



Search for QCD Critical Point: Higher Moments of Net-proton Multiplicity Distributions at RHIC



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Outline

➤ Introduction:

- 1) QCD Phase Diagram.
- 2) Higher Moments Method.
- 3) Baseline Studies.

➤ Data Analysis:

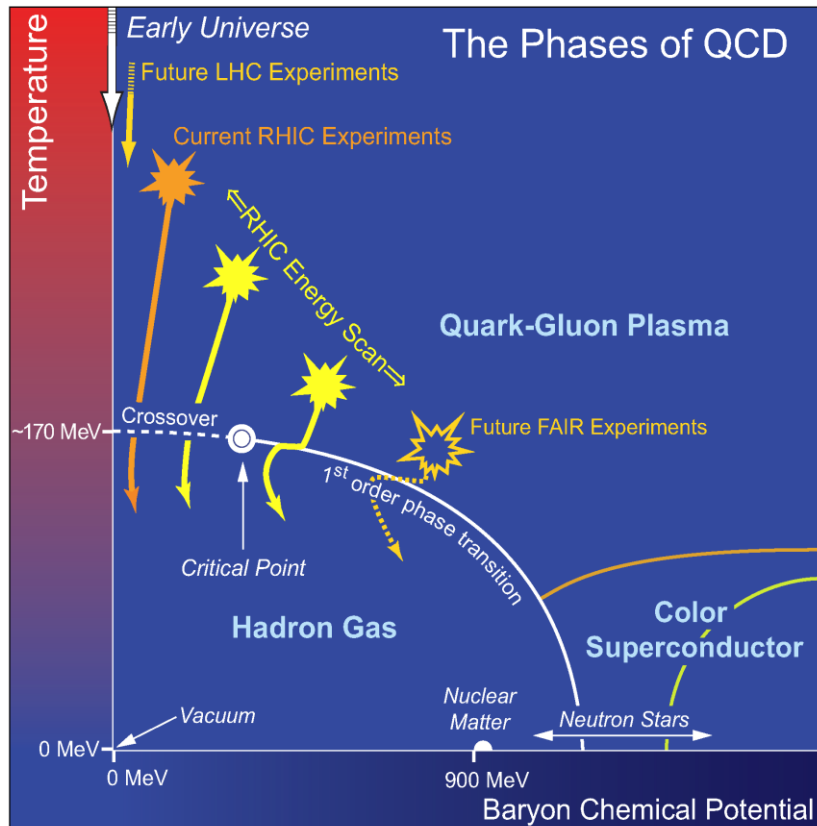
- 1) System: p+p, d+Au, Cu+Cu and Au+Au.
Energy: 19.6, 22.4, 39, 62.4 and 200 GeV.
- 2) Centrality Dependence.
- 3) Energy Dependence.

➤ Comparison with Lattice QCD and Thermal Model.

➤ Summary and Outlook.

QCD Phase Diagram

Shows condition at which thermodynamically distinct phases can occur at equilibrium.



Lattice QCD:

➤ Crossover at $\mu_B = 0$, 1st order phase transition at large μ_B .

➤ QCD Critical Point: The end point of first order phase transition boundary.

Y. Aoki et al., Nature 443:675-678, 2006

Exploring the phase structure.

➤ Map the QCD Phase Boundary.

➤ Search for the QCD Critical Point (CP).

Where is the QCD Critical Point?

Theoretical Calculations:

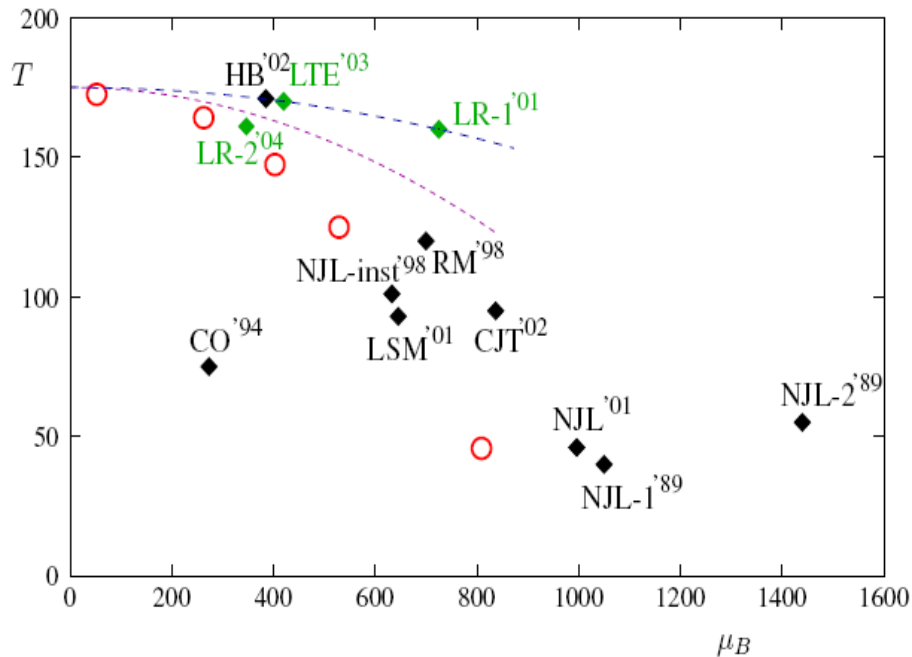
- Lattice QCD
- QCD Based Models

Vs

Experimental measurements

- Sensitive Observable.

Large uncertainties of theoretical calculation.



Approach to QCD Critical Point (CP):

- Diverge of the Correlation length (ξ)
- Non-Gaussian fluctuations.



Non-monotonic signal expected around CP.

