

Results from the BES program at STAR

Hanna Zbroszczyk for the STAR Collaboration
hanna.zbroszczyk@pw.edu.pl

Warsaw University of Technology

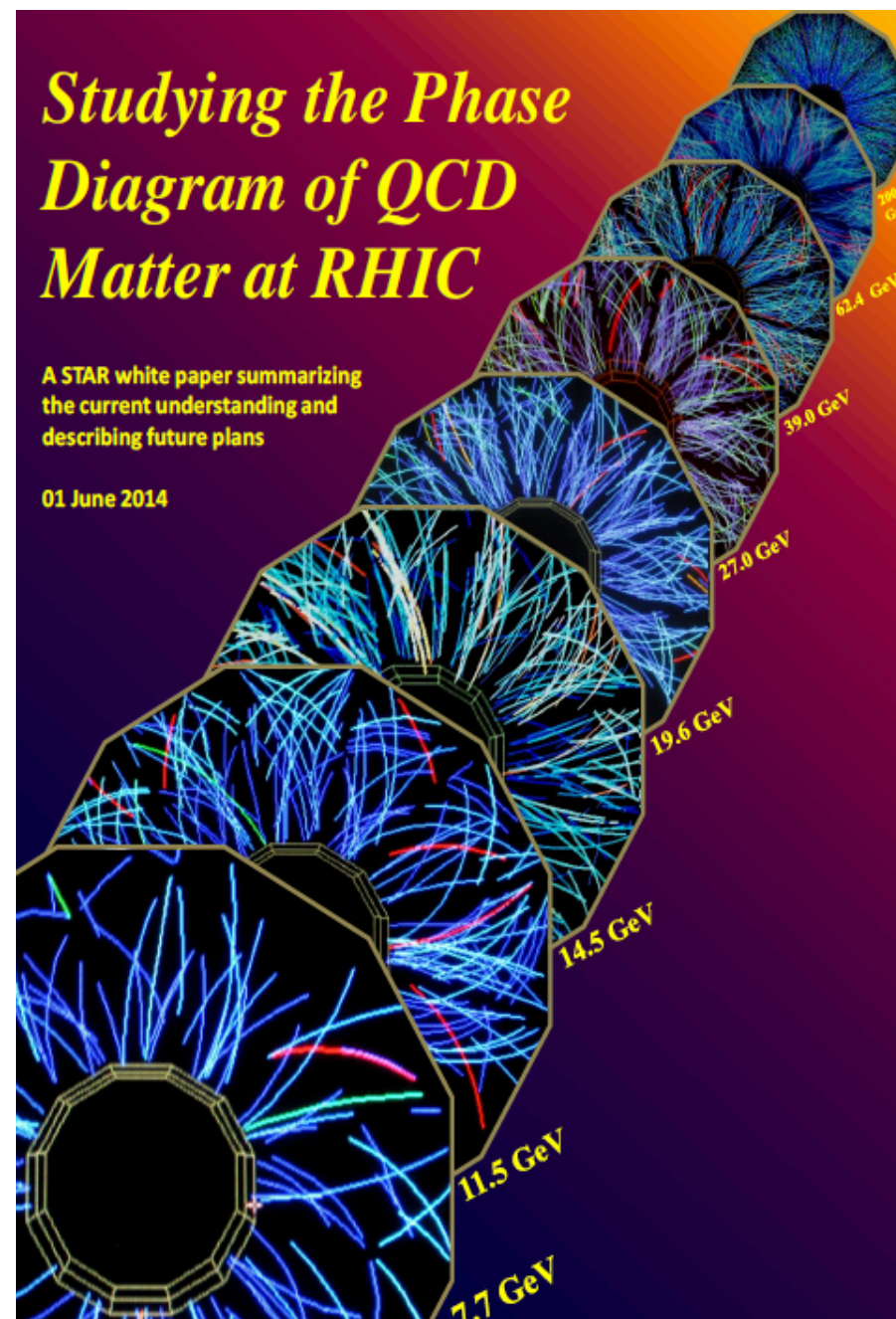


