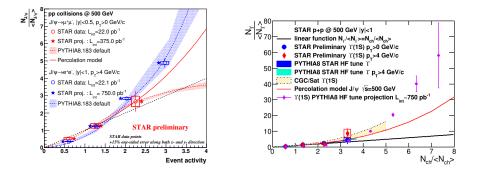
Upgrades and plans for 2022+



Future plans for STAR

- iTPC already running improved momentum resolution
- Forward upgrade new detectors
- High integrated luminosity for precision quarkonium production studies
- And more!

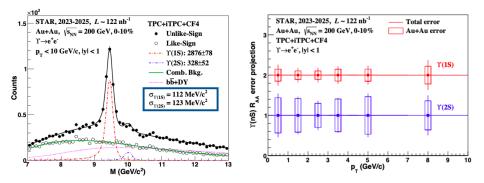
J/ψ and Υ in p+p projections 2017+2022



Projections 2017+2022

- High precision measurement of J/ψ and Υ dependence on normalized N_{ch}
- Very high integrated luminosity $\mathcal{L}_{int} \sim 750 \ \mathrm{pb}^{-1}$ for BHT e and $\mathcal{L}_{int} \sim 375 \ \mathrm{pb}^{-1}$ for $\mu\mu$ triggers
- Possible to discriminate different models

ΥR_{AA} projections 2023+



Projections 2023+

- High precision measurement:
 - High integrated luminosity $\mathcal{L}_{int} \sim 122 \; \mathrm{nb}^{-1}$
 - Improved momentum resolution
 - Low material budget less background
- R_{AA} of $\Upsilon(1S)$ and $\Upsilon(2S)$ vs:
 - centrality, p_T

Summary

arY in p+p collisions at $\sqrt{s}=200~{ m GeV}$ and $\sqrt{s}=500~{ m GeV}$

- $\varUpsilon(1S)$ data reasonably described by CEM model, while overestimated by CGC+NRQCD
- Rapidity spectrum flatter at $\sqrt{s}=500~{
 m GeV}$ than at $\sqrt{s}=200~{
 m GeV}$
- Similar trends for J/ψ and Υ vs. $N_{ch}/< N_{ch}>$ at RHIC and LHC
 - $\bullet\,$ Qualitatively reproduced by the models may be discriminated with 2017+2022 datasets

\varUpsilon in p+Au collisions at $\sqrt{s_{NN}}=200~{\rm GeV}$

• Indication of $\Upsilon(1S+2S+3S)$ suppression $R_{pAu}|_{|y|<0.5} = 0.82 \pm 0.10(stat.)^{+0.08}_{-0.07}(syst.) \pm 0.10(glob.)$

Υ in Au+Au collisions at $\sqrt{s_{NN}}=200~{ m GeV}$

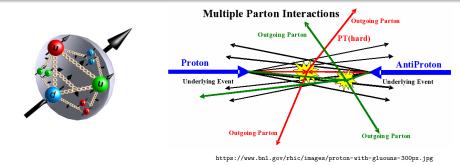
- Υ R_{AA} measured in dielectron and dimuon channels combined for better precision
- Similar suppression of $\Upsilon(1S)$ at RHIC and LHC
- Stronger suppression of $\Upsilon(2S+3S)$ than $\Upsilon(1S)$ in central collisions
 - Sequential suppression
 - Hint of smaller $\Upsilon(2S+3S)$ suppression at RHIC than at LHC
- $\Upsilon(1S)$, $\Upsilon(2S+3S)$ R_{AA} consistent with model calculations
- High precision $\Upsilon(1S)$, $\Upsilon(2S)$ R_{AA} measurements vs. p_T and centrality beyond 2023+

STAR presentations at ICHEP 2020

- 976. Measurements of J/ψ photoproduction in ultra-peripheral collisions at RHIC
 Jaroslav Adam, 29 July 2020 (Wednesday), 19:18
- 1052. Central exclusive production of charged particle pairs in proton-proton collisions at $\sqrt{s}=200~{\rm GeV}$ with the STAR detector at RHIC
 - Rafal Sikora, 30 July 2020 (Thursday), 10:25
- 537. Measurements of open charm hadrons in Au+Au collisions at $\sqrt{s_{NN}}=200~{\rm GeV}$ by the STAR experiment
 - Lukáš Kramárik, 30 July 2020 (Thursday), 12:12
- \bullet 414. Production of D^\pm mesons in Au+Au collisions at $\sqrt{s_{NN}}=200~{\rm GeV}$ at the STAR experiment
 - Jan Vaněk(poster), 30 July 2020 (Thursday), 13:39
- 611. Geometry and Dynamics in Heavy-ion Collisions Seen by the Femtoscopy Method in the STAR experiment
 - Prof. Hanna Zbroszczyk, 31 July 2020 (Friday), 8:30
- 686. Study of the central exclusive production of $\pi^+\pi^- K^+K^-$ and $p\bar{p}$ pairs in proton-proton collisions at $\sqrt{s} = 510 \text{ GeV}$ with the STAR detector at RHIC
 - Tomáš Truhlář(poster), 31 July 2020 (Friday), 13:30