M. Stephanov, Phys. Rev. Lett. 102, 032301 (2009)

M. Stephanov, Acta Phys.Polon.B35:2939-2962,2004

PoSLAT2006:024,2006 (hep-lat/0701002)

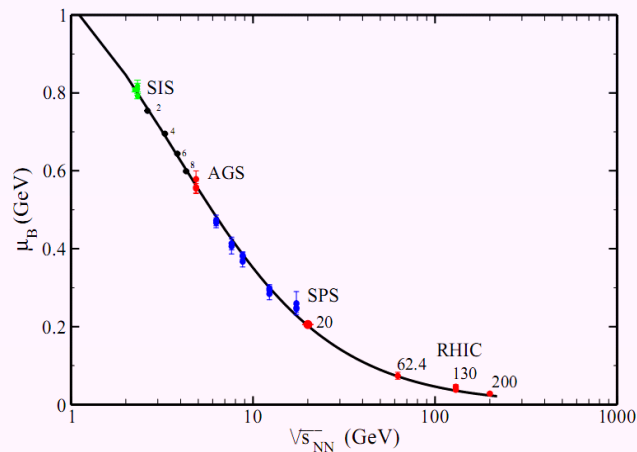
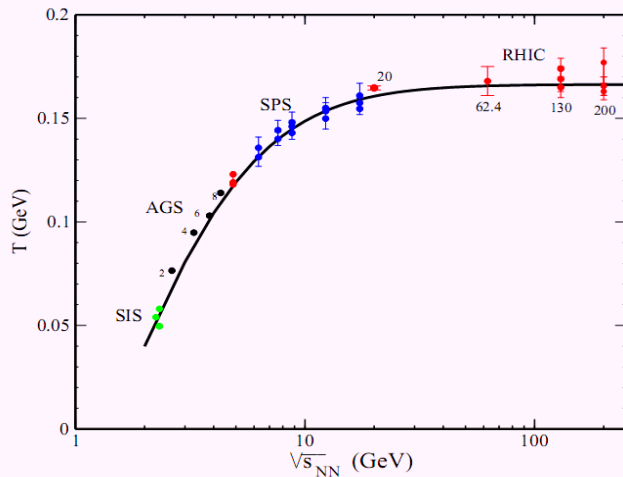
Z. Fodor, S. D. Katz, hep-lat/0106002, hep-lat/0401023.

hep-lat/0402006.

- Particle ratio fit with Thermal Model:
Chemical freeze out temperature (T)
and baryon chemical potential (μ_B).

J. Cleymans et al, Phys. Rev. C73 (2006) 034905

- Varying the colliding energy, we can access different regions (T, μ_B) on the QCD phase diagram.



- **RHIC Beam Energy Scan (BES) Program.**

Year	$\sqrt{s_{NN}}$ (GeV)
2010	7.7, 11.5, 39, 62.4, 200 (μ_B Coverage : 20~420 MeV)
2011	5 (Test Run) , 18, 27, 200

- **STAR Detector : Large Uniform Acceptance.**

Good opportunities to search for CP !

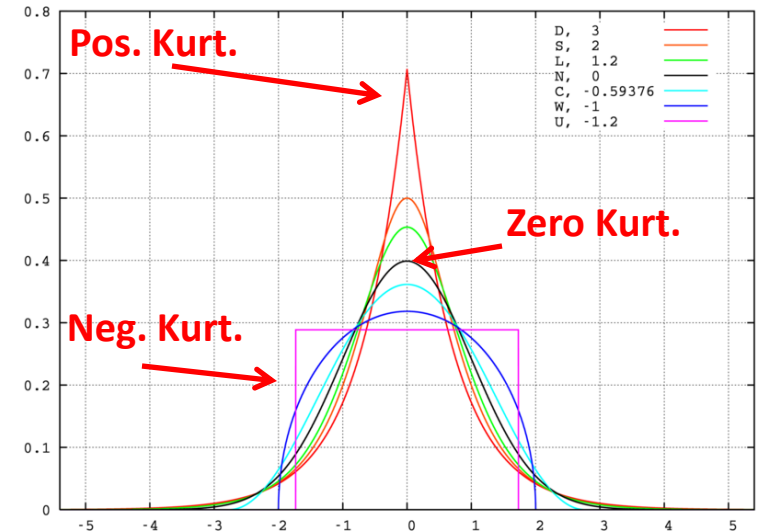
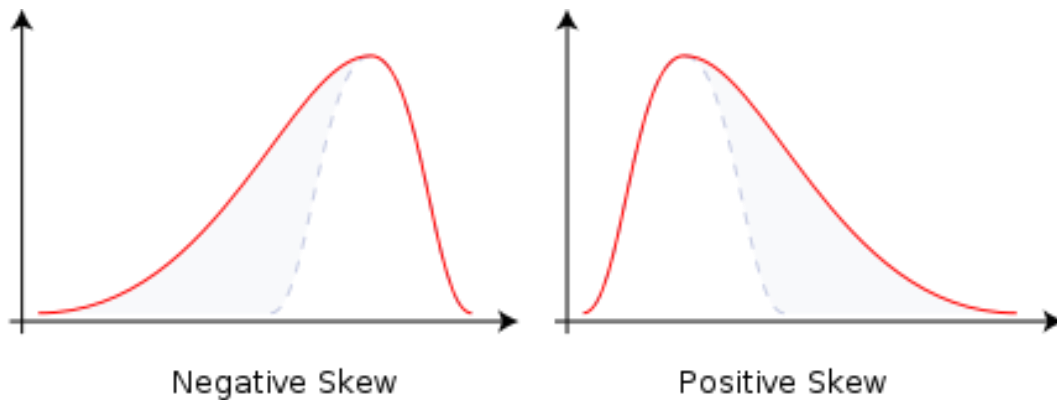
Definition : **N: Event by Event Multiplicity Distribution**

Mean: $Y = \langle N \rangle$

St. Deviation: $\sigma = \sqrt{\langle (N - \langle N \rangle)^2 \rangle}$

Skewness: $s = \frac{\langle (N - \langle N \rangle)^3 \rangle}{\sigma^3}$

Kurtosis: $\kappa = \frac{\langle (N - \langle N \rangle)^4 \rangle}{\sigma^4} - 3$



➤ For Gaussian distribution, the skewness and kurtosis are equal to zero. **Ideal probe of the non-Gaussian fluctuations at CP.**

