

# QCD Critical Point and Net-Proton Number Fluctuations at RHIC-STAR

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In the search of QCD phase boundary and critical point, higher-order cumulants of conserved quantities are proposed as promising observables and have been studied extensively both experimentally and theoretically. In this paper we present cumulant ratios up-to 6<sup>th</sup>-order of net-proton number distributions in Au+Au collisions at  $\sqrt{s_{NN}} = 7.7 - 200$  GeV from STAR Beam Energy Scan program phase I and  $\sqrt{s} = 200$  GeV  $p + p$  collisions. The results are compared with various models and Lattice QCD calculations.

## 1 Introduction

Quantum chromodynamics (QCD) is the theory which describes the strong interactions between quarks and gluons. It is predicted that under a very high temperature and baryon density a deconfined quark-gluon plasma (QGP) phase can be created. Studying the QCD phase structure is one of the main goals in heavy-ion collision physics. A phase diagram in terms of temperature and baryon chemical potential ( $\mu_B$ ) is usually used to explore the QCD phase structure. Regarding the phase transition between the QGP phase and hadronic phase, first principle Lattice QCD calculation<sup>1</sup> at  $\mu_B = 0$  MeV suggests that the phase transition is a smooth crossover. While at large  $\mu_B$ , various QCD-based models predict first order phase transition<sup>2</sup>. Thermodynamically there should be an end point of the first order phase boundary which is called QCD critical point. The possible QCD critical point and first order phase boundary have been investigated both experimentally and theoretically.

Higher-order cumulants of conserved quantities like net-baryon number, net-charge number and net-strangeness number are proposed as promising observables to search for the QCD critical point and the first order phase boundary. Higher-order cumulants are sensitive to the correlation length ( $\xi$ )<sup>3,4</sup> and are directly related to susceptibility ( $\chi$ ) of the system<sup>5</sup>. It is predicted that the fourth-order fluctuations will exhibit a non-monotonic energy dependence<sup>6,7,8</sup> when passing through the critical region. For 5th- and 6th-order cumulants recent calculations from Lattice QCD<sup>9</sup> and the functional renormalisation group approach (FRG)<sup>10</sup> show that they will be negative due to the crossover transition between QGP and hadronic phase. At high baryon density region, on the other hand, they are also sensitive to the 1<sup>th</sup>-order phase boundary<sup>11,12</sup>.

## 2 Cumulants

This section discusses cumulant definition. Let  $N$  represent conserved quantity like net-proton number from data sample. The deviation from its mean value ( $\langle N \rangle$ ) is defined as  $\delta N = N - \langle N \rangle$ . Then cumulants up-to 6<sup>th</sup>-order can be written as:

$$\begin{aligned} C_1 &= \langle N \rangle, \\ C_2 &= \langle (\delta N)^2 \rangle, \\ C_3 &= \langle (\delta N)^3 \rangle, \\ C_4 &= \langle (\delta N)^4 \rangle - 3\langle (\delta N)^2 \rangle^2, \\ C_5 &= \langle (\delta N)^5 \rangle - 10\langle (\delta N)^3 \rangle \langle (\delta N)^2 \rangle, \\ C_6 &= \langle (\delta N)^6 \rangle - 15\langle (\delta N)^4 \rangle \langle (\delta N)^2 \rangle - 10\langle (\delta N)^3 \rangle^2 + 30\langle (\delta N)^2 \rangle^3. \end{aligned} \tag{1}$$

Various cumulant ratios like  $C_3/C_2$ ,  $C_4/C_2$ ,  $C_5/C_1$  and  $C_6/C_2$  are constructed to cancel volume effects and to readily make comparison with ratios of susceptibility ( $\chi_n$ ). The cumulant ratios

31  $C_3/C_2$ ,  $C_4/C_2$  are also named  $S\sigma$  and  $\kappa\sigma^2$  respectively.

