

Recent results of Υ production measured with the STAR experiment

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EUROPEAN UNION
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Development and Education



MINISTRY OF EDUCATION,
YOUTH AND SPORTS

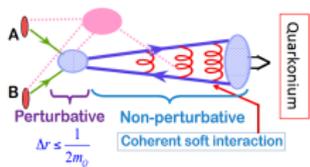
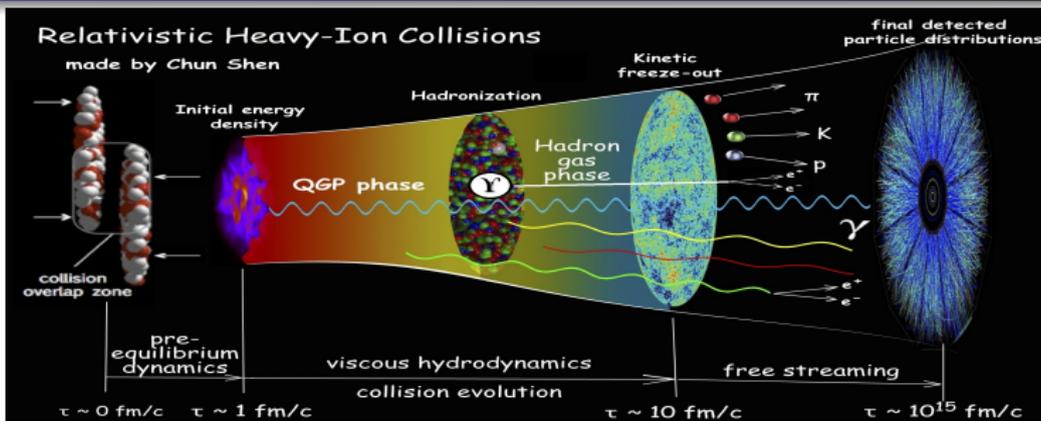


This work was also supported from European Regional Development Fund-Project "Center of Advanced Applied Science" No. CZ.02.1.01/0.0/0.0/16-019/0000778 and by the grant LTT18002 of Ministry of Education, Youth and Sports of the Czech Republic.

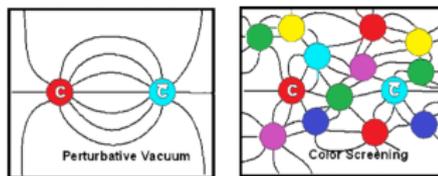
- 1 Introduction
- 2 STAR experiment
- 3 Υ in p+p
 - $\sqrt{s} = 200$ GeV
 - $\sqrt{s} = 500$ GeV
- 4 Υ production in p+Au
 - $\sqrt{s_{NN}} = 200$ GeV
- 5 Υ production in Au+Au
 - $\sqrt{s_{NN}} = 200$ GeV
- 6 Summary



Υ - a probe of quark-gluon plasma



[E. Ferreira, SQM 2019]



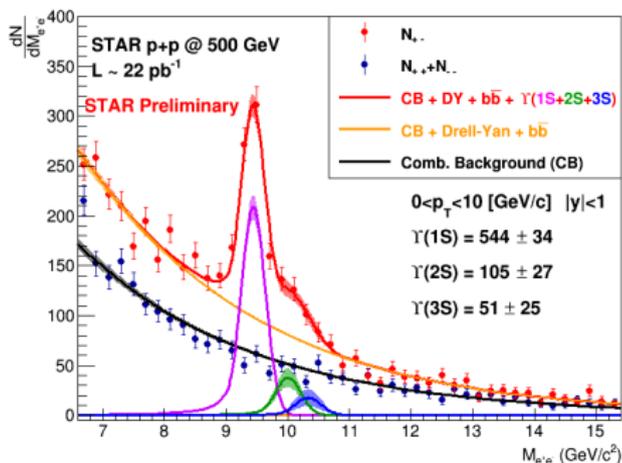
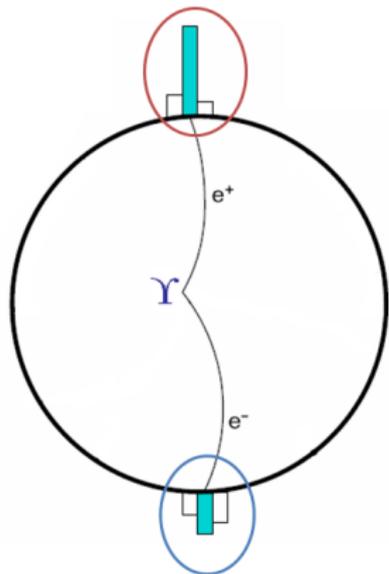
[A. Rothkopf, Hard Probes 2012]

Quark-gluon plasma studies with Υ states

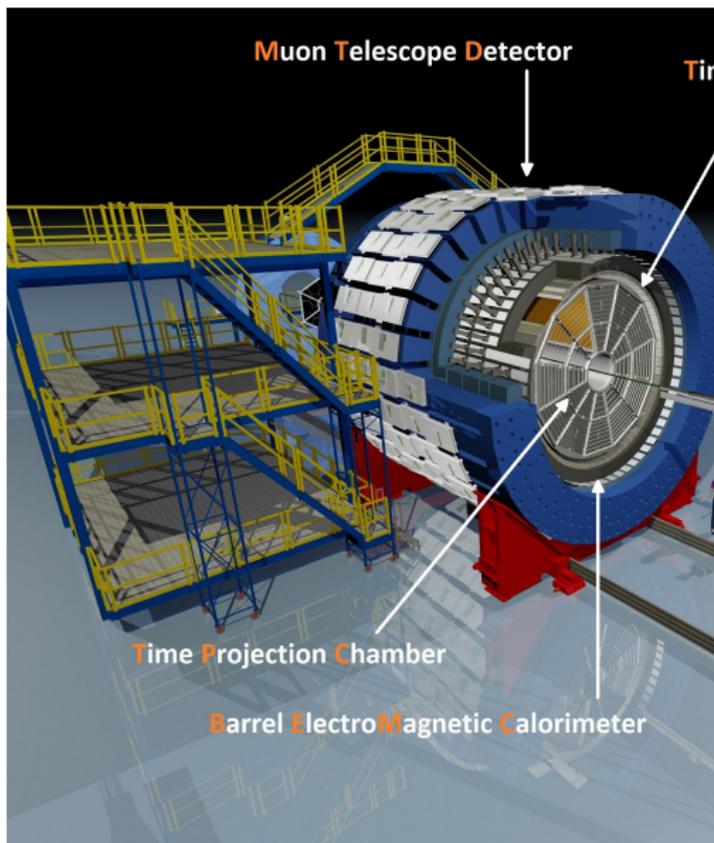
- QGP can be created in heavy-ion collisions and probed using Υ states
- $\Upsilon(nS)$ states $\Upsilon = b\bar{b}$ ($m_{u,d} \ll m_b$):
 - Contain heavy quarks, created at the early stages of the collision
 - Dissociate at high T in QGP via Debye-like screening
[T.Matsui, H.Satz, Phys.Lett.B 178(4),416-422(1986)]
 - Similar to J/ψ

- 1 Upsilon mesons are reconstructed through $\Upsilon \rightarrow e^+e^-$ ($\Upsilon \rightarrow \mu^+\mu^-$) decays
- 2 Detect and select electrons and calculate invariant mass m_{ee} in order to obtain the Υ peak

$$m_{ee} = \sqrt{(E_1 + E_2)^2 - (\vec{p}_1 + \vec{p}_2)^2}$$

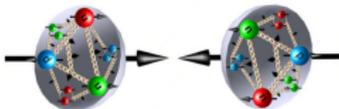


- 3 Apply efficiency corrections to e^+e^- ($\mu^+\mu^-$) obtained from simulated $\Upsilon \rightarrow e^+e^-$ ($\Upsilon \rightarrow \mu^+\mu^-$) decays