➤ **Link to Thermodynamic Susceptibilities in Lattice QCD and Hadron Resonance Gas (HRG) Model:**

$$\chi_B^{(n)} = \left. \frac{\partial^n (P/T^4)}{\partial (\mu_B/T)^n} \right|_T$$

M.Cheng et al, Phys. Rev. D 79, 074505 (2009)

F. Karsch and K. Redlich, Phys. Lett. B 695, 136 (2011)

$$\chi_B^2 = \frac{1}{VT^3} \langle \delta N_B^2 \rangle$$

$$\chi_B^3 = \frac{1}{VT^3} \langle \delta N_B^3 \rangle$$

$$\chi_B^4 = \frac{1}{VT^3} (\langle \delta N_B^4 \rangle - 3 \langle \delta N_B^2 \rangle^2)$$

$$\chi_B^4 / \chi_B^2 = (\kappa \sigma^2)_B$$

$$\chi_B^3 / \chi_B^2 = (S \sigma)_B$$

Experimental measurable net-proton numbers fluctuations can reflect baryon and charge number fluctuations.

Y. Hatta et al, PRL 91, 102003 (2003)

➤ **Sensitive to Correlation Length (ξ) : QCD Based Model Calculation.**

Due to finite size, finite time effects.
in heavy ion collisions. $\xi \sim 2-3$ fm.

$$\langle (\delta N)^2 \rangle \approx \xi^2$$

$$\langle (\delta N)^3 \rangle \approx \xi^{4.5}$$

$$\langle (\delta N)^4 \rangle - 3 \langle (\delta N)^2 \rangle^2 \approx \xi^7$$

M. A. Stephanov, Phys. Rev. Lett. 102, 032301 (2009)

C. Athanasiou, M. Stephanov, K. Rajagopal, Phys. Rev. D 82, 074008 (2010)

Baseline (I): Thermal Fluctuations

With the Boltzmann approximation, thermodynamic pressure in the HRG model (Grand Canonical Ensemble):

$$\frac{P}{T^4} = \frac{1}{\pi^2} \sum_i d_i (m_i / T)^2 K_2(m_i / T) \cosh[(B_i \mu_B + S_i \mu_S + Q_i \mu_Q) / T]$$

F. Karsch and K. Redlich, Phys. Lett. B 695, 136 (2011)

Consider the net-proton fluctuations reflect the net-baryon fluctuations.

$$\kappa \sigma^2 = \frac{\chi_B^{(4)}}{\chi_B^{(2)}} = 1$$

$$S \sigma = \frac{\chi_B^{(3)}}{\chi_B^{(2)}} = \tanh(\mu_B / T) < 1$$

$$\mu_S \ll \mu_B$$

$$\mu_Q \ll \mu_B$$

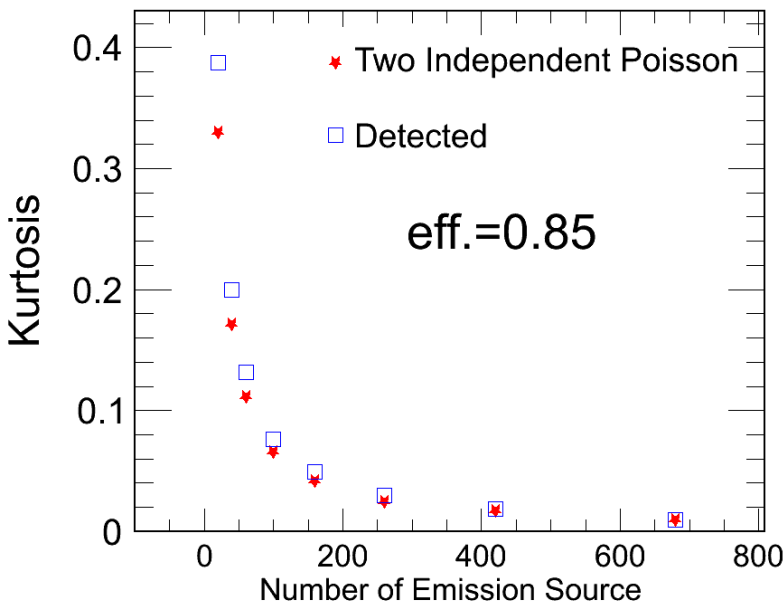
➤ $\kappa \sigma^2$ is unity and $S \sigma$ is related to μ_B / T ratio of the thermal system.

Baseline (II): Detector Efficiency Effect

- **Binomial Process of Detected Particles:** With the total produced multiplicity N and the detector efficiency ε .

$$B(n; N, \varepsilon) = \frac{N!}{n!(N-n)!} \varepsilon^n (1-\varepsilon)^{N-n} \quad \Rightarrow \quad T(k) = \sum_N B(k; N, \varepsilon) P(N)$$

