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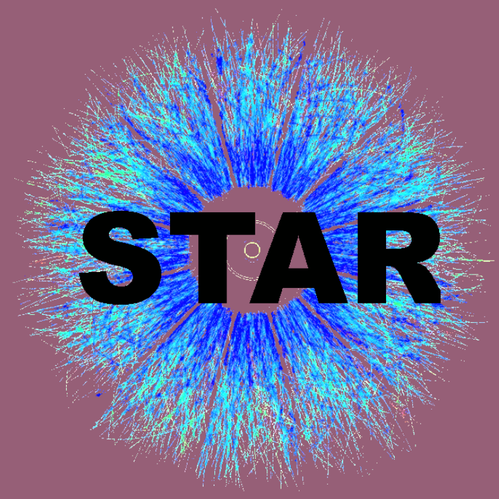

Charged and neutral kaons femtoscopy measured by the STAR experiment

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**15th Workshop on Particle Correlations and Femtoscopy,
East Lansing, Michigan, USA**

18 – 22 July, 2022

Partially funded by:

U.S. DEPARTMENT OF ENERGY | Office of Science

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Femtoscopy - introduction

statistical **model** **experiment**

$$CF(\vec{q}) = \frac{P_{12}(\overrightarrow{p_1},\overrightarrow{p_2})}{P_1(\overrightarrow{p_1})P_2(\overrightarrow{p_2})} = \int d^3r\, S(\vec{q},\vec{r})|\Psi(\vec{q},\vec{r})|^2 = \frac{A(\vec{q})}{B(\vec{q})}$$

$\overrightarrow{p_1}, \overrightarrow{p_2}$ - single particle momentum

$S(\vec{q},\vec{r})$ – emission function

$\Psi(\vec{q},\vec{r})$ – pair wave function

$\vec{q} = |\overrightarrow{p_1} - \overrightarrow{p_2}|$ - relative momentum

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Kaon correlation functions are sensitive to:

$\mathbf{K^{\pm}\,K^{\pm}}$	$\mathbf{K^0_s K^0_s}$	$\mathbf{K^0_s K^{\pm}}$
Quantum Statistical effects (QS) Final State Interaction (FSI) - Coulomb interaction (COUL)	Quantum Statistical effects (QS) Final State Interaction (FSI) - strong interaction (SI)	Final State Interaction (FSI) - strong interaction (SI)

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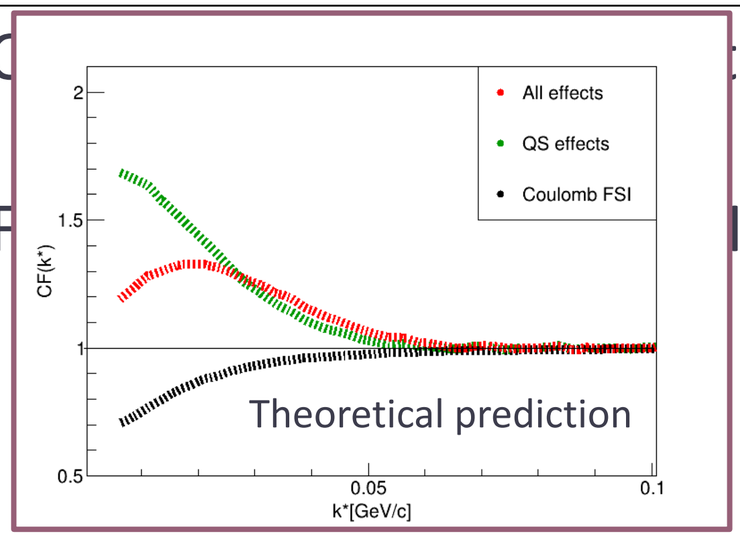
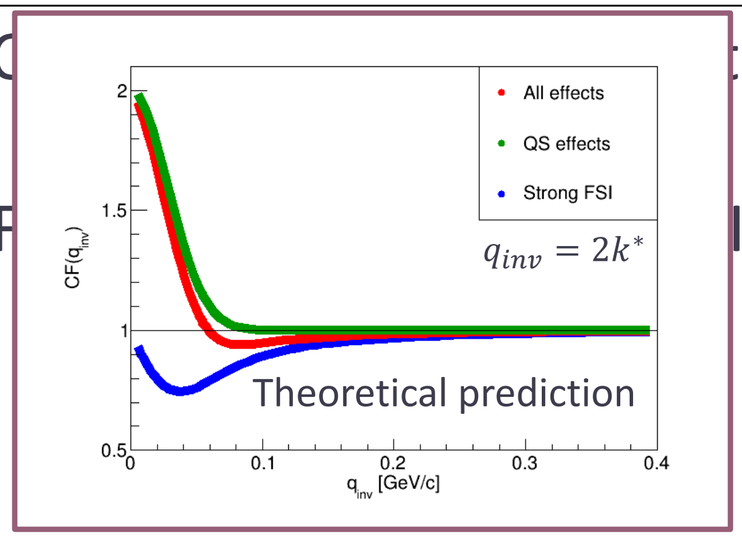
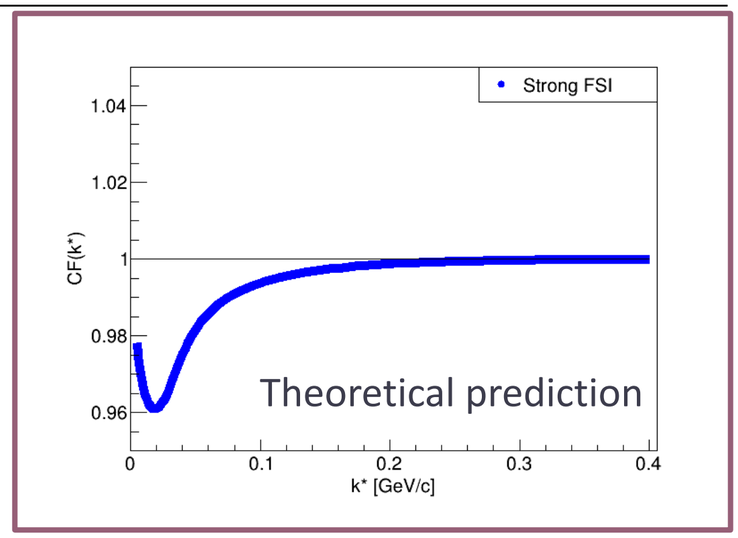


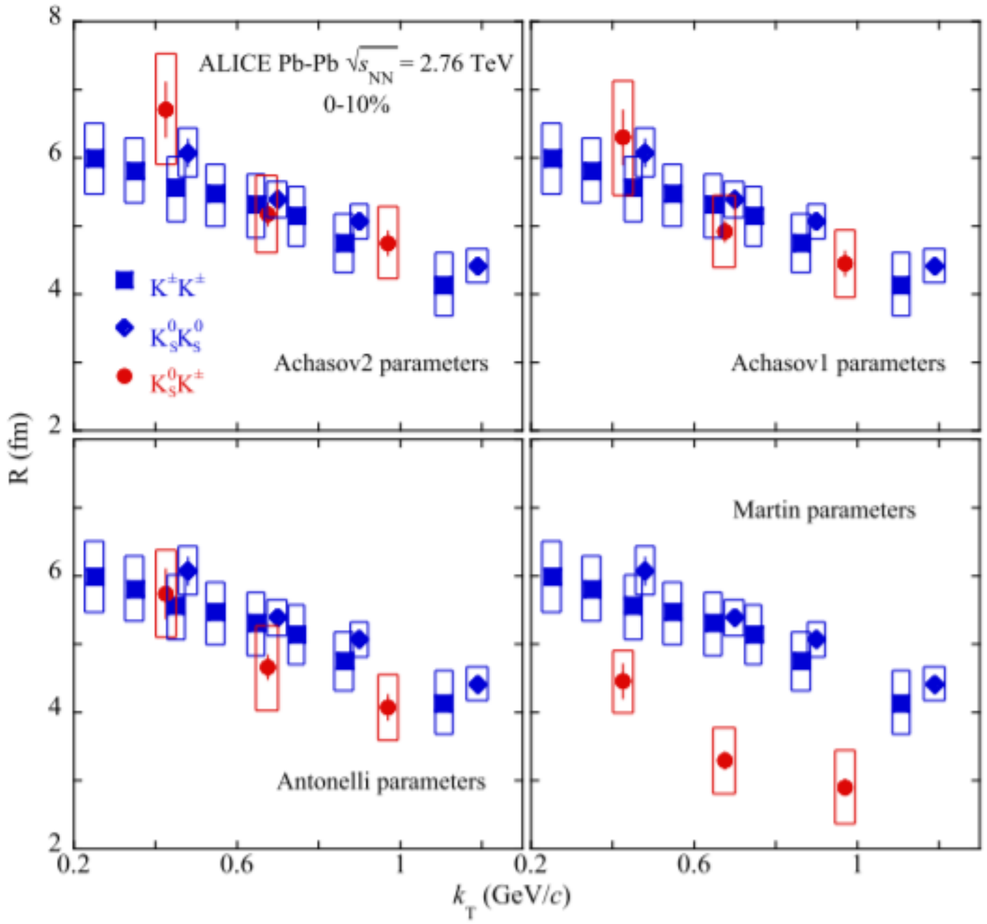




k^*- first particle momentum in PRF

Motivation

ALICE, Physics Letters B **774**(2017) 64



$$k_T = \frac{p_{T1} + p_{T2}}{2}$$

Kaons provide complementary information to pions:

- contain strange quarks (larger production of strange particles is one of the signatures of QGP)
- less affected by the feed-down from resonance decays
- smaller cross section on reaction with the hadronic matter

Very interesting:

- compare femtoscopic results for all possible kaon combination ($K^{\pm}K^{\pm}$, $K^0_SK^0_S$, $K^0_SK^{\pm}$)
- $K^0_SK^{\pm}$ - a_0 could be a 4-quark state (a tetraquark)

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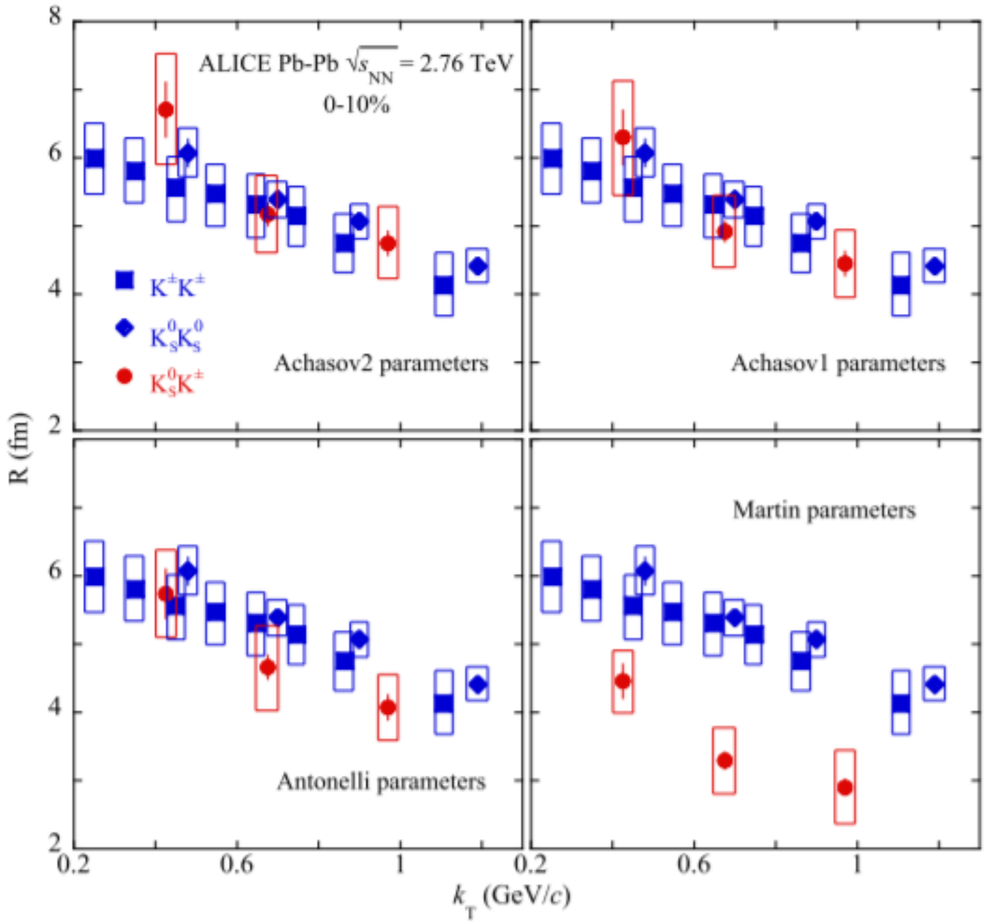