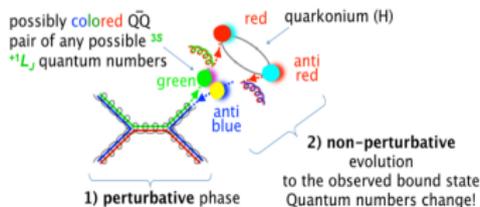
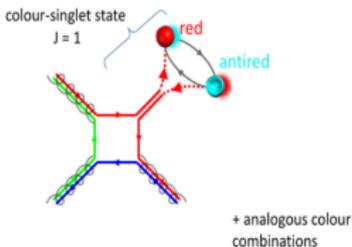


Detectors used for quarkonium studies

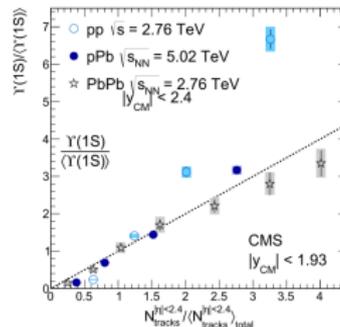
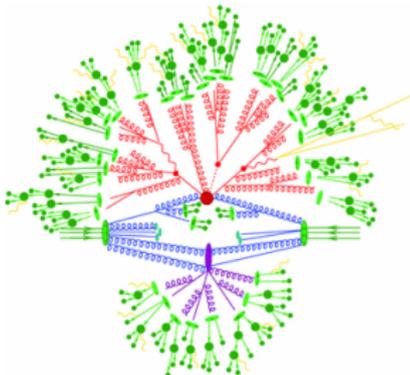
- TPC $|\eta| < 1, 0 \leq \phi < 2\pi$
 - Tracking - momentum measurement
 - Particle identification based on energy loss $\frac{dE}{dx}$
- TOF $|\eta| < 1, 0 \leq \phi < 2\pi$
 - Particle identification based on time-of-flight
 - Fast detector used to remove pile-up for N_{ch} determination
- BEMC $|\eta| < 1, 0 \leq \phi < 2\pi$
 - Trigger on high- p_T electrons
 - Electron identification via E/p and EM shower shape
- MTD $|\eta| < 0.5, 45\%$ in ϕ
 - Magnet used as hadron absorber
 - Dimuon trigger
 - Muon identification utilizing position and time-of-flight information
 - μ - less bremsstrahlung than e



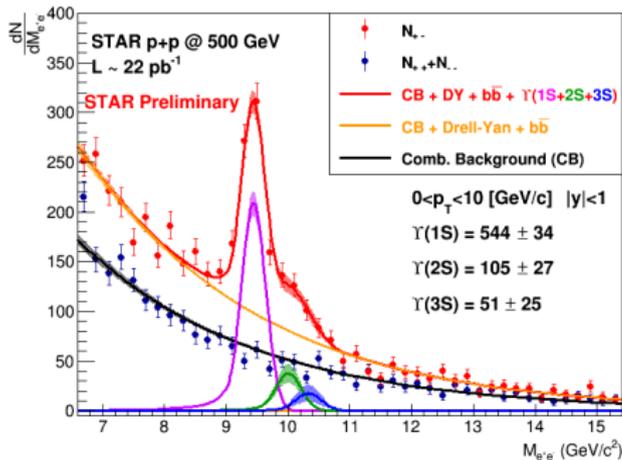
- Study of production mechanism: Color Singlet vs. Color Octet channels



- Charged particle multiplicity N_{ch} dependence
 - Studied with $\Upsilon / \langle \Upsilon \rangle$ vs $N_{ch} / \langle N_{ch} \rangle$
 - Interesting stronger-than-linear increase observed at CMS
 - Investigates interplay between hard and soft processes

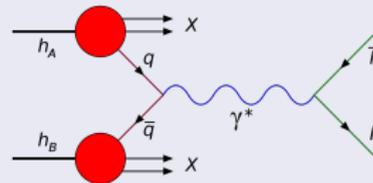


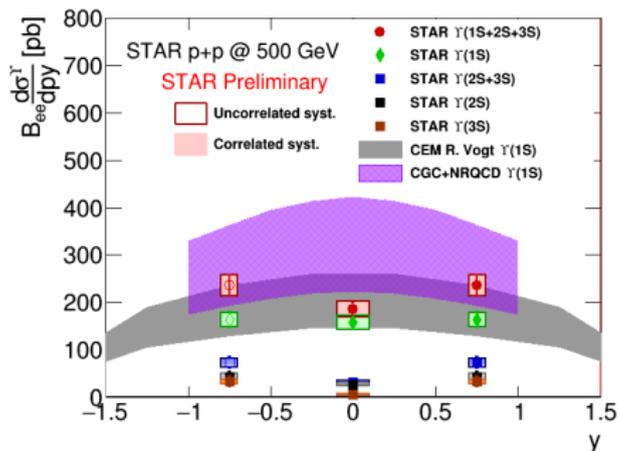
[L.Adamczyk, et al., J.Phys.Lett.B 735(2014)127]



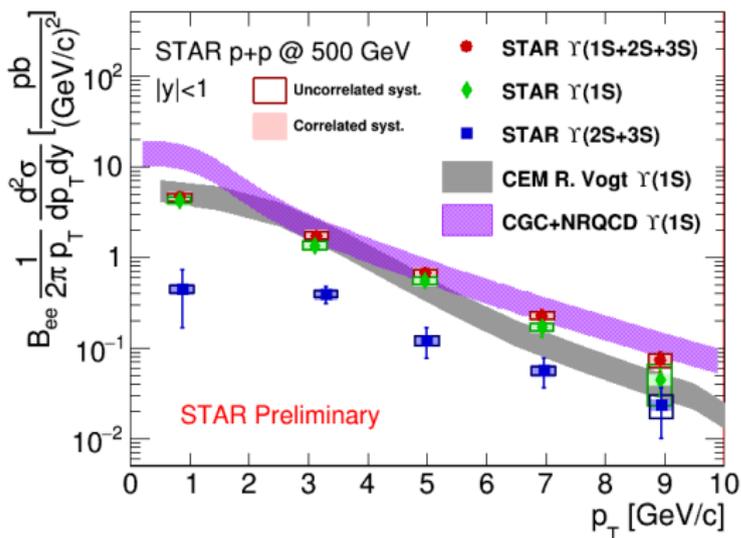
Signal extraction

- Challenging to extract individual $\Upsilon(nS)$ yields!
- Fit m_{ee} distributions with:
 - Signal lineshapes from STAR detector simulation
 - Backgrounds:
 - combinatorial
 - $b\bar{b} \rightarrow B\bar{B} \rightarrow e^+e^-$
 - Drell-Yan $q\bar{q} \rightarrow e^+e^-$





- Data reflected symmetrically around $y = 0$ ($y = \frac{1}{2} \ln \frac{E+p_{LC}}{E-p_{LC}}$)
- Separate $\Upsilon(1S)$ and $\Upsilon(2S)$ (NEW!), $\Upsilon(3S)$ (NEW!) spectra.
- Flatter rapidity spectrum compared to $\sqrt{s} = 200$ GeV (see backup: 27)
- Dip at mid-rapidity for $\Upsilon(2S+3S) \approx 2\sigma$ level from flat
- CEM model (inclusive) consistent with the measurement for $\Upsilon(1S)$
 [R. Vogt, *Phys.Rev.C* 92 034909(2015)]
- CGC+NRQCD predictions for direct $\Upsilon(1S)$ are above the data for inclusive $\Upsilon(1S)$
 [H. Han, et al., *Phys.Rev.D* 94, 014028(2016)]; [Y. Ma, R. Venugopalan, *Phys.Rev.Lett.* 113, 192301(2014)]



- Separate $\Upsilon(1S)$ and $\Upsilon(2S+3S)$ spectra.
- CEM calculation for inclusive $\Upsilon(1S)$

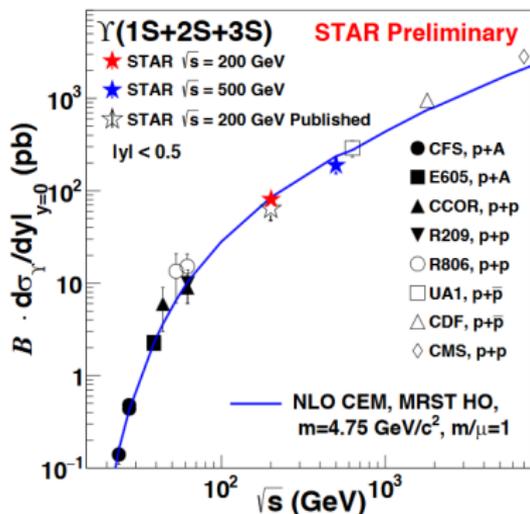
[R. Vogt, Phys.Rev.C 92 034909(2015)]

- Agree with data reasonably well

- CGC+NRQCD for direct Υ

[H. Han, Phys.Rev.D 94, 014028(2016)] [Y. Ma, Phys.Rev.Lett. 113, 192301(2014)]

- $\Upsilon(1S)$: model calculation is above the data points. Caveat: additional corrections (Sudakov resummation) are needed at low p_T according to authors.



