

Measurements of low- p_T e^+e^- pairs and J/ψ in heavy-ion collisions at STAR

Shuai Yang (for the STAR Collaboration)

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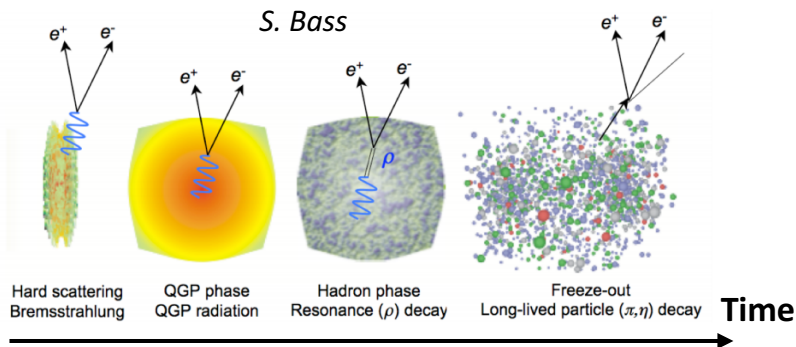
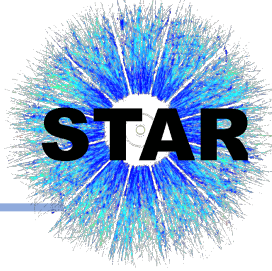


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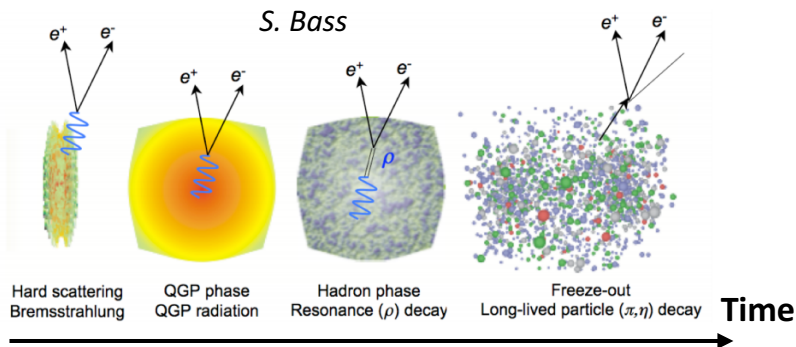
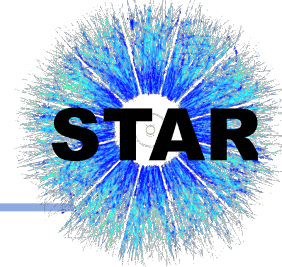


Dileptons - electromagnetic probe

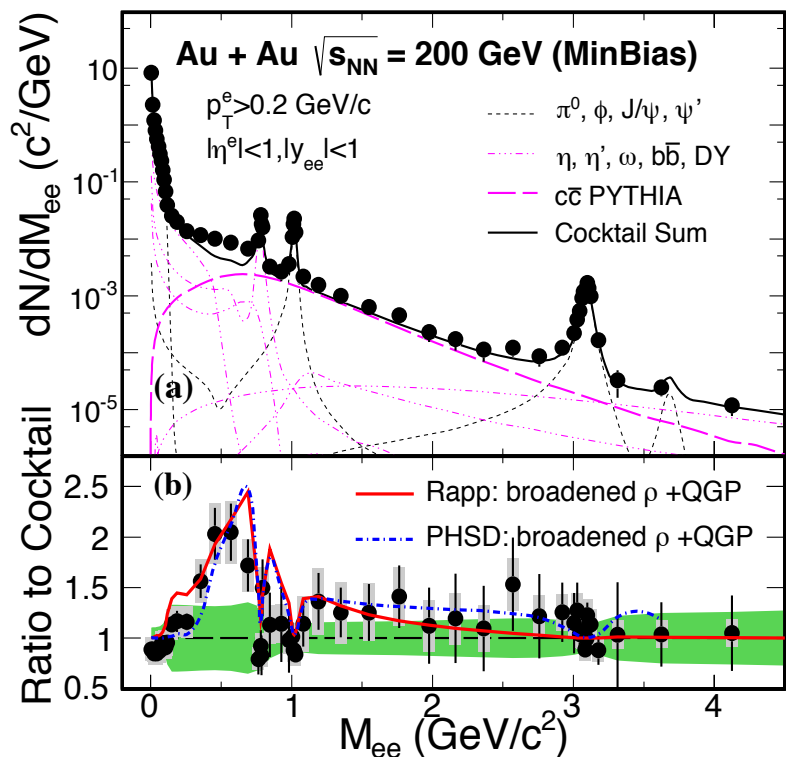


- Penetrating probe
- Direct information about the medium created in heavy-ion collisions

Dileptons - electromagnetic probe

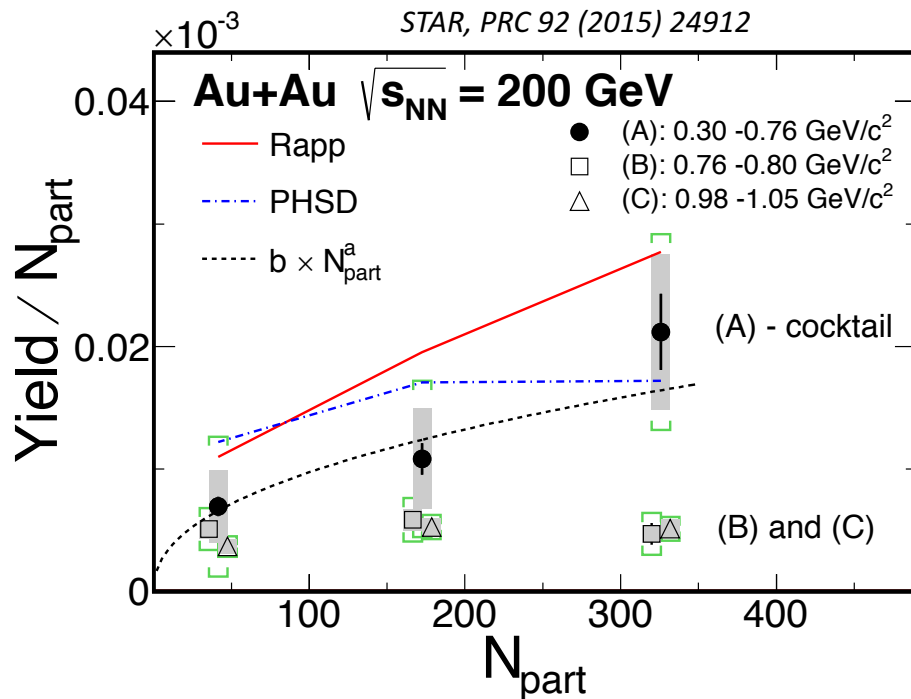


- Penetrating probe
- Direct information about the medium created in heavy-ion collisions



STAR, PRL 113 (2014) 022301

Shuai Yang

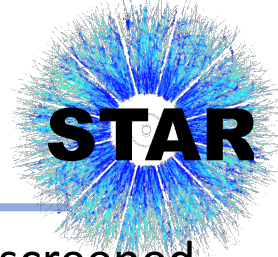


- Consistent with in-medium broadened ρ model calculation [R. Rapp, PRC 63 (2001) 054907]

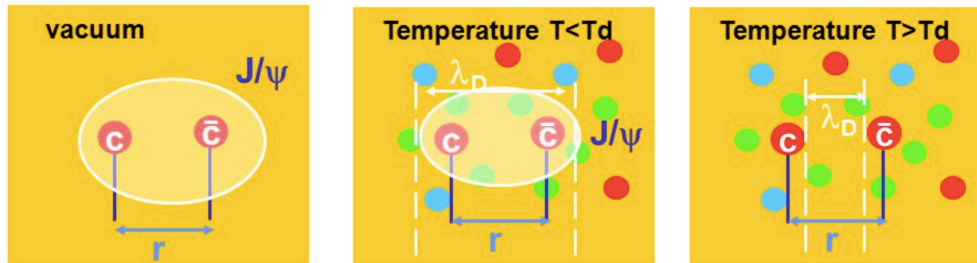
➤ $\propto N_{part}^{1.44 \pm 0.1}$, sensitive to hot medium dynamics

WWND2019, Beaver Creek

Quarkonia - heavy flavor probe

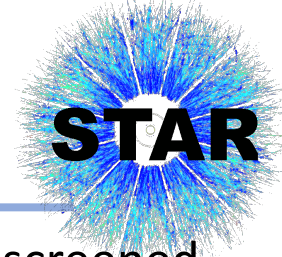


- **Color-screening in QGP:** the quark-antiquark binding potential is screened by the color charges of the surrounding light quarks and gluons -> *dissociation*
- J/ψ suppression was proposed a direct proof of QGP formation [T. Matsui and H. Satz, PLB 178 (1986) 416]

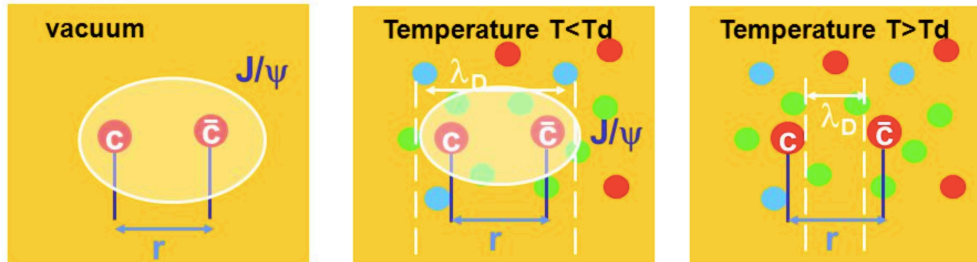


$$r_{q\bar{q}} \sim 1/E_{binding} > r_D \sim 1/T$$

Quarkonia - heavy flavor probe



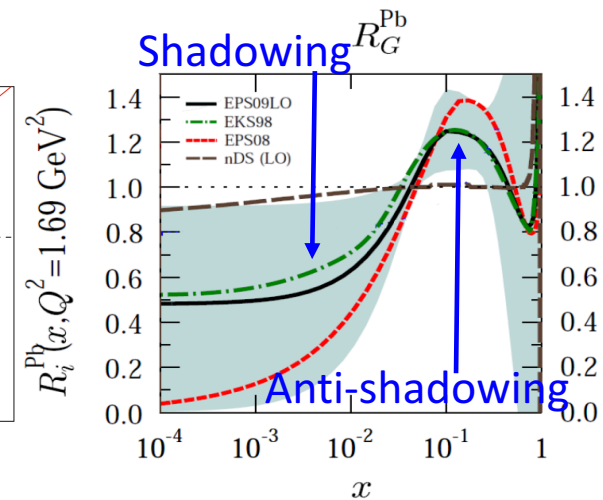
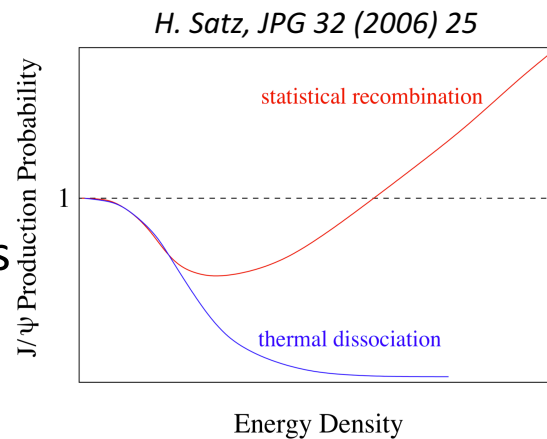
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$$r_{q\bar{q}} \sim 1/E_{binding} > r_D \sim 1/T$$

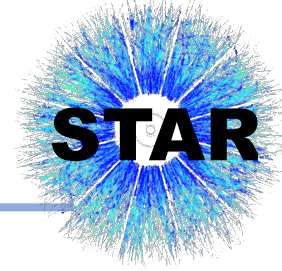
➤ Other effects

- Hot medium effects
 - ✓ regeneration, ...
- Cold nuclear matter effects
 - ✓ nPDF, ...
- Feed-down

