

Neutral kaon femtoscopy in STAR*

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Properties of nuclear matter can be studied by relativistic heavy-ion collisions in high-energy experiments like STAR. One of the methods to learn about properties of nuclear matter is femtoscopy, which relies on information carried by particles produced in the collisions. Using femtoscopic observables, space-time characteristics of the source can be extracted. During heavy-ion collisions mostly pions are produced, and therefore pion femtoscopy is a particularly useful tool. High statistics data sets from RHIC have also made it possible to study the strange particle correlations. The lightest strange hadrons are charged and neutral kaons. The strong interaction, which conserves the strangeness quantum number, is responsible for the kaon production. It is possible to study the neutral kaons, K_S^0 , which can be measured through their decay products to charged pions.

In these proceedings, one-dimensional correlation functions of neutral kaon pairs in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV measured by the STAR experiment at RHIC are presented.

1. Introduction

To study the Quantum Chromodynamics (QCD) phase diagram, the comprehensive program called Beam Energy Scan (BES) was started at Relativistic Heavy Ion Collider (RHIC) in 2010 [1]. The Au+Au collisions at $\sqrt{s_{NN}} = 7.7, 11.5, 14.5, 19.6, 27, 39$ and 62.4 GeV were collected during the first phase of the BES program. The main aims of the BES program include:

- to study the phase diagram at different energies and find areas in which QGP signatures are turned off,
- to search for the first-order phase transition and a critical point,
- to study first order phase transition between Hadron Gas (HG) and Quark-Gluon Plasma (QGP).

* Presented at XIV Workshop on Particle Correlation and Femtoscopy

32 The Solenoidal Tracker at RHIC (STAR) during the first phase of BES
 33 program gathered data, which were used for femtoscopic analysis. This
 34 method uses measurements of momentum of emitted particles in order to
 35 study properties of system created in heavy-ion collision.

36 1.1. The method of femtoscopy

37 In 1954, Robert Hanbury Brown
 38 and Richard Q. Twiss created the
 39 method to measure the angular
 40 sizes of the astronomical objects
 41 using the Michelson interferometry [2]. Femtoscopy method ap-
 42 plied to two particles originates
 43 from the HBT technique used in as-
 44 tronomy. It aims to examine the
 45 particle emitting source (sizes of or-
 46 der 10^{-15} m and life times 10^{-23} s)
 47 through measurements of particle's
 48 momentum distributions. Figure 1
 49 shows the two-particle interference mechanisms used in astronomy (left) and
 50 particle physics (right).
 51

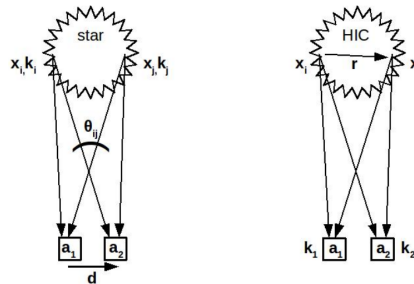


Fig. 1: Schematic diagram of intensity interferometry used in astronomy (left) and particle physics (right)[3].

52 1.2. Correlation function

53 The correlation function (CF) is described as a ratio between probability
 54 of observing two particles with momenta \vec{p}_1 and \vec{p}_2 (P_2) and the probability
 55 of observing these two particles separately (where P_1 is the probability of
 56 observation of single particle):

$$CF(\vec{p}_1, \vec{p}_2) = \frac{P_2(\vec{p}_1, \vec{p}_2)}{P_1(\vec{p}_1)P_1(\vec{p}_2)} \quad (1)$$

57 Theoretical correlation function is determined by the emission function,
 58 $S(\vec{q}, \vec{r})$, which contains all space-time characteristics of the effective source
 59 and pair wave function, $\psi(\vec{q}, \vec{r})$, which includes information about statistical
 60 effects and interactions:

$$CF(\vec{q}) = \int d^3r S(\vec{q}, \vec{r}) |\psi(\vec{q}, \vec{r})|^2 \quad (2)$$

61 where \vec{q} is a difference between momenta \vec{p}_1 and \vec{p}_2 , and \vec{r} is a difference
 62 between the position of the first and second particles in the pair. Through
 63 measurements of quantum statistical effects and interactions one can learn
 64 source parameters, like its size [4].