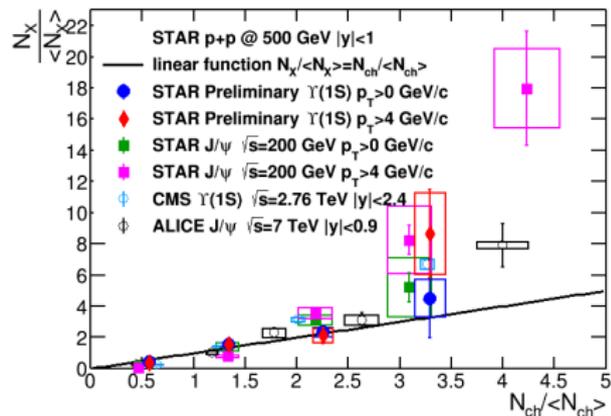
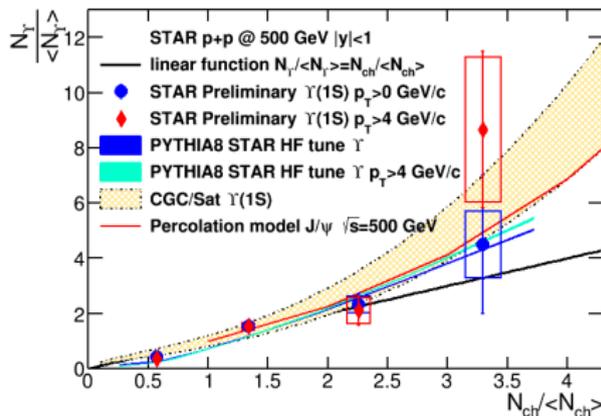
STAR [L.Adamczyk, Phys.Lett.B 735,127–137(2014)]
 CDF [D. Acosta et al., Phys.Rev.Lett. 88,161802(2002)]
 CMS [V. Khachatryan et al., Phys.Rev.D 83,112004(2010)]
 CFS [W. R. Innes et al., Phys.Rev.Lett. 39,1240-1242(1977)]
 CFS [J. K. Yoh et al., Phys.Rev.Lett. 41,684–687(1978)]
 CFS [K. Ueno et al., Phys.Rev.Lett. 42,486–489(1979)]
 CFS [S. Childress et al., Phys.Rev.Lett. 55,1962–1964(1985)]
 E605 [G. Moreno et al., Phys.Rev.D 43,2815–2835(1991)]
 E605 [T. Yoshida et al., Phys.Rev.D 39,3516(1989)]
 CCOR [A.L.S.Angelis et al., Phys.Lett.B 87,398–402(1979)]
 L. Camilleri, T.B.W. Kirk, H.D.I. Abarbanel (Eds.)
 E866 [L. Y. Zhu et al., Phys.Rev.Lett. 100,062301(2008)]
 ISR [C.Kourkoumelis et al., Phys.Lett.B 91,481-486(1980)]

Integrated cross section

- $B_{ee} \frac{d\sigma}{dy} |_{|y|<0.5} = 81 \pm 5(stat) \pm 8(syst)$ pb in p+p collisions at $\sqrt{s} = 200$ GeV
- $B_{ee} \frac{d\sigma}{dy} |_{|y|<0.5} = 186 \pm 14(stat) \pm 33(syst)$ pb in p+p collisions at $\sqrt{s} = 500$ GeV
- STAR results follow the world data trend
- Consistent with the Color Evaporation Model calculation

[A.D.Frawley et al., Phys.Rep. 462, pp.125–175(2008)]

$$\Upsilon \rightarrow e^+e^-$$



- Distributions of N_{ch} fully corrected using unfolding procedure. See backup: 26
- PYTHIA8 and String Percolation models reproduce the trend in the data
[E. G. Ferreira, C. Pajares, *Phys.Rev.C*, 86, 034903(2012)]
- CGC/Saturation model describes the data within large uncertainties
[E. Levin, M. Siddikov, *EPJC*, 97(5), 376(2019)], [M. Siddikov, et al., *arXiv:1910.13579 [hep-ph]*]
- Similar trends at RHIC and LHC for Υ and J/ψ
- Suggest Υ production in MPI or saturation effects
- Prospects: large data sets coming from runs 2017+2022 - increased precision

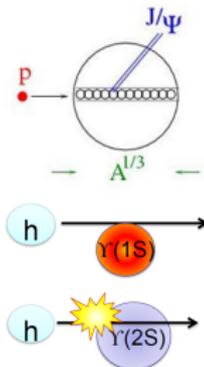


p+A collisions (Cold Nuclear Matter (CNM) effects - assume no QGP effects):

- Nuclear absorption: σ_{abs}
 - $\Upsilon + A \rightarrow X$
- Comover interactions - very small for $\Upsilon(1S)$

[Z. Lin, C.M. Ko, Phys.Lett.B 503, 104(2001)]

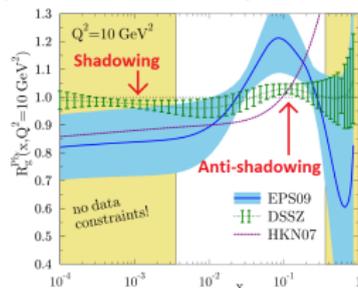
 - $\Upsilon + h \rightarrow X$



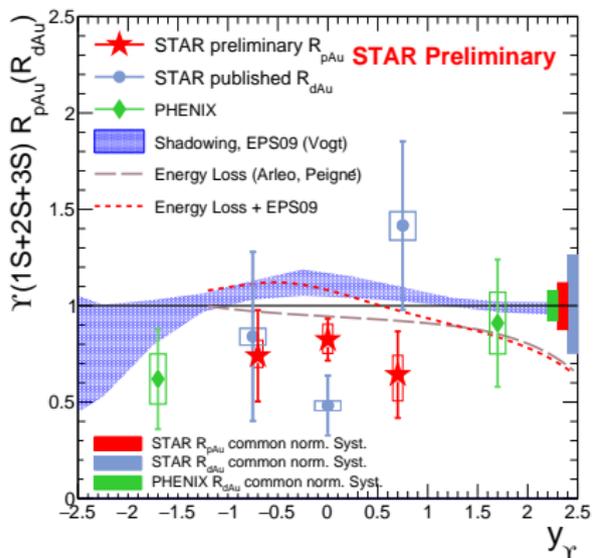
- Modified nuclear PDFs with respect to free nucleons:
 - shadowing
 - anti-shadowing
- Studied by measuring Nuclear Modification

$$\text{Factor: } R_{pA} = \frac{\sigma_{inel}}{\langle N_{coll} \rangle} \frac{d^2 N_{p+Au}/dp_T dy}{d^2 \sigma_{pp}/dp_T dy}$$

[L. Grandchamp, LBNL 2005]



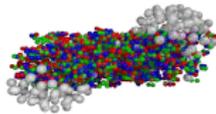
[H. Paukkunen, Nucl.Phys.A 926 24-33(2014)]



[L. Adamczyk et al., J.Phys.Lett.B 735(2014)127],
 [A. Adare et al., Phys. Rev. C 87, 044909],
 [F. Arleo, S. Peigné, JHEP 03, 122(2013)]

$\Upsilon(1S+2S+3S)$

- Improved precision over published results of R_{dAu}
 - $\sim 50\%$ smaller statistical uncertainty vs. y
 - $R_{pAu}|_{|y|<0.5} = 0.82 \pm 0.10(stat.)_{-0.07}^{+0.08}(syst.) \pm 0.10(glob.)$
- Indication of $\Upsilon(1S+2S+3S)$ suppression in p+Au collisions

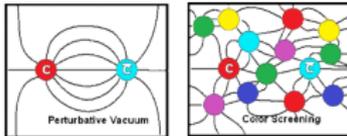


Heavy ion (A+A) collisions (QGP+CNM effects):

- Quarkonium states dissociate at high temperature in QGP via Debye-like screening

[T.Matsui, H.Satz, Phys.Lett.B 178(4),416-422(1986)]

- Provides information about QGP (interactions, temperature...)