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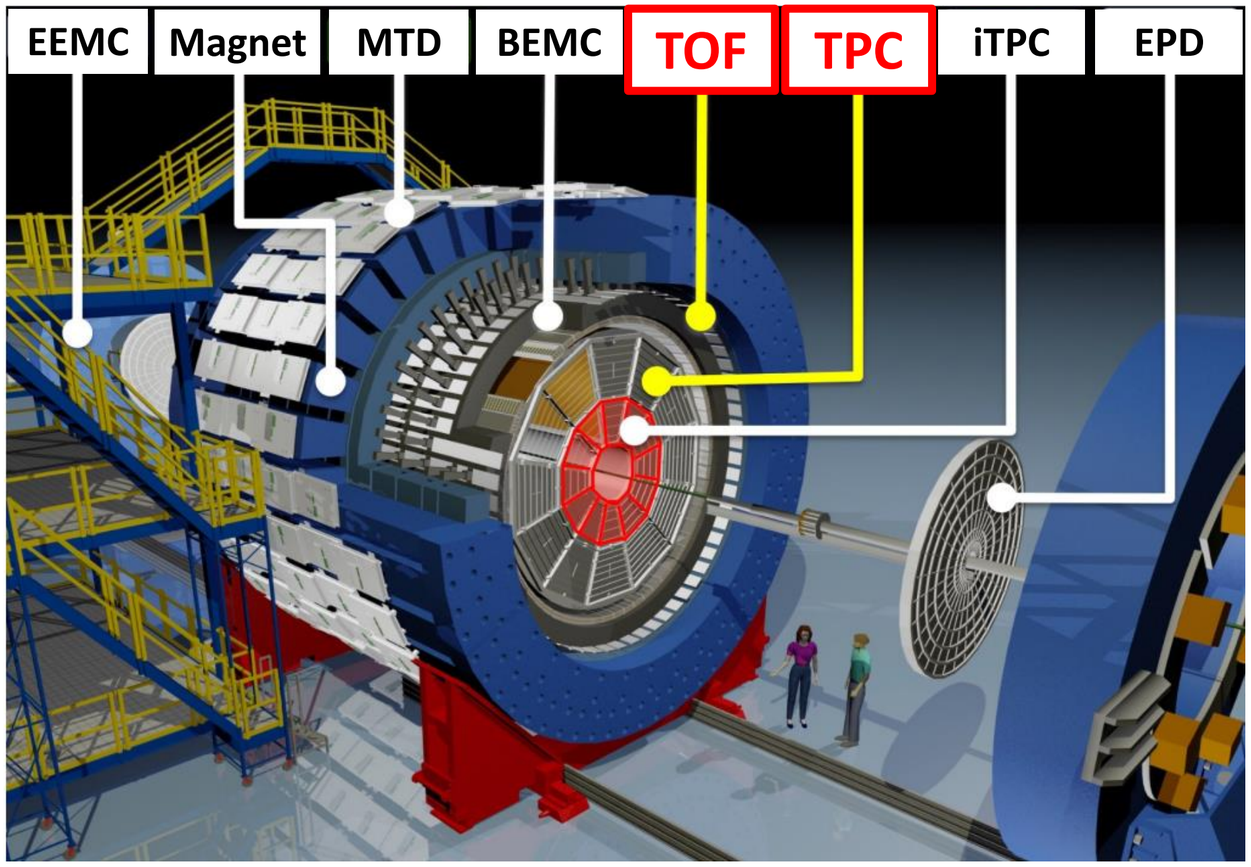
The STAR experiment



→ Exellent particle identification

→ Large, uniform acceptance at mid-rapidity

Time Projection Chamber

PID: dE/dx

Tracking

$0 < \phi < 2\pi$, $|\eta| < 1$

Time-Of-Flight

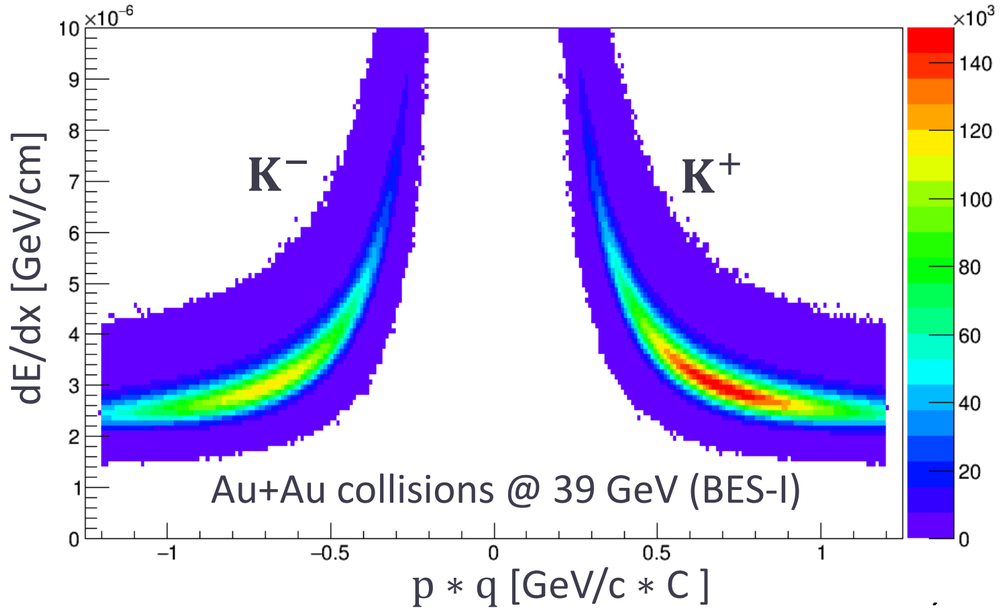
Time resolution < 80 ps

PID: m^2

Datasets

Energy $\sqrt{s_{NN}}$	Year	Mode	Statistics (M)
39 GeV	2010	Collider	~83
200 GeV	2010	Collider	~260

Au+Au collisions

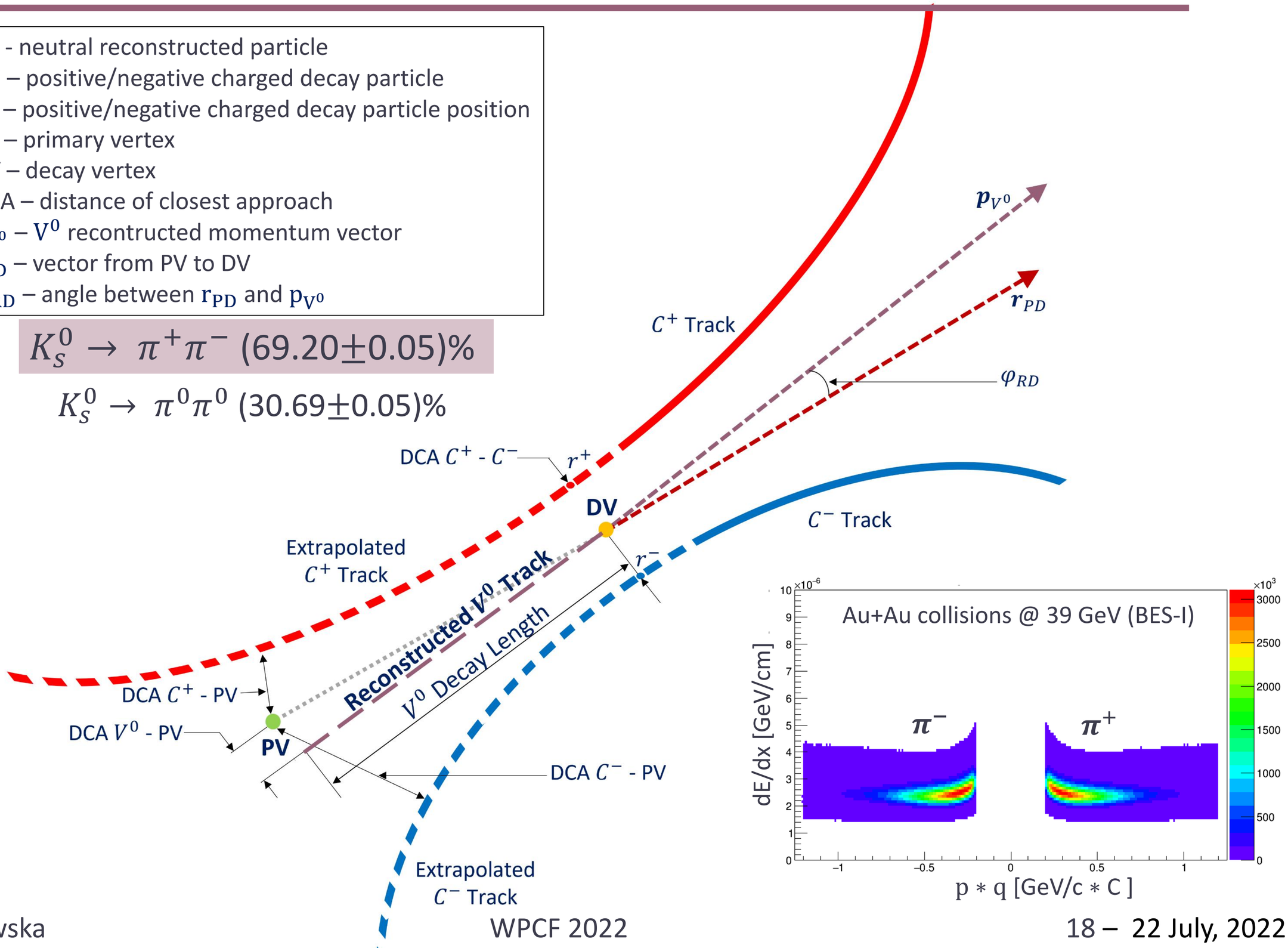


Neutral kaon reconstruction

V^0 - neutral reconstructed particle
$C^{\pm}$ – positive/negative charged decay particle
$r^{\pm}$ – positive/negative charged decay particle position
PV – primary vertex
DV – decay vertex
DCA – distance of closest approach
p_{V^0} – V^0 recontructed momentum vector
r_{PD} – vector from PV to DV
φ_{RD} – angle between r_{PD} and p_{V^0}

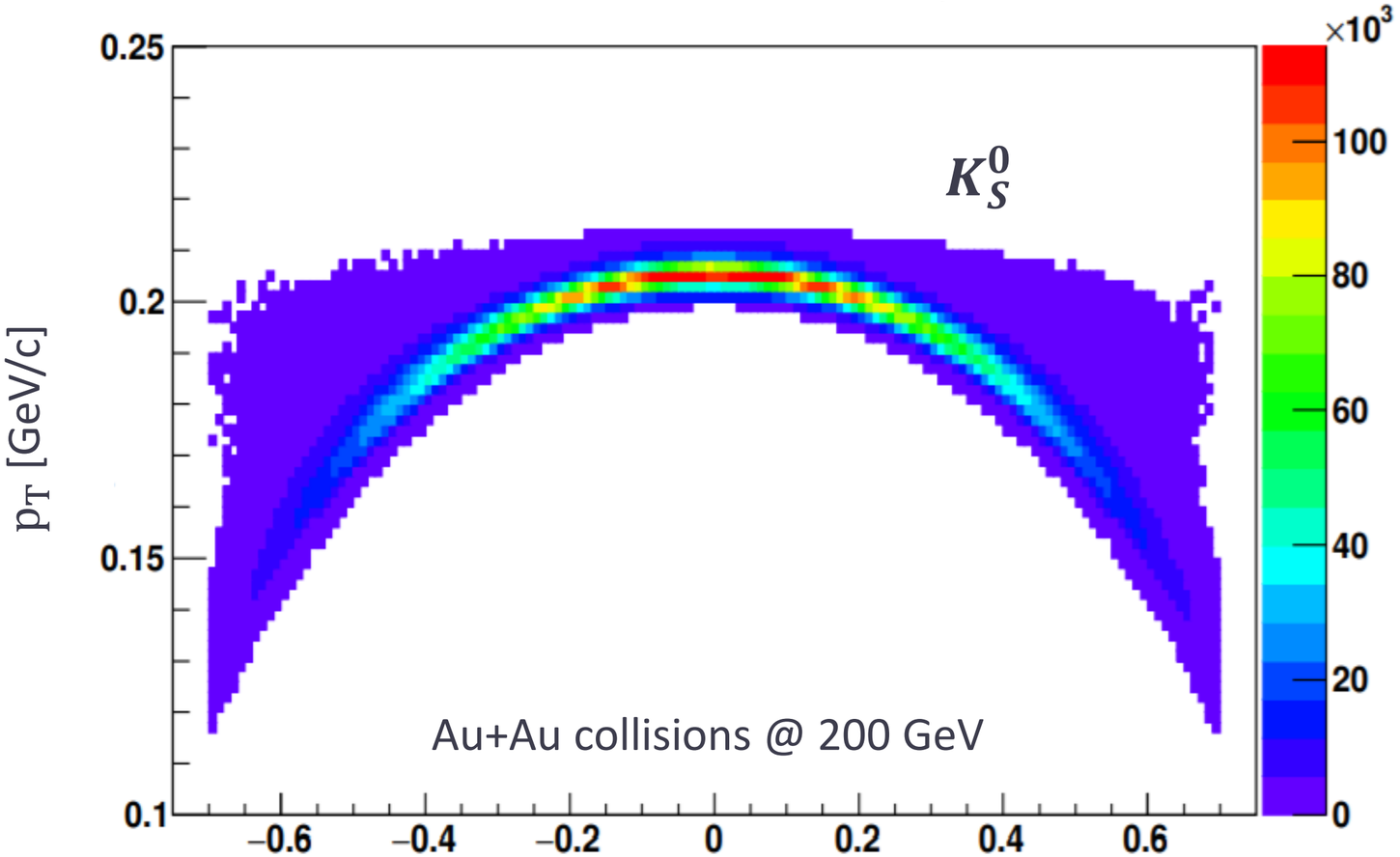
$K_s^0 \rightarrow \pi^+\pi^-$ (69.20±0.05)%

