



# $\Upsilon$ measurements in p+p collisions at $\sqrt{s}=500$ GeV with the STAR experiment



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ZIMÁNYI SCHOOL 2014

December 1-5, 2014

Wigner Research Center for Physics, Budapest, Hungary



EUROPEAN UNION  
EUROPEAN  
SOCIAL FUND



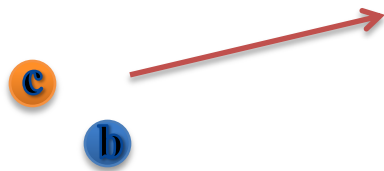
# Outline

- Why Upsilon at RHIC?
- Experimental approach in STAR
- Measured yields and expected precision
  - Efficiency studies
- Prospects for Upsilon measurements in STAR

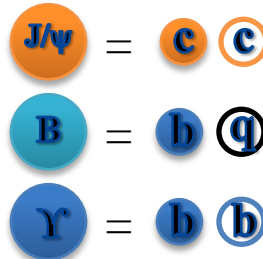
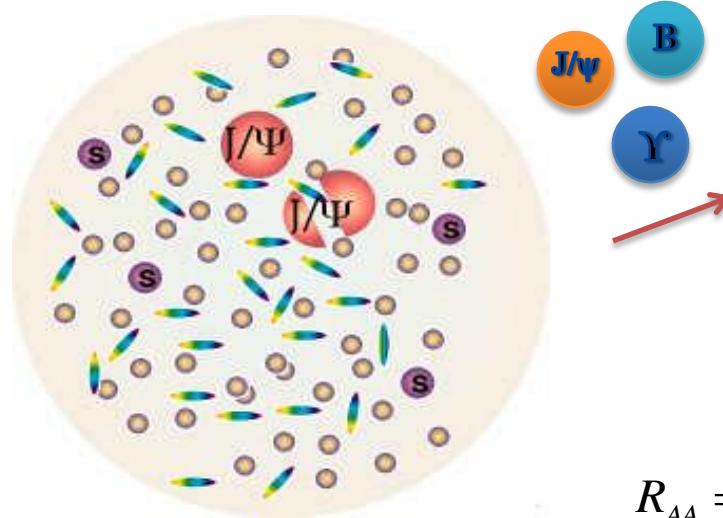


# Physics introduction

Heavy quarks (b,c) are produced in an early stage of the collision (before quark-gluon plasma (QGP) formation)



Heavy quarks are „probes”, which interact with QGP



Hadrons containing heavy quarks can be observed and their modification in QGP measured.

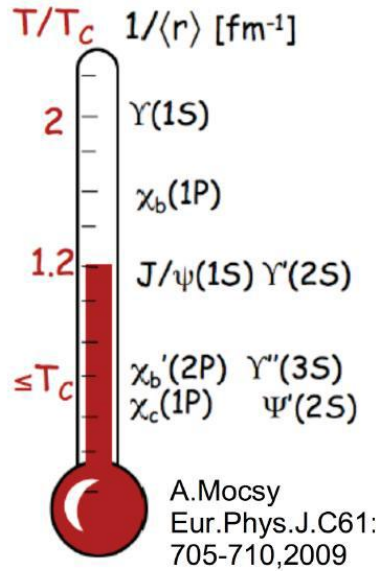
$$R_{AA} = \frac{1}{N_{coll}} \frac{(\text{invariant yield})_{AA}}{(\text{invariant yield})_{pp}}$$

From this we can learn about QGP and interactions between quarks and gluons.

•Suppression of  $J/\psi$  in A+A is a signature of QGP formation.

T.Matsui and H.Satz 1986 Phys. Lett. B 178:416, 1986

# Quarkonium suppression

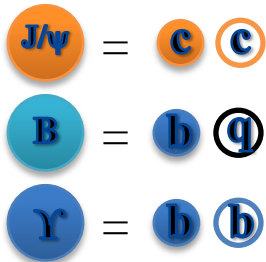


- Upsilon states provide a cleaner probe than  $J/\psi$ , because of smaller regeneration and cold nuclear matter effects
- Sequential suppression of more tightly bound states with increasing temperature – indirect QGP temperature measurement

Á. Mócsy, P. Petreczky, Phys. Rev. D77, 014501 (2008)



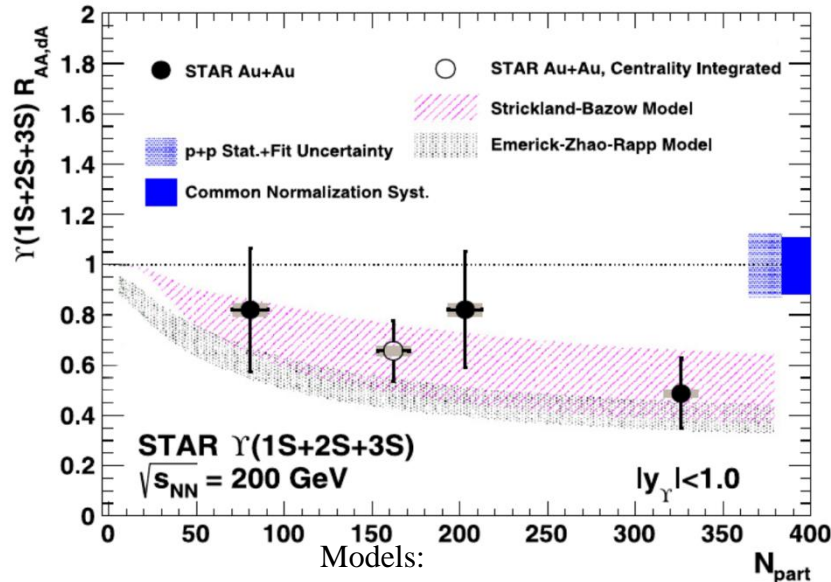
Requires measurements of suppression of different quarkonium states



# Why Upsilon at RHIC?

STAR measured  $R_{AA}$  for Upsilon 1S+2S+3S and 1S as a function of  $N_{part}$ .

• Shows significant suppression for integrated  $p_T$ .

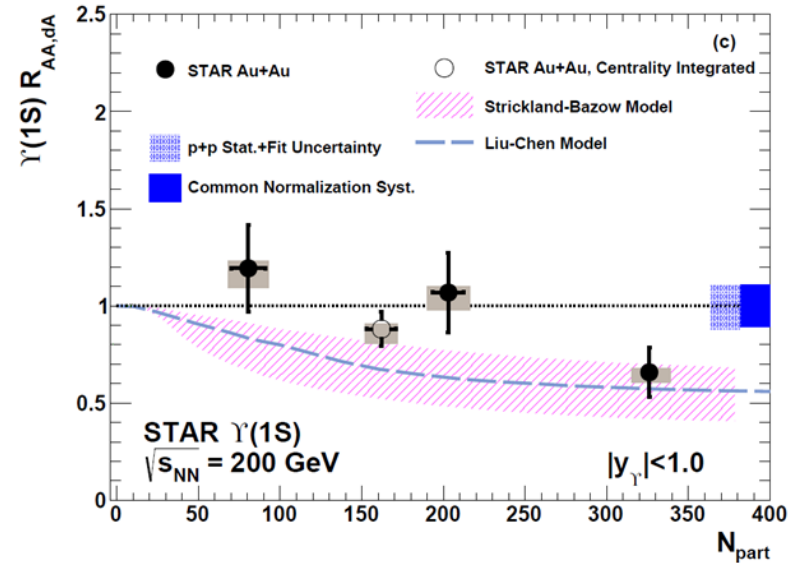


Models:

Hot and cold nuclear matter effects: Strickland, Bazov, Nucl. Phys. A879, 25 (2012)

Hot nuclear matter effects only: arXiv:1009.2585

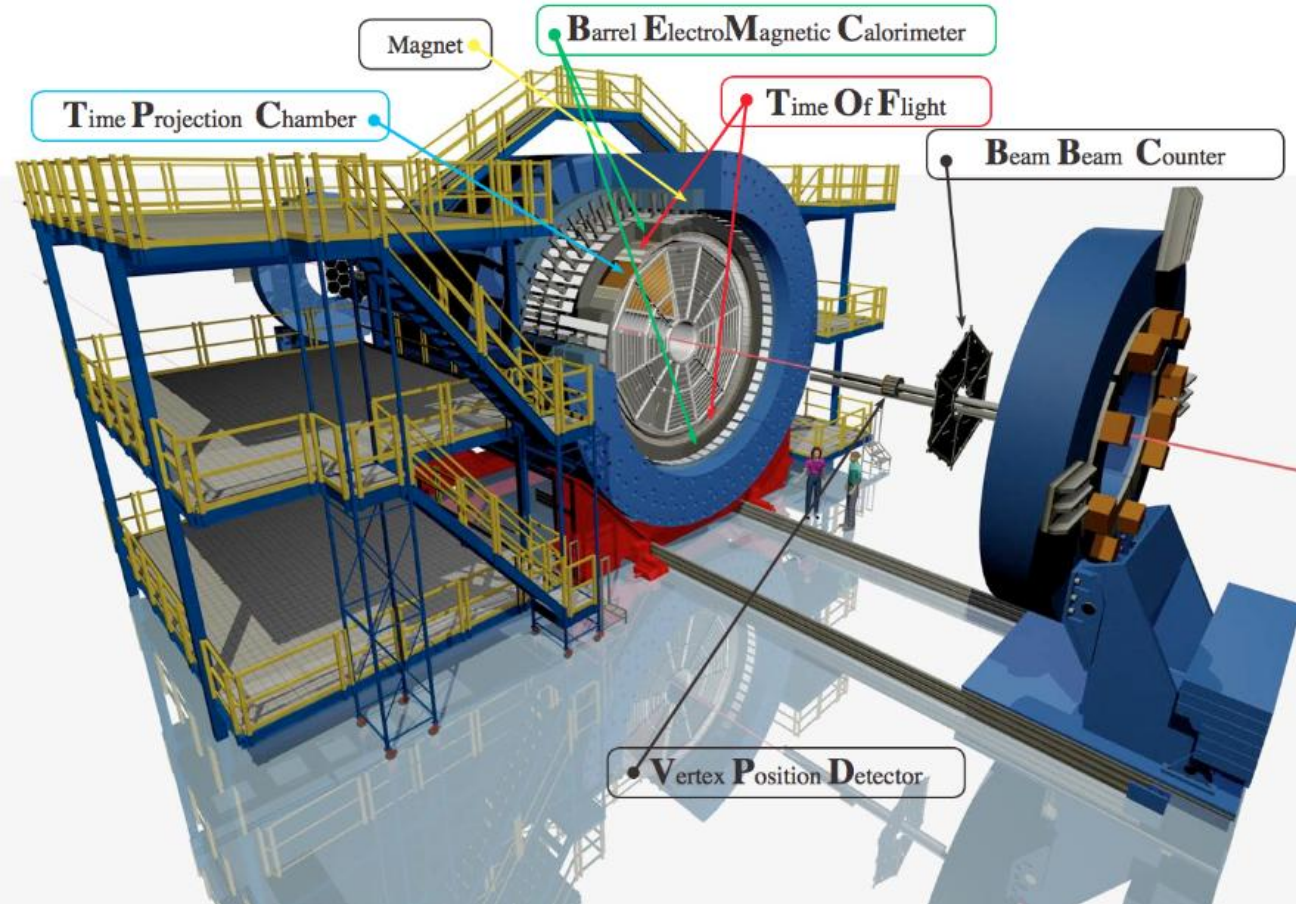
Phys.Lett. B735 (2014) 127



- We need  $p_T$  spectrum from p+p as a baseline for  $R_{AA}(p_T)$  to study dynamical aspects of Upsilon interaction with QGP.
- STAR p+p 500 GeV provides much better statistical precision (larger  $b\bar{b}$  cross section than at 200 GeV, and higher integrated luminosity)
- Will provide an opportunity to measure ratios of different states in production yields.
- May provide constraints for quarkonium production models

# STAR experiment

Solenoidal Tracker At RHIC :  $-1 < \eta < 1, 0 < \phi < 2\pi$



- TPC: Particle identification via  $dE/dx$  and tracking

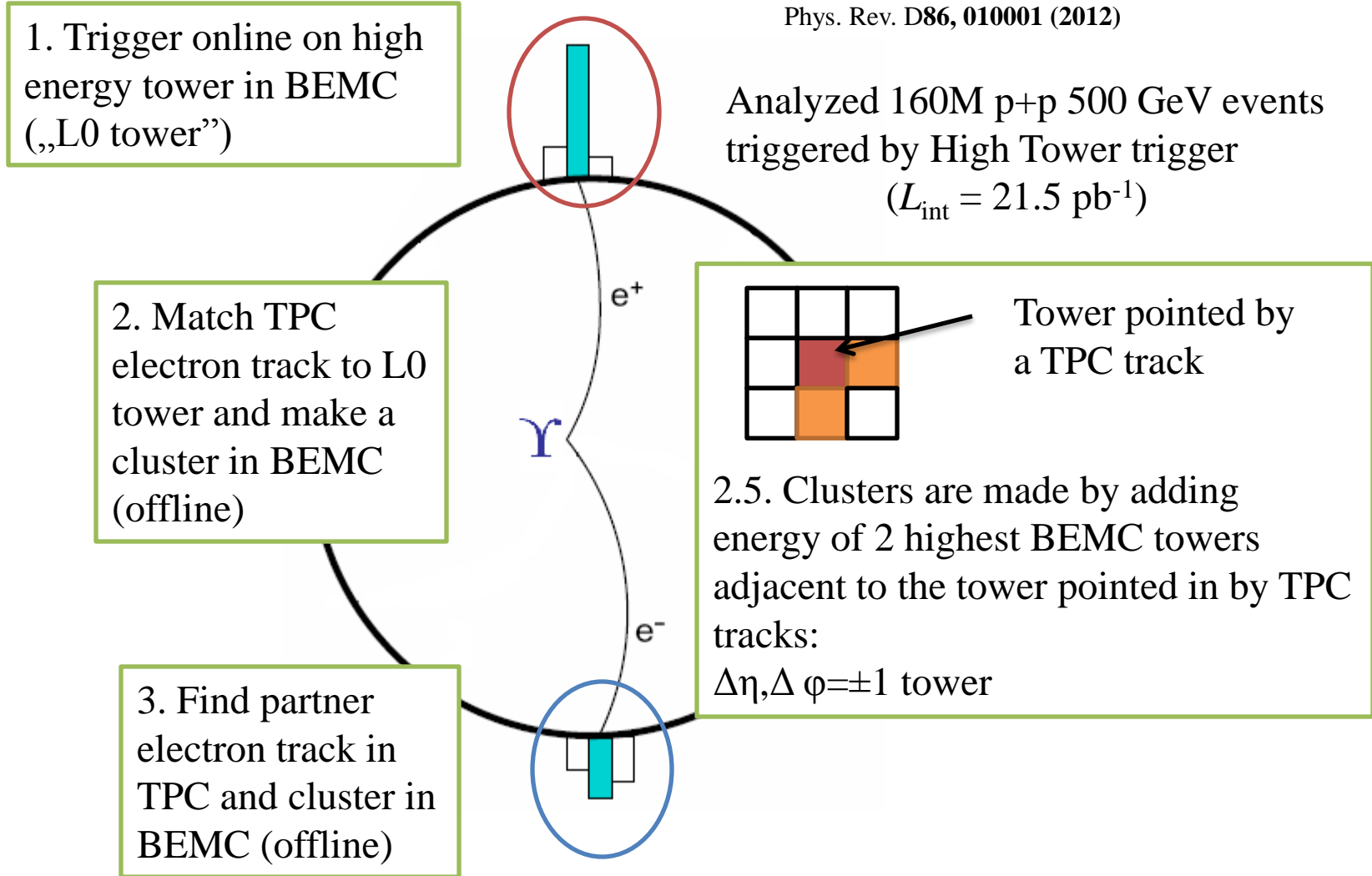
- BEMC: Particle identification via  $E/p$  and shower shape, also provides online High Tower trigger