### 32 3 Analysis details

33 The data are collected in STAR Beam Energy Scan program phase I Au+Au collisions at  $\sqrt{s_{NN}} = 7.7 - 200$  GeV and  $p + p$  collisions at  $\sqrt{s} = 200$  GeV. Protons and antiprotons are identified by the Time Projection Chamber (TPC) and Time of flight (TOF) detectors at rapidity window 35  $-0.5 < y < 0.5$  and transverse momentum window  $0.4 < p_T < 2.0$  GeV/c. At  $p_T < 0.8$  GeV/c 37 only TPC particle identification cuts are used and at  $p_T > 0.8$  GeV/c additional TOF particle 38 identification cuts are used to ensure proton purity.

39 The centrality is determined using charged particle multiplicity within  $|\eta| < 1.0$  excluding 40 protons and antiprotons to avoid auto-correlation effect<sup>14</sup>. The centrality bin width correction 41<sup>14</sup> is applied to suppress initial volume fluctuation effect. Cumulants are calculated at each 42 multiplicity bin and then their weighted averages are taken for each centrality bin. The weight 43 is number of events at the corresponding multiplicity bin. Detector efficiency correction<sup>13</sup> in 44 cumulant calculations are done by assuming binomial detector efficiency. Statistical uncertainties 45 of cumulants are estimated by Bootstrap and Delta methods<sup>15</sup>.

### 46 4 Results

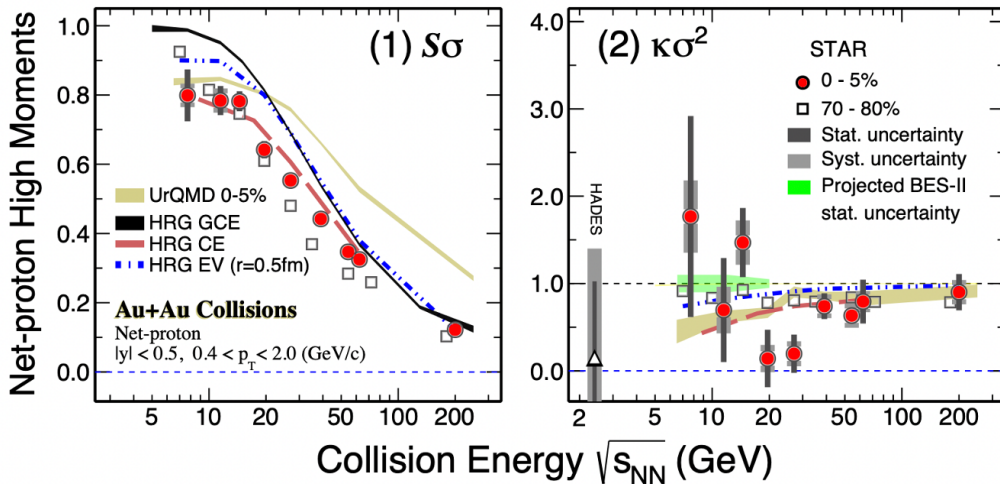


Figure 1 – Energy dependence of  $S\sigma$  and  $\kappa\sigma^2$  of net-proton number distributions in Au+Au collisions at  $\sqrt{s_{NN}} = 7.7 - 200$  GeV. The calculations from different variants (GCE, EV, CE) of hadron resonance gas model (HRG) and the hadronic transport UrQMD model are shown as black, red, blue bands and a gold band respectively.

47 Figure 1 shows energy dependence of  $S\sigma$  and  $\kappa\sigma^2$  of net-proton distributions from 0-5% and 48 70-80% centrality bins within  $|y| < 0.5$  and  $0.4 < p_T < 2.0$  GeV/c in Au+Au collisions at  $\sqrt{s_{NN}} = 7.7 - 200$  GeV.<sup>18,19</sup> The  $S\sigma$  (left panel) shows a decreasing trend with the increase of collision 49 energy both in central and peripheral collisions. The decreasing trend can be qualitatively de- 50 scribed by HRG<sup>16</sup> and UrQMD<sup>17</sup> models. The  $\kappa\sigma^2$  (right panel) shows a non-monotonic energy 51 dependence in central collisions while peripheral collisions shows no non-monotonic energy de- 52 pendence. The non-monotonic trend in central collisions is not qualitatively described by HRG 53 and UrQMD models.