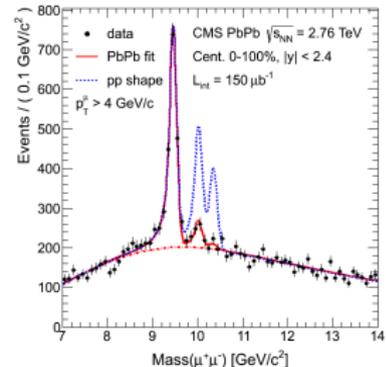
- Sequential suppression due to each $\Upsilon(nS)$ state dissociating at different $T \rightarrow$ estimate of

T [S. Digal, Phys.Rev.D 64, 094015(2001)]

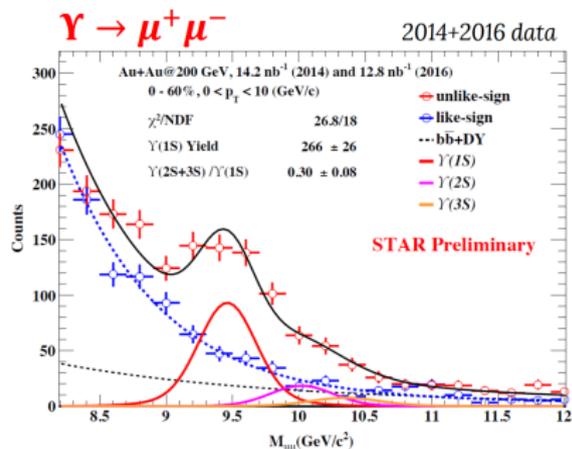
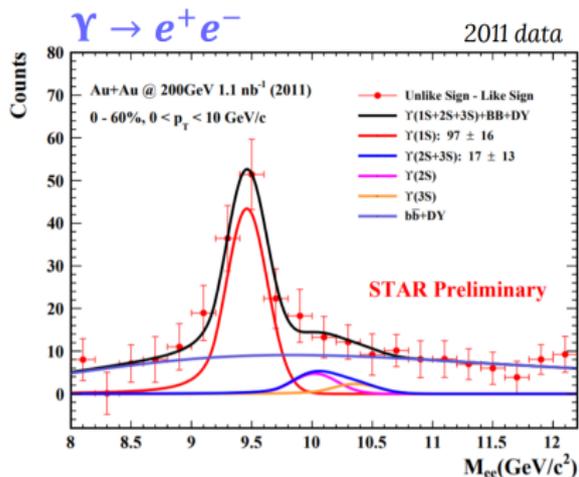
- Modified feed-down pattern



[A. Rothkopf, Hard Probes 2012]

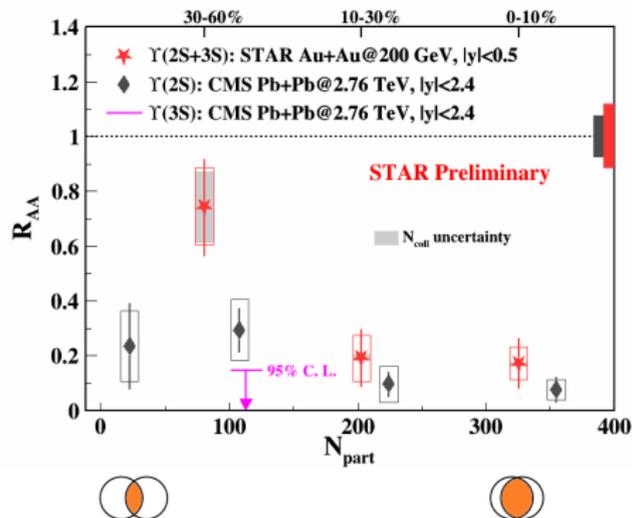
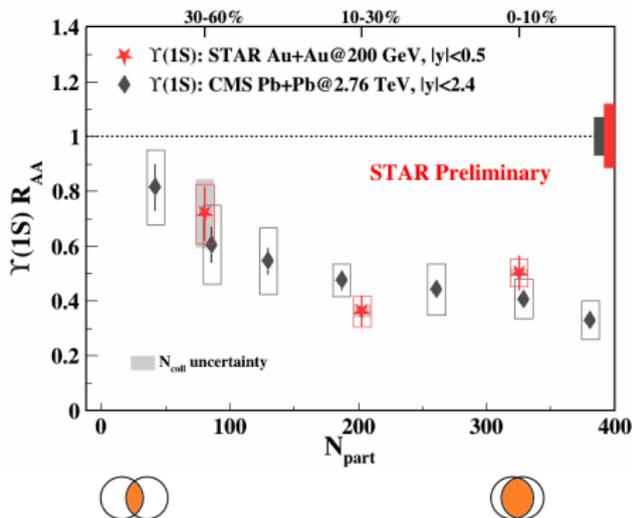


[S. Chatrchyan, et. al, Phys.Rev.Lett. 109(22), 222301]



Υ signal

- Υ measured in both e^+e^- and $\mu^+\mu^-$
- Combined R_{AA} for better precision

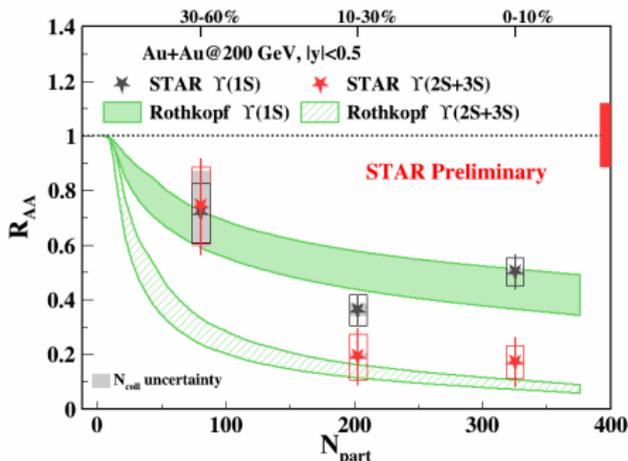


CMS: [V. Khachatryan et al., Phys.Lett.B 770, 357-379(2017)]

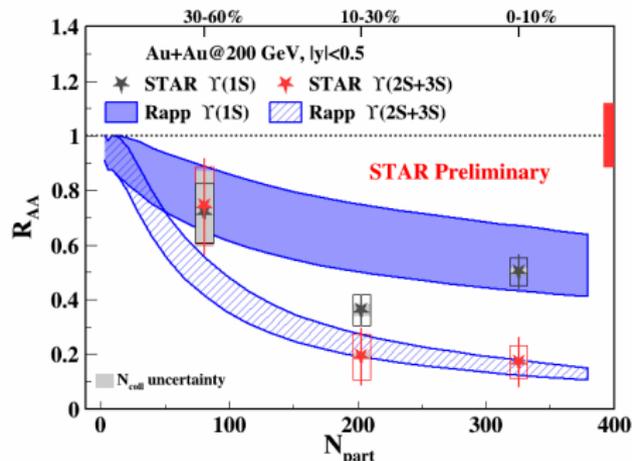
STAR vs. CMS

- Similar suppression for $\Upsilon(1S)$, despite higher medium temperature at the LHC
 - $R_{AA}(1S) \approx 0.4$ could be largely due to suppression of excited states contribution
 - Sequential suppression
 - Regeneration? Larger at LHC than at RHIC
 - CNM effects - need better constraints
- Indication of smaller suppression for $\Upsilon(2S+3S)$ at RHIC than at LHC