- **Monte Carlo:** Input two Independent Poisson Distribution: $N=N_1-N_2$



$$\sigma_k^2 = \varepsilon \sigma_N^2$$

$$S_k = \frac{S_N}{\sqrt{\varepsilon}}$$

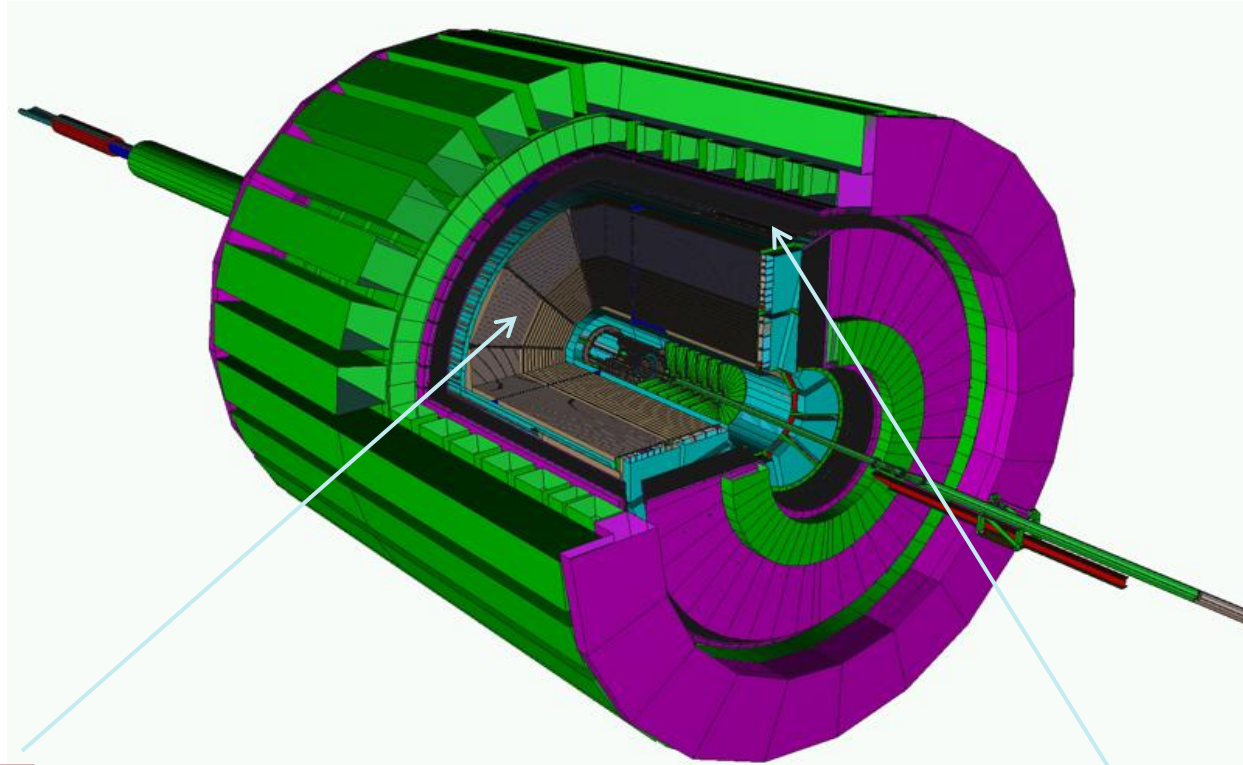
$$K_k = \frac{K_N}{\varepsilon}$$

$$S_k \sigma_k = \frac{S_N}{\sqrt{\varepsilon}} * \sqrt{\varepsilon} \sigma_N = S_N \sigma_N$$

$$K_k \sigma_k^2 = \frac{K_N}{\varepsilon} * \varepsilon \sigma_N^2 = K_N \sigma_N^2$$

- For this case, the efficiency effects are simple and will be cancelled out for moment products.

STAR Detector



TPC

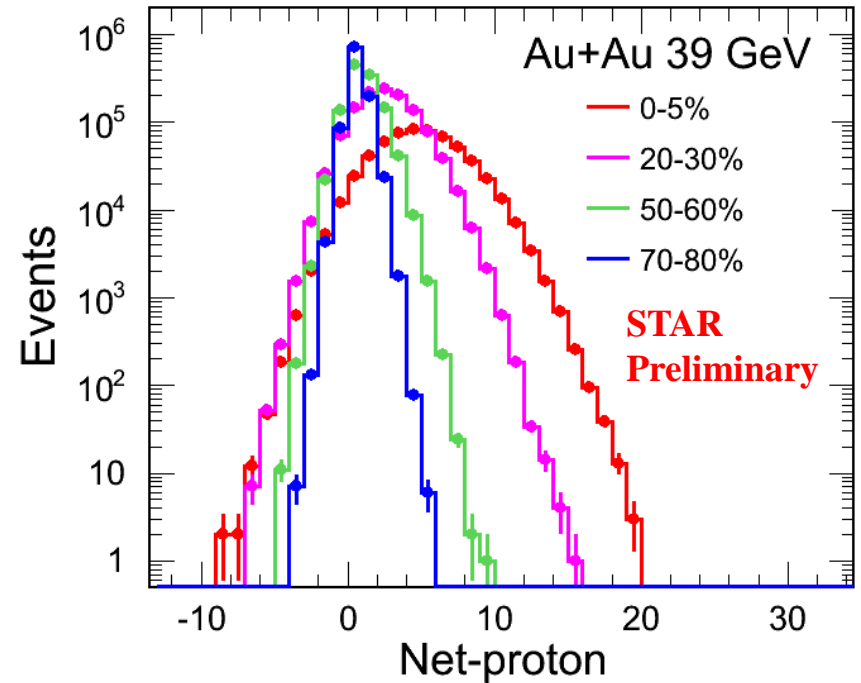
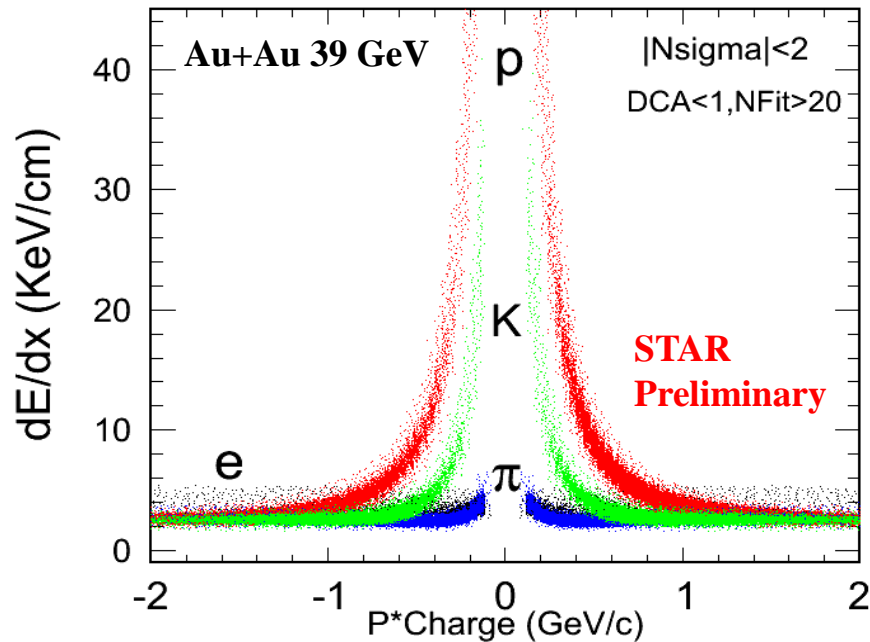
TOF