# Experimental approach in STAR

$$Y \rightarrow e^+e^- \text{ BR} = 2.38 \pm 0.11\%$$

Phys. Rev. D86, 010001 (2012)

Analyzed 160M p+p 500 GeV events  
triggered by High Tower trigger  
( $L_{\text{int}} = 21.5 \text{ pb}^{-1}$ )



# Offline analysis and particle identification

Track which fired the trigger:

Match to L0 tower required

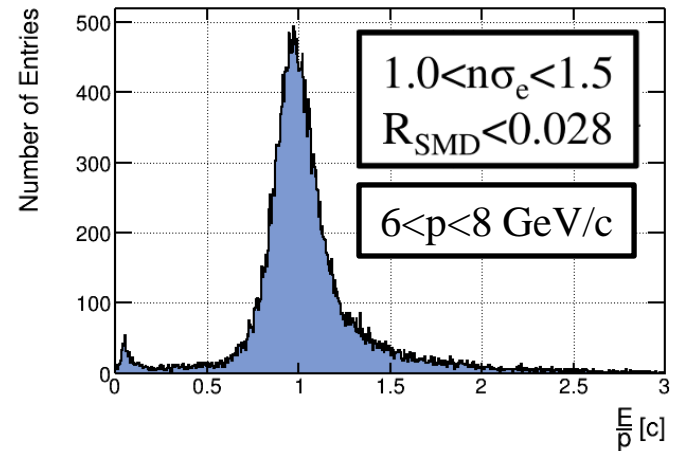
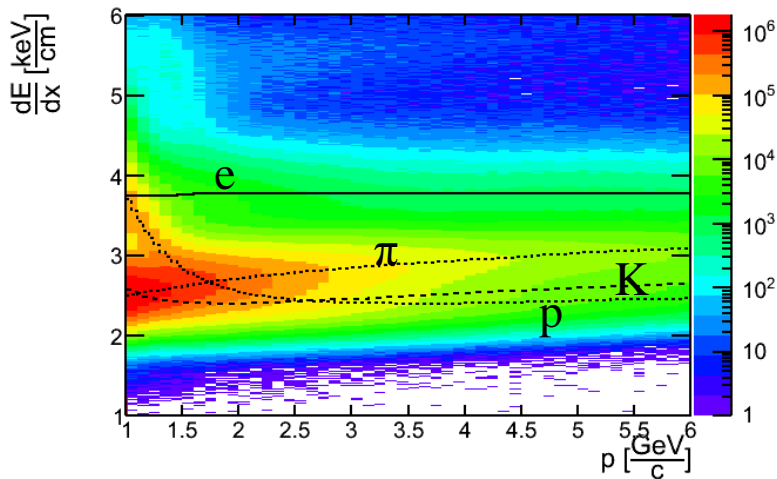
Partner track:

$p > 1 \text{ GeV}/c$

Both tracks:

$-1.2 < n\sigma_e < 3.0$   
 $E_{\text{tower}}/E_{\text{cluster}} > 0.5$   
 $0.55 < E_{\text{cluster}}/pc < 1.45$   
 $R_{\text{SMD}} < 0.028$

dEdx vs p all



Normalized  $dE/dx$ :  $n\sigma_e = \log\left(\frac{dE/dx_e}{dE/dx|_{\text{Bichsel}}}\right) / \sigma$

Distance between a track projection on BEMC SMD and center of a cluster:  $R_{\text{SMD}} = \sqrt{\Delta\eta^2 + \Delta\phi^2}$



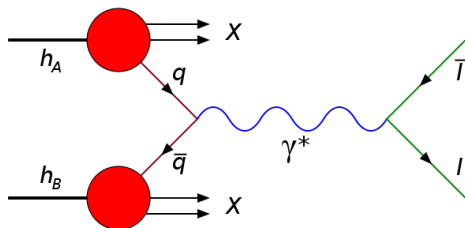
# Upsilon signal and backgrounds

Invariant mass histogram of  $e^+e^-$  contains:

- Upsilon signal
  - Combinatorial background
- described by like sign sum:

$$N_{e^+e^+} + N_{e^-e^-}$$

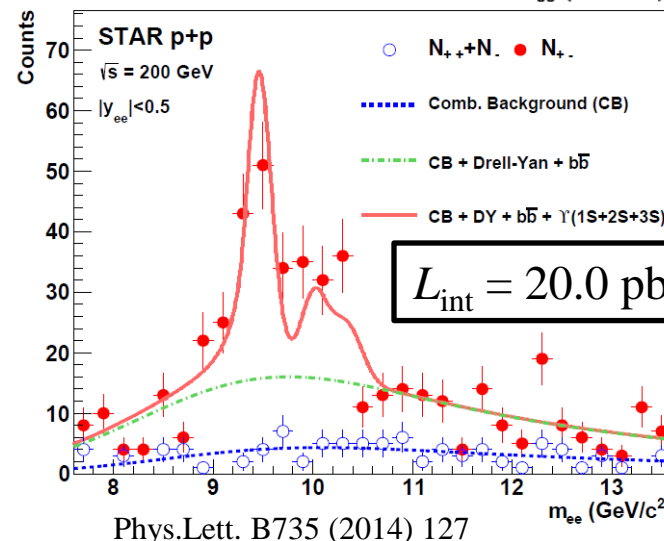
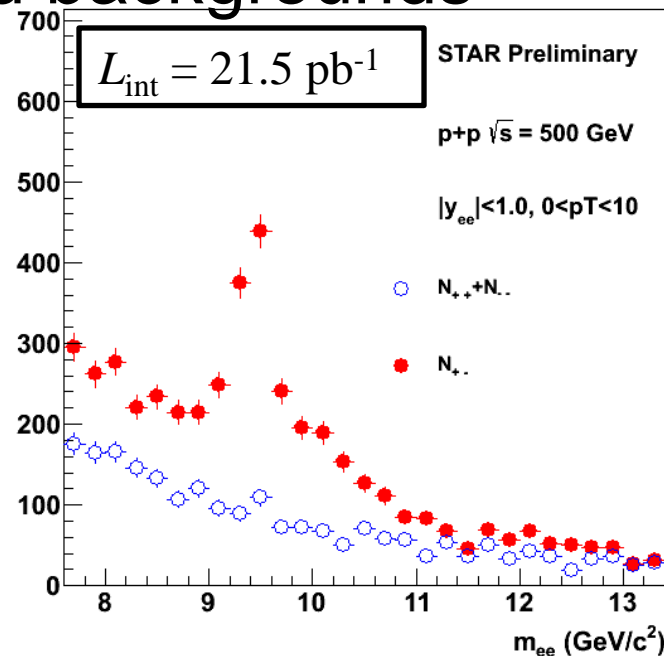
- Drell-Yan background:  $q\bar{q} \rightarrow l\bar{l}$