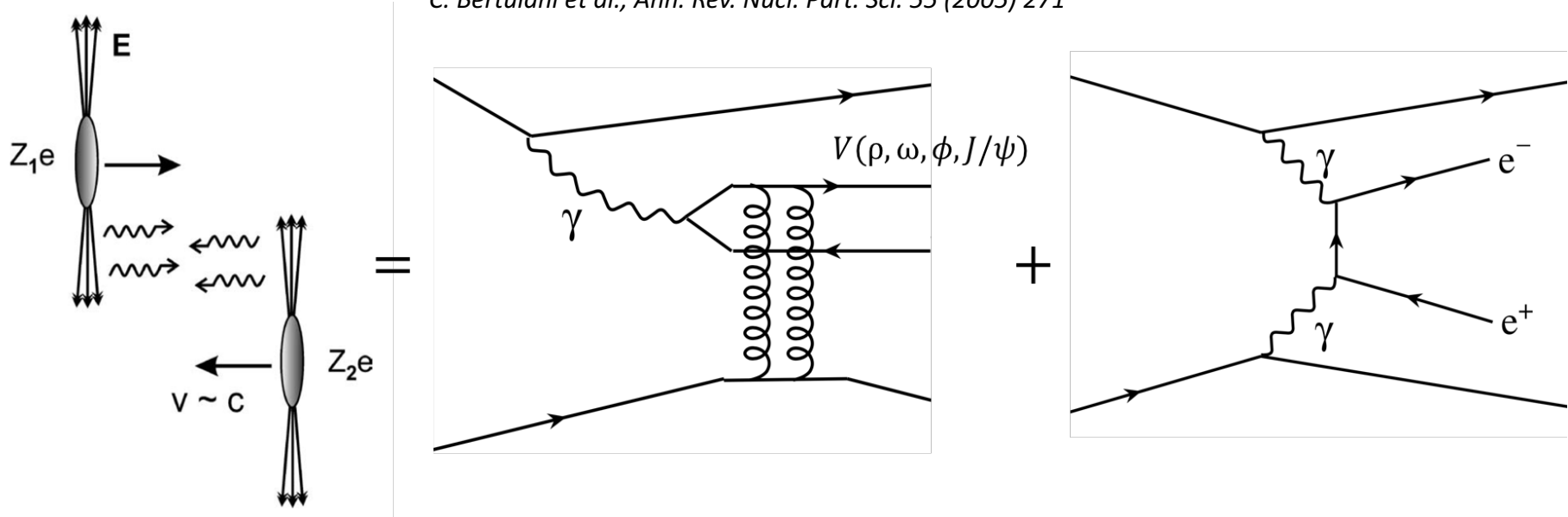


Ferreiro et al., PRC 81 (2010) 064911

Photon interactions

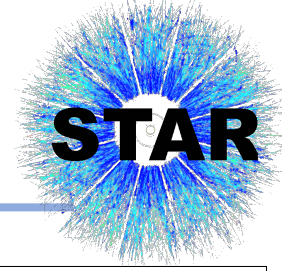


C. Bertulani et al., Ann. Rev. Nucl. Part. Sci. 55 (2005) 271



- Large quasi-real photon flux $\propto Z^2$
- Photon interactions
 - Photon-photon interaction (dilepton...) $\propto Z^4$
 - Photonuclear interaction (vector mesons) $\propto Z^2$
 - ✓ Coherent & Incoherent
- Conventionally only studied in UPC ($b > 2R_A$)

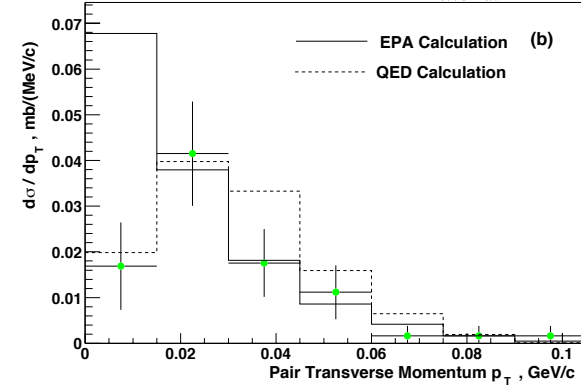
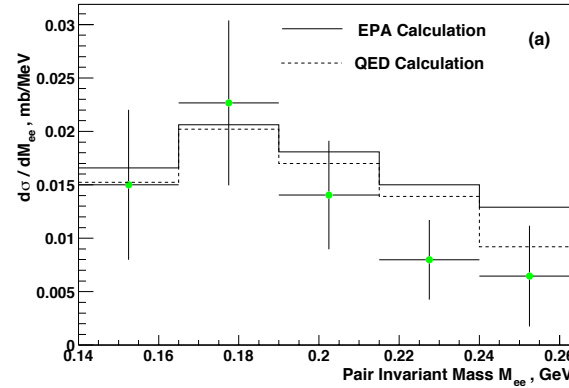
Features of photon interactions



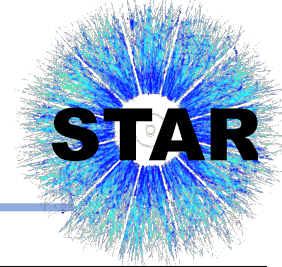
Two-photon interaction

- Continuous mass spectrum
- Concentrated at low p_T

STAR, PRC 70 (2004) 031902



Features of photon interactions



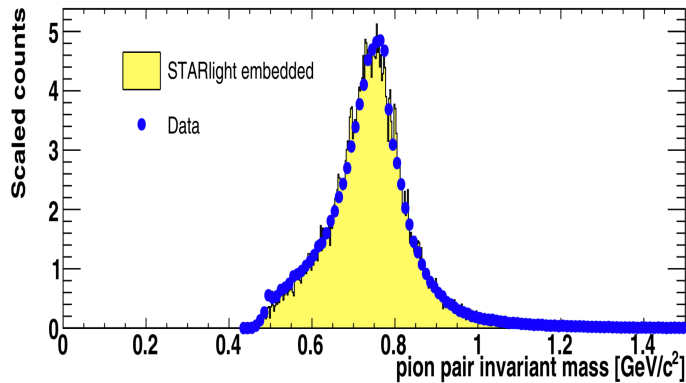
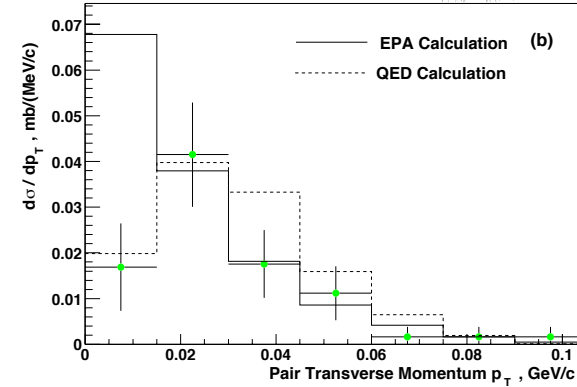
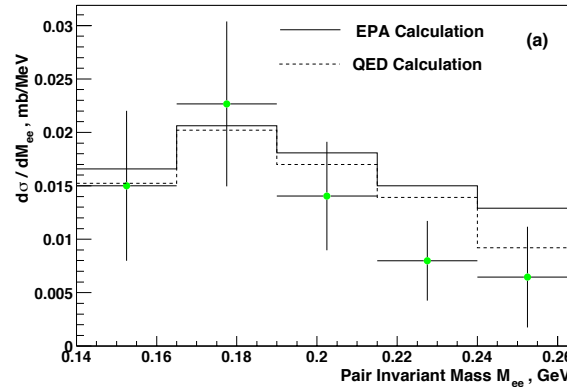
Two-photon interaction

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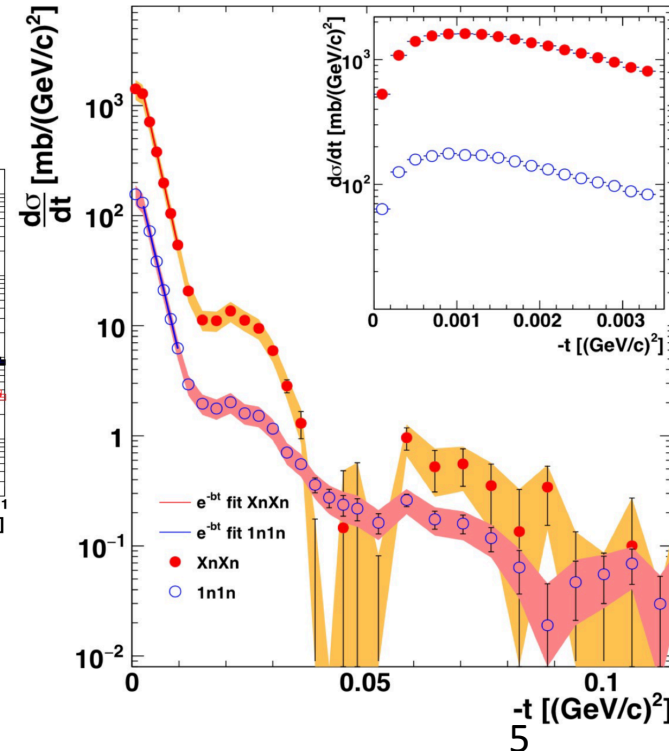
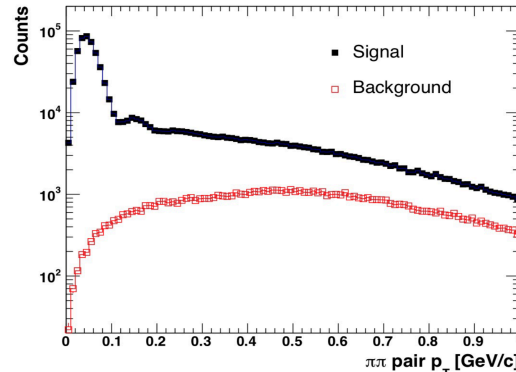
Coherent photonuclear interaction

- Vector meson production
- Concentrated at low p_T
- Interference structure ($p_T \ll 1/b$)