Time Projection Chamber:

- Acceptance: $-1 < \eta < 1, 0 < \varphi < 2\pi$
- Tracking: Particle momentum and trajectory.
- PID: Ionization Energy Loss (dE/dx).
 $(\pi, K) : p_T < 0.7$, proton : $p_T < 1 \text{ GeV}/c$

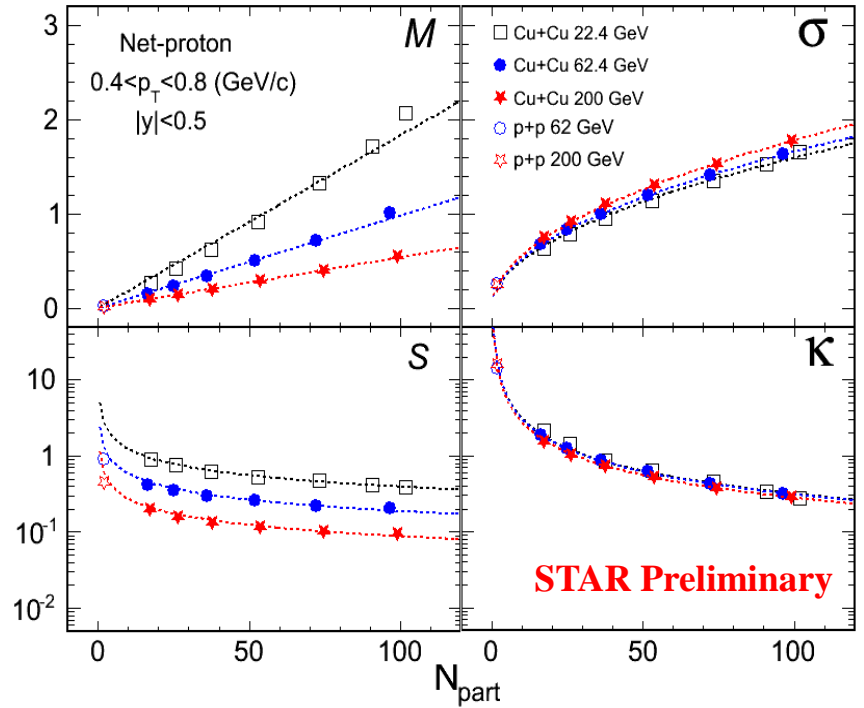
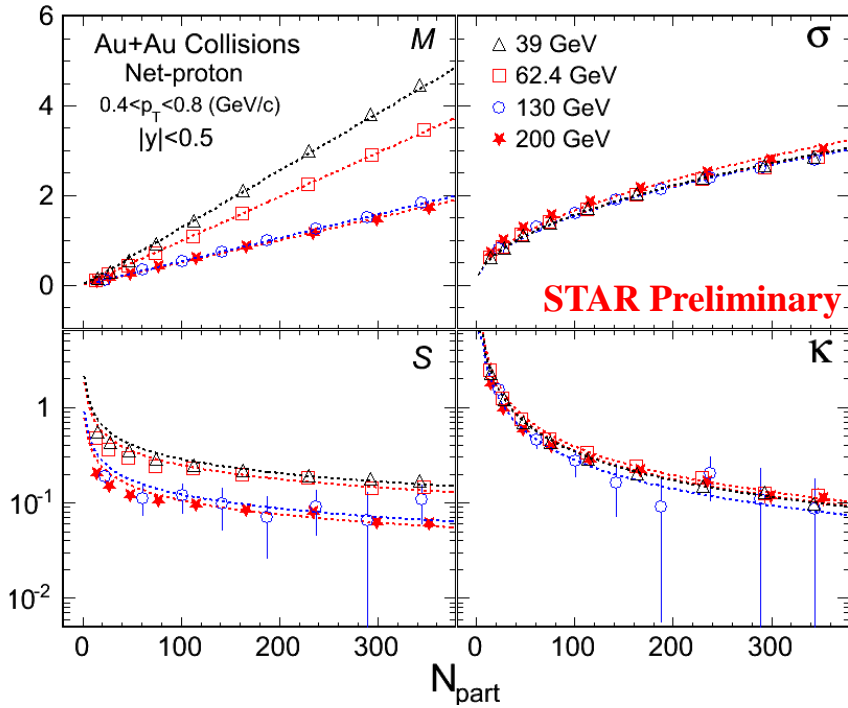
Time Of Flight:

- Acceptance: $-0.9 < \eta < 0.9, 0 < \varphi < 2\pi$
- Timing Resolution $< 100\text{ps}$.
- PID: $(\pi, K) : p_T < 1.6$, proton: $p_T < 3 \text{ GeV}/c$



- Clean Proton and antiproton identification with TPC dE/dx . for $0.4 < p_T < 0.8$ (GeV/c) and $|y| < 0.5$.
- The event-by-event net-proton distributions are more symmetrical in central collision than peripheral.