$K_s^0 \rightarrow \pi^0\pi^0$ (30.69±0.05)%



Neutral kaon reconstruction

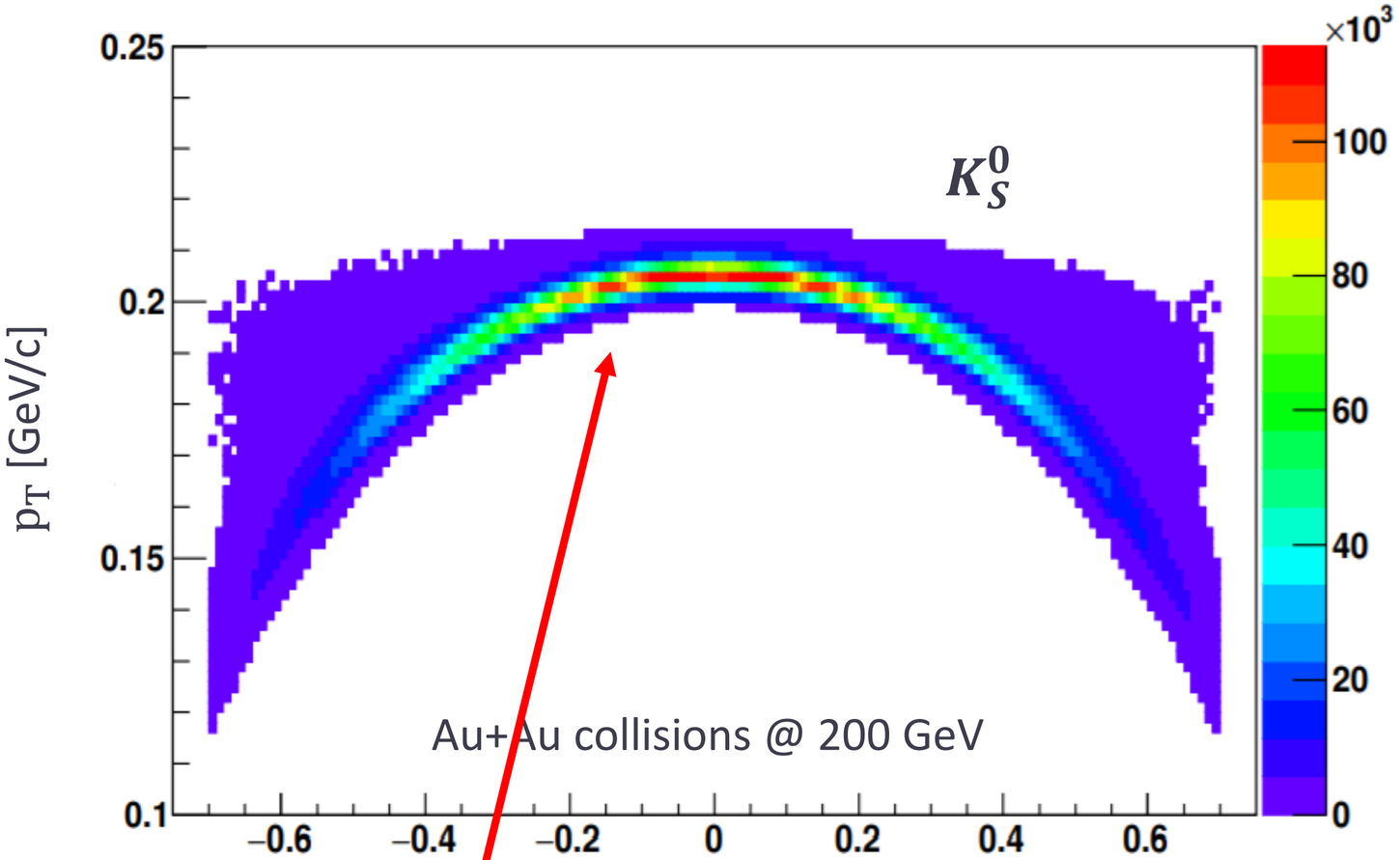
Armenteros-Podolański plot



$$\alpha = \frac{p_l^+ - p_l^-}{p_l^+ + p_l^-}$$

Neutral kaon reconstruction

Armenteros-Podolański plot



$$\alpha = \frac{p_l^+ - p_l^-}{p_l^+ + p_l^-}$$

Decay products of the neutral kaon have the same mass and therefore their momenta are distributed symmetrically on average

Parametrization - $K^0_s K^0_s$

Gaussian density distribution (includes only **QS effects**): $CF(q_{inv}) = 1 + \lambda e^{[-R^2_{inv} q^2_{inv}]}$

λ - the correlation strength, R_{inv} - the size of the particle-emitting source.

Lednicky & Lyuboshitz model includes **strong FSI**: [Sov.J.Nucl.Phys. 35, 770 (1982)]

$$CF(q_{inv}) = 1 + \lambda\left(e^{[-R^2_{inv} q^2_{inv}]} + \frac{1}{2}\left[\left|\frac{f(k^*)}{R_{inv}}\right|^2 + \frac{4\Re f(k^*)}{\sqrt{\pi} R_{inv}} F_1(q_{inv} R_{inv}) - \frac{2\Im f(k^*)}{\sqrt{\pi} R_{inv}} F_2(q_{inv} R_{inv})\right]\right)$$

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QS effect

strong FSI through the $f_0(980)$ and $a_0(980)$ resonances

$$F_1(z) = \int_0^z dx \frac{e^{x^2} - x^2}{z}, \quad F_2(z) = \frac{1 - e^{x^2}}{z}$$

$$f(k^*) = \frac{1}{2}[f_0(k^*) + f_1(k^*)], \quad f_I(k^*) = \frac{\gamma_r}{m_r - s - i\gamma_r k^* - i\gamma'_r k'_r}, \quad s = 4\left(m_K^2 + k^{*2}\right)$$

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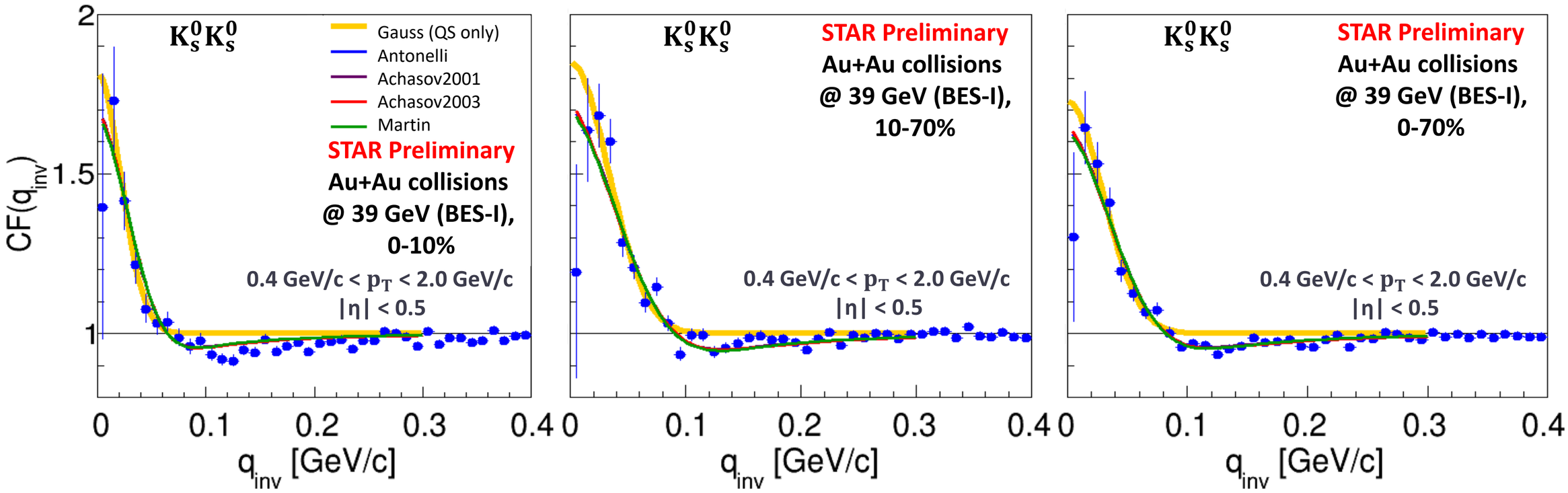
	$m_{f_0}\left[\frac{GeV}{c^2}\right]$	$\gamma_{f_0 K\bar{K}}$	$\gamma_{f_0 \pi\pi}$	$m_{a_0}\left[\frac{GeV}{c^2}\right]$	$\gamma_{a_0 K\bar{K}}$	$\gamma_{a_0 \pi\pi}$
Antonelli [1]	0.973	2.763	0.5283	0.985	0.4038	0.3711
Achasov2001 [2]	0.996	1.305	0.2684	0.992	0.5555	0.4401
Achasov2003 [3]	0.996	1.305	0.2684	1.003	0.8365	0.4580
Martin [4]	0.978	0.792	0.1990	0.974	0.3330	0.2220

[1] eConf C020620, THAT06 (2002), [2] Phys. Rev. D 63, 094007 (2001)
[3] Phys. Rev. D 68, 014006 (2003), [4] Nucl. Phys. B 121, 514–530 (1977)

$K^0_S K^0_S$ femtoscopy at 39 GeV & 200 GeV



$q_{inv} = \sqrt{(\overrightarrow{p_1} - \overrightarrow{p_2})^2 - (E_1 - E_2)^2}$

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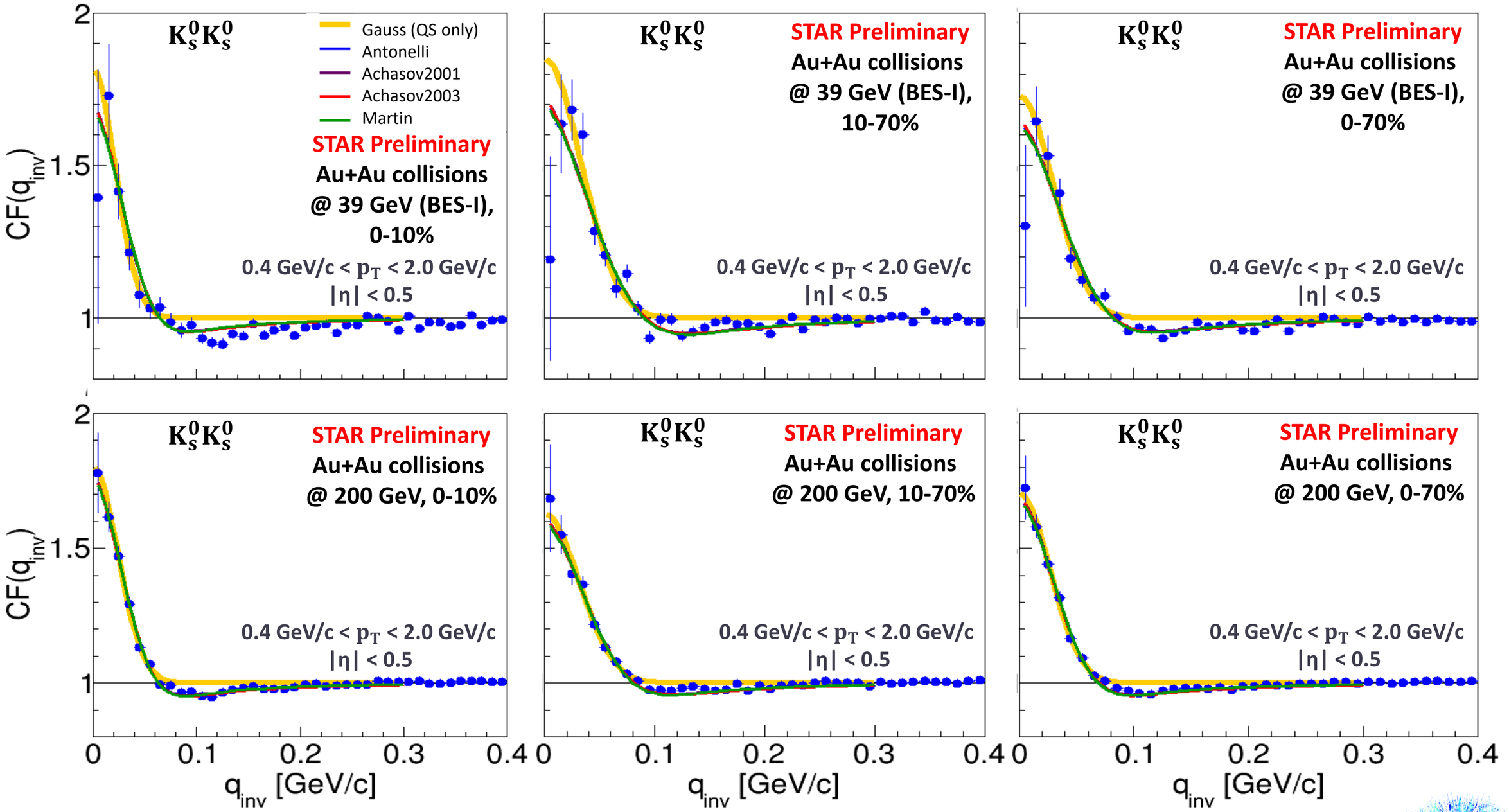