STAR, PRC 70 (2004) 031902

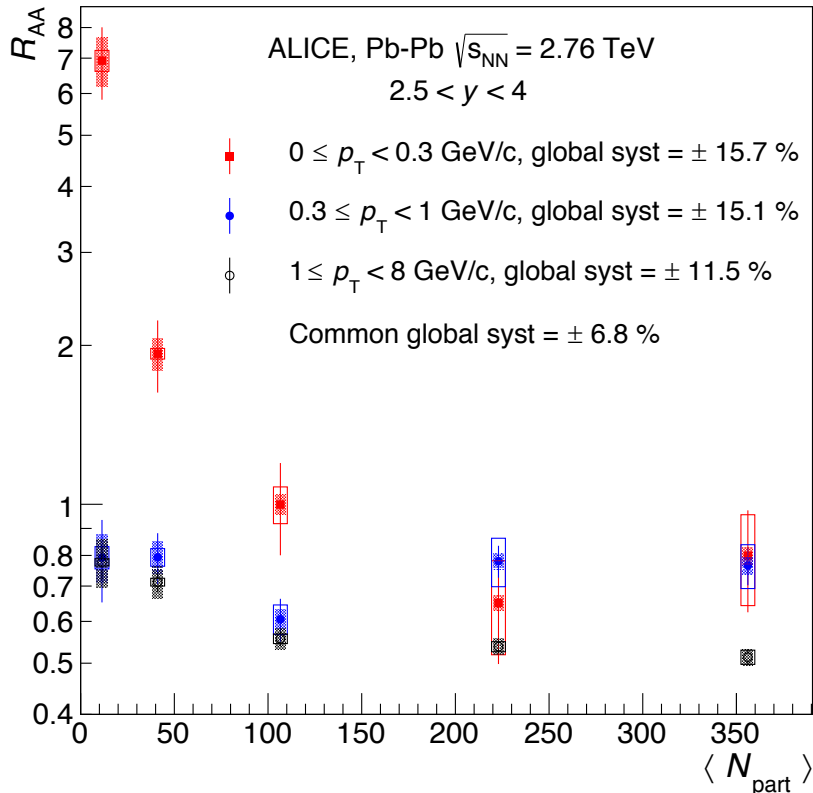


STAR, PRC 96 (2017) 54904



Anomalous J/ψ enhancement at LHC

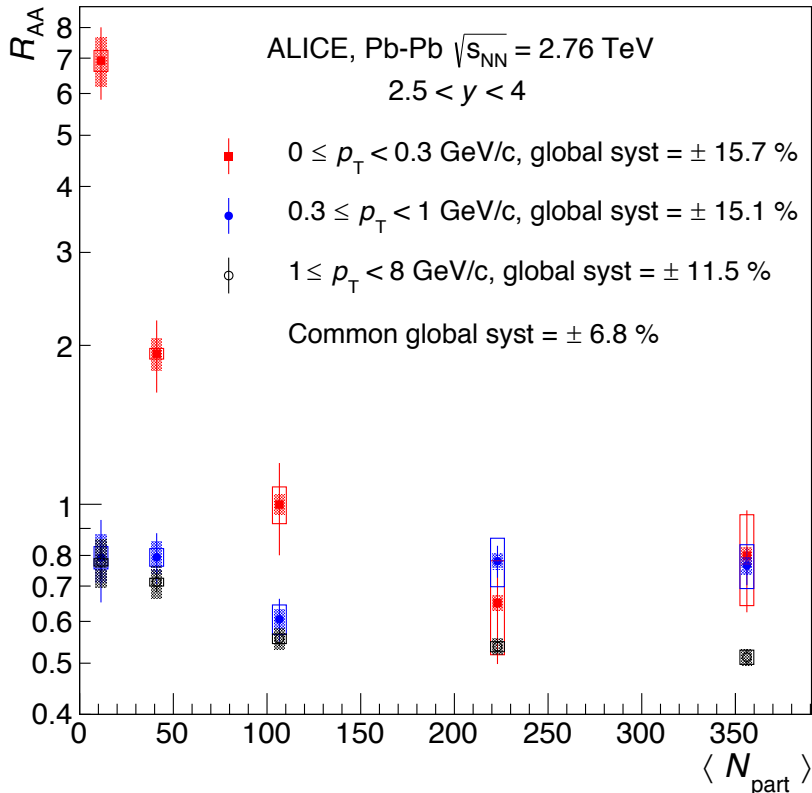
ALICE, PRL 116 (2016) 222301



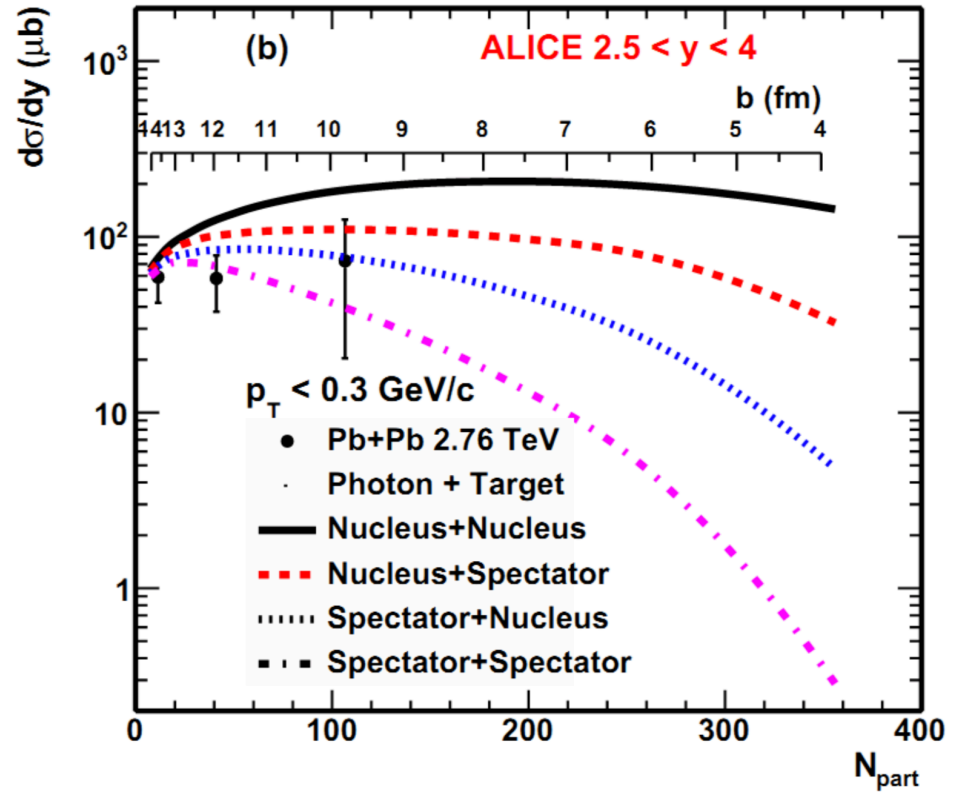
➤ Significant enhancement at low p_T in peripheral Pb+Pb collisions

Anomalous J/ψ enhancement at LHC

ALICE, PRL 116 (2016) 222301

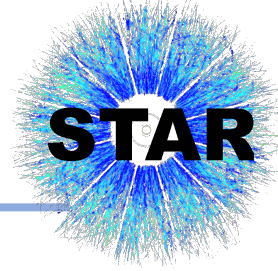


W. Zha et al., PRC 97 (2018) 044910

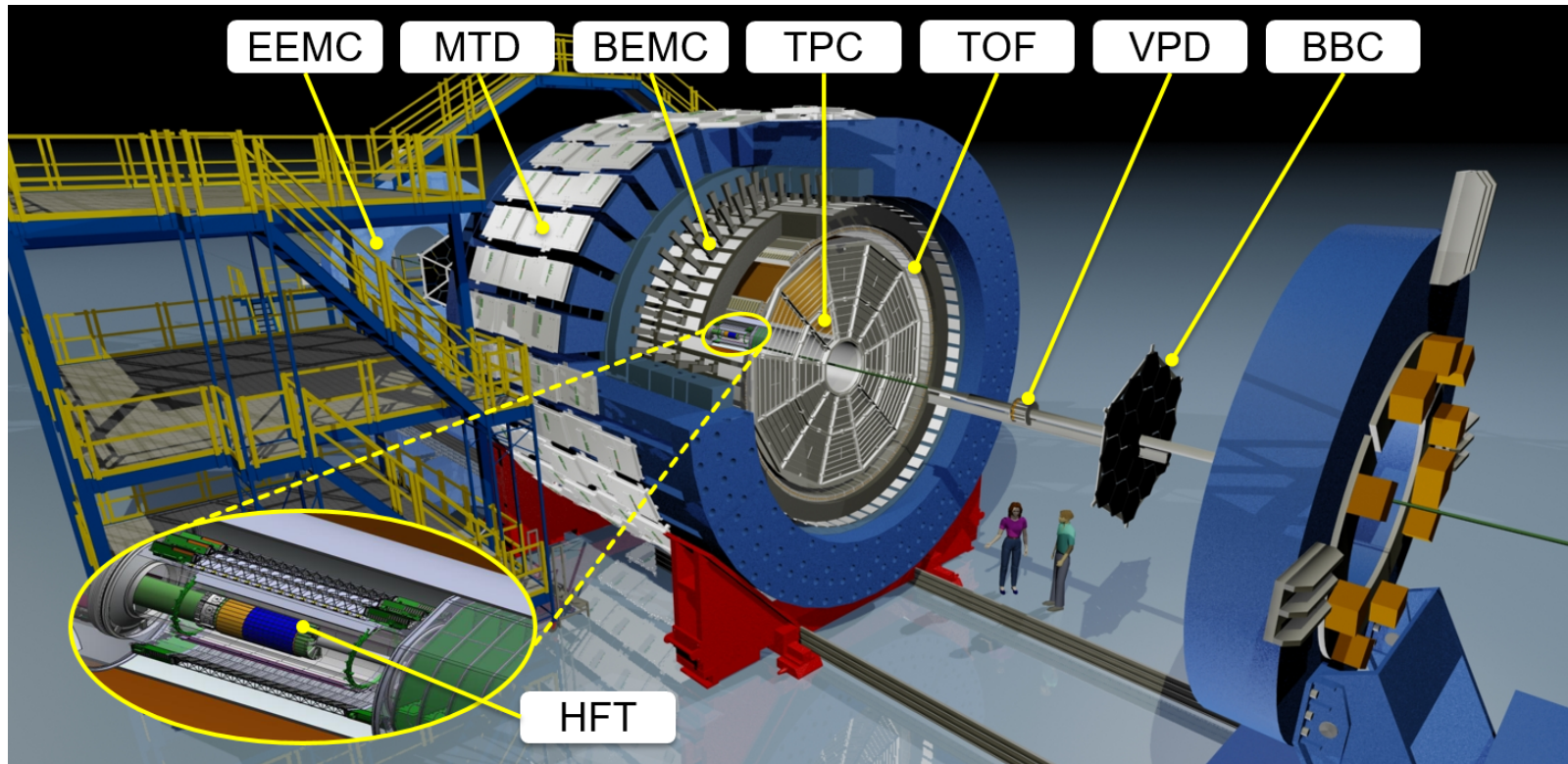


- Significant enhancement at low p_T in peripheral Pb+Pb collisions
- Qualitatively explained by coherent photonuclear production mechanism

The STAR detector



- Midrapidity, large acceptance: $|\eta| < 1, 0 < \phi < 2\pi$

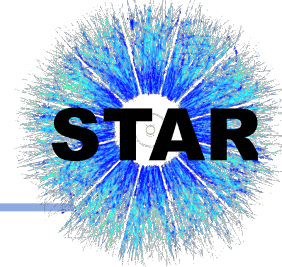


- Time Projection Chamber: tracking, momenta, and energy loss

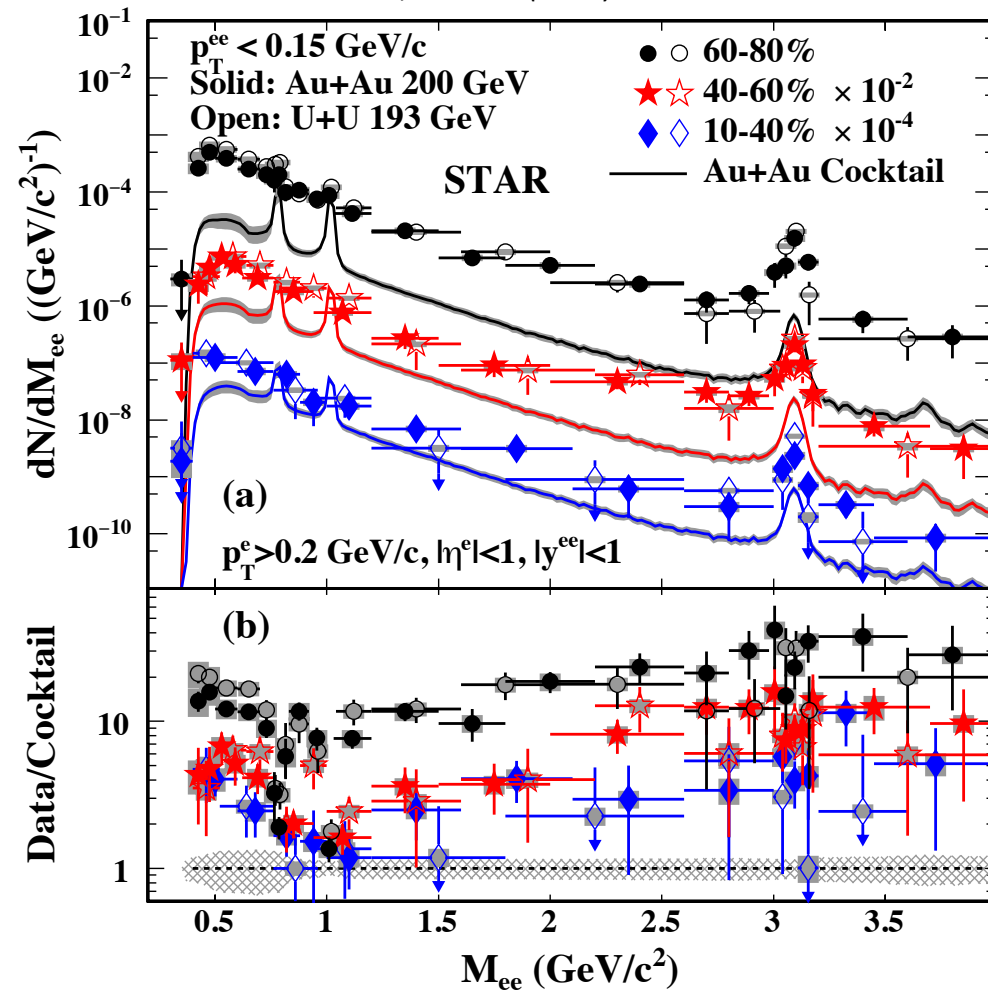
- Barrel Electromagnetic Calorimeter: trigger on and identify high- p_T electrons

- Time-Of-Flight: velocity

Low- p_T e^+e^- invariant mass spectra



STAR, PRL 121 (2018) 132301



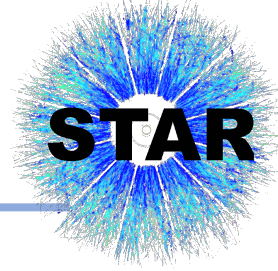
➤ Significant enhancement with respect to the cocktail for $p_T < 0.15$ GeV/c in 60-80% Au+Au and U+U collisions

Ratio of data/cocktail

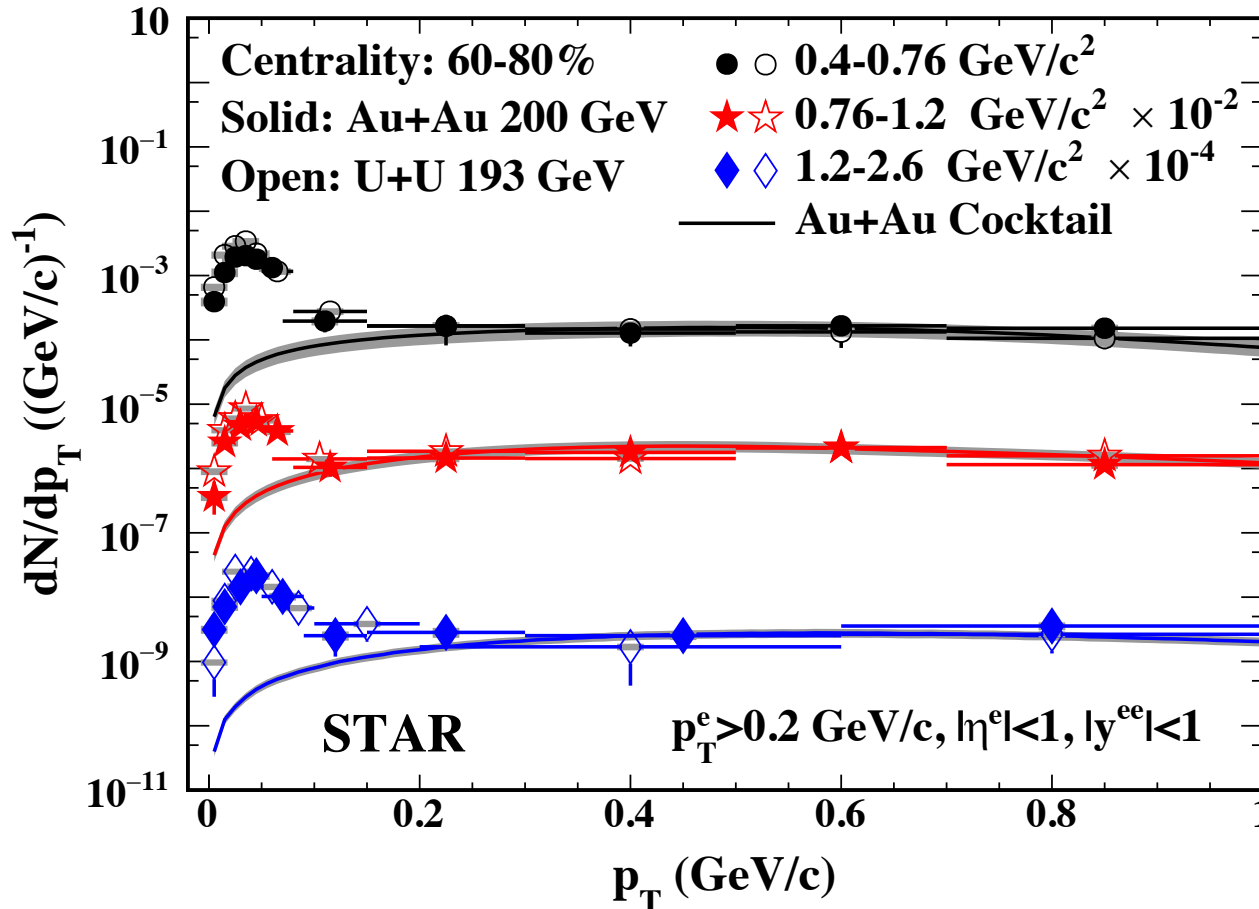
Mass (GeV/c^2)	Au+Au	U+U
0.4 - 0.76	$12.4 \pm 0.7 \pm 2.0 \pm 3.7$	$17.1 \pm 1.0 \pm 2.4 \pm 5.1$
0.76 - 1.2	$3.9 \pm 0.3 \pm 0.6 \pm 0.8$	$4.6 \pm 0.4 \pm 0.5 \pm 0.9$
1.2 - 2.6	$12.6 \pm 1.2 \pm 1.7 \pm 1.9$	$13.8 \pm 1.9 \pm 1.5 \pm 2.1$