54 Figure 2 shows energy dependence of  $C_5/C_1$  and  $C_6/C_2$  of net-proton distributions from 55 0-40% and 70-80% centrality bins within  $|y| < 0.5$  and  $0.4 < p_T < 2.0$  GeV/c in Au+Au 56 collisions at  $\sqrt{s_{NN}} = 7.7 - 200$  GeV. It is suggested from Lattice QCD and FRG calculations 57 that 5<sup>th</sup>- and 6<sup>th</sup>-order cumulants show negative sign while calculations from UrQMD and HRG 58

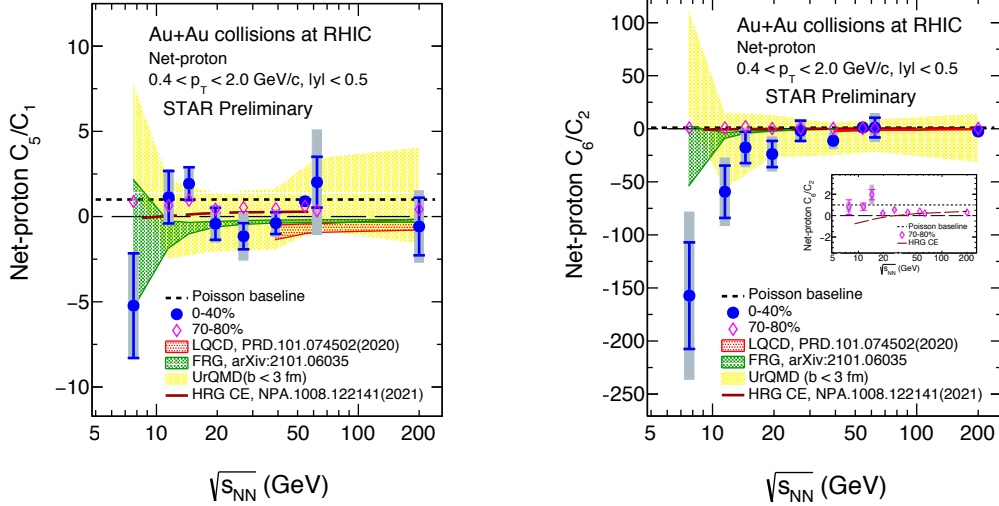


Figure 2 – Energy dependence of  $C_5/C_1$  and  $C_6/C_2$  of net-proton number distributions in Au+Au collisions at  $\sqrt{s_{NN}} = 7.7 - 200$  GeV. The calculations from FRG model, Lattice QCD, hadronic transport UrQMD model and Hadron Resonance Gas (HRG) model are shown as red, green and yellow bands and red lines respectively.

59 models are consistent with either zero or unity. In those models, no phase transition physics  
60 is implemented. The measurements of BES-I data are shown as blue dots for 0-40% and red  
61 diamonds for 70-80%. The cumulant ratio  $C_5/C_1$  (left panel) shows negative sign at  $\sqrt{s_{NN}} =$   
62 7.7, 19.6, 27, 39 and 200 GeV and shows positive sign at  $\sqrt{s_{NN}} = 11.5, 14.5, 54.4$  and 62.4 GeV.  
63 In peripheral collisions  $C_5/C_1$  shows positive sign for all energies. The ratio  $C_6/C_2$  (right panel)  
64 for 0-40% is increasingly negative with decreasing energy with less than  $2\sigma$  significance while it  
65 shows positive sign in peripheral collisions (70-80%) for all energies.