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Outline

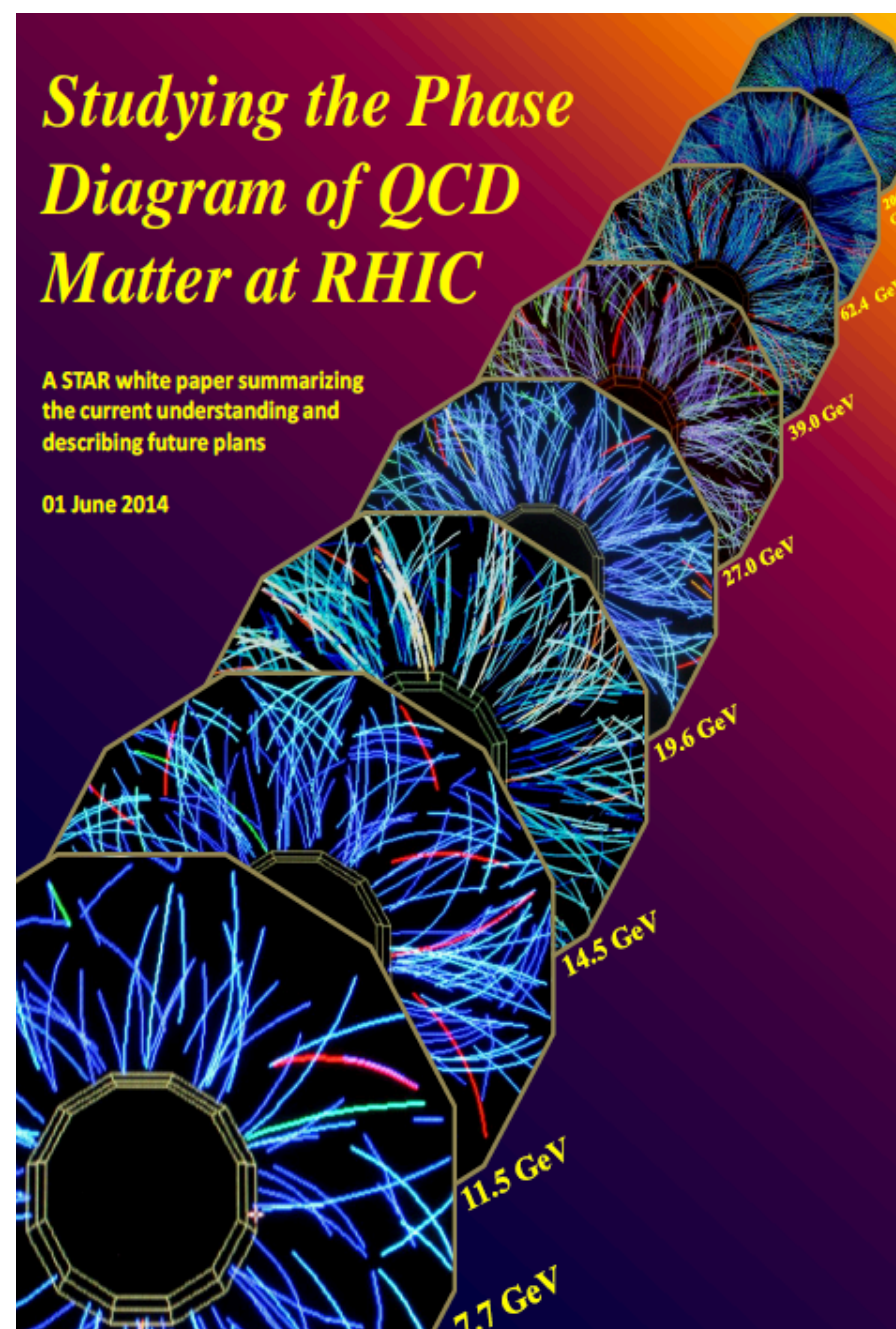
- Beam Energy Scan Program
- QCD phase diagram
- Experiments
- Main questions
- Summary and perspectives



