

1 Higher-order Proton Cumulants in Au+Au Collisions at
2 $\sqrt{s_{\text{NN}}} = 3$ GeV from RHIC-STAR*

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9 In these proceedings, we present the higher-order cumulants of proton
10 multiplicity distributions of the fixed-target (FXT) run in Au+Au collisions
11 at $\sqrt{s_{\text{NN}}} = 3.0$ GeV. The cumulant ratios are presented as a function of
12 centrality and collision energy. The proton cumulant ratio C_4/C_2 is consistent
13 with fluctuations driven by baryon number conservation and indicates
14 an energy regime dominated by hadronic interactions. These data imply
15 that the QCD critical point could exist at energies higher than 3 GeV if
16 created in heavy-ion collisions.

17 **1. Introduction**

18 Experimental evidences [1] at RHIC and the LHC have demonstrated
19 the formation of Quark-Gluon Plasma (QGP) in ultra-relativistic heavy-ion
20 collisions at small baryon chemical potential ($\mu_B \approx 0$ MeV) where the phase
21 transition from the hadronic matter to QGP is suggested to be a crossover
22 from state-of-the-art Lattice QCD calculations [2]. It has been conjectured
23 that there is a first-order phase transition and a QCD critical point at the
24 finite μ_B region in the QCD phase diagram. In the search for the possible
25 QCD critical point, higher-order cumulants of conserved quantities
26 such as net-baryon number, net-strangeness number, and net-charge number
27 are sensitive observables to locate its position [3, 4, 5, 6]. Experimentally
28 net-proton and net-kaon numbers are used as a proxy for net-baryon and
29 net-strangeness numbers due to the difficulty to detect neutral particles in
30 the experiment. Recent results from the STAR experiment on net-proton
31 fourth-order cumulant ratio have shown intriguing non-monotonic energy

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dependence with 3.1σ significance in the most central Au+Au collisions at $\sqrt{s_{\text{NN}}} = 7.7 - 200$ GeV [7, 8] while there are still large statistical uncertainties for energy $\sqrt{s_{\text{NN}}} < 19.6$ GeV. These proceedings reports proton cumulants and cumulant ratios up to 4th-order in $\sqrt{s_{\text{NN}}} = 3$ GeV Au+Au collisions from the STAR fixed-target experiment. The relevant analysis details and correction methods will also be shortly discussed. To understand the collision dynamics in the absence of the critical behavior, we have carried out simulations with a microscopic transport model UrQMD [9] for Au+Au collisions at $\sqrt{s_{\text{NN}}} = 3$ GeV. Connections between experimental data and physics implications in the high baryon density region will be discussed.

2. Experimental Observables

This section shows the definitions of cumulants and cumulant ratios. Let N represent net-proton number. The deviation from its mean value ($\langle N \rangle$) is defined as $\delta N = N - \langle N \rangle$. Then cumulants up to 4th-order can be written as:

$$\begin{aligned} C_1 &= \langle N \rangle, C_2 = \langle (\delta N)^2 \rangle, C_3 = \langle (\delta N)^3 \rangle, \\ C_4 &= \langle (\delta N)^4 \rangle - 3\langle (\delta N)^2 \rangle^2. \end{aligned} \quad (1)$$

The cumulants are related to the various moments as

$$M = C_1, \quad \sigma^2 = C_2, \quad S = \frac{C_3}{(C_2)^{3/2}}, \quad \kappa = \frac{C_4}{C_2^2}, \quad (2)$$

where M , σ^2 , S , and κ are mean, variance, skewness, and kurtosis, respectively. Various cumulant ratios like C_2/C_1 , C_3/C_2 , and C_4/C_2 are constructed to cancel volume dependence:

$$\frac{C_2}{C_1} = \sigma^2/M, \quad \frac{C_3}{C_2} = S\sigma, \quad \frac{C_4}{C_2} = \kappa\sigma^2. \quad (3)$$

3. Analysis Details

The analysis used around 140 million $\sqrt{s_{\text{NN}}} = 3$ GeV Au+Au collisions events which are collected by the dedicated physics fixed-target run of the STAR experiment in the year 2018. The centrality is determined using charged particle reference multiplicity excluding protons and light nuclei within $-2 < \eta < 0$ where η is pseudo-rapidity in the lab frame. As shown in Fig. 1 protons are identified by comparing the energy loss measured by the Time Projection Chamber (TPC) with theoretical predictions (Fig. 1(a)). At high momentum ($p_{\text{lab}} > 2$ GeV/c), due to the contamination from other particles, the mass square measured by Time of Flight (TOF) is used to

61 ensure proton purity (Fig. 1(b)). The anti-protons are negligible ($\bar{p}/p <$
 62 10^{-6}) at $\sqrt{s_{\text{NN}}} = 3$ GeV thus the proton cumulants are measured in the
 63 analysis. Figure 1(c) shows proton acceptance with the combination of TPC
 64 and TOF. The red dashed box indicates the acceptance window used in this
 65 analysis.