➤ Enhancement factor (data/cocktail) decreases from peripheral to central collisions

p_T spectra in 60-80% collisions

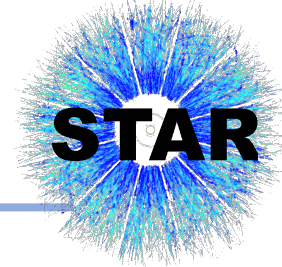


STAR, PRL 121 (2018) 132301

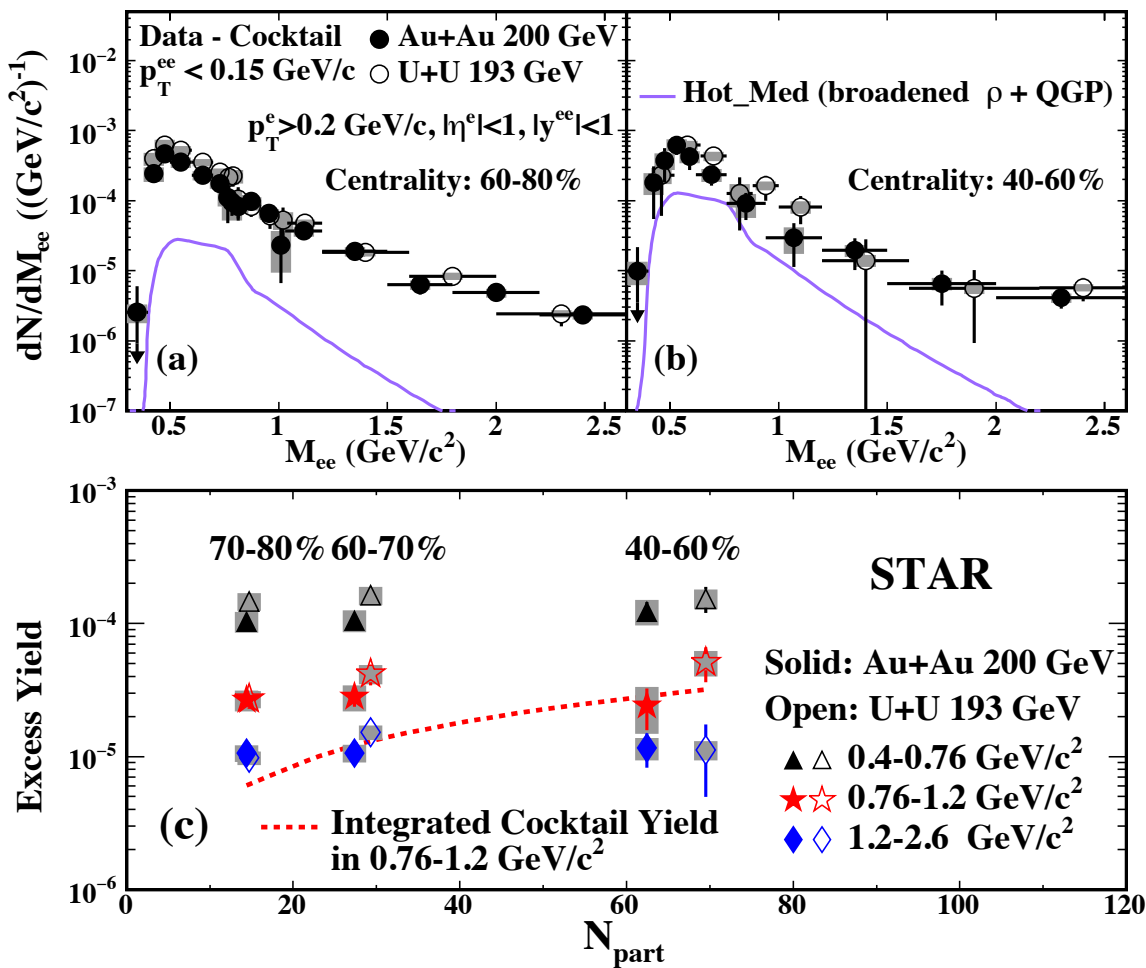


- Excess concentrated below $p_T \approx 0.15$ GeV/c
- Data are consistent with hadronic cocktail for $p_T > 0.15$ GeV/c

Origin of the low- p_T enhancement

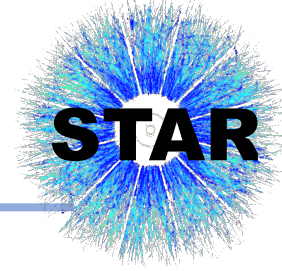


STAR, PRL 121 (2018) 132301
R. Rapp, PRC 63 (2001) 054907



- Can not be explained by in-medium broadened ρ model
- Compared to hadronic production, excess yield exhibits a much weaker centrality dependence
- Need additional source(s)

Models of two-photon interaction



➤ Equivalent Photon Approximation (EPA) method

- Photon is treated as real
- Weizsäcker–Williams method to estimate photon flux
- No impact parameter dependence of p_T spectrum for the dilepton from initial photon-photon interaction

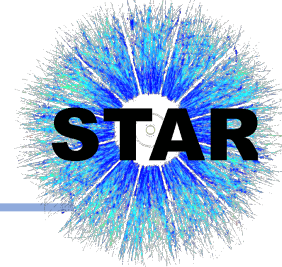
➤ Models based on EPA method

- Model by Zha *et al.* [W. Zha *et al.*, *PLB* 781 (2018) 182]
 - ✓ Use Woods-Saxon charge distribution in nucleus for photon flux estimation
 - ✓ Consider dilepton production inside nucleus
- STARlight [S. Klein, *PRC* 97 (2018) 054903]
 - ✓ Ignore dilepton production inside nucleus
- STARlight with next-to-leading order correction and hot medium effects - Coulomb scattering [S. Klein *et al.*, *arxiv*: 1811.05519]

➤ Model based on external classical field approach [M. Vidovic *et al.*, *PRC* 47 (1993) 2308]

- Model by Zha *et al.* [W. Zha *et al.*, *arxiv*: 1812.02820]
 - ✓ Consider impact parameter dependence of p_T spectrum for the dilepton from initial photon-photon interaction

Origin of the low- p_T enhancement

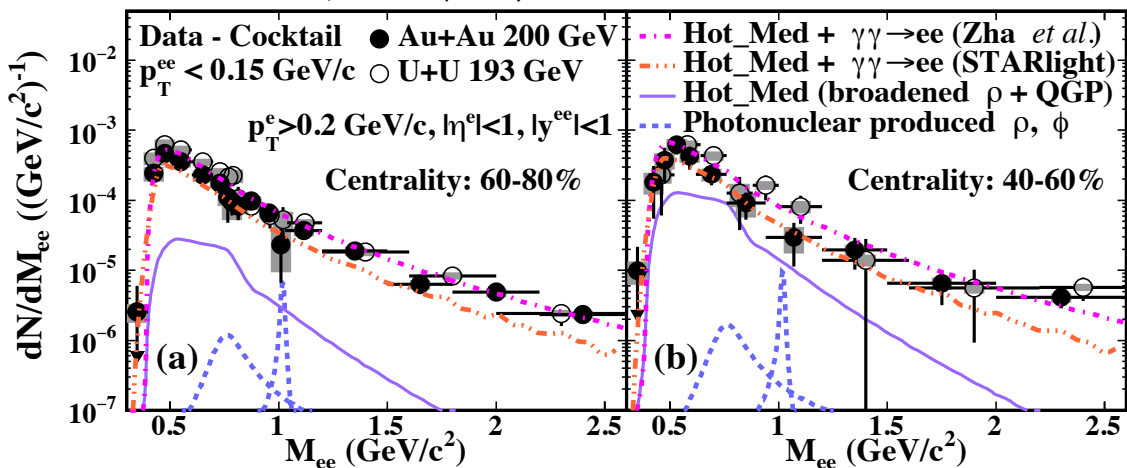


STAR, PRL 121 (2018) 132301

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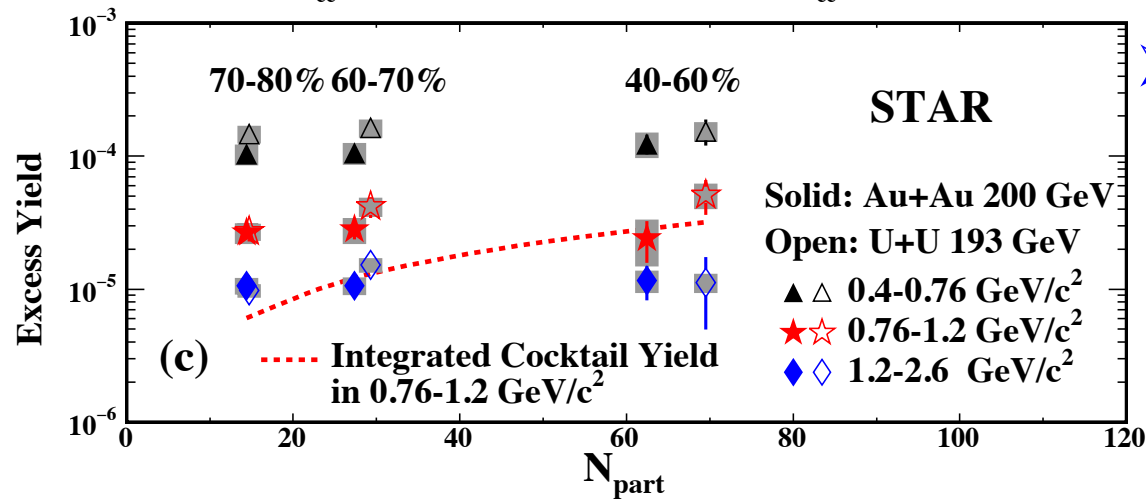
W. Zha et al., PRC 97 (2018) 044910, PLB 781 (2018) 182

S. Klein, PRC 97 (2018) 054903



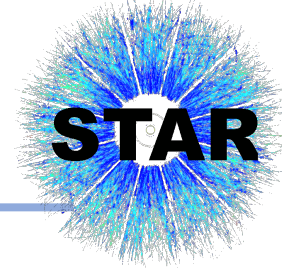
➤ Qualitatively described by models including two-photon interaction

- Data from 60-80% collisions favors model by Zha et al.