Υ : STAR vs. models



Rothkopf: [B. Krouppa et al., Phys.Rev.D 97, 016017(2018)]



Rapp: [X. Du et al., Phys.Rev.C 96,054901(2017)]

Models

- Krouppa, **Rothkopf**, Strickland
 - QGP modeled by hydrodynamics + Debye-like screening
 - No regeneration, no CNM effects
- De, He, **Rapp**
 - QGP modeled by hydrodynamics + Debye-like screening
 - Includes both regeneration and CNM effects
- Both models agree with STAR $\Upsilon(1S)$ data
- Indication that Rothkopf's model underestimates the STAR $\Upsilon(2S+3S)$ results for 30 – 60% centrality

Υ in p+p collisions at $\sqrt{s} = 200$ GeV and $\sqrt{s} = 500$ GeV

- $\Upsilon(1S)$ data reasonably described by CEM model, while overestimated by CGC+NRQCD
- Similar trends for J/ψ and Υ vs. $N_{ch}/\langle N_{ch} \rangle$ at RHIC and LHC

Υ in p+Au collisions at $\sqrt{s_{NN}} = 200$ GeV

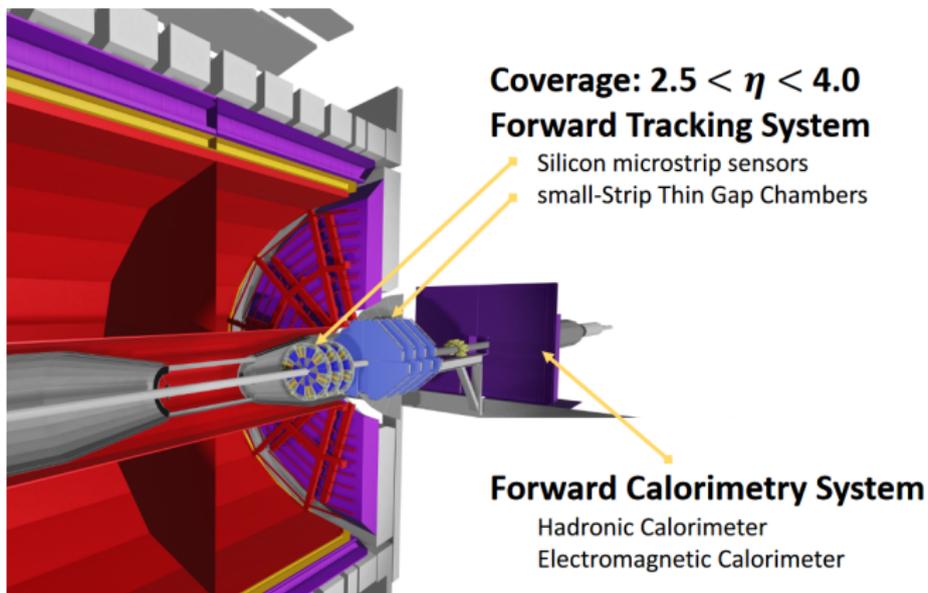
- Indication of $\Upsilon(1S+2S+3S)$ suppression

Υ in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV

- ΥR_{AA} measured in dielectron and dimuon channels - combined for better precision
- Similar suppression of $\Upsilon(1S)$ at RHIC and LHC
- Stronger suppression of $\Upsilon(2S+3S)$ than $\Upsilon(1S)$ in central collisions
 - Sequential suppression
- $\Upsilon(1S)$, $\Upsilon(2S+3S)$ R_{AA} consistent with model calculations

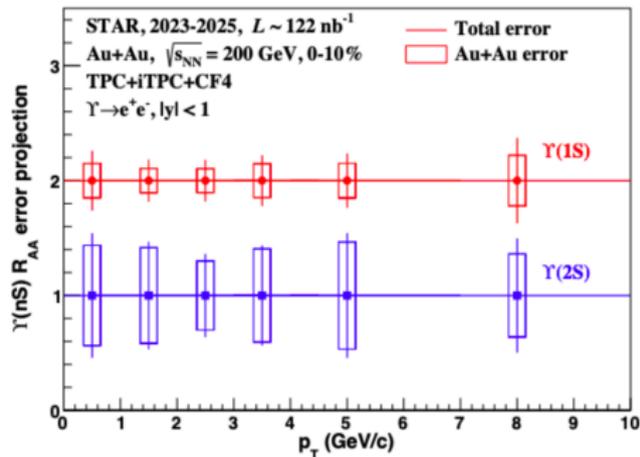
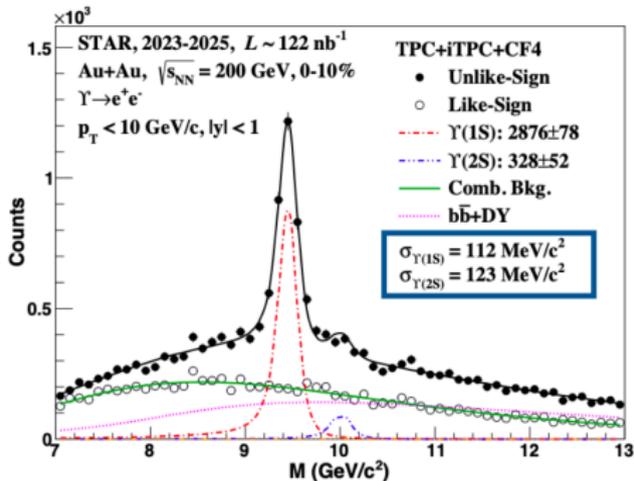
Thank you for your attention!

BACKUP



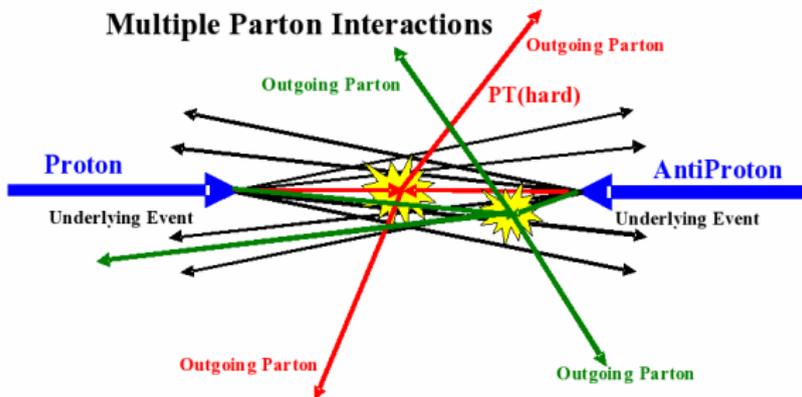
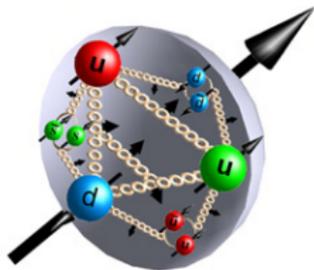
Future plans for STAR

- iTPC already running - improved momentum resolution
- Forward upgrade - new detectors
- High integrated luminosity for precision quarkonium production studies
- And more!