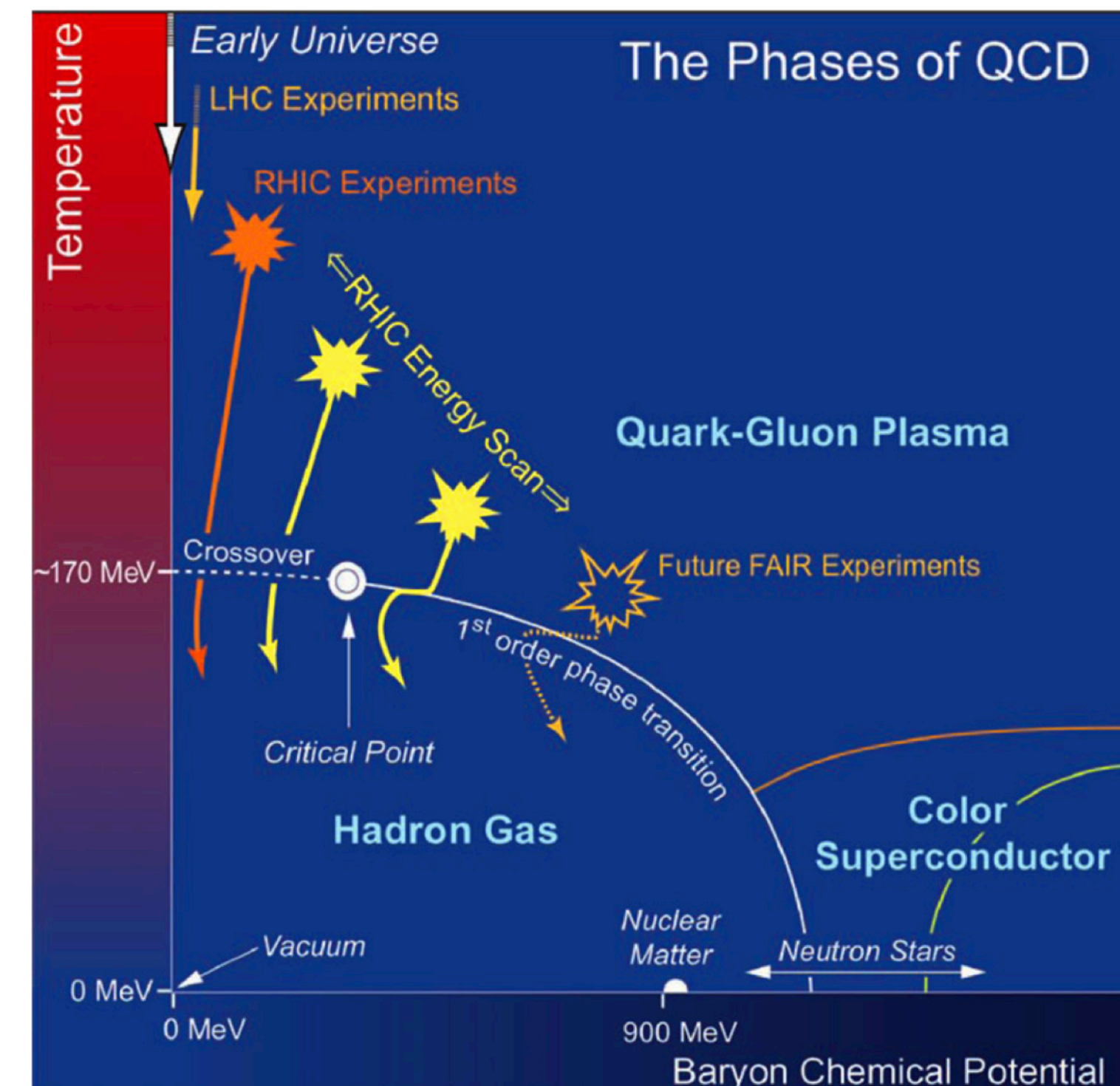
Outline

- Beam Energy Scan Program
- QCD phase diagram
- Experiments
- **Main questions**
- Summary and perspectives