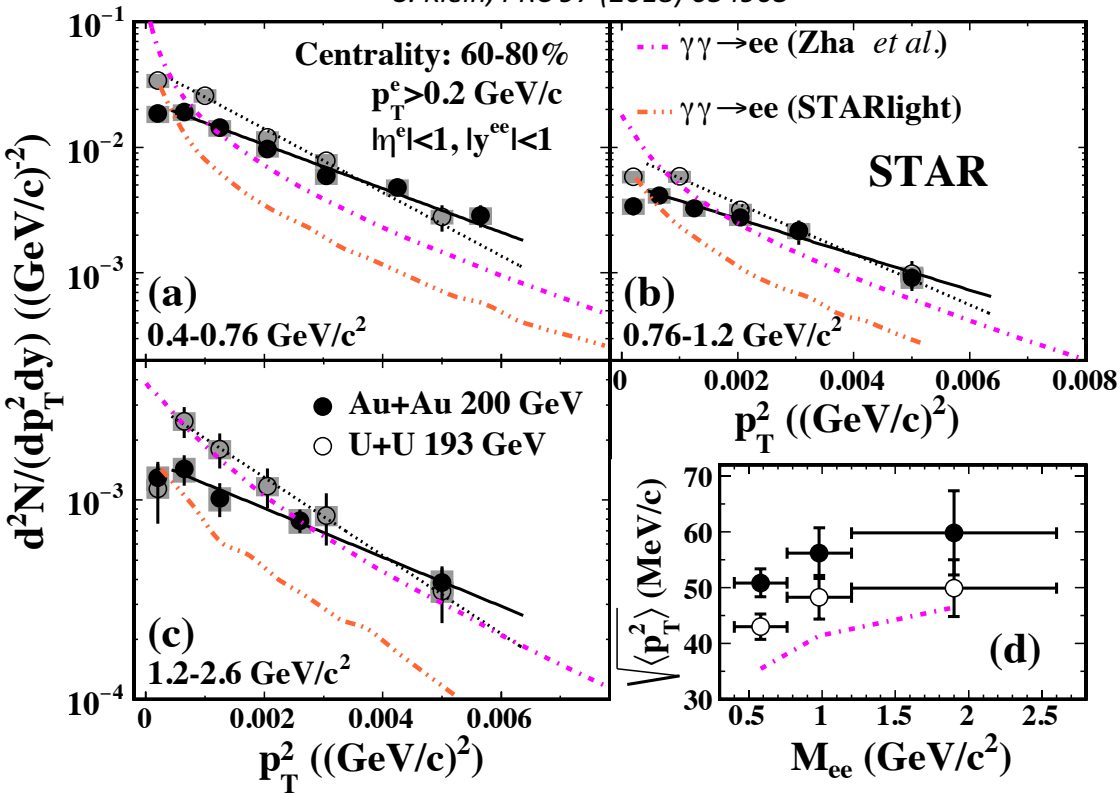


➤ Negligible contributions from photoproduced vector mesons

p_T^2 distributions in 60-80% collisions



STAR, PRL 121 (2018) 132301
 W. Zha et al., PLB 781 (2018) 182
 S. Klein, PRC 97 (2018) 054903

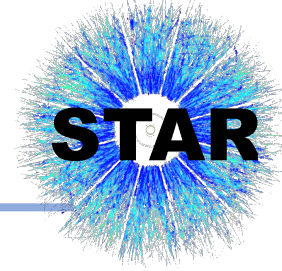


➤ Models fail to describe p_T^2 distributions

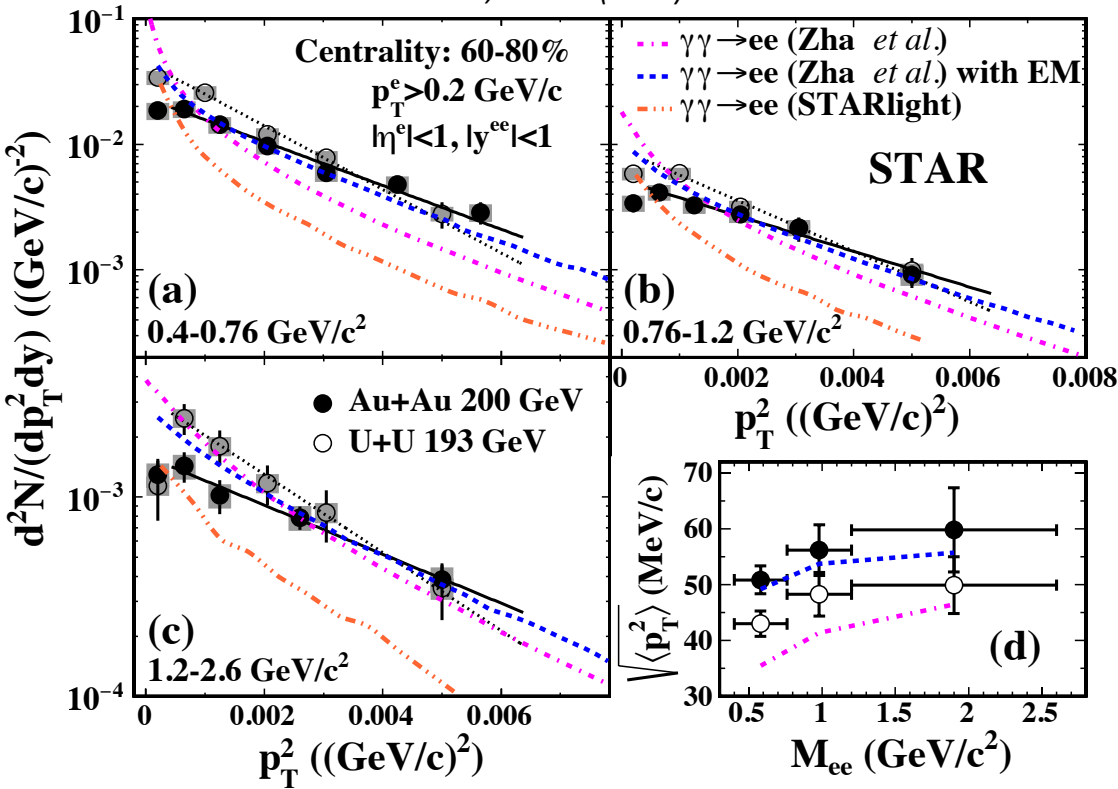
➤ Employ $\sqrt{\langle p_T^2 \rangle}$ to quantify the discrepancy between data and models

- Mass and collision species dependence
- Data are systematically higher than models

p_T^2 distributions in 60-80% collisions

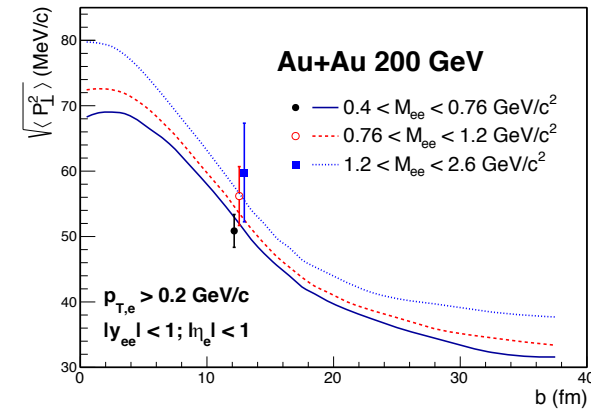
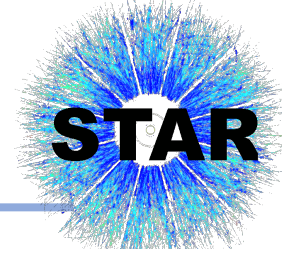


STAR, PRL 121 (2018) 132301
 W. Zha et al., PLB 781 (2018) 182
 S. Klein, PRC 97 (2018) 054903



- Models fail to describe p_T^2 distributions
- Employ $\sqrt{\langle p_T^2 \rangle}$ to quantify the discrepancy between data and models
 - Mass and collision species dependence
 - Data are systematically higher than models
- Model of Zha describes data when including effects of magnetic field on the produced pairs
 - Indication the existence of strong magnetic field trapped in QGP?

Impact parameter dependence of initial two-photon interaction [W. Zha et al., arxiv: 1812.02820]

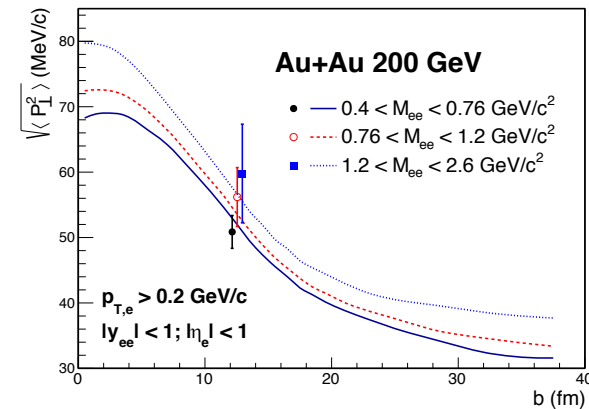
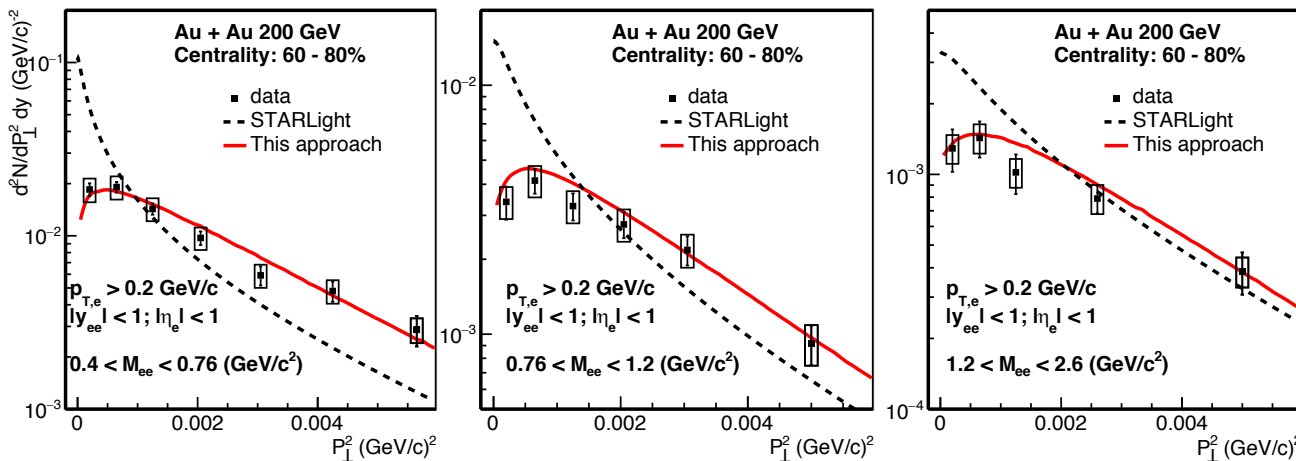
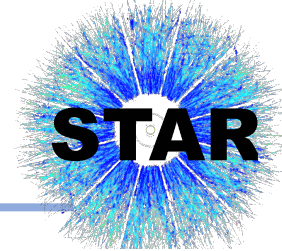


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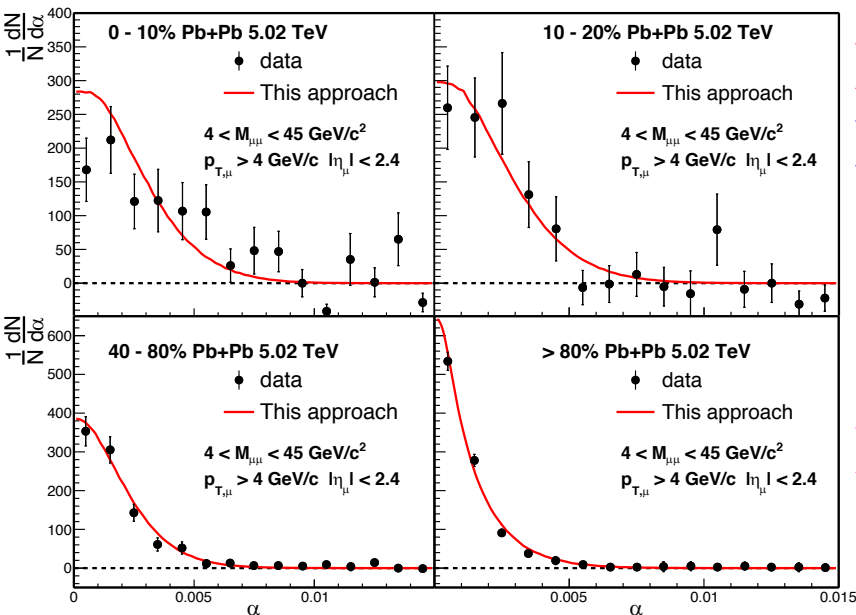
➤ Strong impact parameter dependence

Impact parameter dependence of initial two-photon interaction

[W. Zha et al., arxiv: 1812.02820]

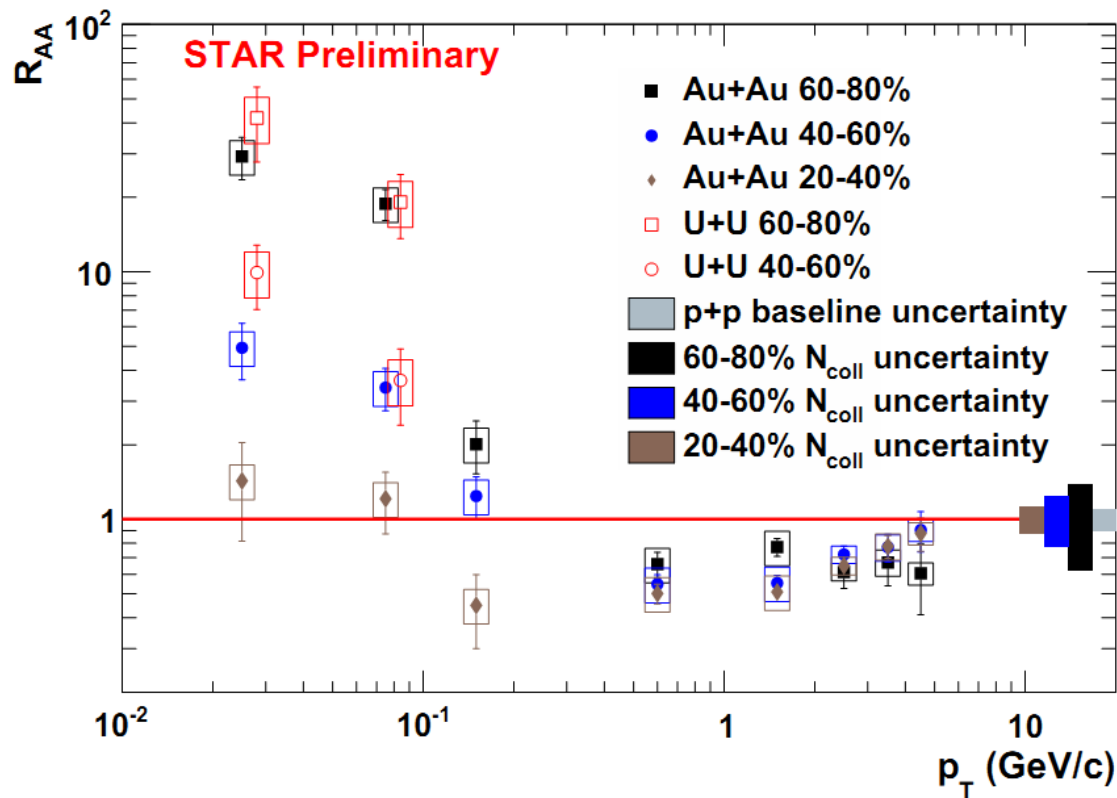
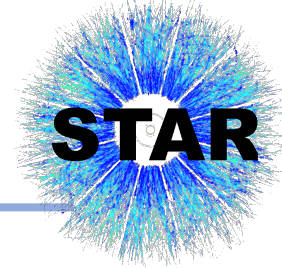