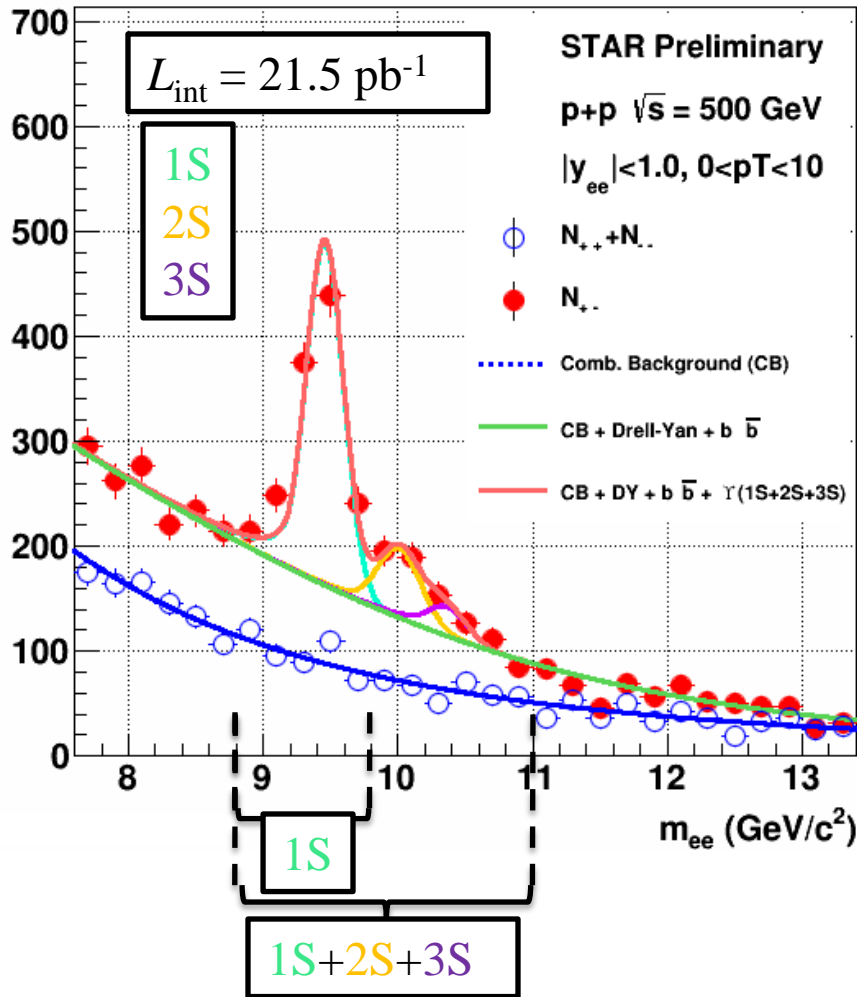
- $b\bar{b}$  background:  $b\bar{b} \rightarrow e^+e^-$

Signal extraction strategy:

- describe background by a function and do a simultaneous fit to signal and combinatorial background
- integrate counts in an Upsilon mass range and subtract background



# Upsilon signal



Signal calculated by integrating counts and subtracting backgrounds from fits:

$$N_Y = N_{e+e-} - N_{CB}^{Fit} - N_{DY}^{Fit} - N_{b\bar{b}}^{Fit}$$

Background fit functions

$$f_{b\bar{b}} = N \frac{m^A}{\left(1 + \frac{m}{B}\right)^C} \quad f_{DY} = N \frac{1}{\left(1 + \frac{m}{A}\right)^B}$$

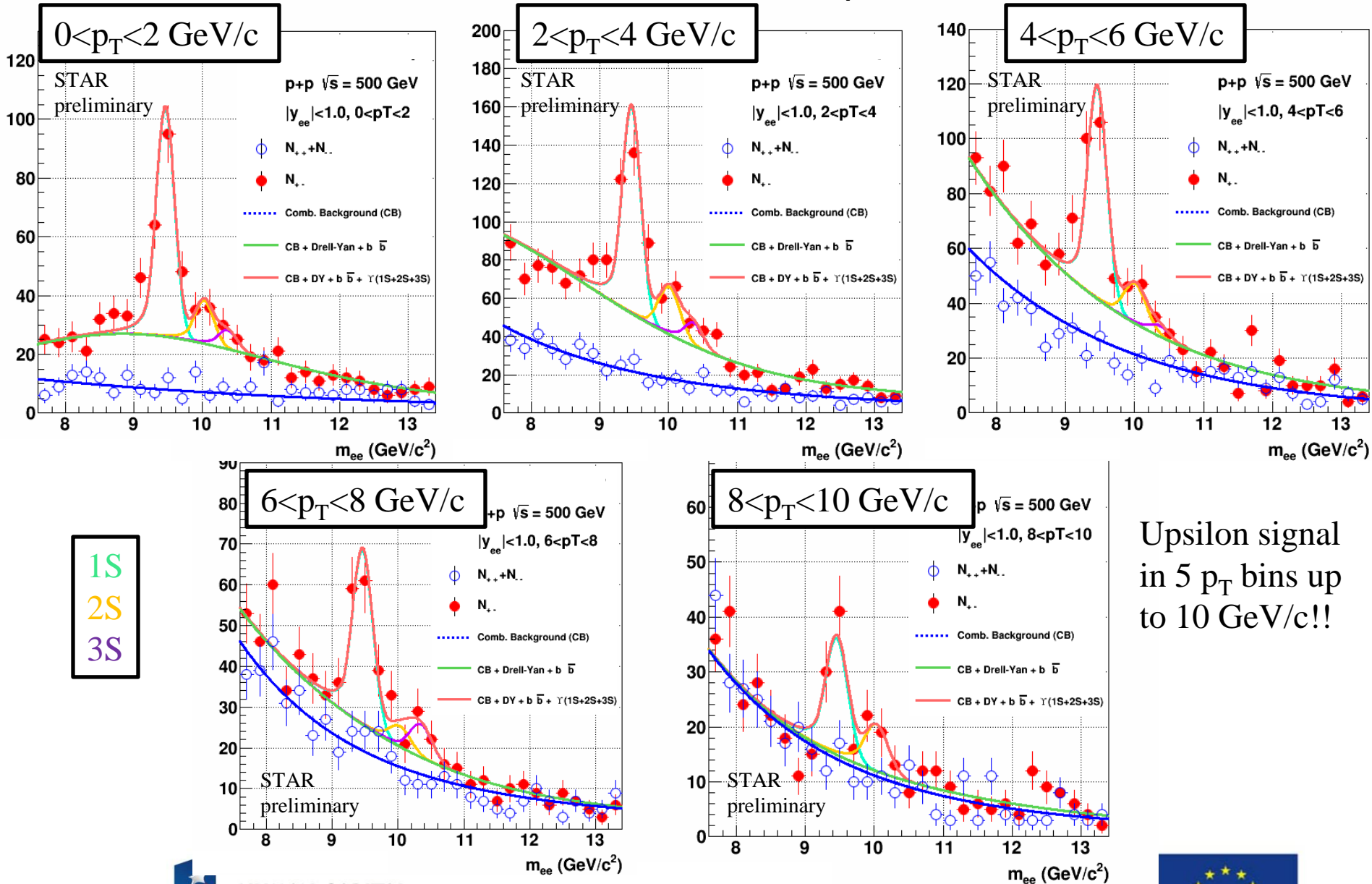
$$f_{CB} = N \cdot A \cdot \exp\left(\frac{-m}{B}\right) + N \cdot (1-A) \cdot \exp\left(\frac{-m}{C}\right)$$