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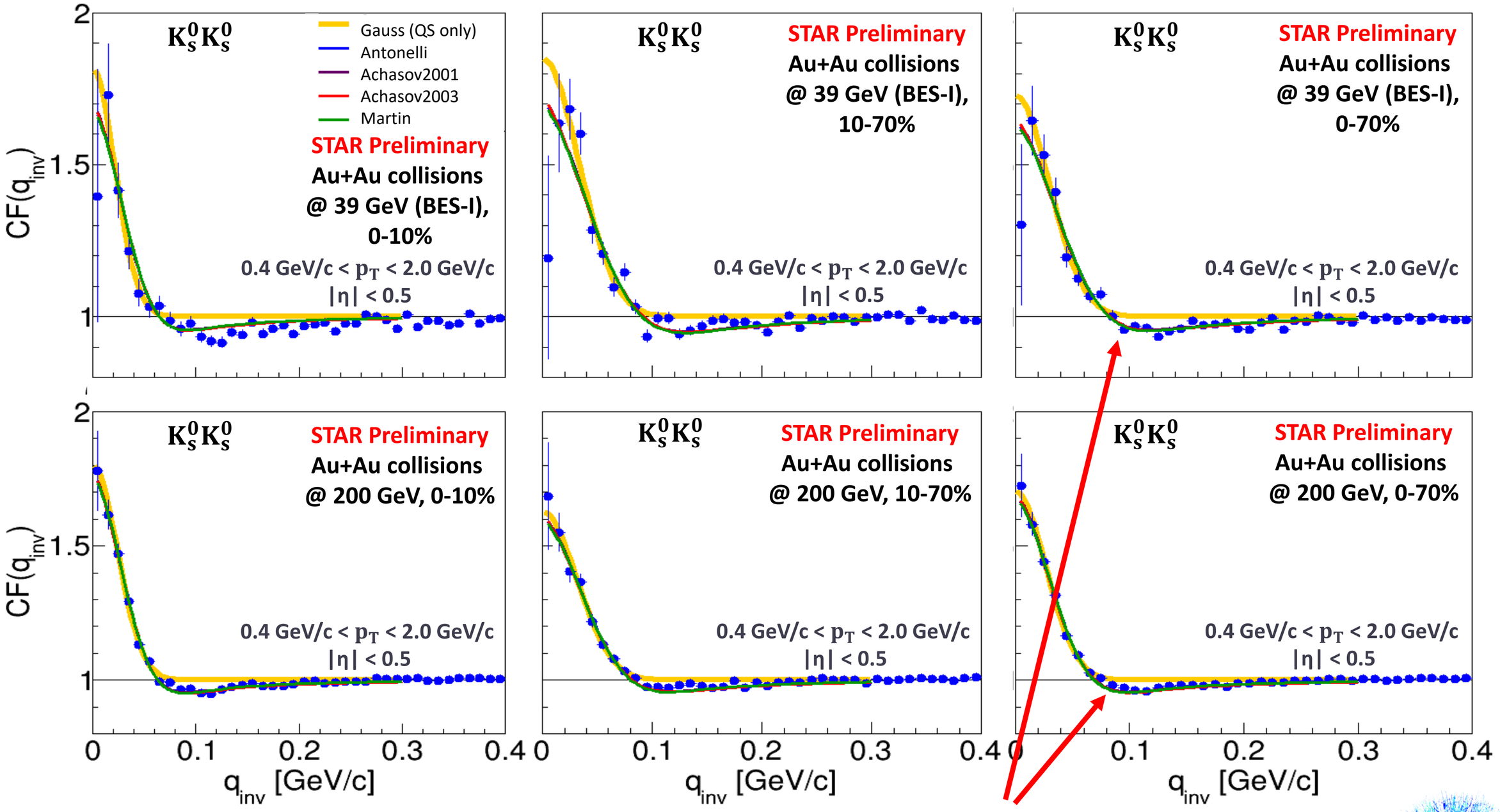

FSI is needed to reproduce the dip structure

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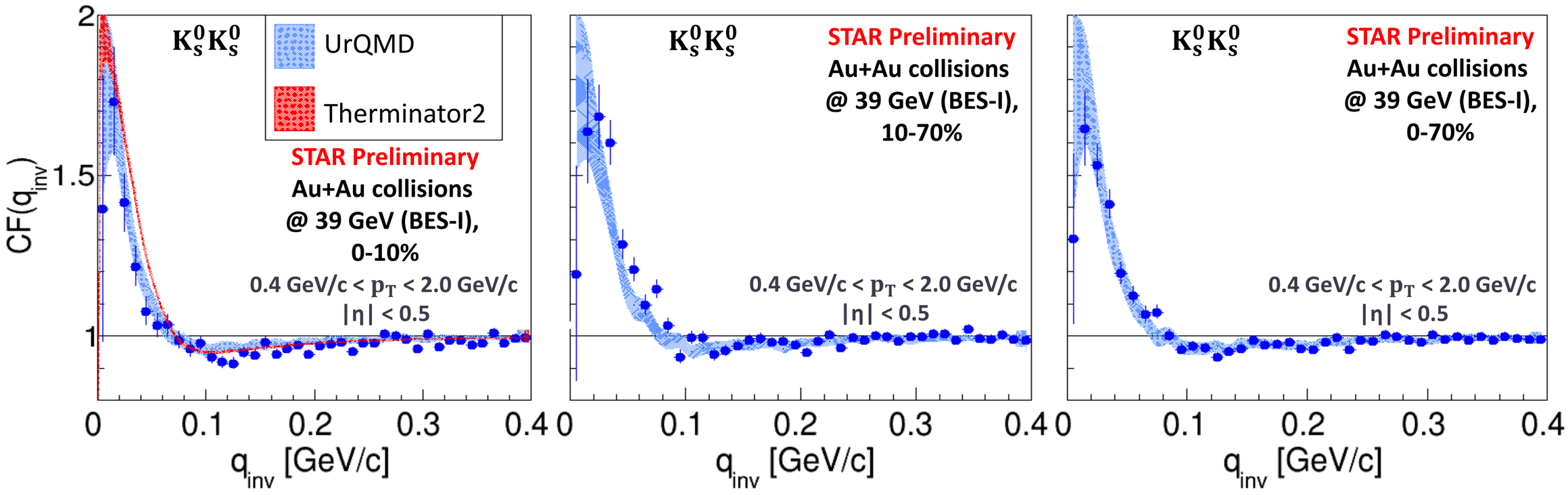
$K^0_S K^0_S$ femtoscopy at 39 GeV & 200 GeV



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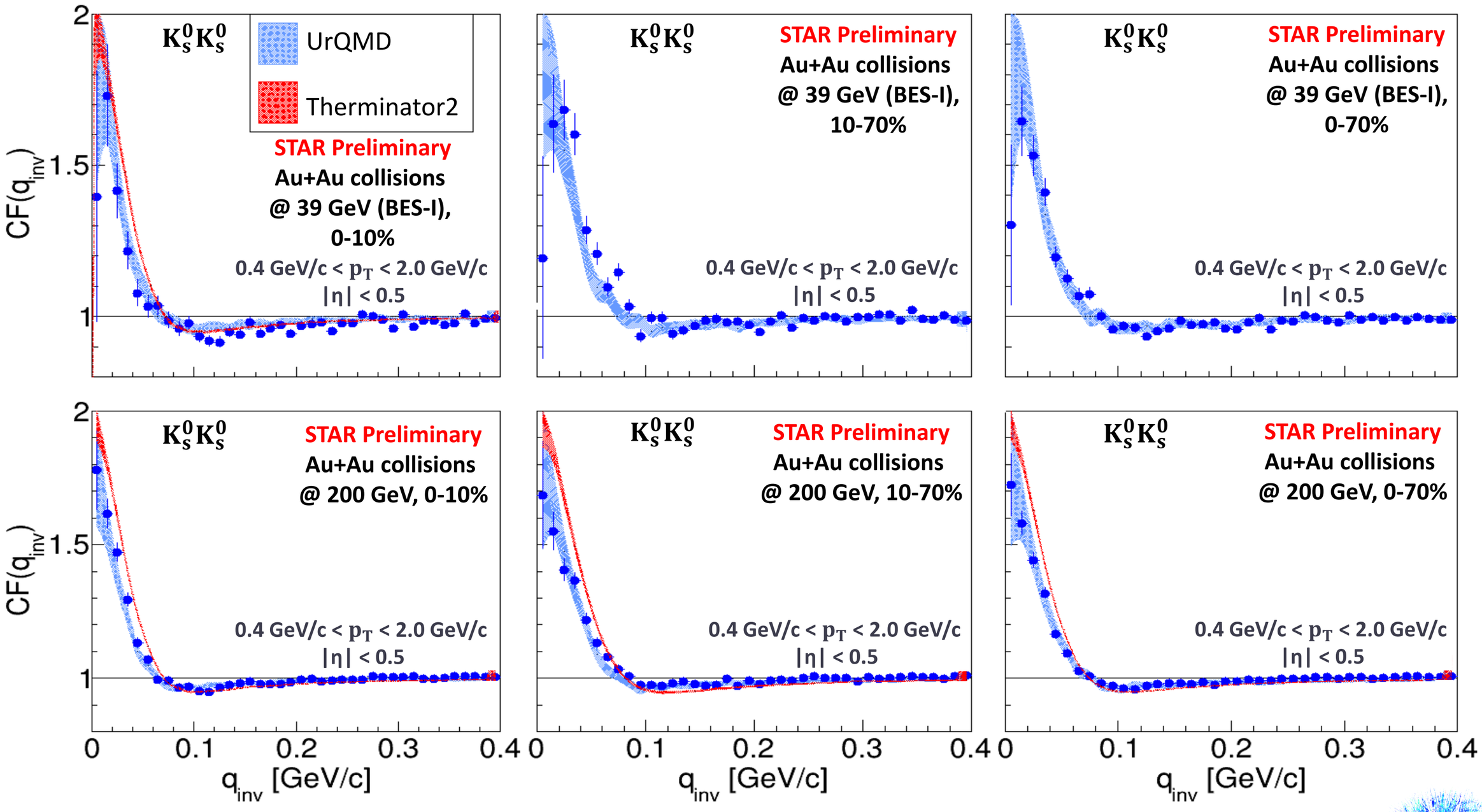