Centrality Dependence (I): Higher Moments



Central Limit Theorem (CLT)

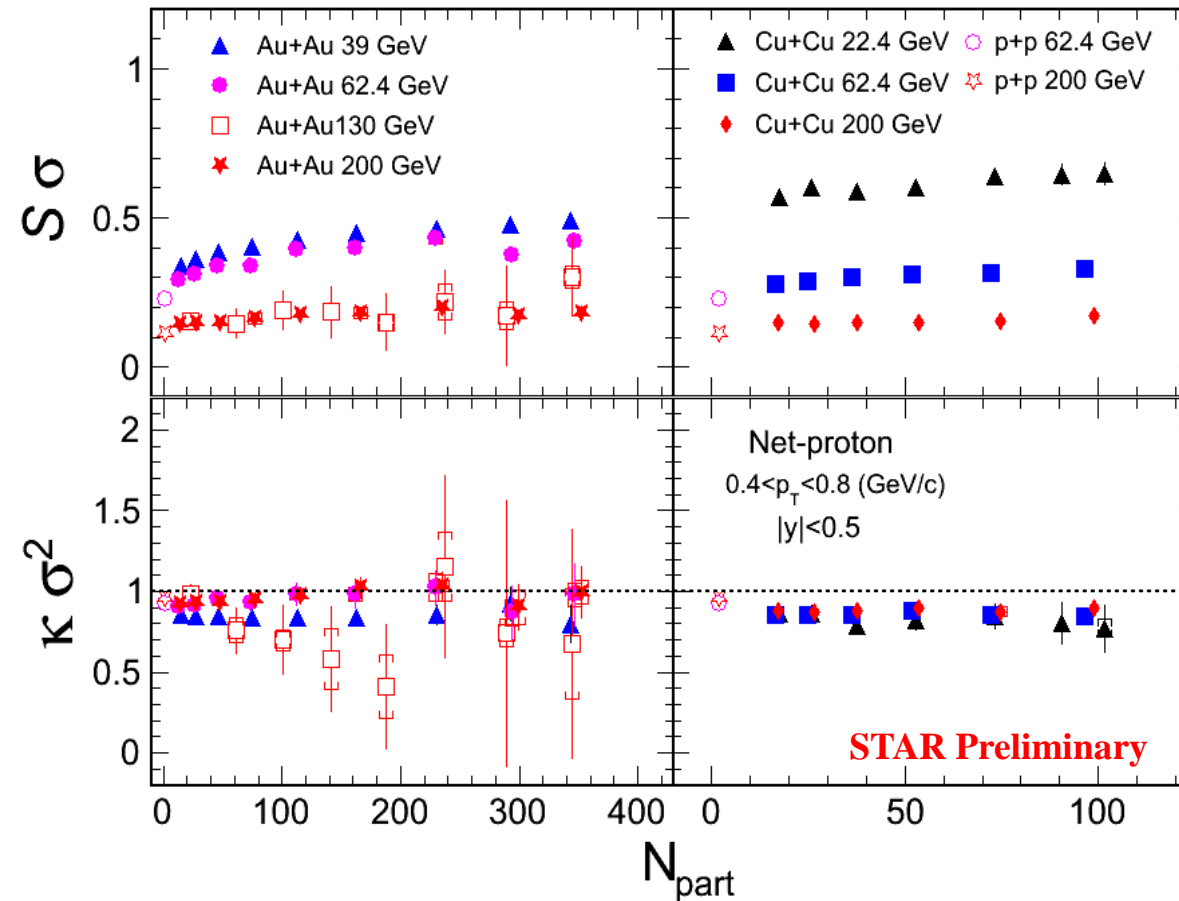
$$M_i = M_x \times C \times N_{part}, \sigma_i^2 = \sigma_x^2 \times C \times N_{part}$$

$$S_i = \frac{S_x}{\sqrt{C \times N_{part}}}, \kappa_i = \frac{\kappa_x}{(C \times N_{part})}$$

Consistent with CLT Expectations (lines).

Indicates many identical, independent particle emission sources.

Centrality Dependence (II): Moment Products



Related to baryon number susceptibility ratio:

$$(S\sigma)_B = \chi_B^3 / \chi_B^2$$

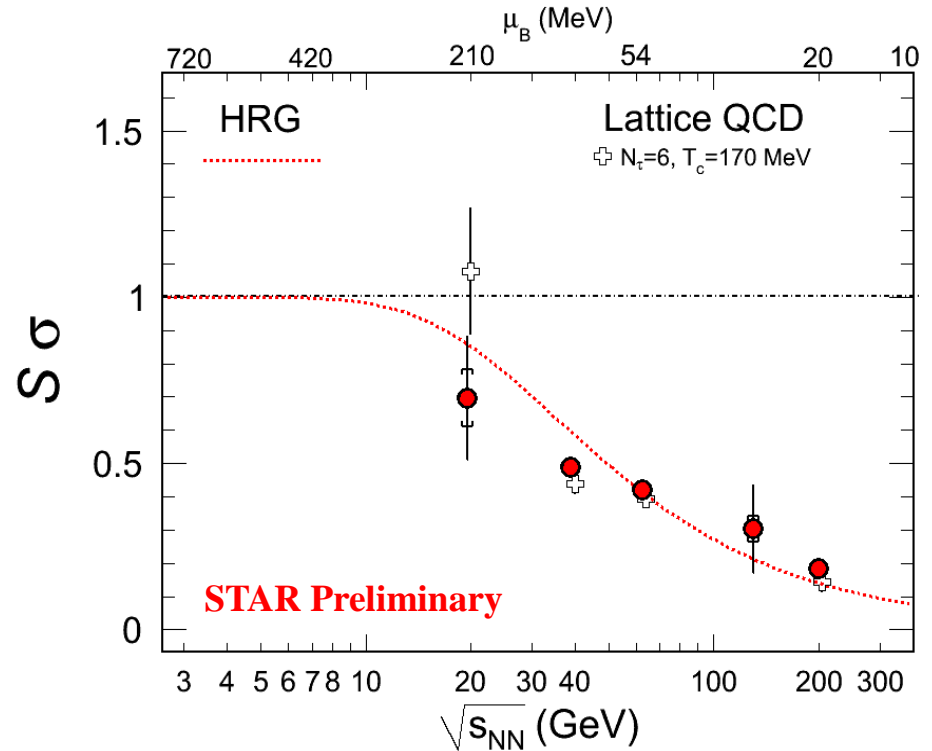
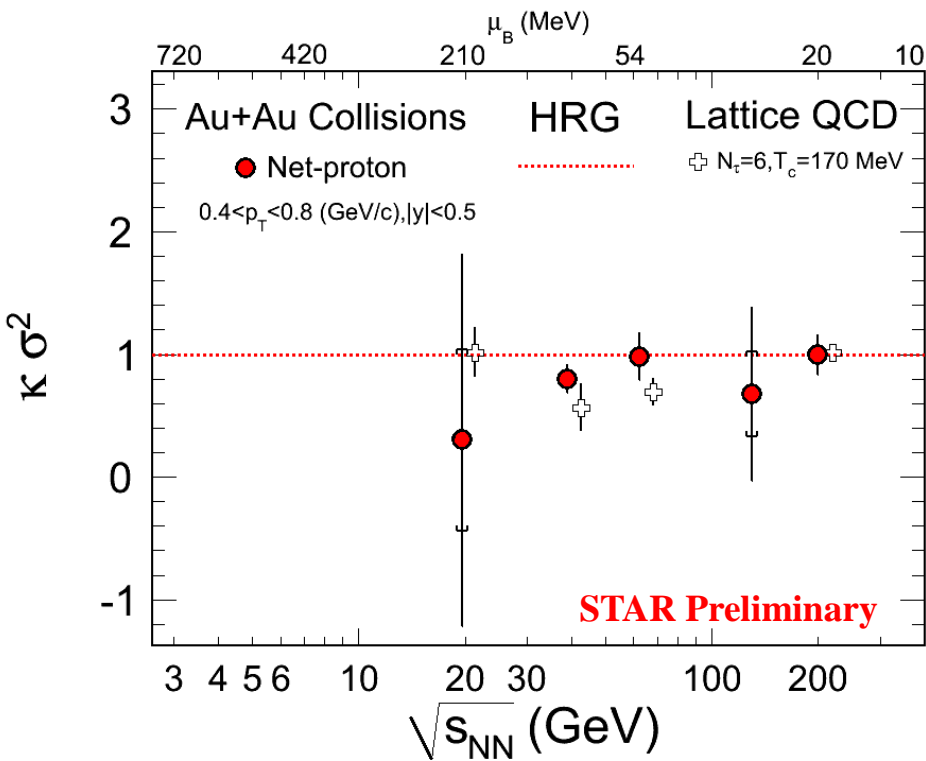
$$(\kappa\sigma^2)_B = \chi_B^4 / \chi_B^2$$