66 Cumulants are corrected for detector efficiencies by a track-by-track
 67 method [10, 11]. The rapidity (y) and transverse-momentum (p_{T}) depen-
 68 dences of detector efficiency are considered. To correct the pileup effect due
 69 to the finite thickness of the gold target, a pileup correction method [12, 13]
 70 is used. As seen in our model simulation, there is a large initial volume
 71 fluctuation effect when calculating cumulants at $\sqrt{s_{\text{NN}}} = 3$ GeV, thus we
 72 tested an initial volume fluctuation correction method [14]. We measured
 73 cumulants as a function of reference multiplicity, and then obtained central-
 74 ity binned results by the Centrality Bin Width Correction (CBWC) [15].
 75 The statistical uncertainties of cumulants are estimated by the bootstrap
 76 method. The systematic uncertainties are estimated by varying analysis cuts
 77 related to centrality, pileup effect, track quality, and detector efficiency.

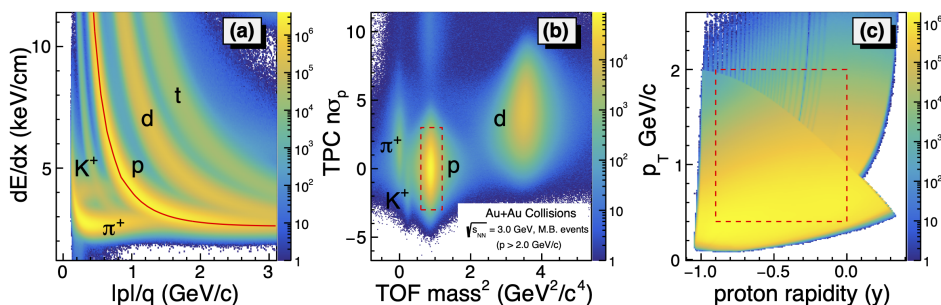


Fig. 1. Panel (a): TPC track energy loss (dE/dx (keV/cm)) vs. momentum; pion, kaon, deuteron and triton are labeled. The proton Bethe-Bloch curve is plotted with red line. Panel (b): TPC $n\sigma_p$ vs. TOF mass². Panel (c): Transverse momentum (p_{T}) vs. proton rapidity.

78 4. Results

79 Figure 2 shows the centrality dependence of the proton cumulant ratios
 80 C_2/C_1 , C_3/C_2 , and C_4/C_2 within $-0.5 < y < 0$ and $0.4 < p_{\text{T}} < 2.0$ GeV/ c .
 81 The 3 GeV data shown with black open squares are corrected for detector
 82 efficiency and pileup effect and then the CBWC was applied to obtain cen-
 83 trality binned results. The red circles and blue triangles are additionally

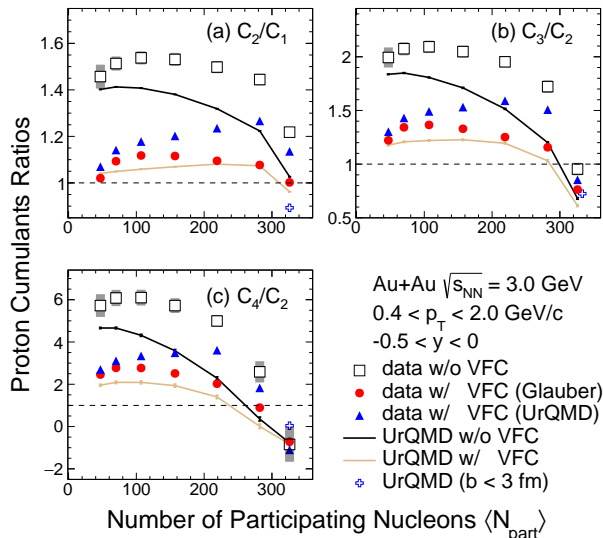


Fig. 2. Centrality dependence of proton cumulants and cumulant ratios up to 4th-order in Au+Au collisions at $\sqrt{s_{\text{NN}}} = 3$ GeV within kinematic acceptance $-0.5 < y < 0$ and $0.4 < p_{\text{T}} < 2.0$ GeV/c. The black squares are results without volume correction while red circles and blue triangles represent results with volume correction using Glauber and UrQMD model, respectively.