1S yield extracted in 8.8-9.8 GeV/c<sup>2</sup> range.  
 1S+2S+3S extracted in 8.8-11 GeV/c<sup>2</sup> range.

1S+2S+3S yield:  $831 \pm 45$   
 1S yield:  $653 \pm 35$

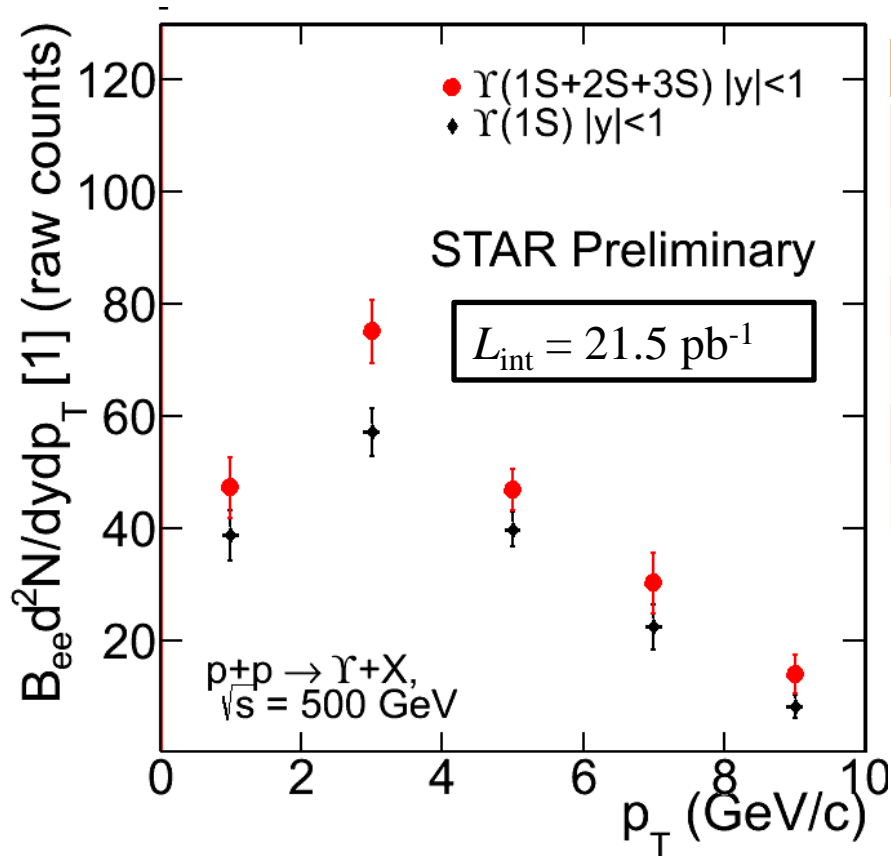
Statistical uncertainty ~5%  
 (~20% stat.+fit uncertainty for published p+p 200 GeV results)

# Upsilon signal – $p_T$ bins



Upsilon signal  
in 5  $p_T$  bins up  
to 10  $\text{GeV}/c$ !!

# Uncorrected $p_T$ spectrum

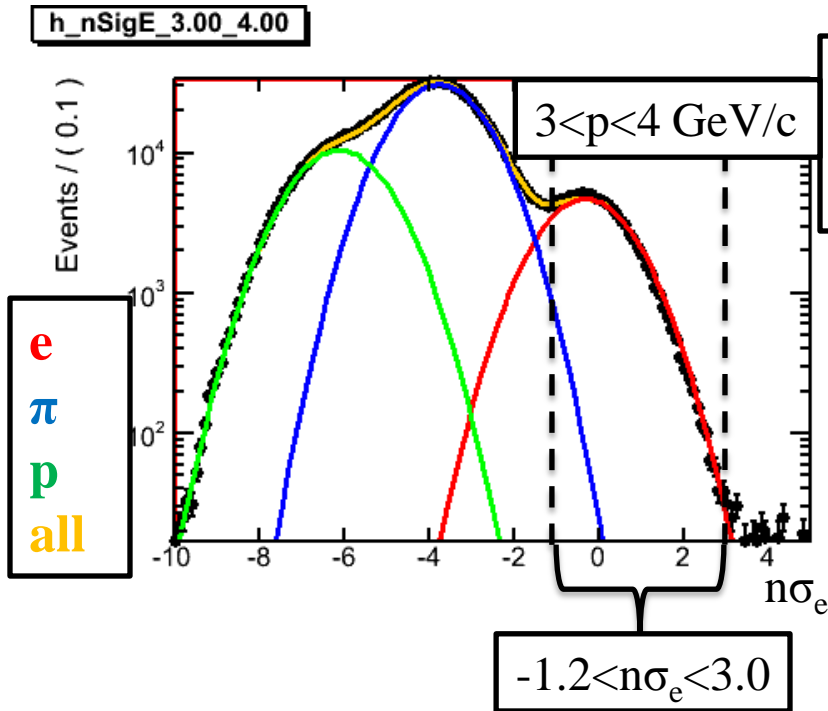


	1S+2S+3S		
$p_T$ [GeV/c]	Yield	Uncertainty	Uncertainty %
0-2	189	22	12
2-4	301	23	8
4-6	188	15	8
6-8	121	22	19
8-10	56	15	27

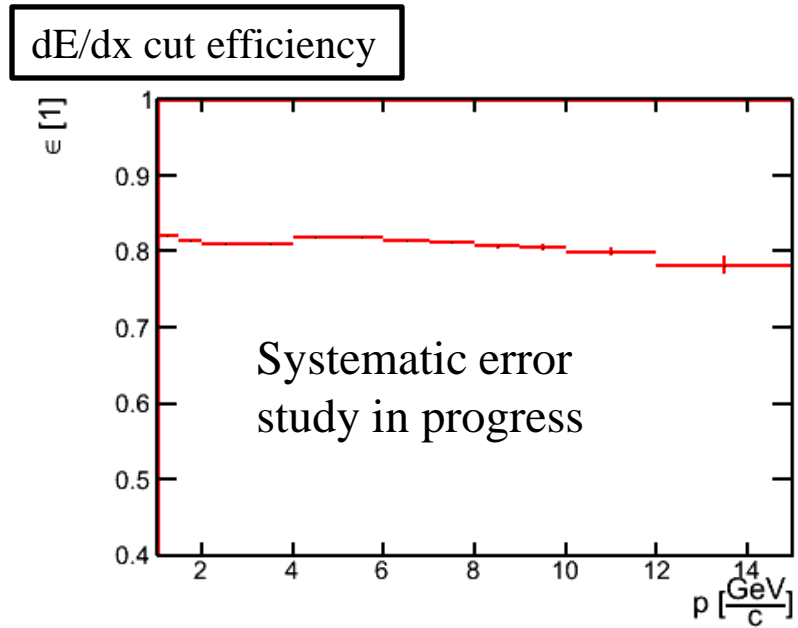
	1S		
$p_T$ [GeV]	Yield	Uncertainty	Uncertainty %
0-2	155	18	12
2-4	229	18	8
4-6	159	13	9
6-8	90	16	18
8-10	32.9	8.4	26