65 Experimentally, in the one-dimensional case, the correlation function is
 66 defined as a ratio of signal (A) to background (B), where signal measures
 67 relative momentum distribution (q_{inv}) of pairs from the same collision and
 68 background measures of pairs from different collisions with similar prop-
 69 erties. For pairs constructed from different events the quantum statistical
 70 effects and final state interactions are absent:

$$CF(q_{inv}) = \frac{A(q_{inv})}{B(q_{inv})} \quad (3)$$

71 The q_{inv} is define as [5]:

$$q_{inv} = \sqrt{(p_1 - p_2)^2 - (E_1 - E_2)^2} \quad (4)$$

72 where E_1 and E_2 are the energies of the 1st and 2nd particle from the pair
 73 respectively.

74 2. Neutral kaon correlation function

75 Figure 2 shows correlation functions for neutral kaons which de-
 76 pend on:
 77

- 78 • Quantum Statistical effects (QS) - Bose-Einstein statistics, which increases the prob-
 79 ability of finding two particles with similar momentum
- 80 • Final State Interactions (FSI):
- 81 - Strong Interaction (SI)
- 82 - Coulomb Force (COUL) in the case of neutral kaons it is absent
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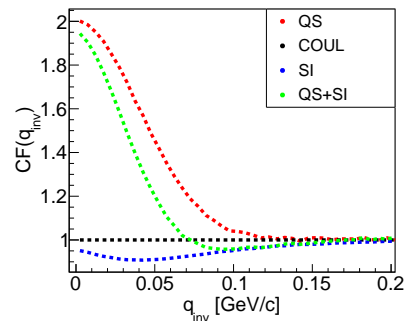


Fig. 2: Correlation function, $CF(q_{inv})$, of the kaon pairs from Therminator model [6] for central Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV.

88 2.1. Parametrization

89 The parametrization was done using Gaussian fit (Eq. 5) taking into
 90 account the QS effects only, where λ is the correlation strength and R_{inv} is
 91 the one-dimensional radius of the particle-emitting source:

$$C(q_{inv}) = 1 + \lambda \exp[R_{inv}^2 q_{inv}^2] \quad (5)$$

3. Results

3.1. Selection criteria

In this analysis, minimum-bias (centrality corresponding to 0 – 80% of the total hadronic cross section of the collision) Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV were used. Kaon identification was performed using information from two detectors, namely the Time Projection Chamber (TPC) and the Time-Of-Flight (TOF). Identification of particles in TPC was done via specific ionization losses (dE/dx). TOF determines velocity of a particle based on its time of passage through the length of the detector. Using information from TOF detector, particles with momentum above threshold are uniquely identified. Table 1 shows criteria for particle's selection for π^+ and π^- (daughter particles of the K_s^0 candidates). Criteria for neutral kaons are presented in Table 2.

Table 1: Daughter criterias of selection

p_T [GeV/c]	0.2-1.2
DCA to the primary vertex [cm]	>1.3
$ n_{\sigma\pi} $	<3
$ n_{\sigma K,p} $	>3
mass [GeV/c ²]	0.07-0.2

Table 2: Neutral kaon selection criteria

p_T [GeV/c]	0.2-1.5
$ \eta $	<0.5
V^0 DCA to the primary vertex [cm]	0-0.3
DCA of daughters [cm]	0-0.3
decay length [cm]	>2
Armenteros q_T [GeV/c]	0.12-0.22
Armenteros $ \alpha $	<0.7
invariant mass range [GeV/c ²]	0.488-0.51

3.2. Correlation function

Figure 3 shows $K_s^0 K_s^0$ correlation functions with Gaussian fit (Eq. 5) before and after applying the pair purity correction. The purity correction was done using the equation:

$$C_{corrected}(q_{inv}) = \frac{C_{measured}(q_{inv}) - 1}{PairPurity(q_{inv})} + 1 \quad (6)$$

where PairPurity is defined as a product of single particle purity. Purity of K_s^0 is calculated as a ratio of number of K_s^0 candidates in the specic range of distribution of invariant mass (Table 2) to number of all considered $\pi^- \pi^+$ pairs in this range.