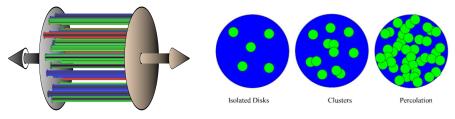
BACKUP

Multiple parton interactions (MPI)



http://www.desy.de/~jung/multiple-interactions/may06/mi-rick.gif

- Protons are complex objects consisting of constituent quarks, sea quarks and gluons.
- Multiple parton interactions (MPI) may happen in *p* + *p* collison implemented in PYTHIA.
 - Besides the main hard process, there may be additional hard and soft processes in MPI.
- As implemented in PYTHIA8, heavy quarks can also be produced during MPI.
- MPI together with initial- (ISR), final-state radiation (FSR) and beam remnants define the event activity, which can be characterized experimentally using the charged particle multiplicity.

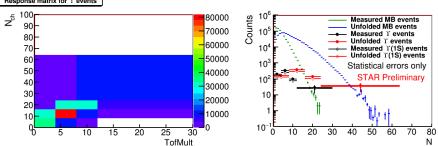


[Ann.Rev.Nucl.Part.Sci.60, 463-489(2010)] [Proc.of SPIE, 100313U(2016)]

- Models particle production originating from strings of color field formed in p + p collisions.
- Soft particle production dampened by interaction of overlapping strings.
- Predicts quadratic dependence of normalized yield for particles from hard processes vs. normalized charged particle multiplicity in high multiplicity events.

$$\frac{N_{hard}}{\langle N_{hard} \rangle} = \langle \rho \rangle \left(\frac{\frac{dN_{ch}}{d\eta}}{\langle \frac{dN_{ch}}{d\eta} \rangle} \right)^2 \text{ [Phys.Rev. C, 86, 034903 (2012)]}$$

Multiplicity distribution via unfolding

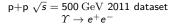


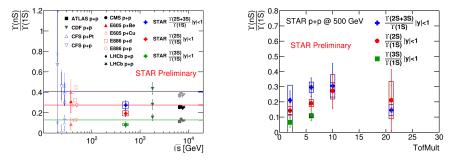
Response matrix for T events

Unfolding method used for multiplicity dependent studies

- A response matrix is obtained using the PYTHIA8 event generator for both min-bias and Υ events taking into account reconstruction efficiency
- 2 The measured distributions are unfolded using their respective response matrices
- This procedure yields the unfolded (true) distribution 3
- 4 Similar procedure used for J/ψ
- Measured N_{ch} distribution obtained from p+p $\sqrt{s} = 500 \text{ GeV}$ 2009 data 6
- Measured distribution of Υ events obtained from p+p $\sqrt{s} = 500 \text{ GeV} 2011 \text{ data}$

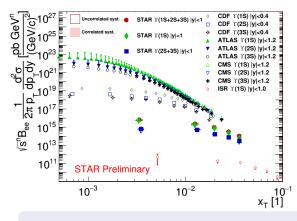
Cross section ratios: $\Upsilon(nS)/\Upsilon(1S)$





[[]W. Zha, et al, Phys.Rev.C 88,067901(2013)]

- Left plot: cross section ratios measured in 500 GeV p+p collisions are slightly below (within 2σ) world data average, shown as solid lines in the left plot.
- Right plot: Ratios vs. TofMult no strong multiplicity dependence observed.
- TofMult: number of tracks matched to TOF within $|\eta| < 1$, $p_T > 0.2 \, {\rm GeV/c}$ (uncorrected)