- Statistical uncertainty < 20% for  $p_T < 8 \text{ GeV}/c$
- Even better than < 10% for  $2 < p_T < 6 \text{ GeV}/c$

# dE/dx cut efficiency



Obtained a cleaner sample for dE/dx cut efficiency calculation and fitted multiple gaussians to  $n\sigma_e$  distributions in momentum slices



Efficiency is a ratio of the integral of the electron gaussian within cut range to the total electron gaussian

$$\text{Normalized } dE/dx: n\sigma_e = \log\left(\frac{dE/dx_e}{dE/dx|_{Bichsel}}\right) / \sigma$$

# BEMC E/p efficiency

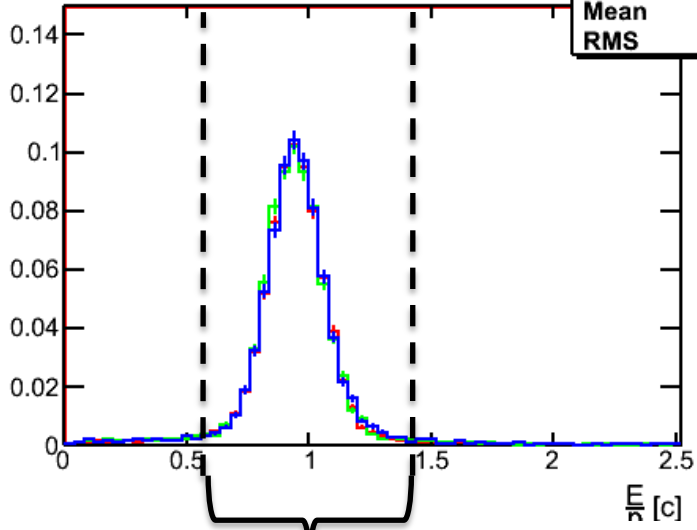
h\_1S\_p\_Eclu

$2 < p < 3 \text{ GeV}$

h_1S_p_Eclu_2_3	
Entries	9745
Mean	0.9398
RMS	0.1826

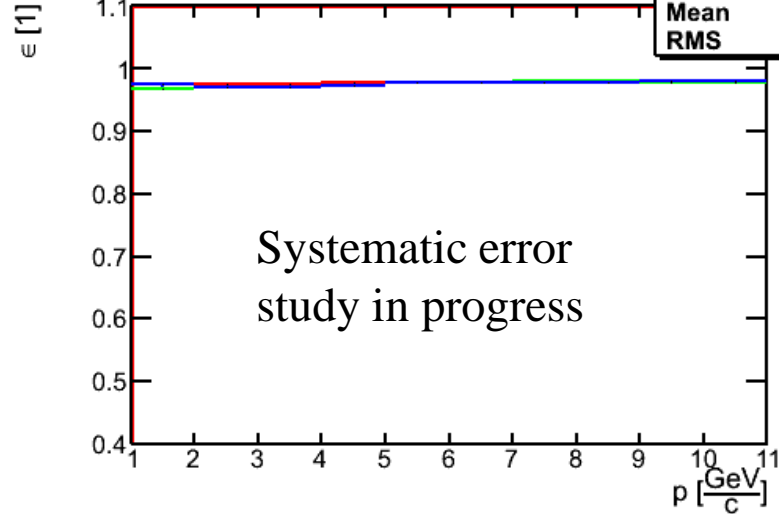
h\_1S\_p\_EoP\_eff

h_1S_p_EoP_eff	
Entries	10
Mean	6.004
RMS	2.872



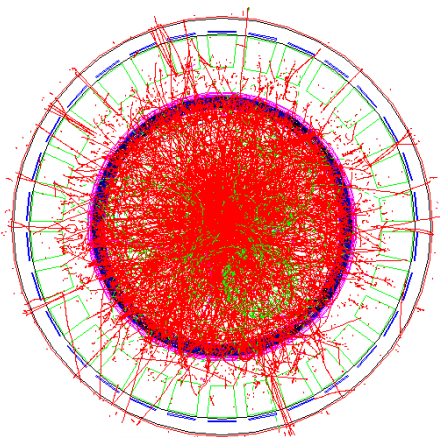
$0.55 < E_{\text{cluster}}/p < 1.45$

Compared  $E_{\text{clu}}/p$  distribution from MC simulation



MC simulation  $\Upsilon(1S)$   
 MC simulation  $\Upsilon(2S)$   
 MC simulation  $\Upsilon(3S)$

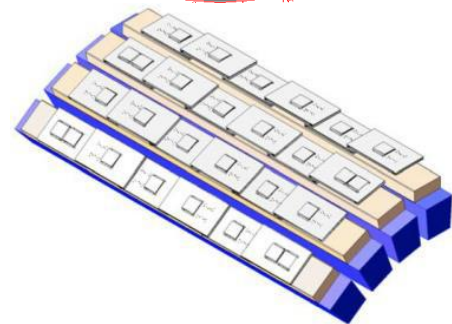
# Muon Telescope Detector



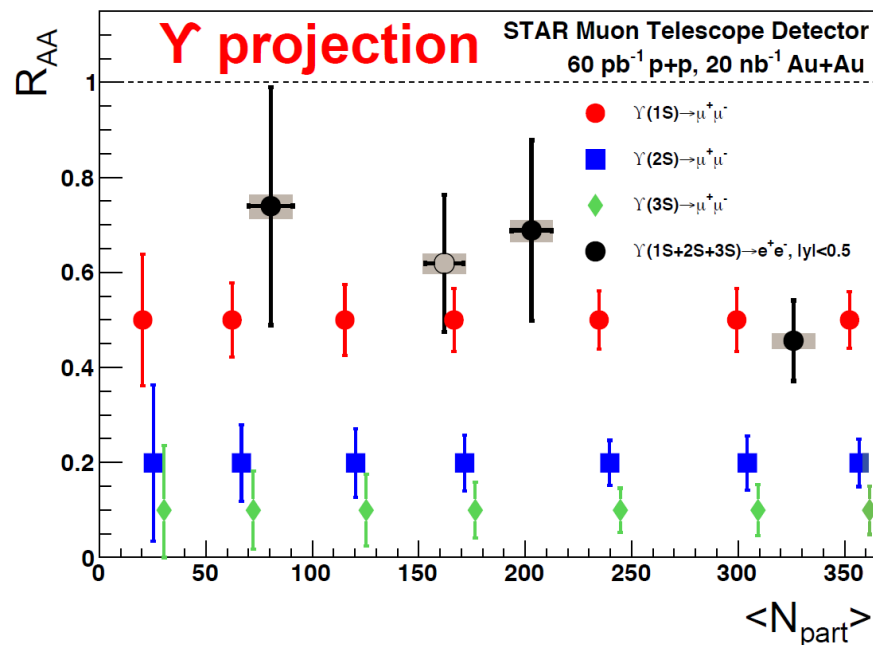
- Multi-gap Resistive Plate Chamber (MRPC) - gas detector outside STAR magnet