STAR, PRL 121 (2018) 132301
 ATLAS, PRL 121 (2018) 212301



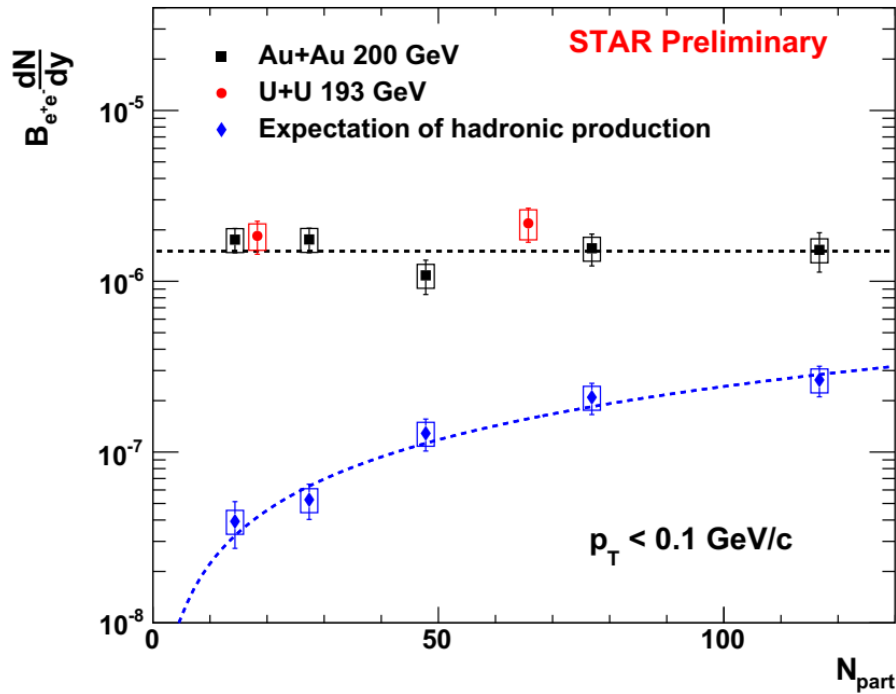
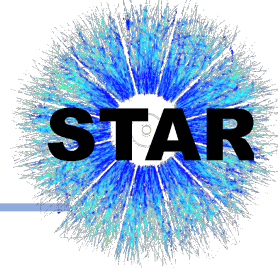
- Strong impact parameter dependence
- Can describe STAR and ATLAS data simultaneously
 - ATLAS data can also be qualitatively described by EPA model incorporating Coulomb scattering [S. Klein et al., arxiv: 1811.05519]
- Critical for the study of possible hot medium effects

Low- p_T J/ψ at STAR



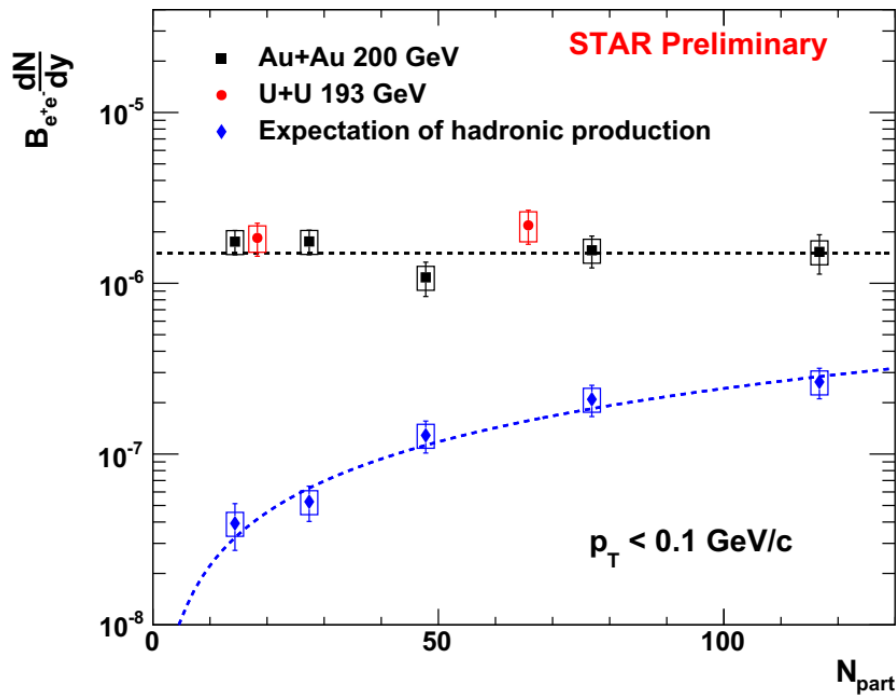
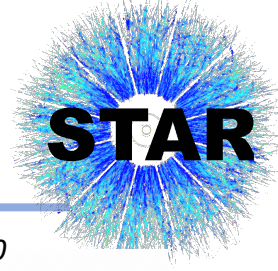
- Significant enhancement at low p_T (< 0.1 GeV/c) in 40-80% collisions
- No significant difference between Au+Au and U+U collisions

Excess yields of low- p_T J/ψ

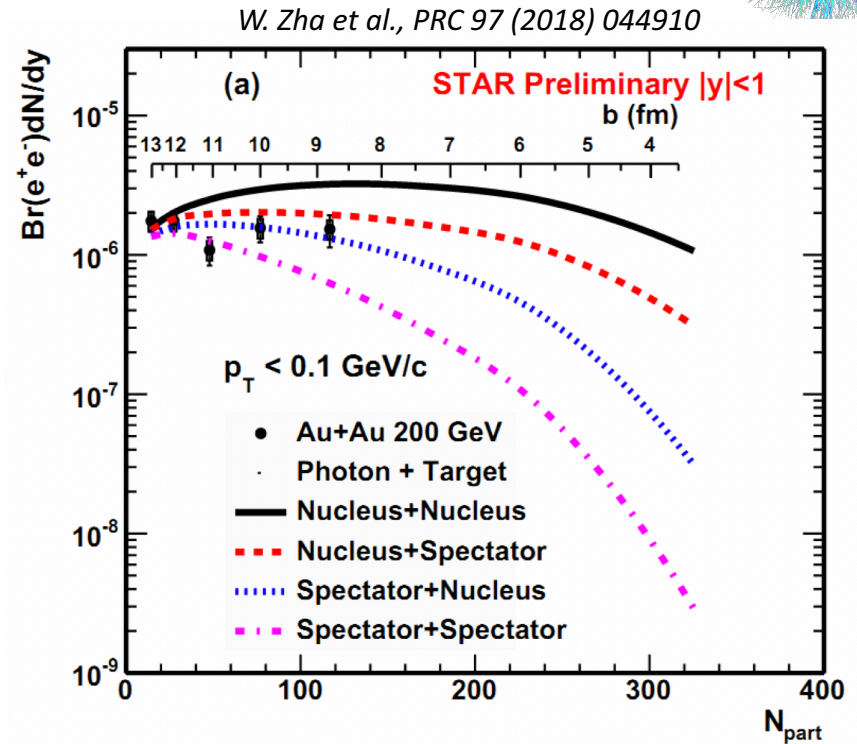


- No significant centrality dependence of the excess yield
- Yield of low- p_T J/ψ from hadronic production is expected to increase dramatically with N_{part}

Excess yields of low- p_T J/ψ

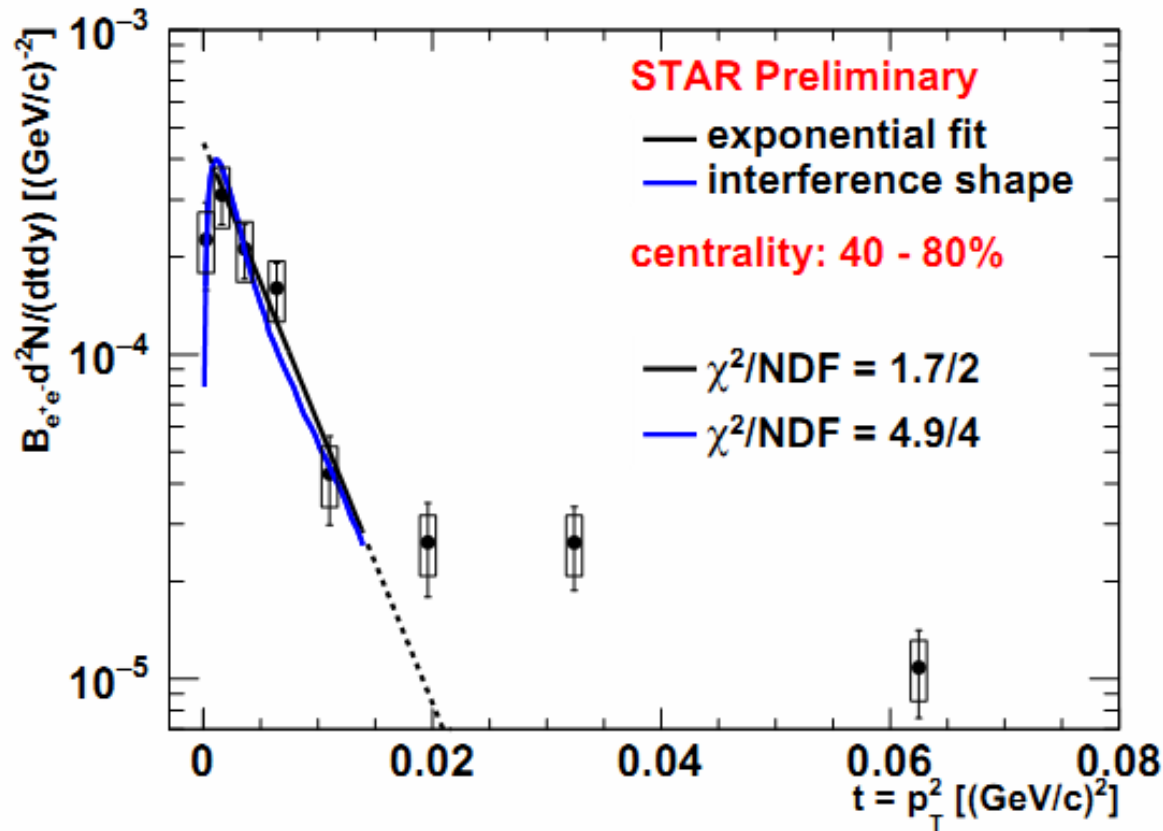
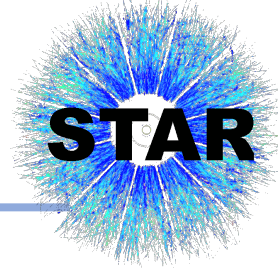


- No significant centrality dependence of the excess yield
 - Yield of low- p_T J/ψ from hadronic production is expected to increase dramatically with N_{part}



- Qualitatively described by photonuclear interaction
 - N+S and S+N scenarios can describe the data reasonably well
 - ✓ Measurements in central collisions are critical

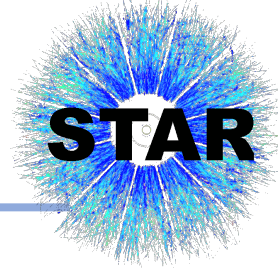
t distribution of J/ψ



➤ Similar structure to that in UPCs

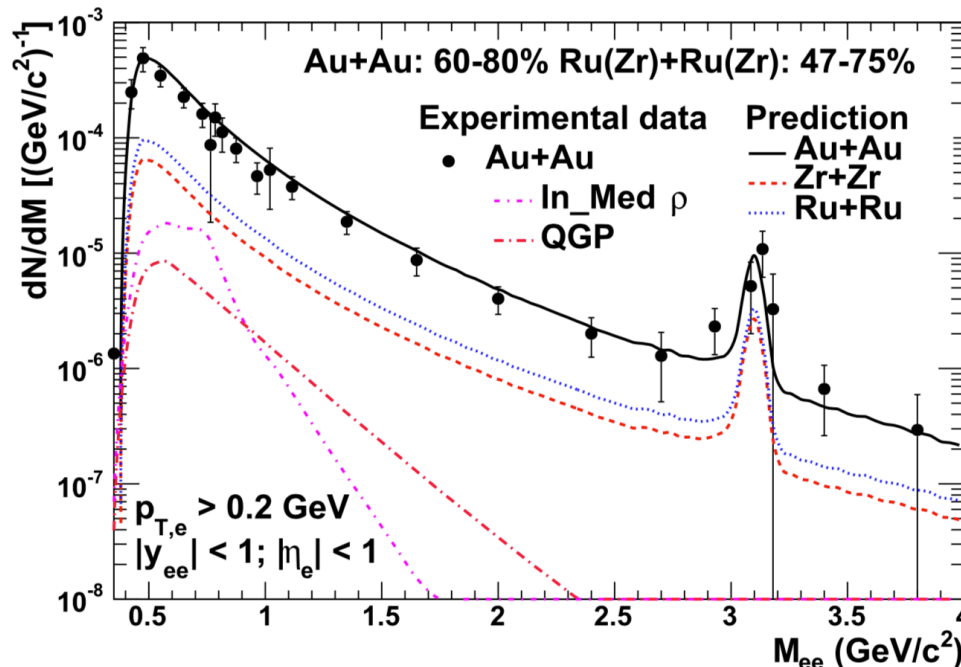
- Indication of interference [S. Klein, PRL 84 (2000) 2330]
- Similar slope parameter for exponential fit
 - ✓ Slope = $196 (\text{GeV}/c)^{-2}$ in UPC case
 - ✓ Slope = $199 \pm 31 (\text{GeV}/c)^{-2}$ (w/o the first point) in 40-80% collisions