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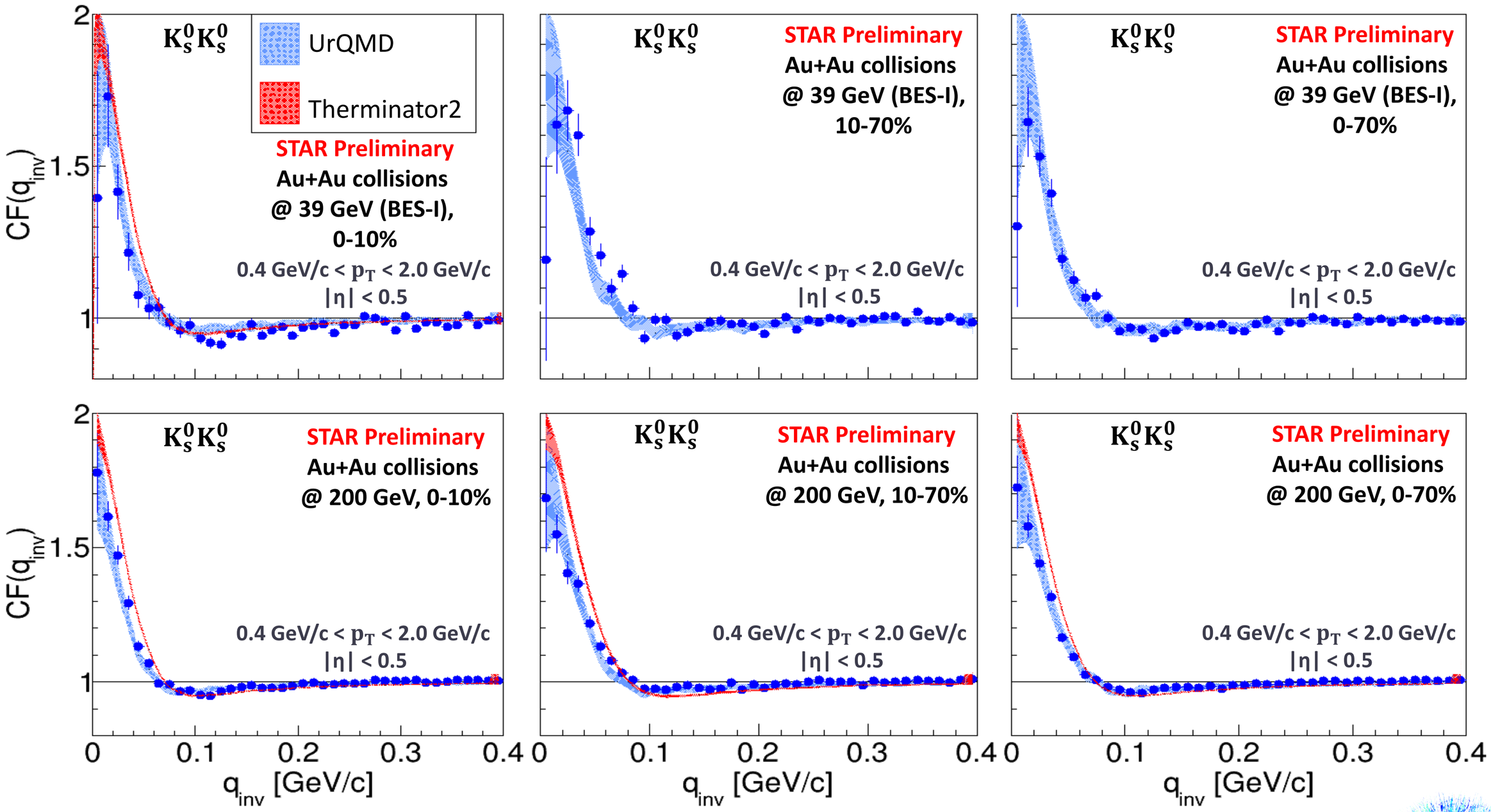
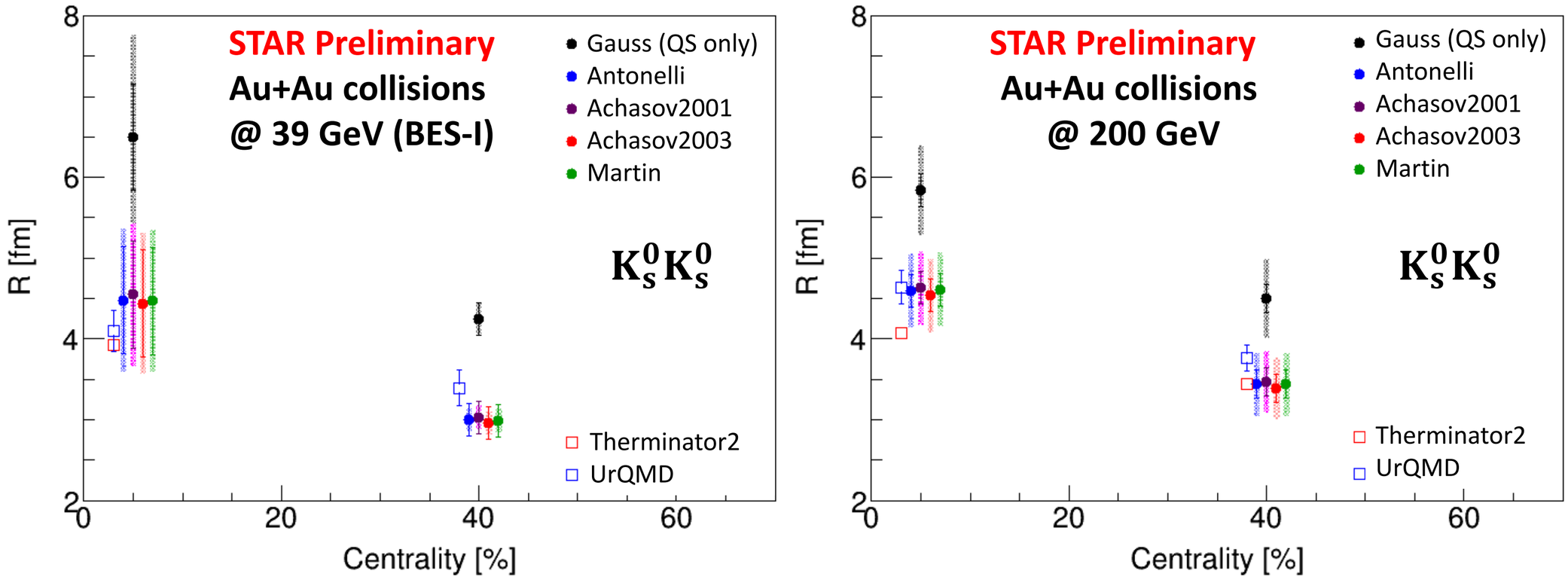


Good agreement of the experimental points with the models

$q_{inv} = \sqrt{(\overrightarrow{p_1} - \overrightarrow{p_2})^2 - (E_1 - E_2)^2}$

$K^0_S K^0_S$ femtoscopy – centrality dependence

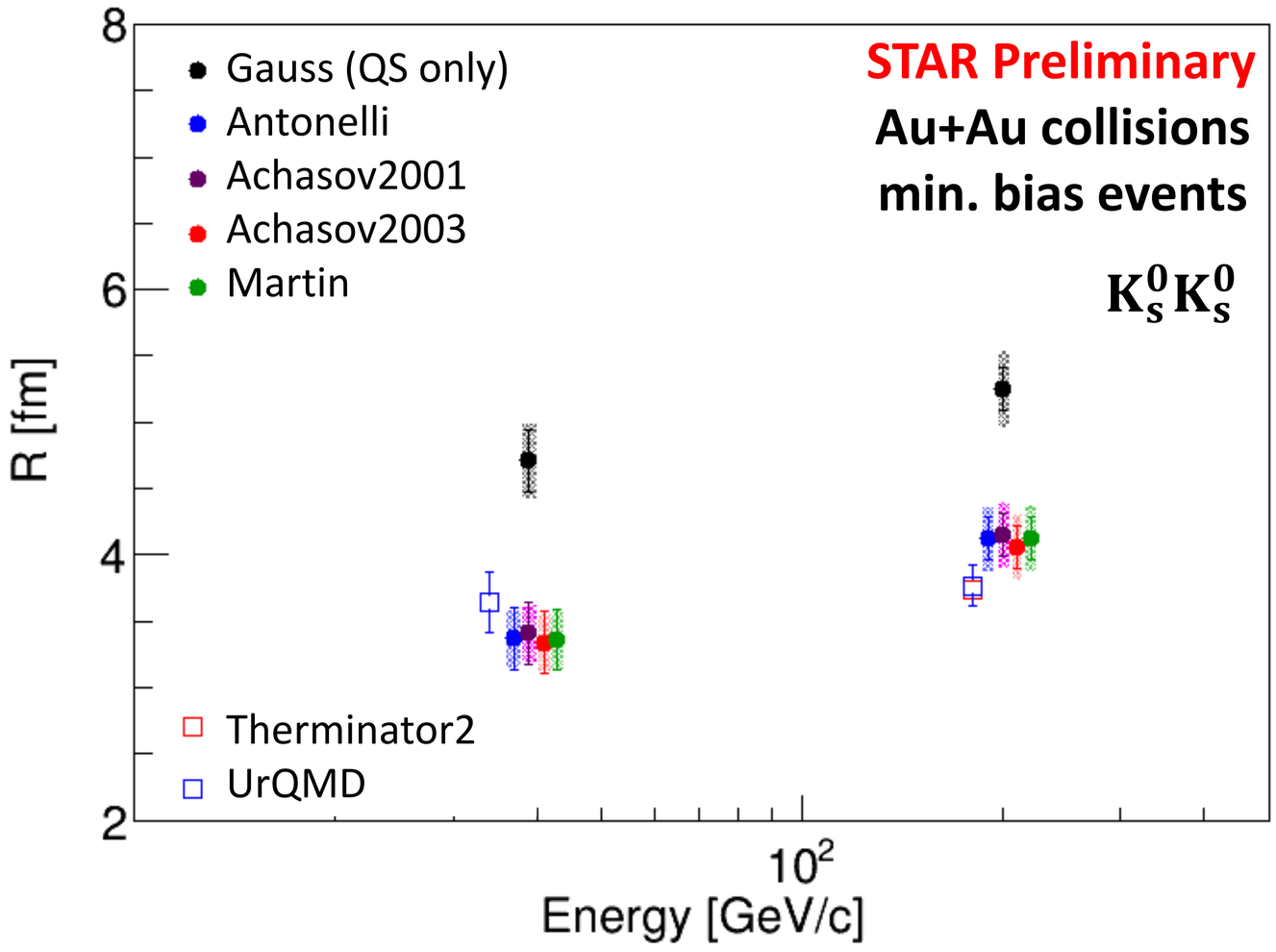


- Visible centrality dependence

$$R_{0-10\%} > R_{10-70\%}$$

- Only Gaussian assumption leads to larger R, showing importance to include SI

$K^0_S K^0_S$ femtoscopy – energy dependence



- Visible energy dependence for both parametrization

$$R_{200\ GeV} > R_{39\ GeV}$$

- Source sizes from models and FSI parametrizations consistent within the range of uncertainty