- Allows precision measurements of heavy quarkonia in dimuon channel (no  $\gamma$  conversions, less bremsstrahlung than e and less dalitz decays)

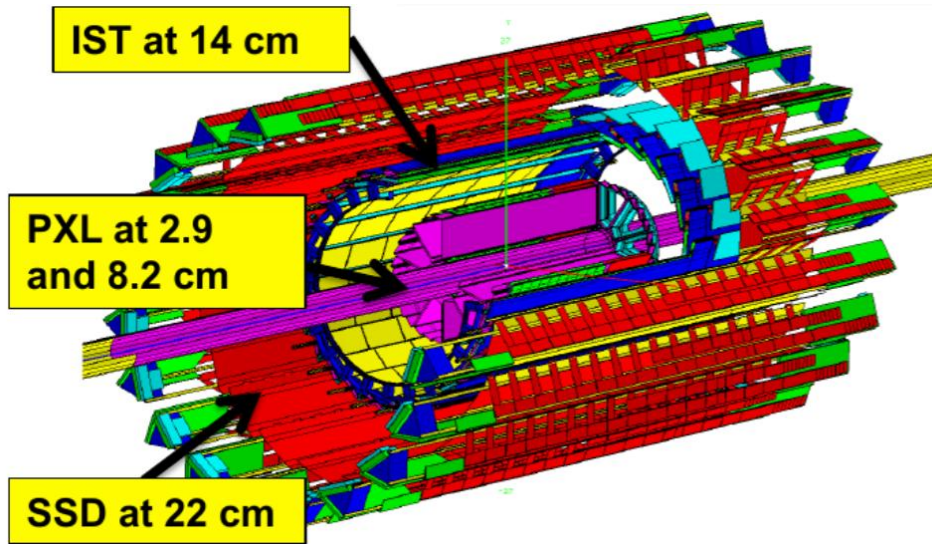
- Complete in 2014, sampled  $\sim 13.8 \text{ nb}^{-1}$  in Au+Au data



Acceptance:  
45% in  $\phi$   
 $|\eta| < 0.5$

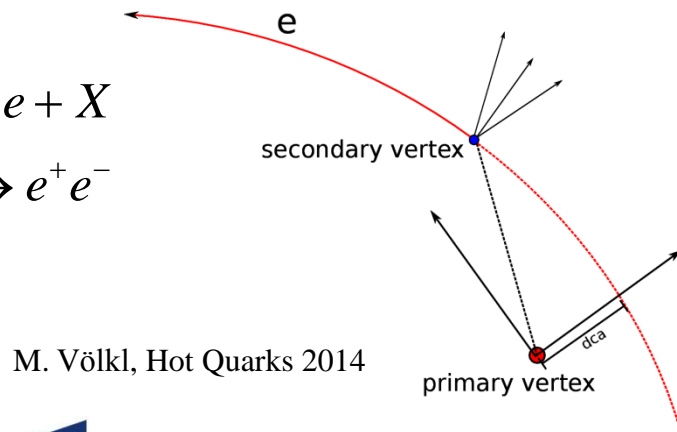
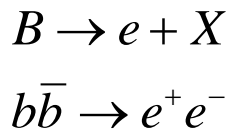


# Heavy Flavor Tracker

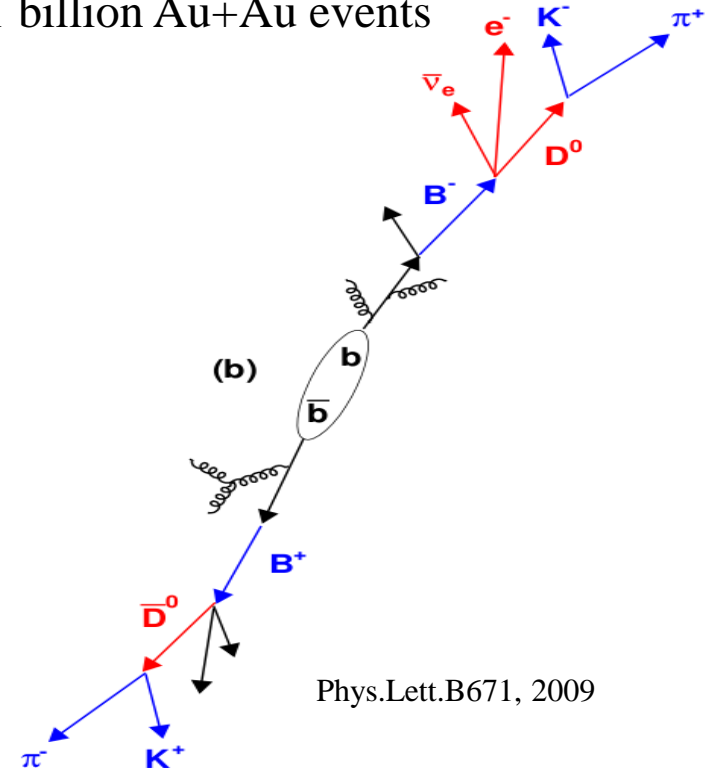


- State-of-the-art silicon tracker for heavy flavor studies
- Allows direct reconstruction of secondary vertices from D and B meson decays
- Fully installed in 2014, recorded ~1 billion Au+Au events

May provide direct  $b\bar{b}$  (background for  $\Upsilon$ ) cross section measurements through:



M. Völkl, Hot Quarks 2014



Phys.Lett.B671, 2009



# Summary

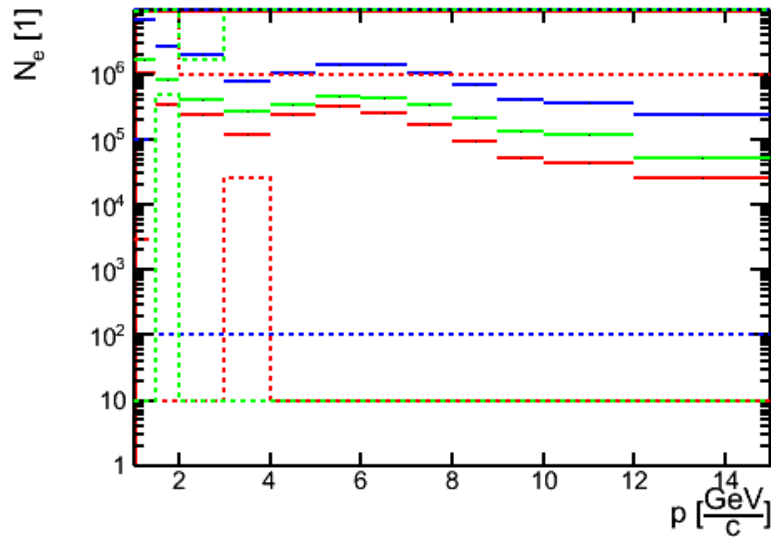
- Upsilon measurement in p+p 500 GeV could provide a better statistical precision baseline for Au+Au 200 GeV after extrapolation
- Data can be divided into 5  $p_T$  bins from 0-10 GeV/c
- With current data ( $L_{\text{int}} = 21.5 \text{ pb}^{-1}$ ), we can expect  $p_T$  spectrum for Upsilon 1S and combined states with statistical uncertainty  $< 20\%$  for  $p_T$  up to 8 GeV/c
- Stay tuned for upcoming results!!