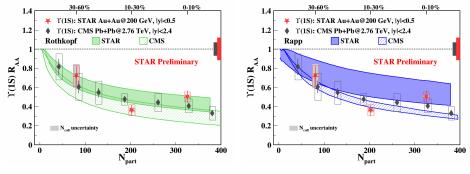
 $\begin{array}{l} {\rm STAR} \; p + p \; \sqrt{s} = 500 \; {\rm GeV} \\ {\rm ATLAS} \; p + p \; \sqrt{s} = 7 \; {\rm TeV} \\ {\rm [Phys, Rev. D87, 052004(2013)]} \\ {\rm CMS} \; p + p \; \sqrt{s} = 7 \; {\rm TeV} \\ {\rm [Phys. Lett. B} \; 749, 14-34(2015)] \\ {\rm CDF} \; p + \bar{p} \; \sqrt{s} = 1.8 \; {\rm TeV} \\ {\rm [Phys. Rev. Lett. \; 88, 161802(2002)] } \\ {\rm ISR} \; p + \bar{p} \; \sqrt{s} = 53, 63 \; {\rm GeV} \\ {\rm [Phys. Lett. B} \; 91, 481-4486(1980)] } \end{array}$

•
$$x_T = \frac{2p_T}{\sqrt{s}}, \sigma^{inv} \equiv E \frac{d^3\sigma}{d^3p} = \frac{F(x_T)}{p_T^{n(x_T,\sqrt{s})}} = \frac{F'(x_T)}{\sqrt{s}^{n(x_T,\sqrt{s})}}$$

[JHEP06,035(2010)]

- pQCD predicts that spectra of hard processes should follow x_T scaling check with n = 5.6 (number of partons taking active part in the process) obtained for J/ψ [*Phys.Rev.C* 80, 041902(2009)]
- No clear scaling observed, some indication for LHC data at high p_T

STAR and CMS $\Upsilon(1S)$ vs. models

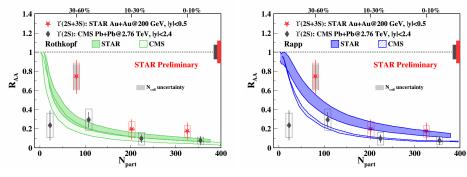


[Phys.Rev.D 97,(2018)016017], [Phys.Rev.C 96,(2017)054901]

$\Upsilon(1S)$ vs. models

• Both models consistent with the data

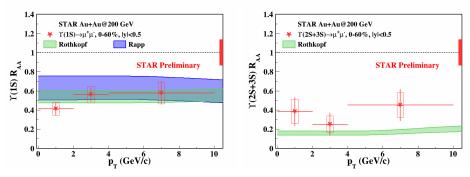
STAR and CMS $\Upsilon(2S+3S)$ vs. models



[Phys.Rev.D 97,016017(2018)], [Phys.Rev.C 96,054901(2017)]

$\Upsilon(2S+3S)$ vs. models

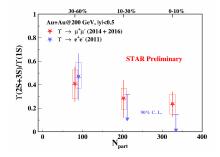
· Both models consistent with the data in central and semi-central collisions



[Phys.Rev.D 97,016017(2018)], [Phys.Rev.C 96,054901(2017)]

R_{AA} vs. p_T vs. models

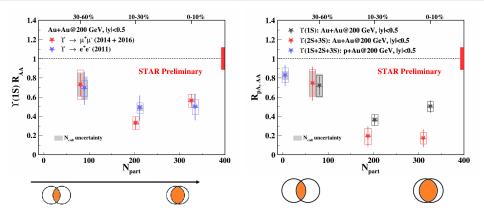
- Both models consistent with the data
- Rothkopf's model slightly lower than $\Upsilon(2S+3S)$
- Flat vs. p_T





• Both channels consistent

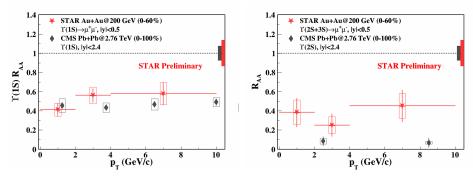
ΥR_{AuAu} vs. N_{part}



R_{AuAu} measured by STAR

- Consistent results from dielectron and dimuon channels
- Both results combined in order to achieve better precision
- Similar level of suppression in peripheral collisions as in p + Au
- Stronger suppression of $\Upsilon(2S+3S)$ than $\Upsilon(1S)$ in central collisions

ΥR_{AA} vs. p_T



CMS: [Phys.Lett.B 770, 357-379(2017)]

Transverse momentum dependence

- Similar suppression for $\Upsilon(1S)$ at RHIC and LHC
- Indication of stronger suppression of high- $p_T \ \Upsilon(2S+3S)$ at LHC than at RHIC
- Both consistent with flat dependence vs. p_T