113 After purity correction a stronger correlation (larger value of the λ pa-
 114 rameter and smaller size of the source) as compared to one before applying
 115 purity correction is observed. Values of radii before and after correction are
 presented in Table 3 and Table 4.

Table 3: Radius of the particle-emitting source before purity correction.

Centrality	Radius [fm]	λ
0-80%	5.08 ± 0.19	0.630 ± 0.051

Table 4: Radius of the particle-emitting source after purity correction.

Centrality	Radius [fm]	λ
0-80%	4.72 ± 0.20	0.701 ± 0.056

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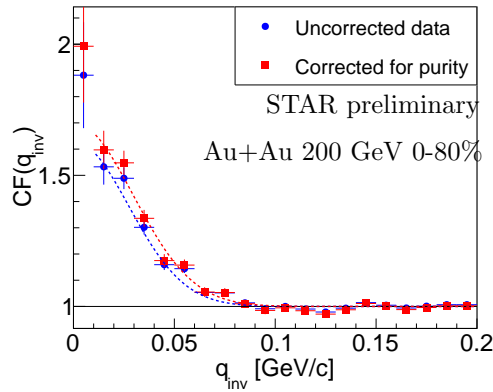
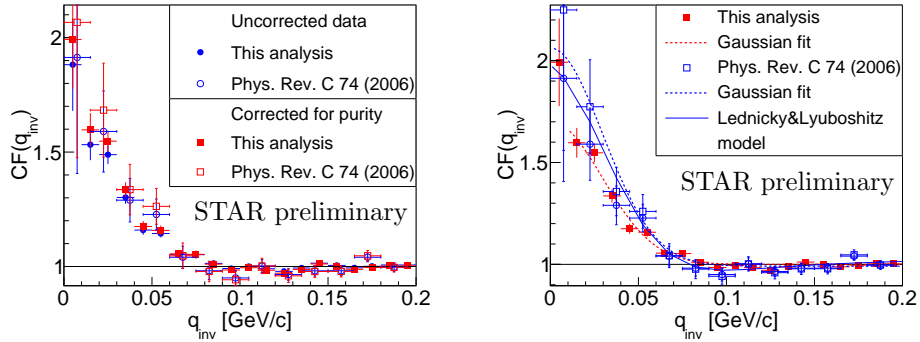


Fig. 3: The $K_S^0 K_S^0$ correlation functions with fits before (blue points) and after (red points) purity corrections. The uncertainties of CF are statistical only.



(a) The $K_S^0 K_S^0$ correlation function from central Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV.

(b) Fits to the $K_S^0 K_S^0$ correlation function.

Fig. 4: Comparison with previous STAR results.

117 *3.3. Comparison with previous result*

118 The first statistically meaningful results from $K_S^0 K_S^0$ femtoscopy in Au+Au
 119 collisions at collision energy $\sqrt{s_{NN}} = 200$ GeV were published by the STAR
 120 Collaboration in 2006 [7]. In that analysis, data from the Zero Degree
 121 Calorimeter (ZDC) and the Central Trigger Barrel (CTB) were used. The
 122 radius of the particle emitting source was obtained using Gaussian fit as
 123 5.02 ± 0.61 fm.

124 Figure 4 presents the comparison between previous results with recent
 125 ones. Both measurements: previous and current ones are consistent with
 126 each other within estimated uncertainties.

127 **4. Summary**

128 In these proceedings, one-dimensional neutral kaon correlation functions
 129 measured by the STAR experiment for minimum-bias (0-80%) Au+Au col-
 130 lisions at $\sqrt{s_{NN}} = 200$ GeV have been presented.

131 To parametrize the experimental correlation functions Gaussian fits were
 132 used, which gave similar values of the radii of the effective source and λ
 133 parameter as in previous STAR results. The purity correction for neutral
 134 kaons slightly reduces the size of the measured particle-emitting source and
 135 increases the statistical uncertainties. As expected, value of λ parameter
 136 increased and source size decreased.

137 **5. ACKNOWLEDGMENTS**

138 This work was supported by the grant: UMO-2017/27/B/ST2/01947
 139 supported by the National Science Centre, Poland (NCN).

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