# BACKUP

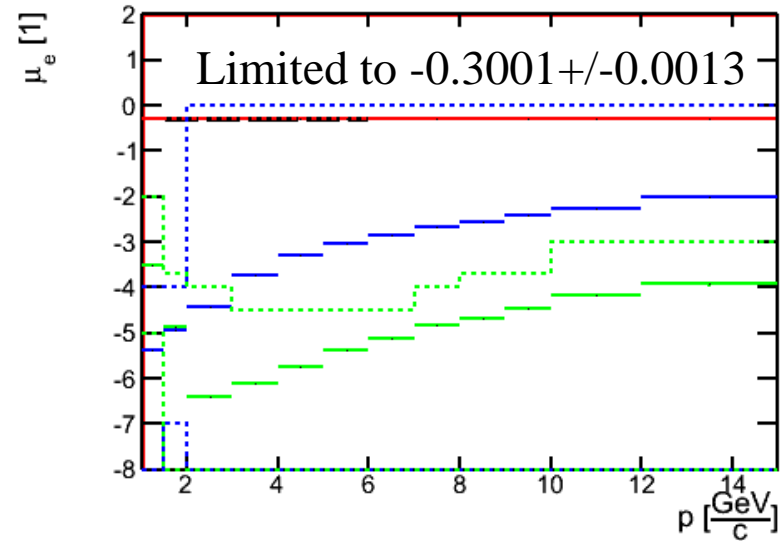


# parameters

Electron yield

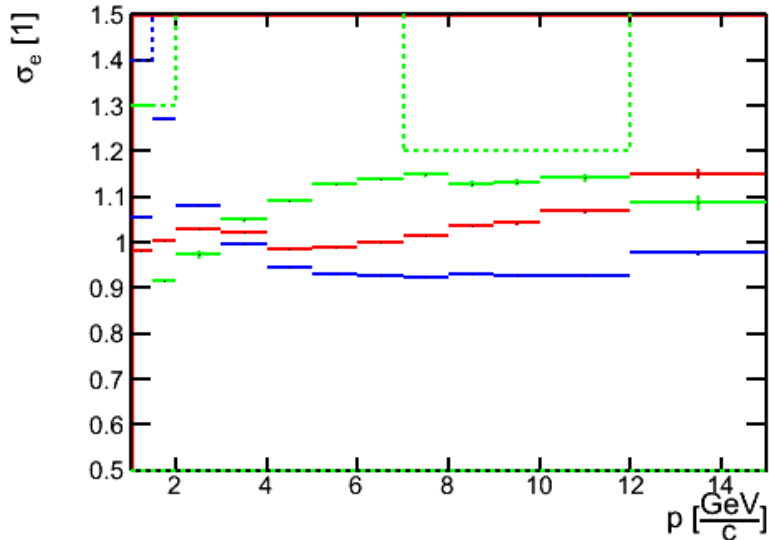


Electron mean



e  
π  
p  
all

Electron sigma



Smooth dependence => good

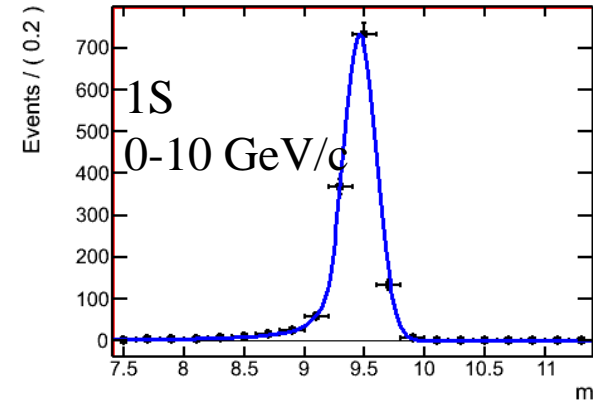
Set limit for electron mean based on  
straight line fit in 1.5-6 GeV:  
-0.3001 +/- 0.0013

Very little constraints (dotted lines)

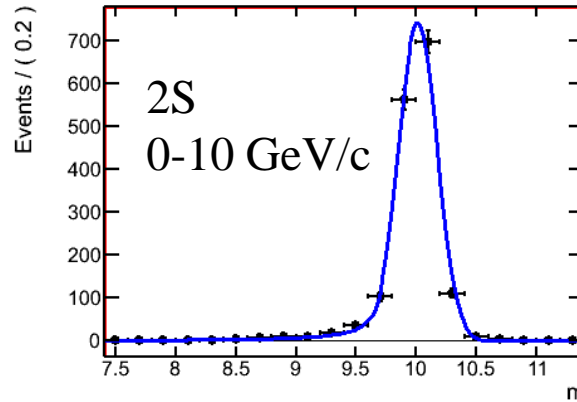


# Upsilon lineshapes

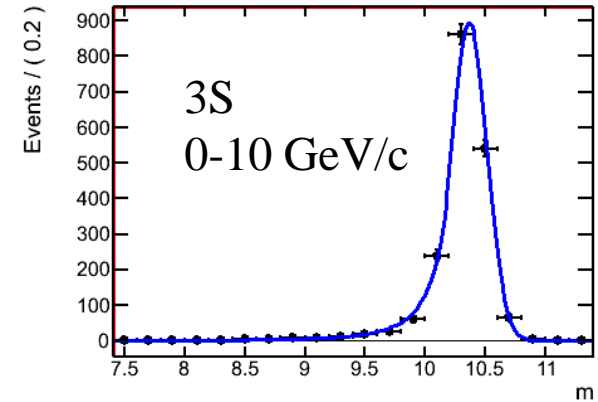
Lineshape\_0.0<pT<10.0\_GeV



Lineshape\_0.0<pT<10.0\_GeV



Lineshape\_0.0<pT<10.0\_GeV



Lineshapes from MC simulation include the reconstruction efficiency.

Fitted Crystal Ball functions to lineshapes

Made likelihood fits with RooFit in each  $p_T$  bin to take into account  $p_T$  dependence

Crystal Ball function:

$$f(x; \alpha, n, \bar{x}, \sigma) = N \begin{cases} \exp\left(-\frac{(x - \bar{x})^2}{2\sigma^2}\right), & \text{for } \frac{x - \bar{x}}{\sigma} > -\alpha \\ A\left(B - \frac{x - \bar{x}}{\sigma}\right)^{-n}, & \text{for } \frac{x - \bar{x}}{\sigma} \leq -\alpha \end{cases}$$

$$A = \left(\frac{n}{|\alpha|}\right)^n \cdot \exp\left(-\frac{|\alpha|^2}{2}\right) \quad C = \frac{n}{|\alpha|} \cdot \frac{1}{n-1} \cdot \exp\left(-\frac{|\alpha|^2}{2}\right)$$

$$B = \frac{n}{|\alpha|} - |\alpha| \quad D = \sqrt{\frac{\pi}{2}} \cdot \left(1 + \operatorname{erf}\left(\frac{|\alpha|}{\sqrt{2}}\right)\right)$$

$$N = \frac{1}{\sigma(C + D)}$$

DESY F31-86-02(1986), Appendix E