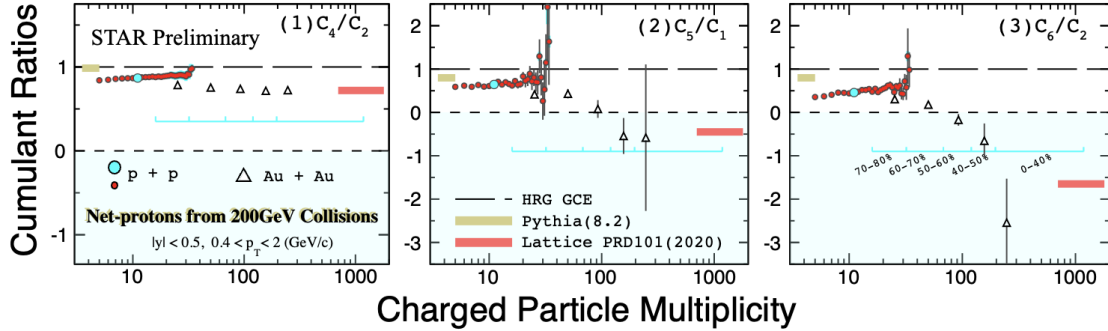


Figure 3 – Multiplicity dependence of  $C_4/C_2$ ,  $C_5/C_1$  and  $C_6/C_2$  of net-proton number distributions in  $p + p$  collisions at  $\sqrt{s} = 200$  GeV. The calculations of HRG model, Pythia (8.2) and Lattice QCD are shown as black dashed line, yellow and red bands respectively.

66 Figure 3 shows multiplicity dependence of net-proton results of  $C_5/C_1$  and  $C_6/C_2$  within  
67  $|y| < 0.5$  and  $0.4 < p_T < 2.0$  GeV/c in  $p + p$  collisions at  $\sqrt{s} = 200$  GeV. We see that the  
68 cumulant ratios ( $C_4/C_2$ , The cumulant ratios  $C_5/C_1$  and  $C_6/C_2$ ) from  $p + p$  collisions fit into  
69 the multiplicity dependence of results from Au+Au collisions which are shown with triangles.  
70  $C_5/C_1$  and  $C_6/C_2$  are negative for 0-40% in Au+Au collisions and positive at peripheral Au+Au  
71 collisions and  $p + p$  collisions. Pythia<sup>20</sup> calculation using version 8.2 of  $p + p$  collisions at  $\sqrt{s} =$   
72 200 GeV is positive as shown with yellow bands. The Lattice QCD calculation<sup>9</sup> at  $\sqrt{s_{NN}} = 200$   
73 GeV is negative as shown with red bands. Compared with calculations from various models, it  
74 is suggested that the negative sign for central Au+Au collisions at  $\sqrt{s_{NN}} = 200$  GeV is due to

75 crossover phase transition between partonic and hadronic phases.

## 76 **5 Summary**

77 In this proceedings, we report the measurements of net-proton cumulant ratios up-to 6<sup>th</sup>-order  
78 in Au+Au at  $\sqrt{s_{NN}} = 7.7 - 200$  GeV and  $p + p$  collisions at  $\sqrt{s_{NN}} = 200$  GeV at STAR. With  
79 results from 200 GeV  $p + p$  collisions and the energy dependence of  $C_4/C_2$ ,  $C_5/C_1$  and  $C_6/C_2$   
80 from the BES-I data sets, and the comparison with LQCD calculations, we conclude:

- 81 1. QCD matter is indeed created in the 200 GeV central (0-5%) Au+Au collisions at RHIC.
- 82 2. Non-monotonic energy dependence of  $C_4/C_2$  is observed from the most central (0-5%)  
83 Au+Au collisions.<sup>18,19</sup>

84 Future results from BES-II and STAR fixed-target experiment  $\sqrt{s_{NN}} = 3$ GeV data sets will allow  
85 to answer if QCD critical point exists in the covered energy region.

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