Parametrization - $K_s^0 K^{\pm}$

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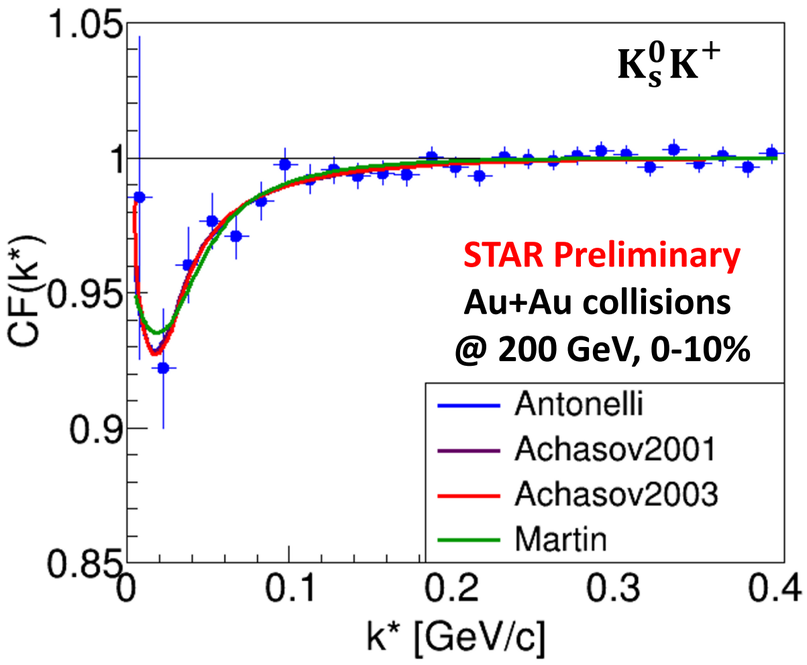
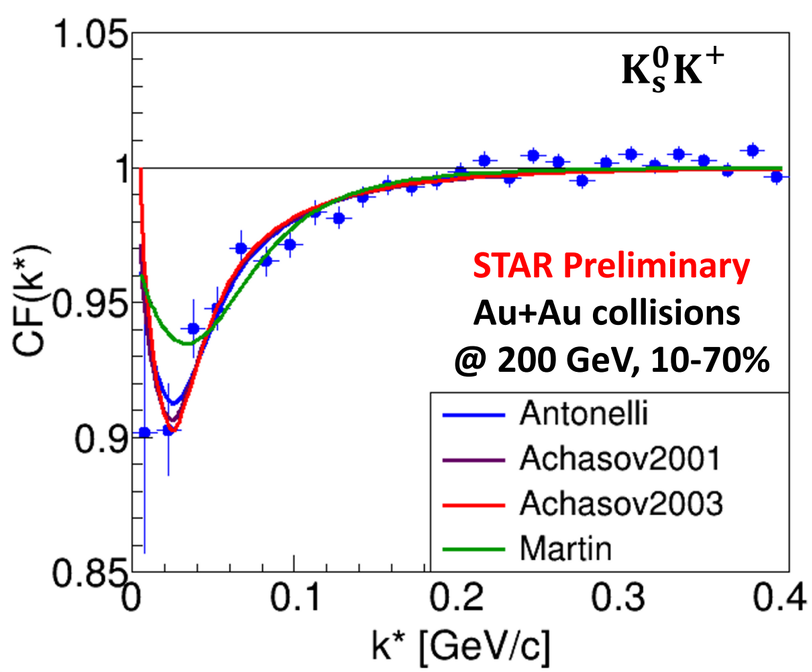
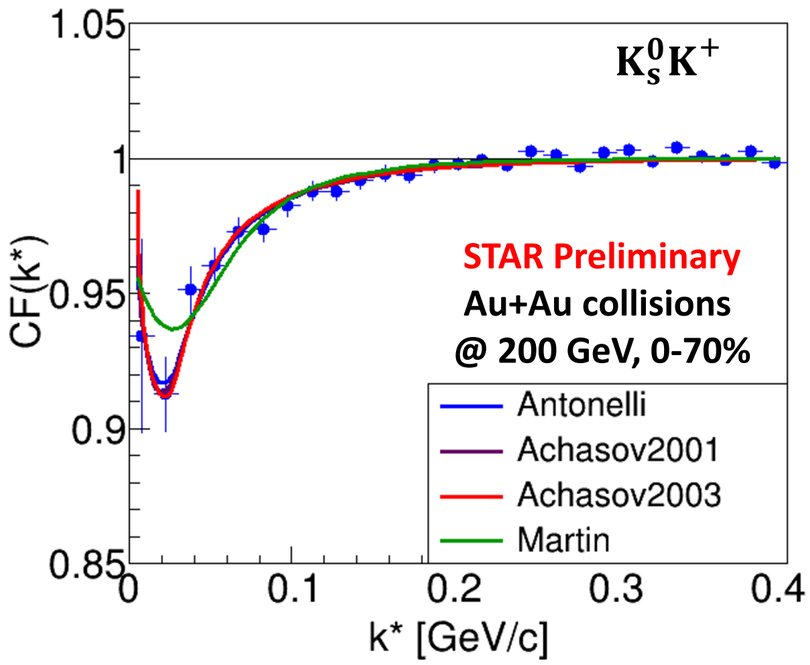
$K^0_S K^+$ femtoscopy at 200 GeV

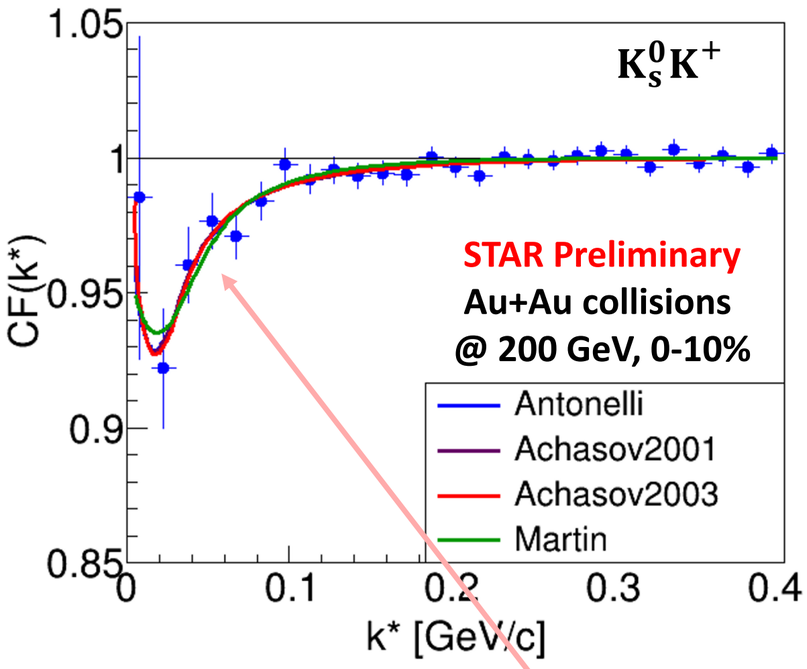


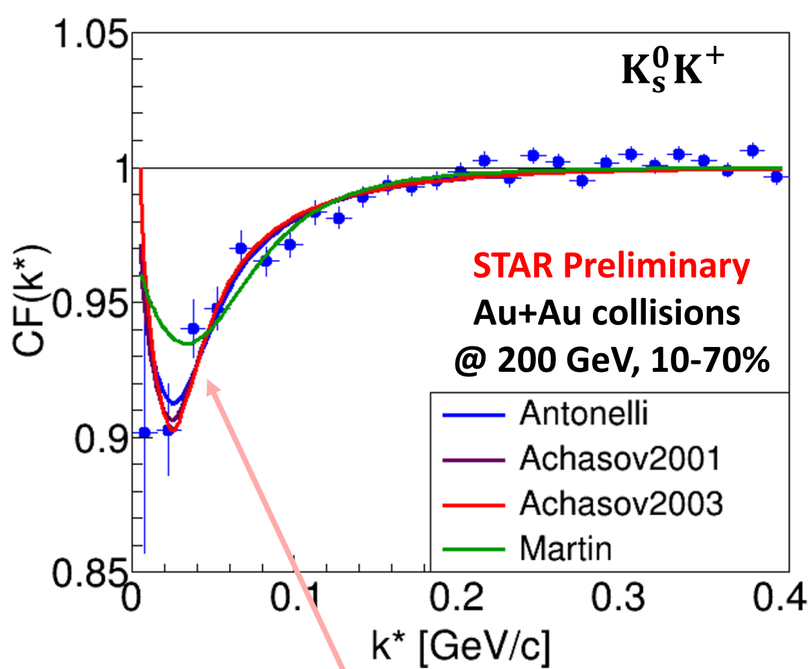


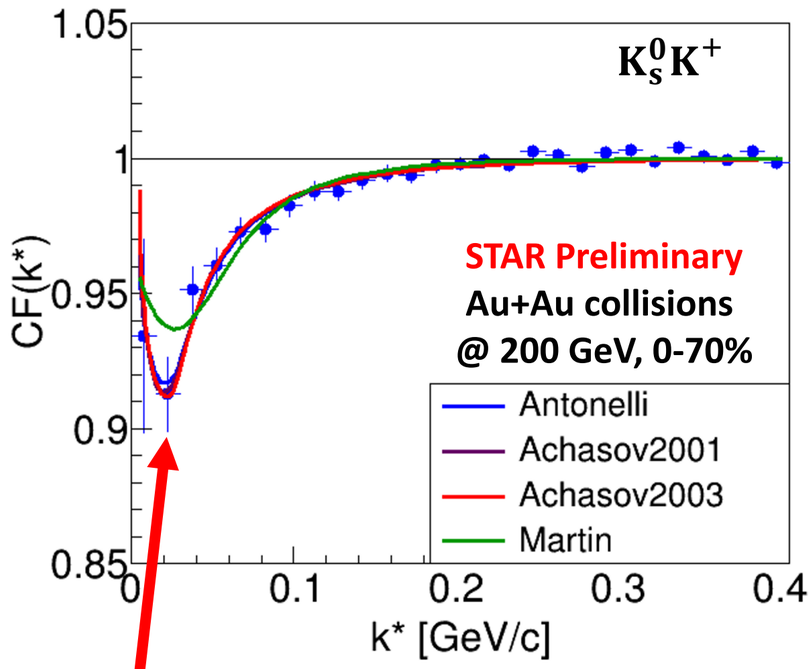


$k^* = |\overrightarrow{p_1}| = |\overrightarrow{p_2}|$

$K_s^0K^+$ femtoscopy at 200 GeV



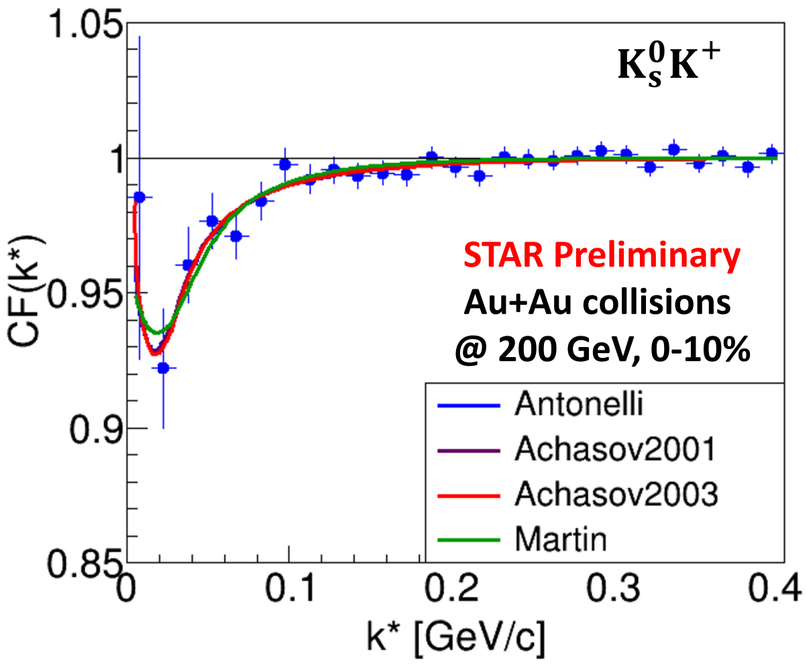


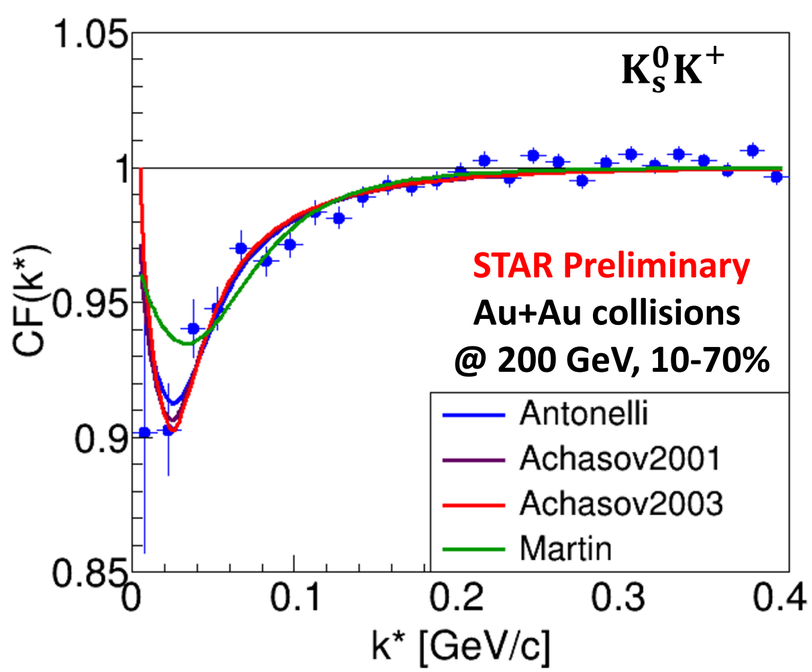


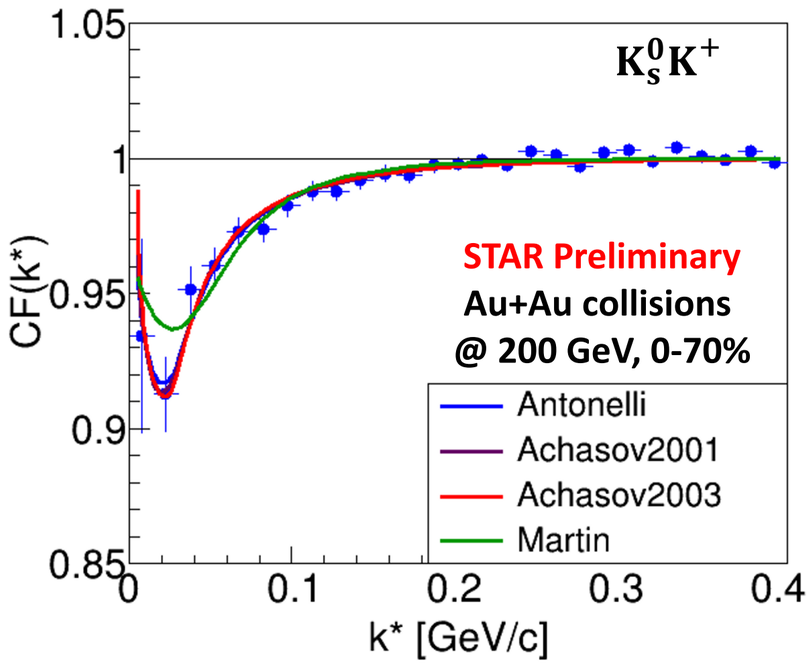
The a_0 FSI parametrization gives an excellent representation of the signal region of the data