M.Cheng et al, Phys. Rev. D 79, 074505 (2009)

F. Karsch and K. Redlich, Phys. Lett. B 695, 136 (2011)

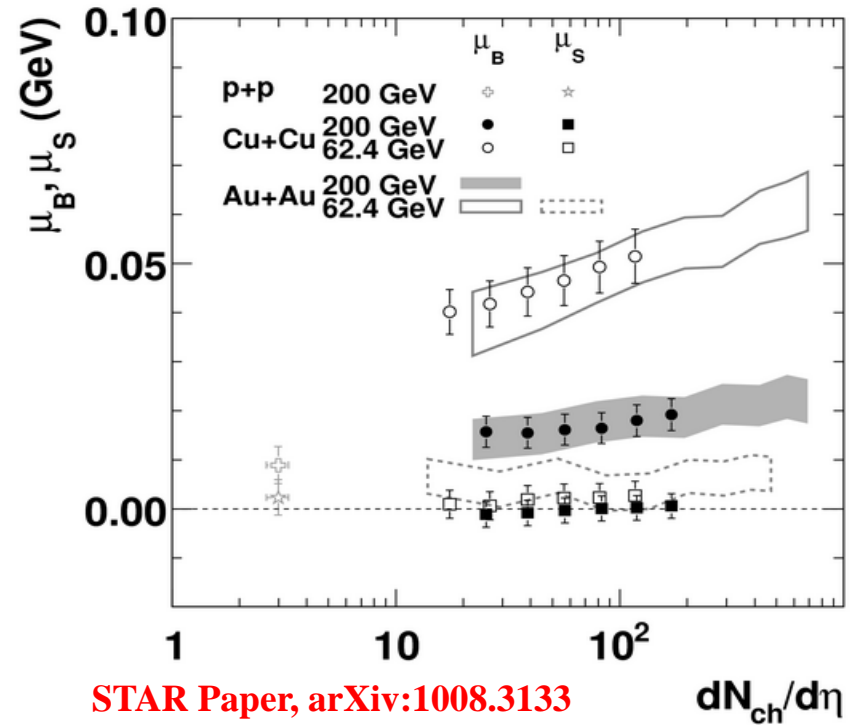
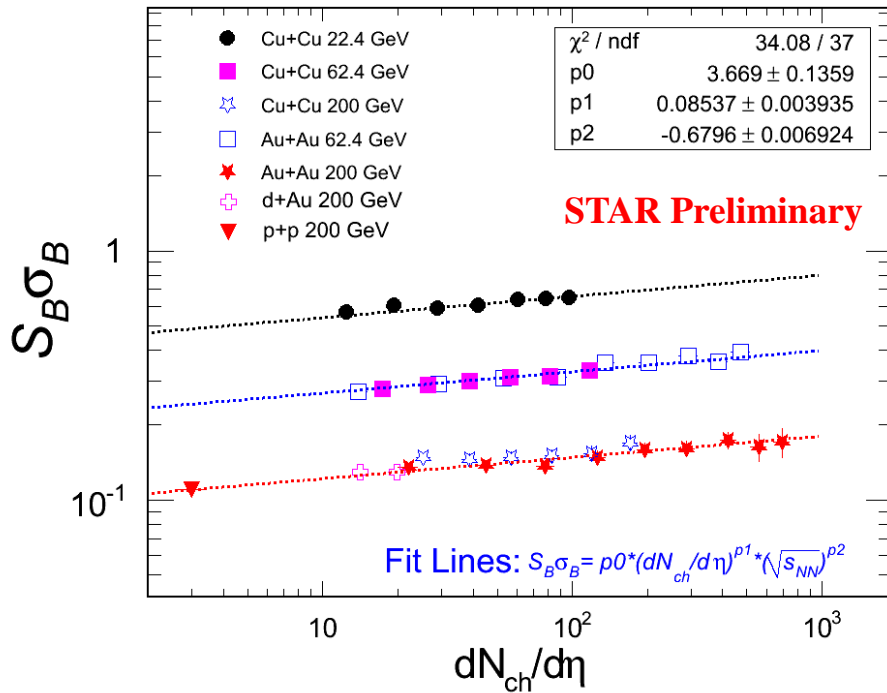
- $S\sigma$: Weak centrality dependence.
- $\kappa\sigma^2$: No centrality dependence.