84 corrected for initial volume fluctuation using Glauber and UrQMD models,
 85 respectively. It is clear that the volume fluctuation correction shows a strong
 86 model dependence and affects the centrality dependence, particularly in pe-
 87 ripheral collisions. The respective dynamics in the UrQMD and Glauber
 88 model for charged hadron production lead to two different mappings from
 89 the measured final charged hadron multiplicity distributions to the initial
 90 geometry. This difference is likely to be the dominant source of the model
 91 dependence in the VFC. On the other hand, one can see in the figure that
 92 the difference between results with and without the VFC is small for higher
 93 order ratios C_3/C_2 and C_4/C_2 in the most central bin.

94 Figure 3 shows the collision energy dependence of cumulant ratio C_4/C_2
 95 of net-proton and proton multiplicity distributions in central Au+Au col-
 96 lisions [16]. As reported in Refs. [8, 7] the net-proton and proton C_4/C_2
 97 show a non-monotonic energy trend in central Au+Au collisions. A min-
 98 imum is seen at around $\sqrt{s_{\text{NN}}} = 20$ GeV and then C_4/C_2 becomes close
 99 to unity with large statistical uncertainty when decreasing collision energy.
 100 The new measurement of proton C_4/C_2 for $\sqrt{s_{\text{NN}}} = 3$ GeV central Au+Au

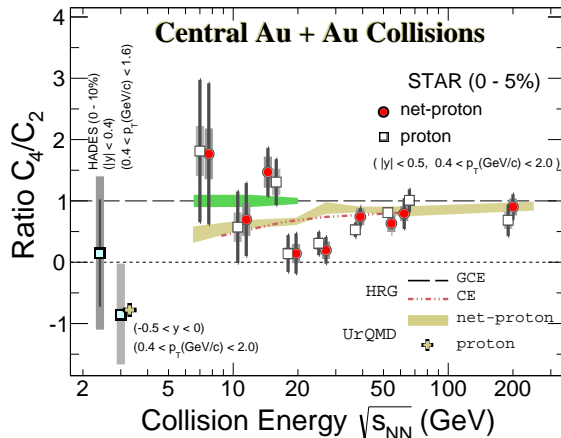


Fig. 3. Collision energy dependence of cumulants ratio C_4/C_2 in central Au+Au collisions within kinematic acceptance cut $0.4 < p_T < 2.0$ GeV/c. UrQMD calculations are shown with gold band for net-proton with rapidity cut $|y| < 0.5$ and gold filled cross for proton with rapidity cut $-0.5 < y < 0$. Statistical and systematic uncertainty are shown with black and grey bars, respectively. The green shaded area indicates the projected statistical uncertainty with BES-II data.

101 collisions is around -1, which is reproduced by the hadronic transport model
 102 UrQMD while at higher energies the non-monotonic energy dependence is
 103 not reproduced by various non-critical models including the UrQMD and
 104 HRG [17] models. Precision data in the energy window of $3 < \sqrt{s_{NN}} < 20$
 105 GeV are needed in order to explore the possibility of critical phenomena.
 106 The HADES collaboration has reported the proton cumulant ratio in $\sqrt{s_{NN}}$
 107 = 2.4 GeV Au+Au collisions within acceptance window $0.4 < p_T < 1.6$
 108 GeV/c and $|y| < 0.4$: $C_4/C_2 = 0.15 \pm 0.9$ (stat.) ± 1.4 (sys.) [18] which is
 109 consistent with 3 GeV result within uncertainty although detailed compar-
 110 ison should be done within same acceptance.

111 5. Summary

112 In this paper, we reported cumulant ratios of proton multiplicity dis-
 113 tributions in $\sqrt{s_{NN}} = 3$ GeV Au+Au collisions by the STAR fixed-target
 114 experiment. The proton C_4/C_2 is observed to be -0.85 ± 0.09 (stat.) ± 0.82
 115 (sys.) in the most central 0-5% centrality at $\sqrt{s_{NN}} = 3$ GeV. Compared to
 116 higher energy results and the transport model calculations, the suppression
 117 in C_4/C_2 is consistent with fluctuations driven by baryon number conser-

118 vation and indicates an energy regime dominated by hadronic interactions,
 119 which implies that the QCD critical point could exist at energies higher than
 120 3 GeV if discovered in heavy-ion collisions. New data sets have been col-
 121 lected during the second phase of the RHIC beam energy scan program for
 122 Au+Au collisions at $\sqrt{s_{NN}} = 3 - 19.6$ GeV. The analysis on those datasets
 123 will be crucial in exploring the QCD phase structure at high baryon density
 124 region and locating the critical point.

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