Projections 2023+

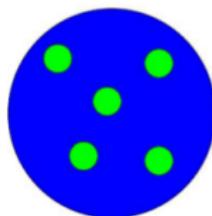
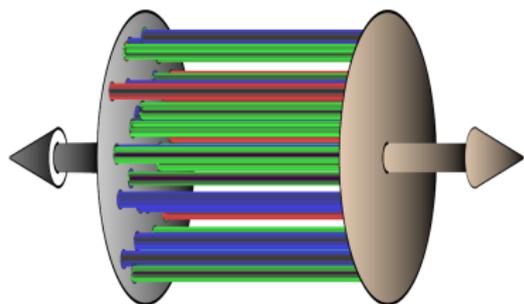
- High precision measurement:
 - High integrated luminosity $\mathcal{L}_{int} \sim 122 \text{ nb}^{-1}$
 - Improved momentum resolution
 - Low material budget - less background
- R_{AA} of $\Upsilon(1S)$ and $\Upsilon(2S)$ vs:
 - centrality, p_T



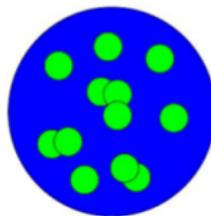
<https://www.bnl.gov/rhic/images/proton-with-gluons-300px.jpg>

<http://www.desy.de/~jung/multiple-interactions/may06/mi-rick.gif>

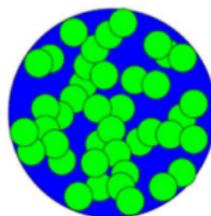
- Protons are complex objects consisting of constituent quarks, sea quarks and gluons.
- Multiple parton interactions (MPI) may happen in $p + p$ collision - implemented in PYTHIA.
 - Besides the main hard process, there may be additional hard and soft processes in MPI.
- As implemented in PYTHIA8, heavy quarks can also be produced during MPI.
- MPI together with initial- (ISR), final-state radiation (FSR) and beam remnants define the event activity, which can be characterized experimentally using the charged particle multiplicity.



Isolated Disks



Clusters



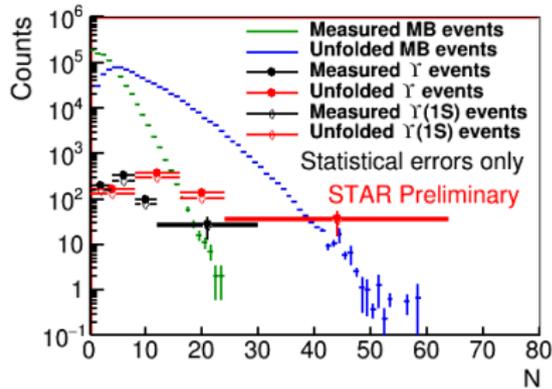
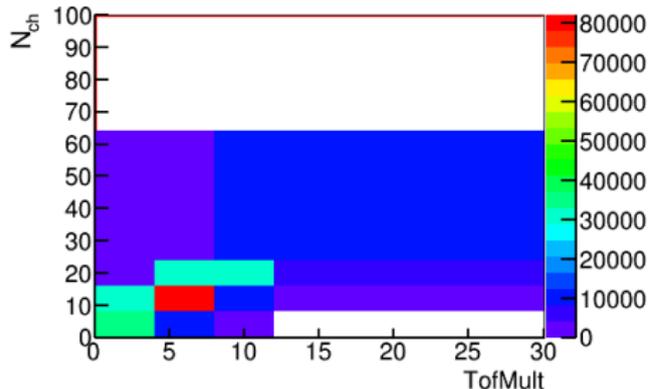
Percolation

[F. Gelis et al., Ann.Rev.Nucl.Part.Sci.60, 463-489(2010)] [L. Kozarzewski, Proc.of SPIE, 100313U(2016)]

- Models particle production originating from strings of color field formed in $p + p$ collisions.
- Soft particle production dampened by interaction of overlapping strings.
- Predicts quadratic dependence of normalized yield for particles from hard processes vs. normalized charged particle multiplicity in high multiplicity events.

$$\frac{N_{hard}}{\langle N_{hard} \rangle} = \langle \rho \rangle \left(\frac{\frac{dN_{ch}}{d\eta}}{\langle \frac{dN_{ch}}{d\eta} \rangle} \right)^2 \quad [E. G. Ferreira, C. Pajares, Phys.Rev. C, 86, 034903 (2012)]$$

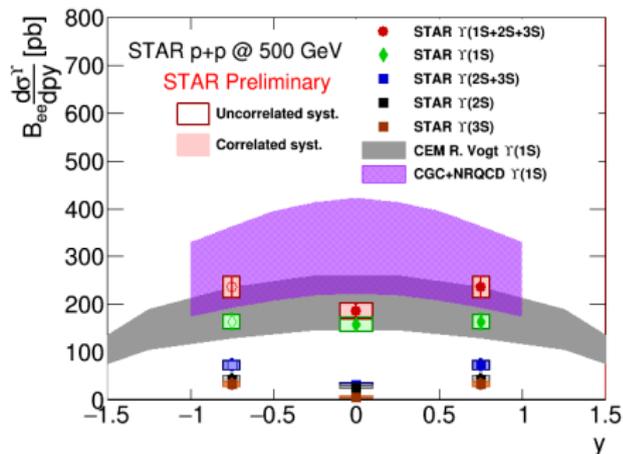
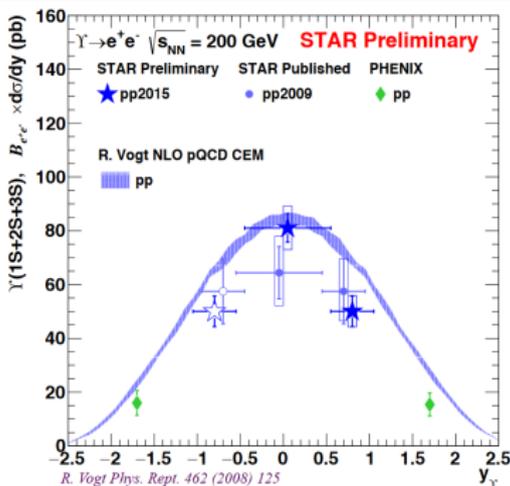
Response matrix for Υ events



Unfolding method used for multiplicity dependent studies

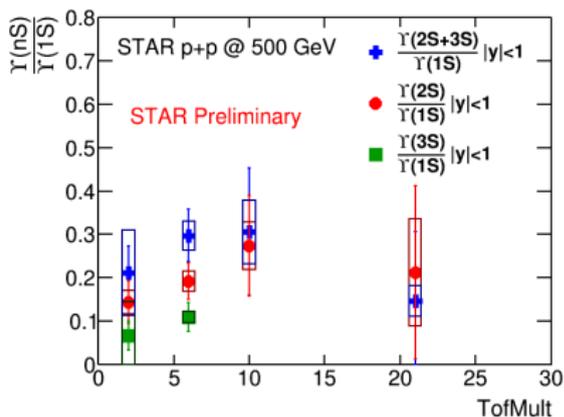
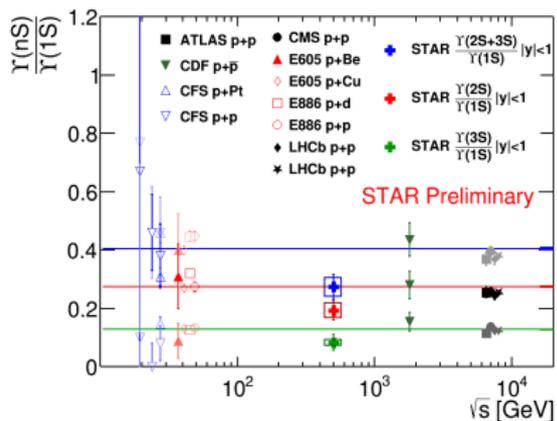
- ① A response matrix is obtained using the PYTHIA8 event generator for both min-bias and Υ events taking into account reconstruction efficiency
- ② The measured distributions are unfolded using their respective response matrices
- ③ This procedure yields the unfolded (true) distribution
- ④ Similar procedure used for J/ψ
- ⑤ Measured N_{ch} distribution obtained from p+p $\sqrt{s} = 500$ GeV 2009 data
- ⑥ Measured distribution of Υ events obtained from p+p $\sqrt{s} = 500$ GeV 2011 data