Isobaric collisions



➤ ${}^{96}_{44}\text{Ru}+{}^{96}_{44}\text{Ru}$ vs. ${}^{96}_{40}\text{Zr}+{}^{96}_{40}\text{Zr}$

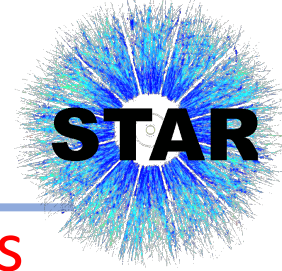
- Charge differs by 10%, everything else is almost the same
- Large statistics taken by STAR in 2018: 3.1B vs. 1.5B (goal) minimum-bias events
- Rapid (daily) switching between Ru and Zr: minimize systematic uncertainty



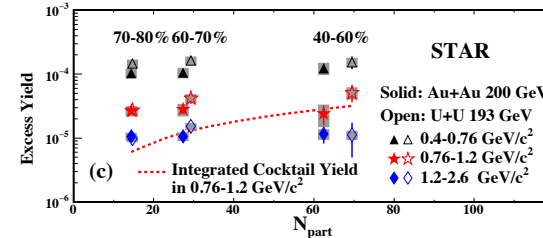
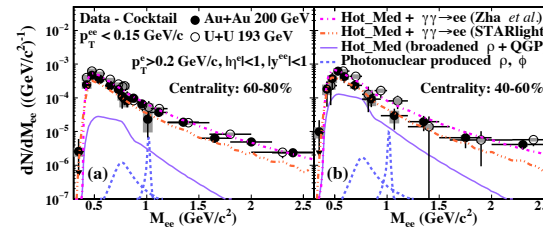
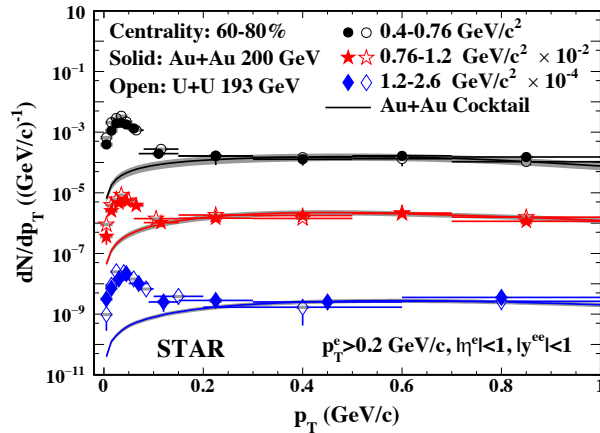
W. Zha et al., PLB 789 (2019) 238
STAR, PRL 121 (2018) 132301

➤ Further constrain the photon interactions and their possible impacts on emerging phenomena in heavy-ion collisions

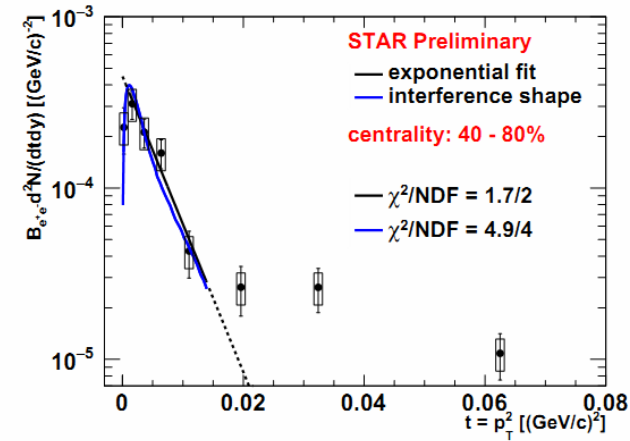
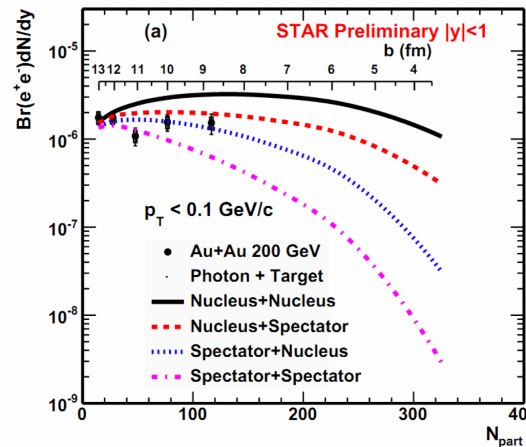
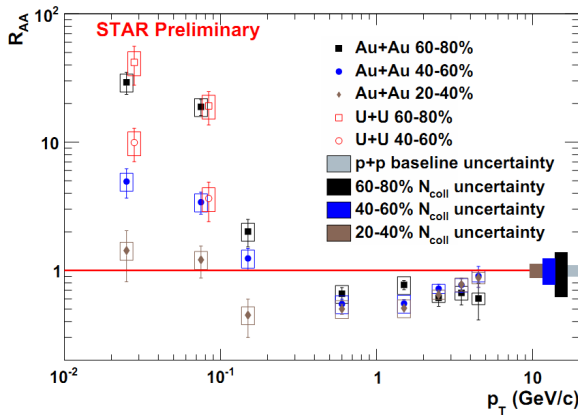
Summary



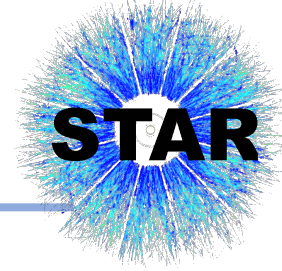
➤ Low- p_T e^+e^- pair production in heavy-ion collisions



➤ Low- p_T J/ψ production in heavy-ion collisions

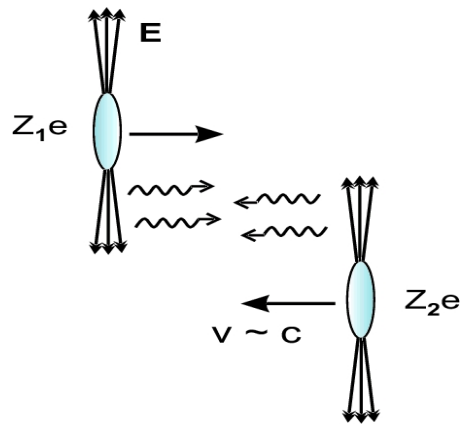
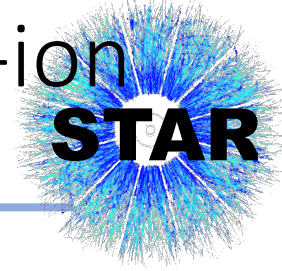


➤ Explore photon interactions in isobaric collisions



Backup

Coherent photons as “partons” in heavy-ion collisions



Coherent limitation: $Q^2 \leq 1/R^2 \Rightarrow$ quasi-real !

Photon four momentum: $q^u = (\omega, \vec{q}_T, \omega/\gamma)$

$$Q^2 = \frac{\omega^2}{\gamma^2} + q_T^2$$

$$\omega \leq \omega_{max} \sim \frac{\gamma}{R}$$

$$q_T \leq 1/R$$

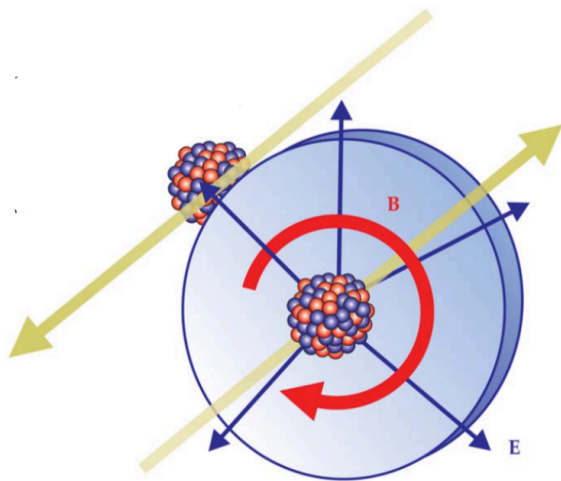
- View photons as “partons” being present with fast moving ions!

The extent of photons swarming about the ions:

The radius of nuclear matter $R_{Nuc} \sim 6.3$ fm (Au)

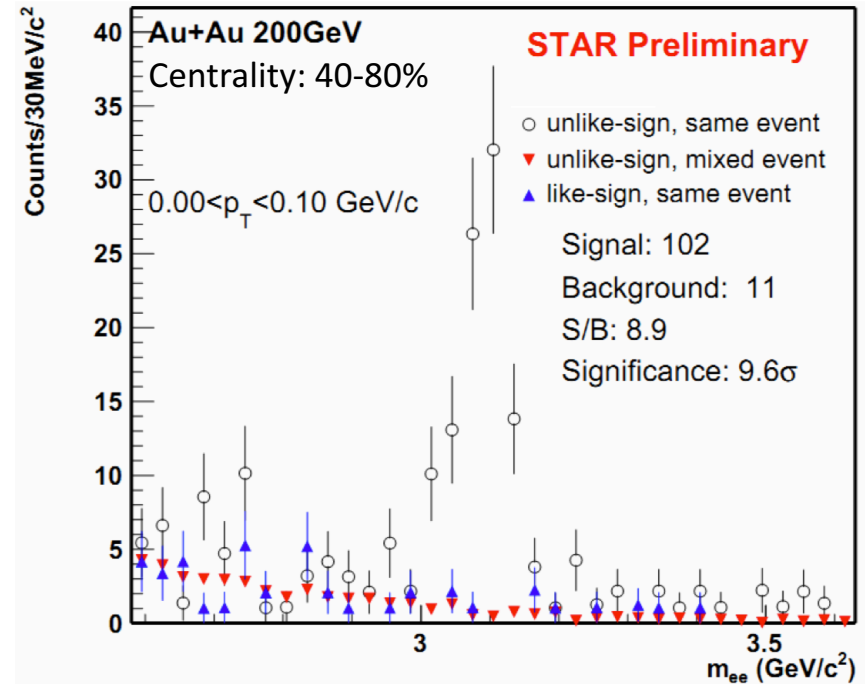
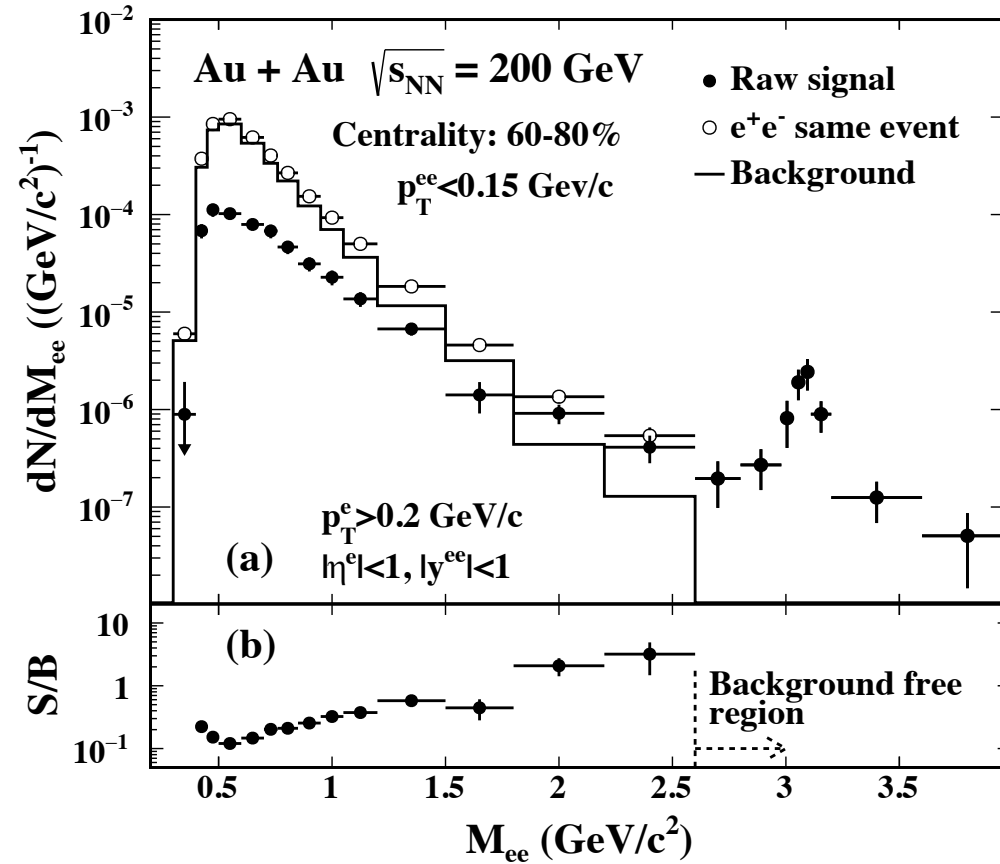
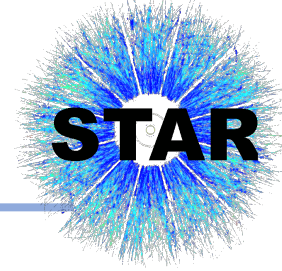
$$R_{photons} \gg R_{Nuc}$$