STAR Energy Dependence: Moment Products (Central Collision)



- 62.4, 130 and 200 GeV data are consistent with Lattice QCD and HRG Model. 39 GeV data starts to deviate from HRG model.
- Non-monotonic signals for CP are not observed at $\mu_B < 200$ MeV region.

Scaling Properties for $S\sigma$



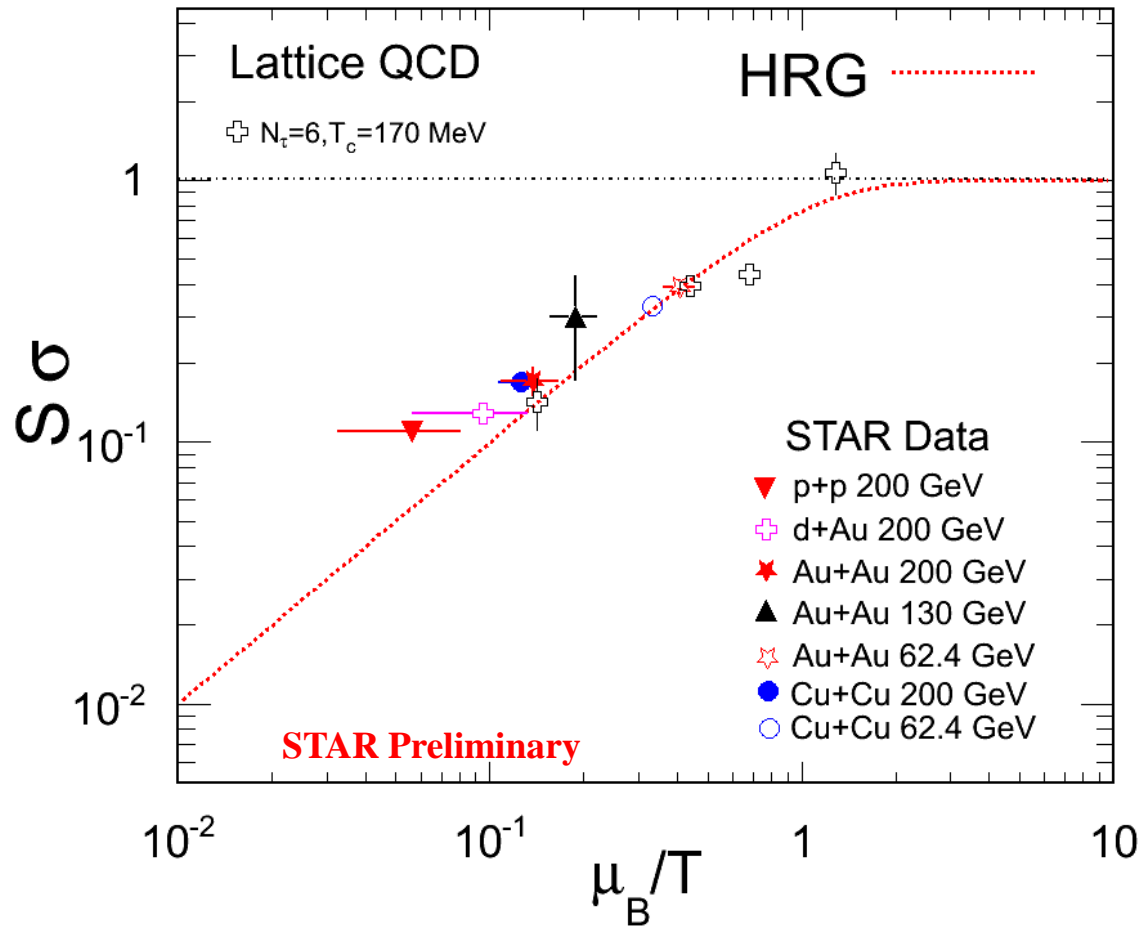
Fit: $S_B \sigma_B = \frac{11}{3} * \left(\frac{1}{s^4} \frac{dN_{ch}}{d\eta} \right)^{\frac{1}{12}}$ **At high energy :**

HRG: $S_B \sigma_B = \tanh\left[\frac{\mu_B}{T}\right]$ $\mu_B/T \ll 1$

$\frac{\mu_B}{T} \approx \tanh\left[\frac{\mu_B}{T}\right] = \frac{11}{3} \left(\frac{1}{s^4} \frac{dN_{ch}}{d\eta} \right)^{\frac{1}{12}}$

➤ Reflect connection between fluctuations, thermodynamic parameter and charged particle density.

s: Square of center of mass energy.
 $dN_{ch}/d\eta$: Charged particle density.



Caveat: p+p and d+Au system may not be described by Grand Canonical Ensemble.

- All HI results are consistent with the thermal model prediction, except the small systems from p+p and d+Au collisions.
- The issue of thermalization is discussed with fluctuation data from high-energy nuclear collisions at RHIC for the first time.

Summary and Outlook

- Higher moments and moment products are expected to be sensitive to the QCD critical point related correlation and fluctuations.
- $S\sigma$ and $\kappa\sigma^2$ are found to be consistent with Lattice QCD and HRG model for high energy data. Non-monotonic signals for QCD critical point are not observed at $\mu_B < 200$ MeV region.
- For the first time, we address the issue of thermalization with high order fluctuation data from high-energy nuclear collisions at RHIC.

Outlook:

1. Run 11 are running. More BES data are coming....
2. TOF will be used for PID in future analysis.