$k^* = |\overrightarrow{p_1}| = |\overrightarrow{p_2}|$

$K_s^0K^+$ femtoscopy at 200 GeV





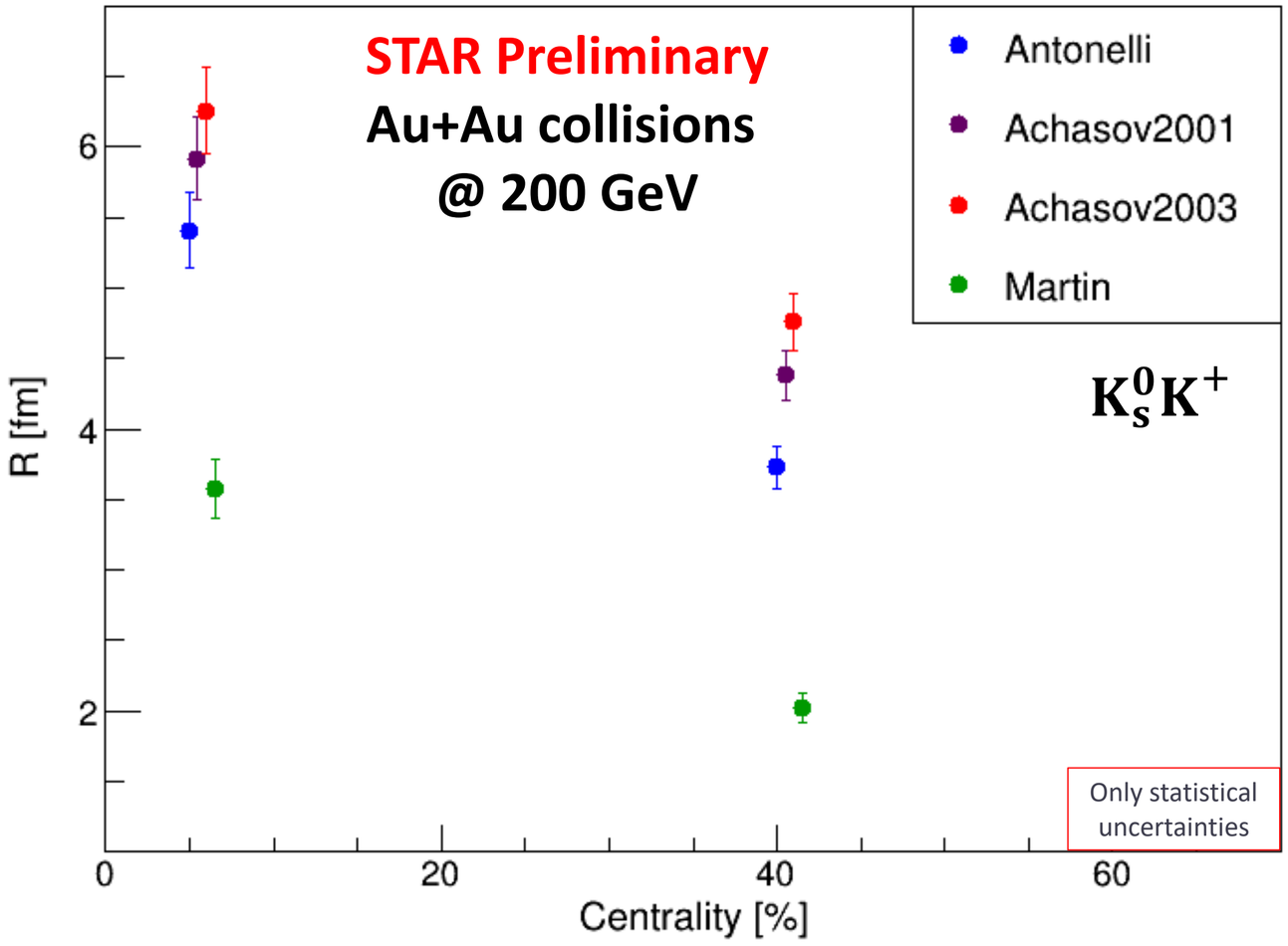


χ^2/NDF

	0-10%	10-70%	0-70%
Antonelli [1]	0.60	1.66	1.04
Achasov2001 [2]	0.59	1.73	1.07
Achasov2003 [3]	0.58	1.85	1.14
Martin [4]	0.65	1.65	1.16

$k^* = |\overrightarrow{p_1}| = |\overrightarrow{p_2}|$

$K_s^0K^+$ femtoscopy – centrality dependence



- Visible centrality dependence $R_{0-10\%} > R_{10-70\%}$
- **Achasov2003** pametrization (the larger a_0 mass) gives the larger size of the souce

Comparison - $K^0_S K^0_S$ and $K^0_S K^+$ - 200 GeV

Antonelli favors a_0 resonance as a tetraquark

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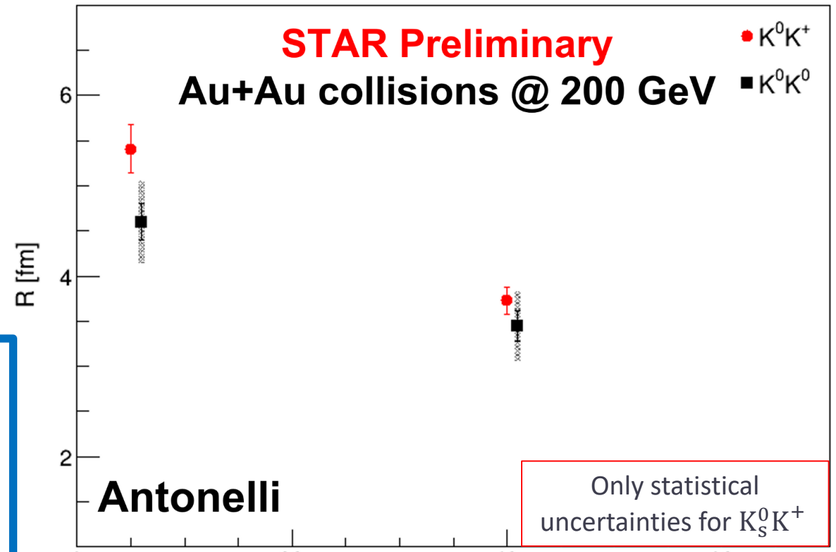