Take the photoproduction of ρ (Au+Au 200 GeV) in UPC as example: $\langle R_{producton} \rangle \sim 40$ fm

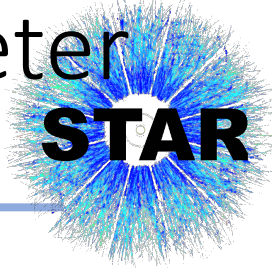


Physics Today **70**, 10, 40 (2017)

Raw e^+e^- and J/ψ signal at low p_T



Model calculation with impact parameter dependence



$$\sigma = 16 \frac{Z^4 e^4}{(4\pi)^2} \int d^2b \int \frac{dw_1}{w_1} \frac{dw_2}{w_2} \frac{d^2k_{1\perp}}{(2\pi)^2} \frac{d^2k_{2\perp}}{(2\pi)^2} \frac{d^2q_{\perp}}{(2\pi)^2} \times \frac{F(-k_1^2)}{k_1^2} \frac{F(-k_2^2)}{k_2^2} \frac{F^*(-k_1'^2)}{k_1'^2} \frac{F^*(-k_2'^2)}{k_2'^2} \boxed{e^{-i\vec{b} \cdot \vec{q}_{\perp}}} \quad (2)$$

$$\times [(\vec{k}_{1\perp} \cdot \vec{k}_{2\perp})(\vec{k}'_{1\perp} \cdot \vec{k}'_{2\perp})\sigma_s(w_1, w_2) + (\vec{k}_{1\perp} \times \vec{k}_{2\perp})(\vec{k}'_{1\perp} \times \vec{k}'_{2\perp})\sigma_{ps}(w_1, w_2)]$$

Integration over b

$$\longrightarrow \sigma = 16 \frac{Z^4 e^4}{(4\pi)^2} \int \frac{dw_1}{w_1} \frac{dw_2}{w_2} \frac{d^2k_{1\perp}}{(2\pi)^2} \frac{d^2k_{2\perp}}{(2\pi)^2} \left| \frac{F(-k_1^2)}{k_1^2} \right|^2 \times \left| \frac{F(-k_2^2)}{k_2^2} \right|^2 k_{1\perp}^2 k_{2\perp}^2 \sigma(w_1, w_2) \quad (6)$$

where the four momenta of photons are

$$k_1 = (w_1, k_{1\perp}, \frac{w_1}{v}), k_2 = (w_2, P_{\perp} - k_{1\perp}, \frac{w_2}{v})$$

$$w_1 = \frac{1}{2}(P_0 + vP_z), w_2 = \frac{1}{2}(P_0 - vP_z) \quad (3)$$

$$k_{2\perp} = P_{\perp} - k_{1\perp}, q_{\perp} = k_{1\perp} - k'_{1\perp}$$

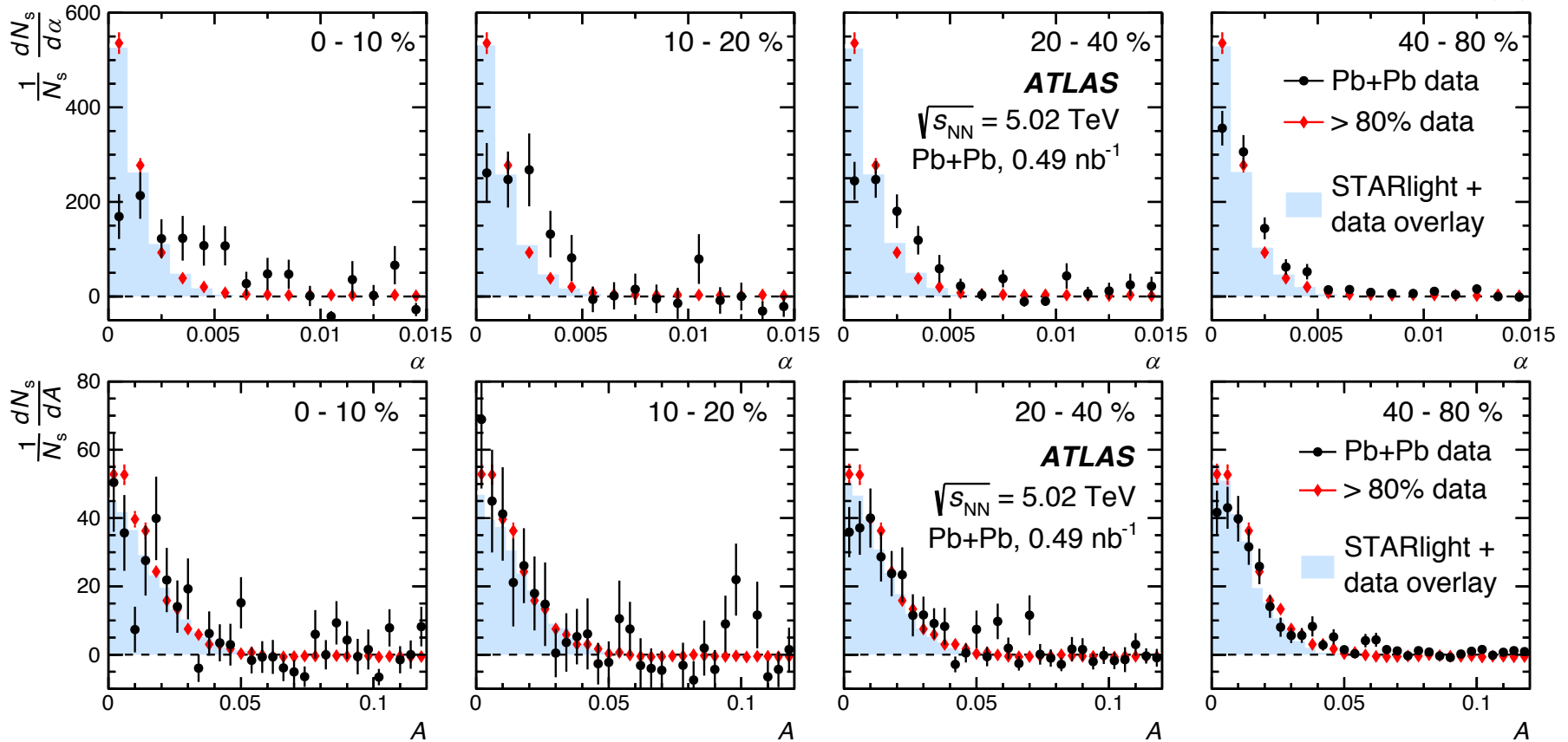
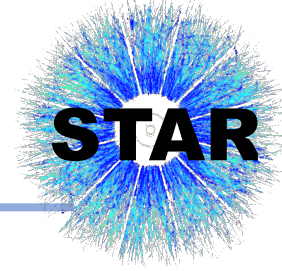
$$k'_1 = (w_1, k_{1\perp} - q_{\perp}, w_1/v)$$

$$k'_2 = (w_2, k_{2\perp} - q_{\perp}, w_2/v)$$

➤ EPA expression commonly used in traditional photon-photon models

W. Zha et al., arxiv: 1812.02820

Dimuon pairs from two-photon interaction at ATLAS



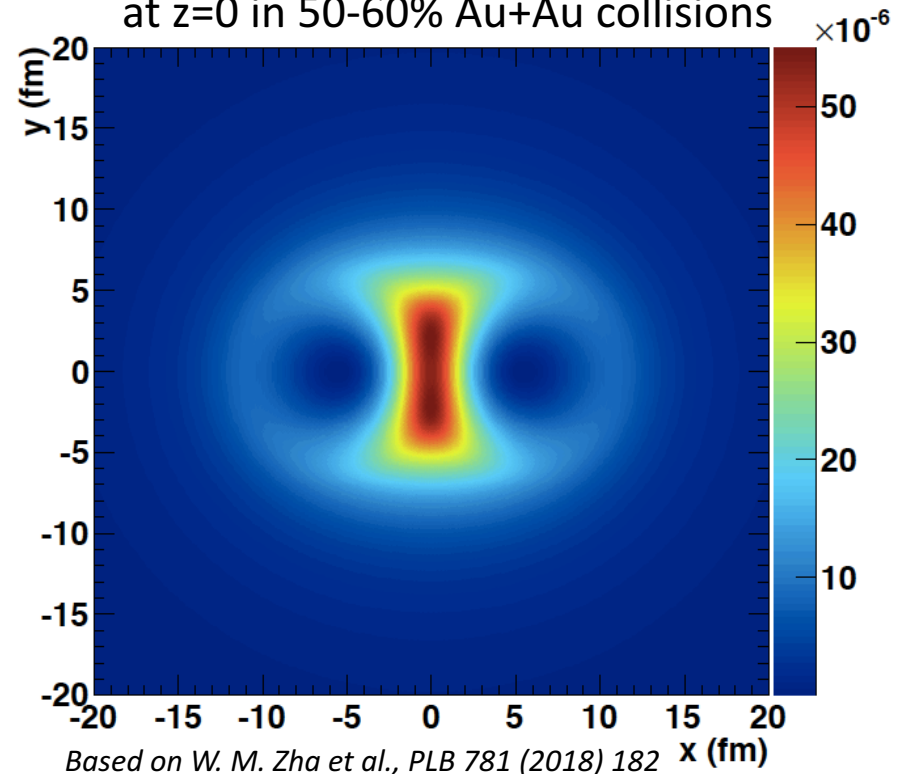
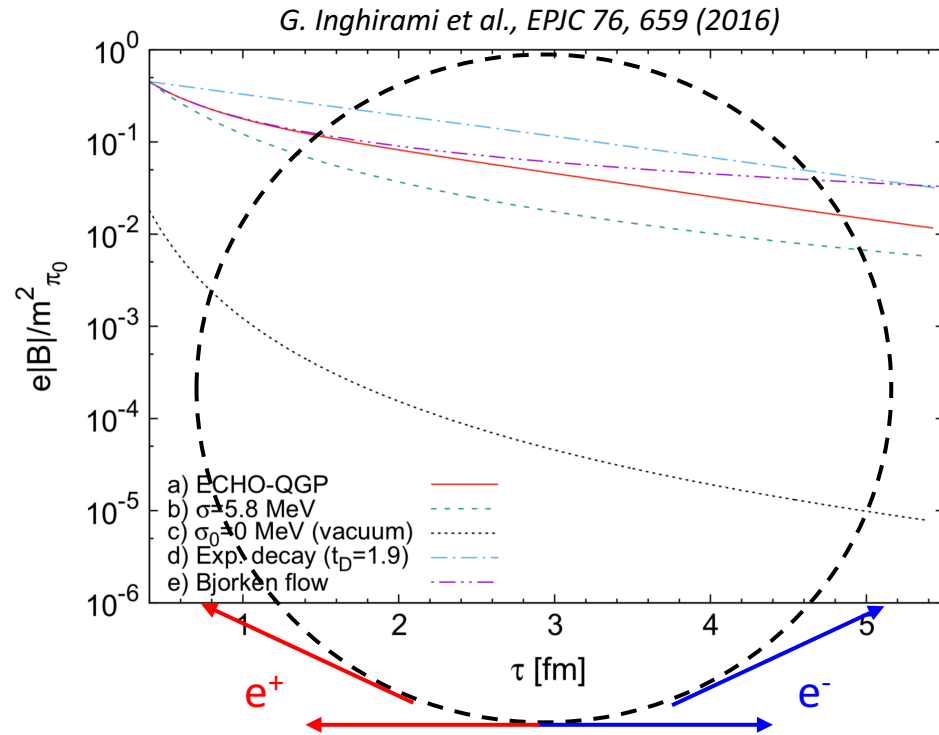
➤ Use >80% collisions as baseline

➤ Indication of Coulomb scattering? [S. Klein et al., arxiv: 1811.05519]

$$\alpha \equiv 1 - \frac{|\phi^+ - \phi^-|}{\pi}, \quad A \equiv \left| \frac{p_T^+ - p_T^-}{p_T^+ + p_T^-} \right|$$

Sensitivity to residual magnetic field?

Example of e^+e^- spatial distributions at $z=0$ in 50-60% Au+Au collisions



➤ To account for the effect of the time-dependent magnetic field on average, the model assumes that all the e^+e^- pairs traverse 1 fm through a magnetic field of 10^{14} T **perpendicular to the beam line**

- The net effect of this approach is close to $\int eB(t)cdt = e\bar{B}L$
- $e\bar{B}L \approx 30$ MeV/c, the extreme pair p_T increase: $2e\bar{B}L \approx 60$ MeV/c