Υ rapidity dependence in p+p at 200 and 500 GeV



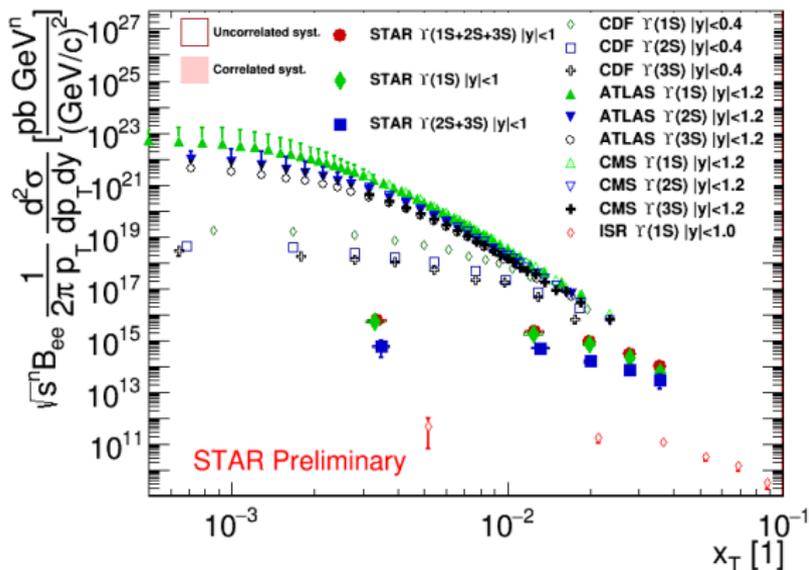
- $\sqrt{s} = 200$ GeV STAR data:
 - Slightly narrower than Color Evaporation Model (CEM)
- $\sqrt{s} = 500$ GeV data:
 - Data reflected symmetrically around $y = 0$ ¹
 - Separate $\Upsilon(1S)$ and $\Upsilon(2S)$ (NEW!), $\Upsilon(3S)$ (NEW!) spectra.
 - Flatter rapidity spectrum compared to $\sqrt{s} = 200$ GeV (see backup)
 - Dip at mid-rapidity for $\Upsilon(2S+3S) \approx 2\sigma$ level from flat
 - CEM model (inclusive) consistent with the measurement for $\Upsilon(1S)$
 [R. Vogt, Phys.Rev.C 92 034909(2015)]
 - CGC+NRQCD predictions for direct $\Upsilon(1S)$ are above the data for inclusive $\Upsilon(1S)$
 [H. Han et al., Phys.Rev.D 94, 014028(2016)], [Y. Ma, R. Venugopalan, Phys.Rev.Lett. 113, 192301(2014)]

p+p $\sqrt{s} = 500$ GeV 2011 dataset
 $\Upsilon \rightarrow e^+e^-$



[W. Zha, et al, Phys.Rev.C 88,067901(2013)]

- Left plot: cross section ratios measured in 500 GeV p+p collisions are slightly below (within 2σ) world data average, shown as solid lines in the left plot.
- Right plot: Ratios vs. TofMult - no strong multiplicity dependence observed.
- TofMult: number of tracks matched to TOF within $|\eta| < 1$, $p_T > 0.2$ GeV/c (uncorrected)



STAR $p + p \sqrt{s} = 500$ GeV
 ATLAS $p + p \sqrt{s} = 7$ TeV
 [G. Aad et al., Phys.Rev.D
 87,052004(2013)]
 CMS $p + p \sqrt{s} = 7$ TeV
 [V.Khachatryan et al., Phys.Lett.B
 749,14-34(2015)]
 CDF $p + \bar{p} \sqrt{s} = 1.8$ TeV
 [D. Acosta et al., Phys.Rev.Lett.
 88,161802(2002)]
 ISR $p + \bar{p} \sqrt{s} = 53, 63$ GeV
 [C.Kourkoumelis et al., Phys.Lett.B
 91,481-486(1980)]

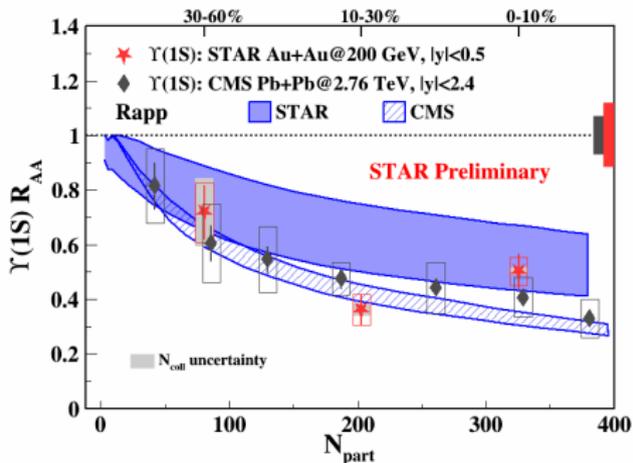
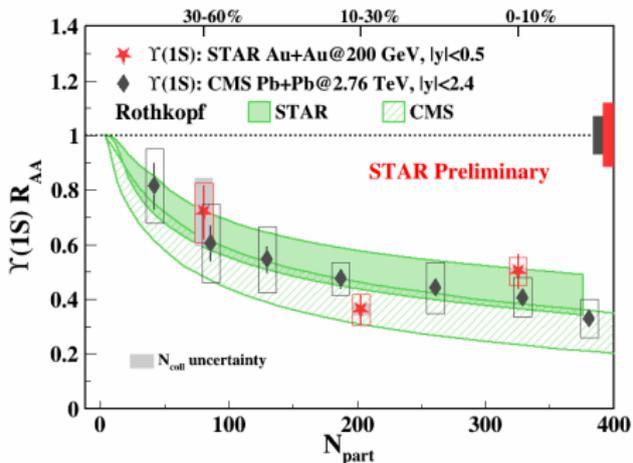
- $$x_T = \frac{2p_T}{\sqrt{s}}, \sigma^{inv} \equiv E \frac{d^3\sigma}{d^3p} = \frac{F(x_T)}{p_T^{n(x_T, \sqrt{s})}} = \frac{F'(x_T)}{\sqrt{s}^{n(x_T, \sqrt{s})}}$$

[F. Arleo et al., JHEP06,035(2010)]

- pQCD predicts that spectra of hard processes should follow x_T scaling - check with $n = 5.6$ (number of partons taking active part in the process) obtained for J/ψ

[L. Adamczyk et al., Phys.Rev.C 80, 041902(2009)]

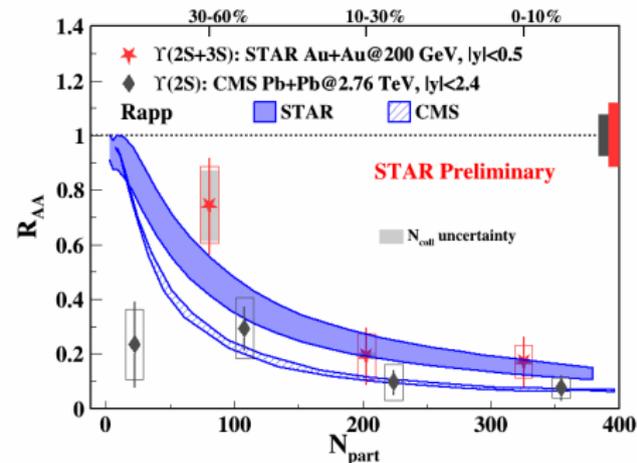
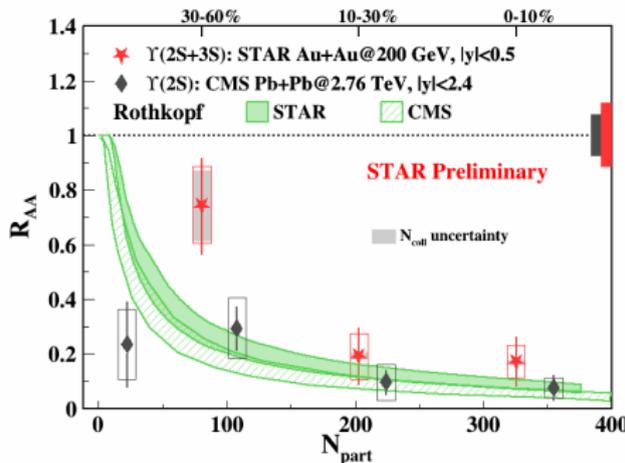
- No clear scaling observed, some indication for LHC data at high p_T