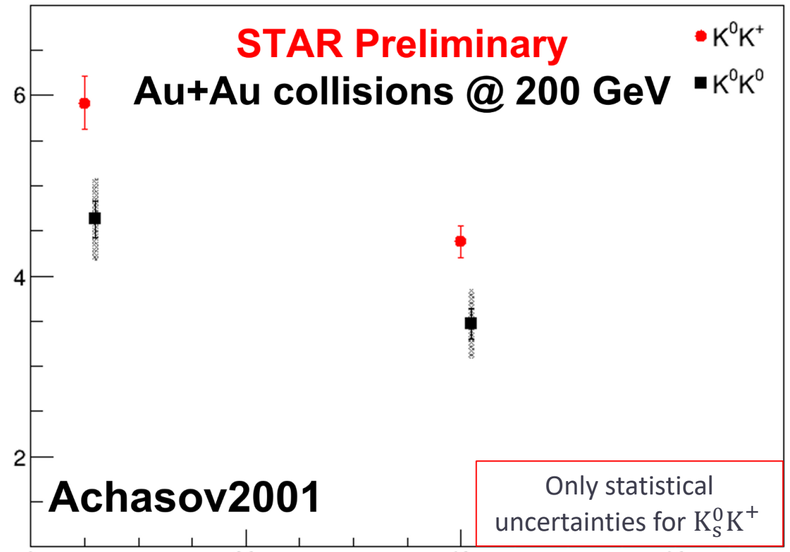

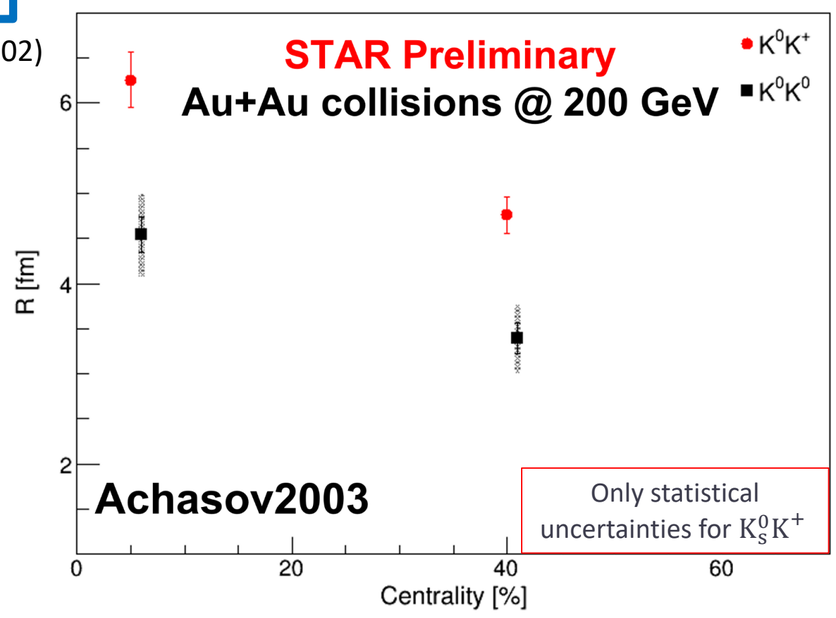

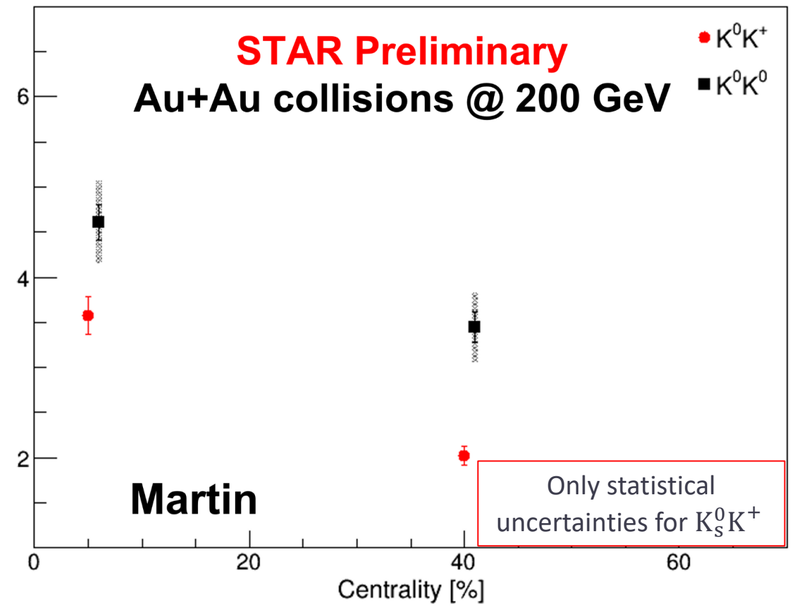

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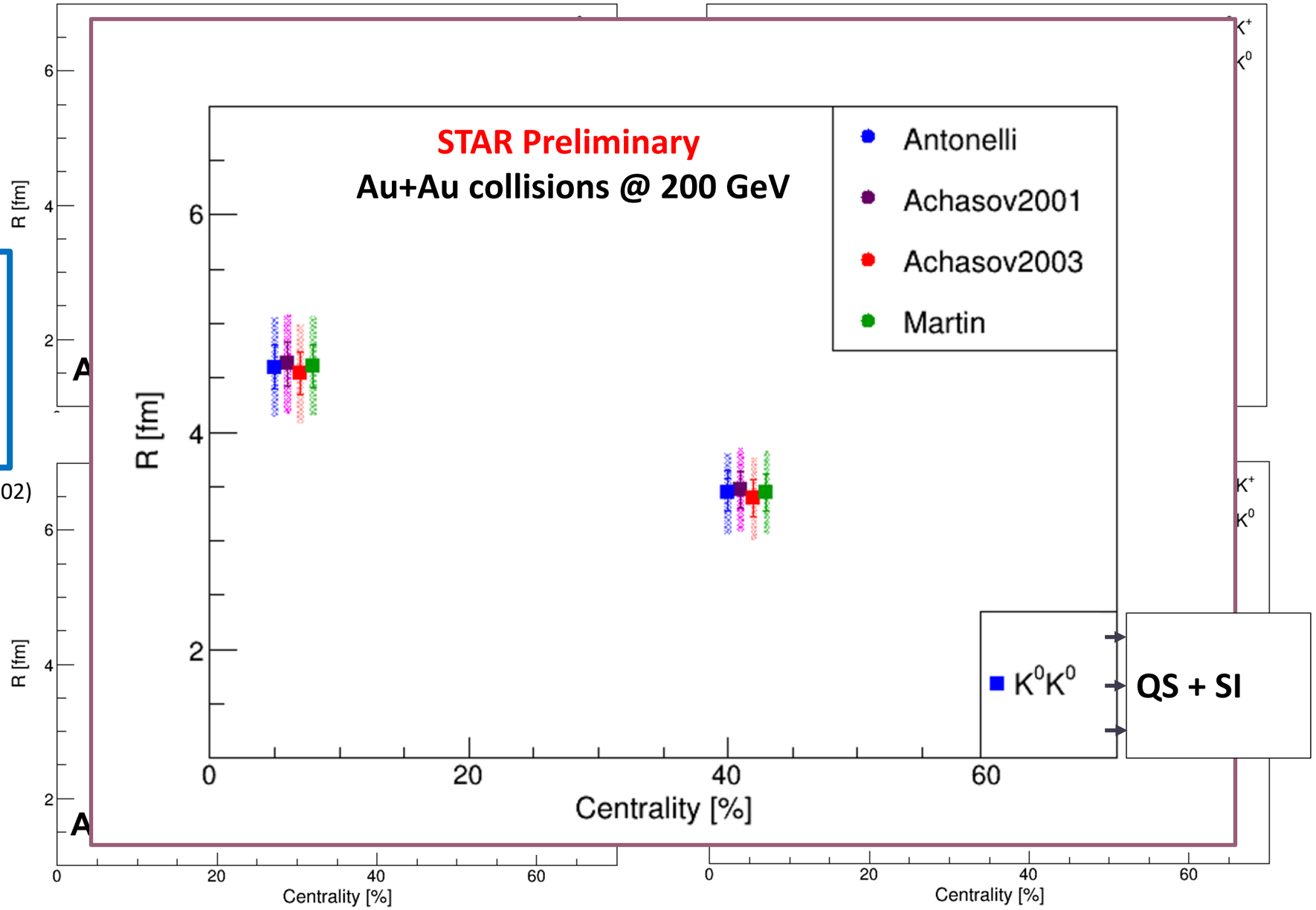

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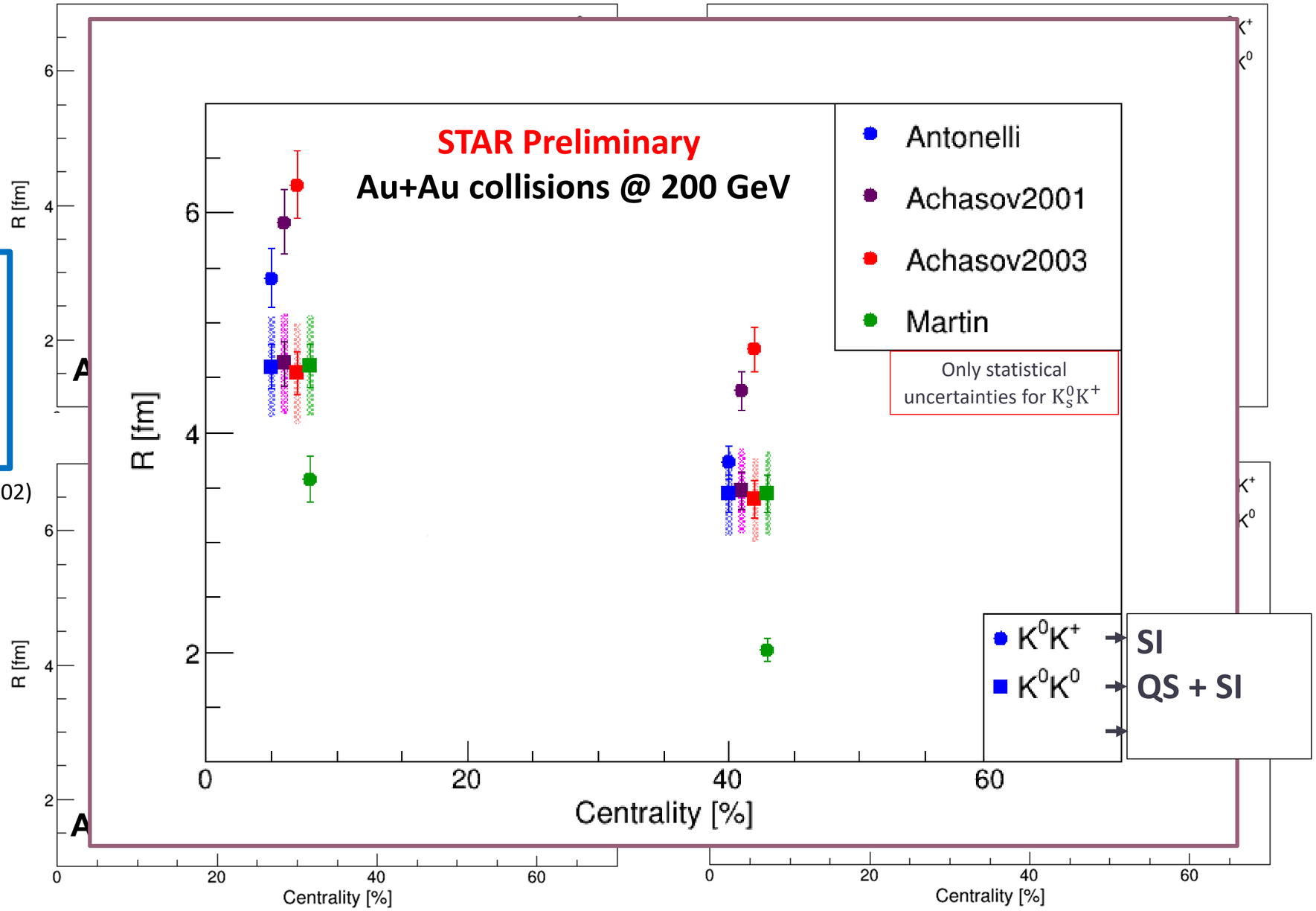

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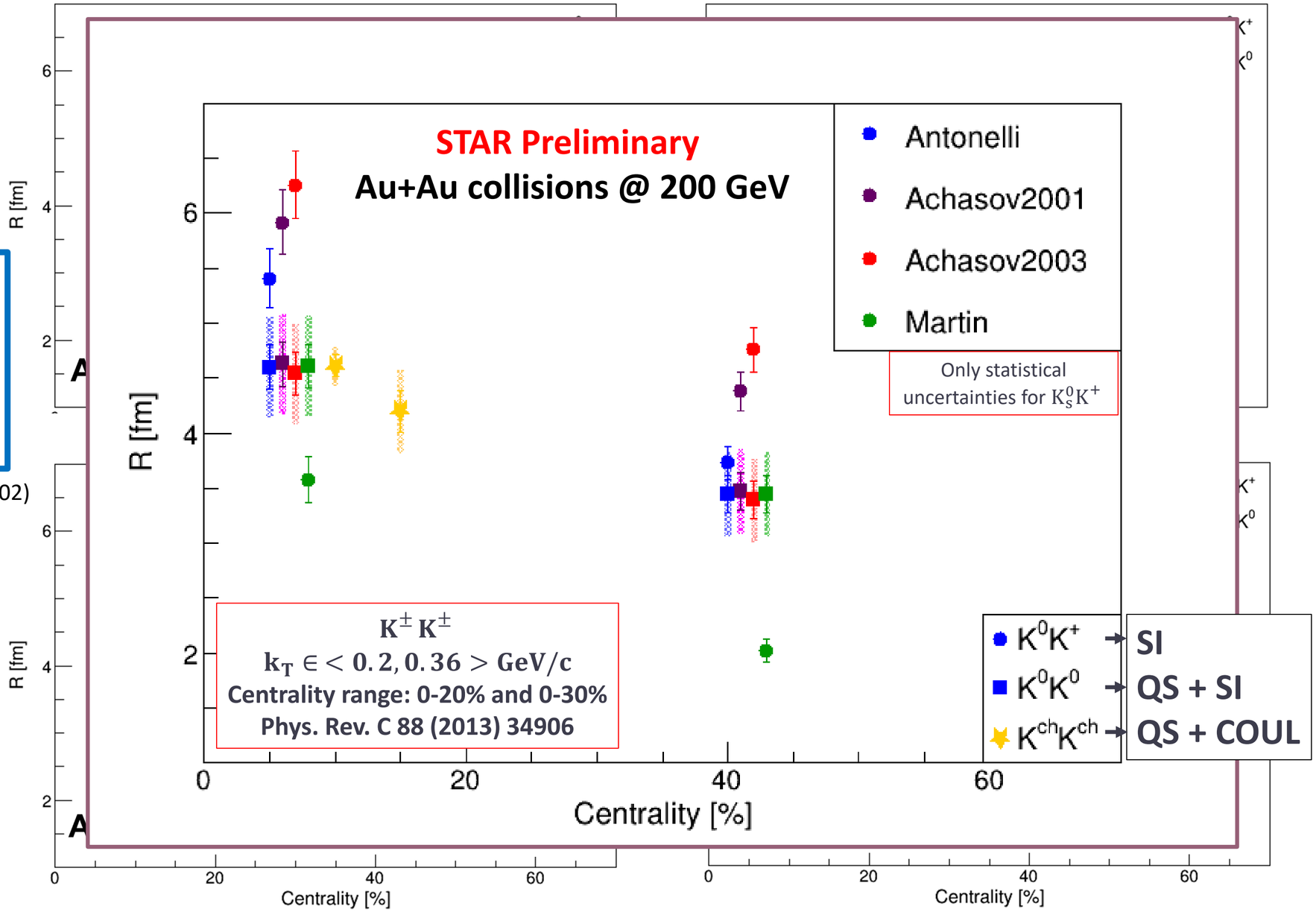

$$R_{K^{\pm}K^{\pm}} = \sqrt{\frac{R^2_{out} + R^2_{side} + R^2_{long}}{3}}$$

Summary

$\mathbf{K_s^0 K_s^0}$ femtoscopy

- The strong final-state interaction has a significant effect on the $\mathrm{K_s^0 K_s^0}$ correlation due to the near-threshold $\mathrm{f_0(980)}$ and $\mathrm{a_0(980)}$ resonances
- The radii of the source depend on centrality and increase with increasing collision energy
- Extracted source radii are comparable to these from models

$\mathbf{K_s^0 K^+}$ femtoscopy

- The $\mathrm{a_0(980)}$ FSI parametrization gives very good representation of the shape of the signal region in CF
- The parametrization with the larger $\mathrm{a_0(980)}$ mass and decay coupling gives larger size of the source
- Comprarison with $\mathrm{K_s^0 K_s^0}$
  - Antonelli parametrization favors $\mathrm{a_0(980)}$ resonance as a tetraquark

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$K_s^0K_s^0$ femtoscopy

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$K_s^0K^+$ femtoscopy

- The $a_0(980)$ FSI parametrization gives very good representation of the shape of the signal region in CF
- The parametrization with the larger $a_0(980)$ mass and decay coupling gives larger size of the source
- Comprarison with $K_s^0K_s^0$
  - Antonelli parametrization favors $a_0(980)$ resonance as a tetraquark

Summary

$K_s^0K_s^0$ femtoscopy

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$K_s^0K^+$ femtoscopy

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- The parametrization with the larger $a_0(980)$ mass and decay coupling gives larger size of the source
- Comprarison with $K_s^0K_s^0$
  - Antonelli parametrization favors $a_0(980)$ resonance as a tetraquark

Compatibility of source sizes for $K_s^0K_s^0$, $K_s^0K^+$ and $K^\pm K^\pm$ pairs in the case of Antonelli

Summary

$K_s^0K_s^0$ femtoscopy

- The strong final-state interaction has a significant effect on the $K_s^0K_s^0$ correlation due to the near-threshold $f_0(980)$ and $a_0(980)$ resonances
- The ra… …se with increas…
- Extract…

$K_s^0K^+$ fem…

- The a_0… …n of the shape …
- The pa… …oupling gives larger size of the source
- Comprarison with $K_s^0K_s^0$
  - Antonelli parametrization favors $a_0(980)$ resonance as a tetraquark

Outlook

High statistic BES-II data give the ability to analyze neutral kaon femtoscopy for energies $\sqrt{s_{NN}}$ < 20 GeV

Compatibility of source sizes for $K_s^0K_s^0$, $K_s^0K^+$ and $K^\pm K^\pm$ pairs in the case of Antonelli

Thank you for your attention!!