[B. KrouppaPhys et al., .Rev.D 97,(2018)016017], [X. Du et al., Phys.Rev.C 96,(2017)054901]

$\Upsilon(1S)$ vs. models

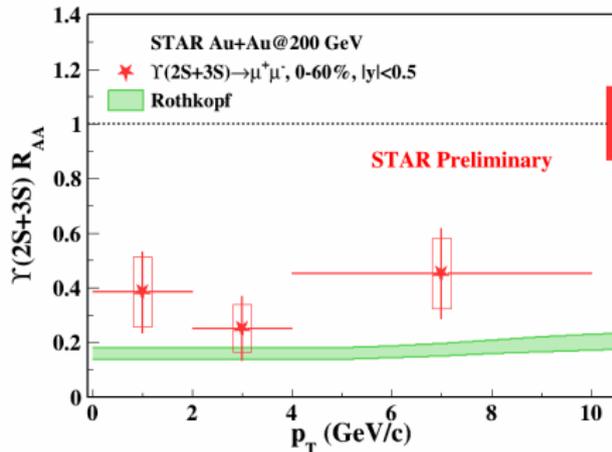
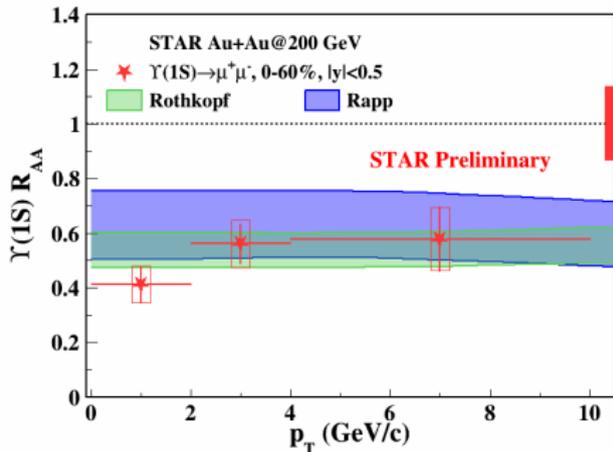
- Both models consistent with the data



[B. Krouppa et al., Phys.Rev.D 97,016017(2018)], [X. Du et al., Phys.Rev.C 96,054901(2017)]

$\Upsilon(2S+3S)$ vs. models

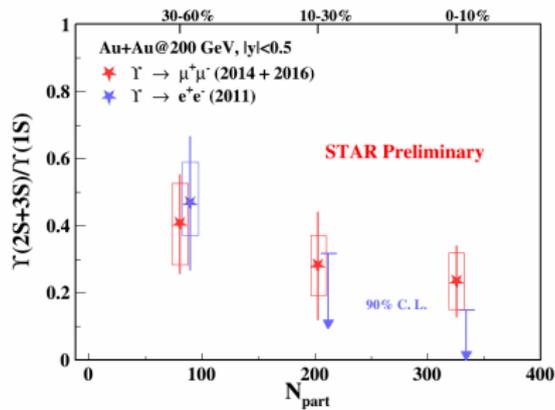
- Both models consistent with the data in central and semi-central collisions



[B. Krouppa et al., Phys.Rev.D 97,016017(2018)], [X. Du et al., Phys.Rev.C 96,054901(2017)]

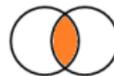
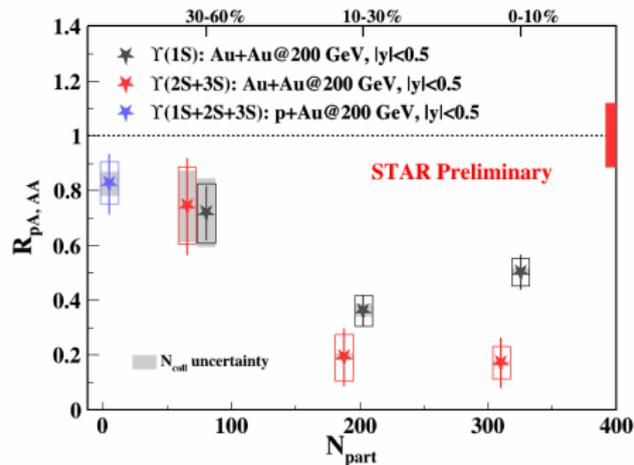
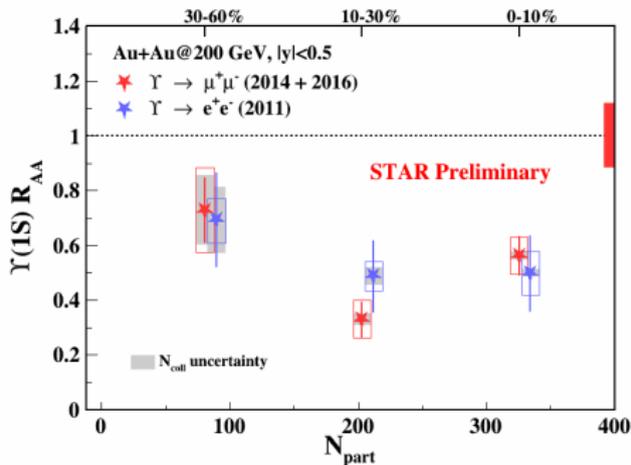
R_{AA} vs. p_T vs. models

- Both models consistent with the data
- Rothkopf's model slightly lower than $Y(2S+3S)$
- Flat vs. p_T



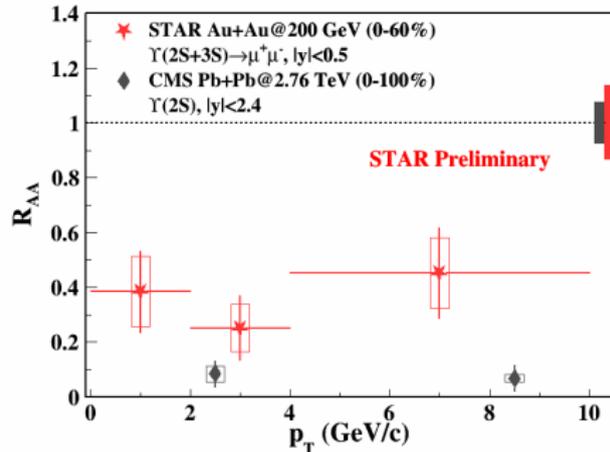
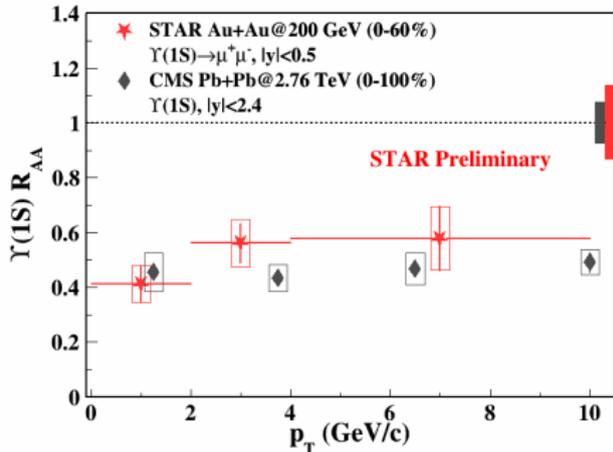
$\frac{\gamma(2S+3S)}{\gamma(1S)}$ vs. N_{part}

- Both channels consistent



R_{AuAu} measured by STAR

- Consistent results from dielectron and dimuon channels
- Both results combined in order to achieve better precision
- Similar level of suppression in peripheral collisions as in $p + Au$
- Stronger suppression of $\Upsilon(2S+3S)$ than $\Upsilon(1S)$ in central collisions



CMS: [V.Khachatryan et al., Phys.Lett.B 770, 357-379(2017)]

Transverse momentum dependence

- Similar suppression for $\Upsilon(1S)$ at RHIC and LHC
- Indication of stronger suppression of high- p_T $\Upsilon(2S+3S)$ at LHC than at RHIC
- Both consistent with flat dependence vs. p_T