- Goals of the Beam Energy Scan Program:**
1. Search for **turn-off** of **QGP** signatures
 2. Search for signals of the **first-order phase transition**
 3. Search for QCD **critical point**
 4. Search for signals of **Chiral symmetry restoration**



<http://drupal.star.bnl.gov/STAR/starnotes/public/sn0598>



Probing QCD Phase Diagram with Heavy-ion collisions

Heavy-ion collision used as a tool to probe QCD phase diagram

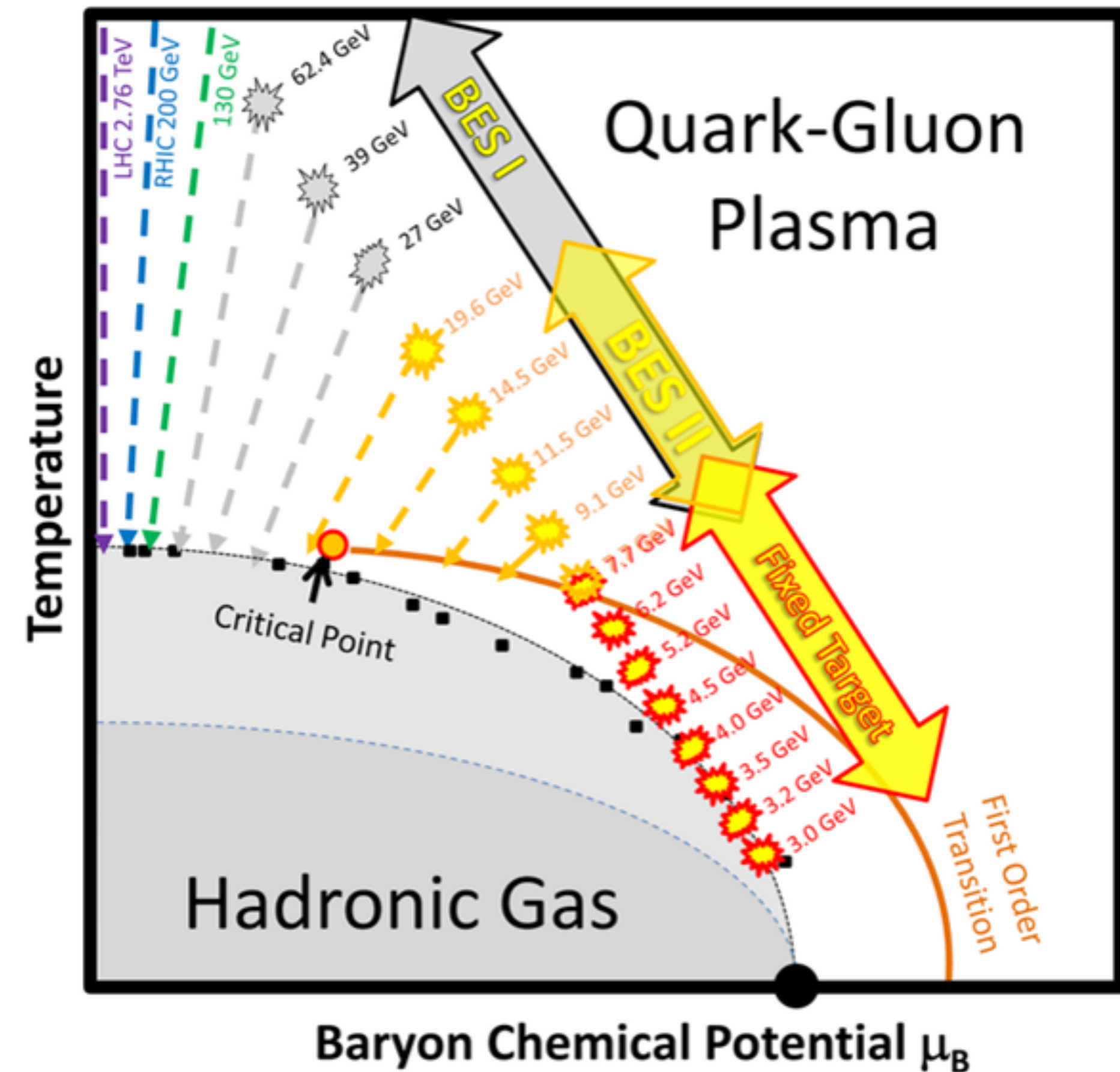
Believed to be understood:

Lattice QCD predicts a smooth **cross-over** transition at large T and $\mu_B \sim 0$

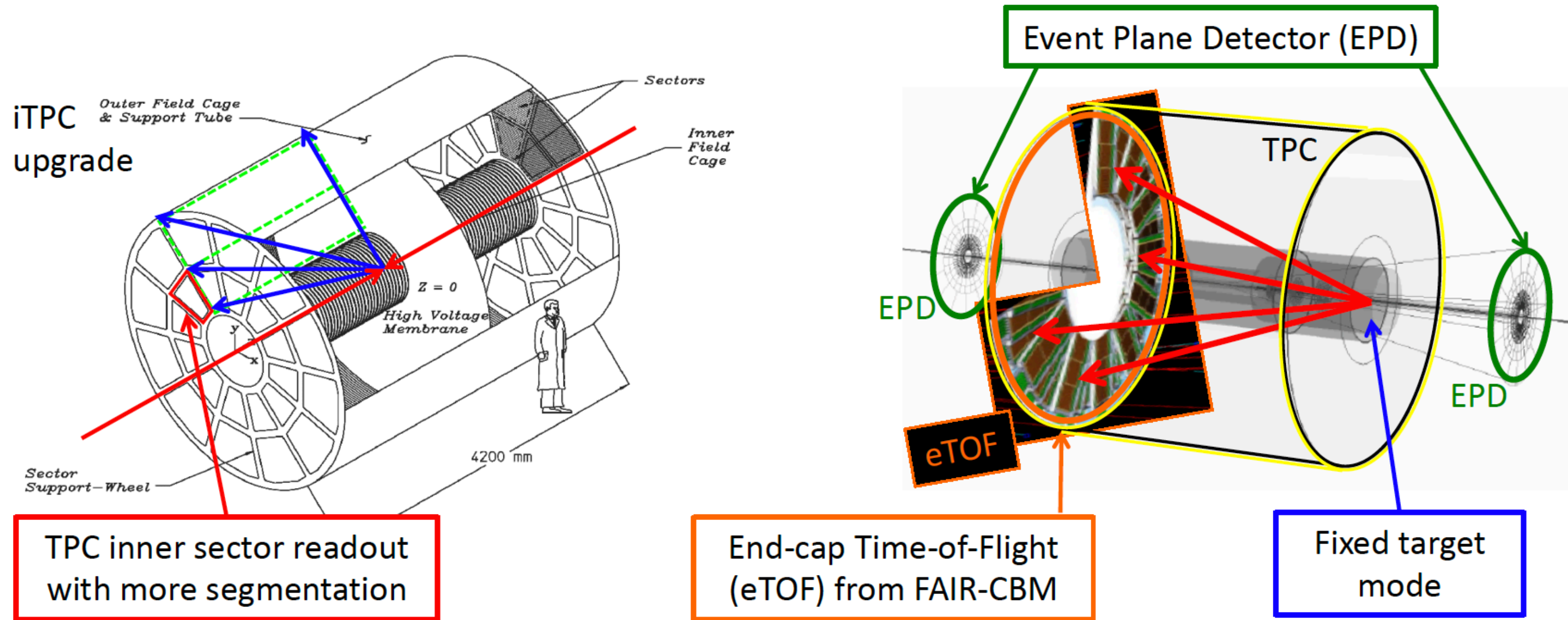
Various models predict **first-order phase transition** at large μ_B

Critical point is believed to exist, but.. where?

Strategy: to map the phase diagram (μ_B, T) using heavy-ion collisions changing their collision energy: BES-I, BES-II (+FXT)



STAR detector system



TPC inner sector readout with more segmentation

End-cap Time-of-Flight (eTOF) from FAIR-CBM

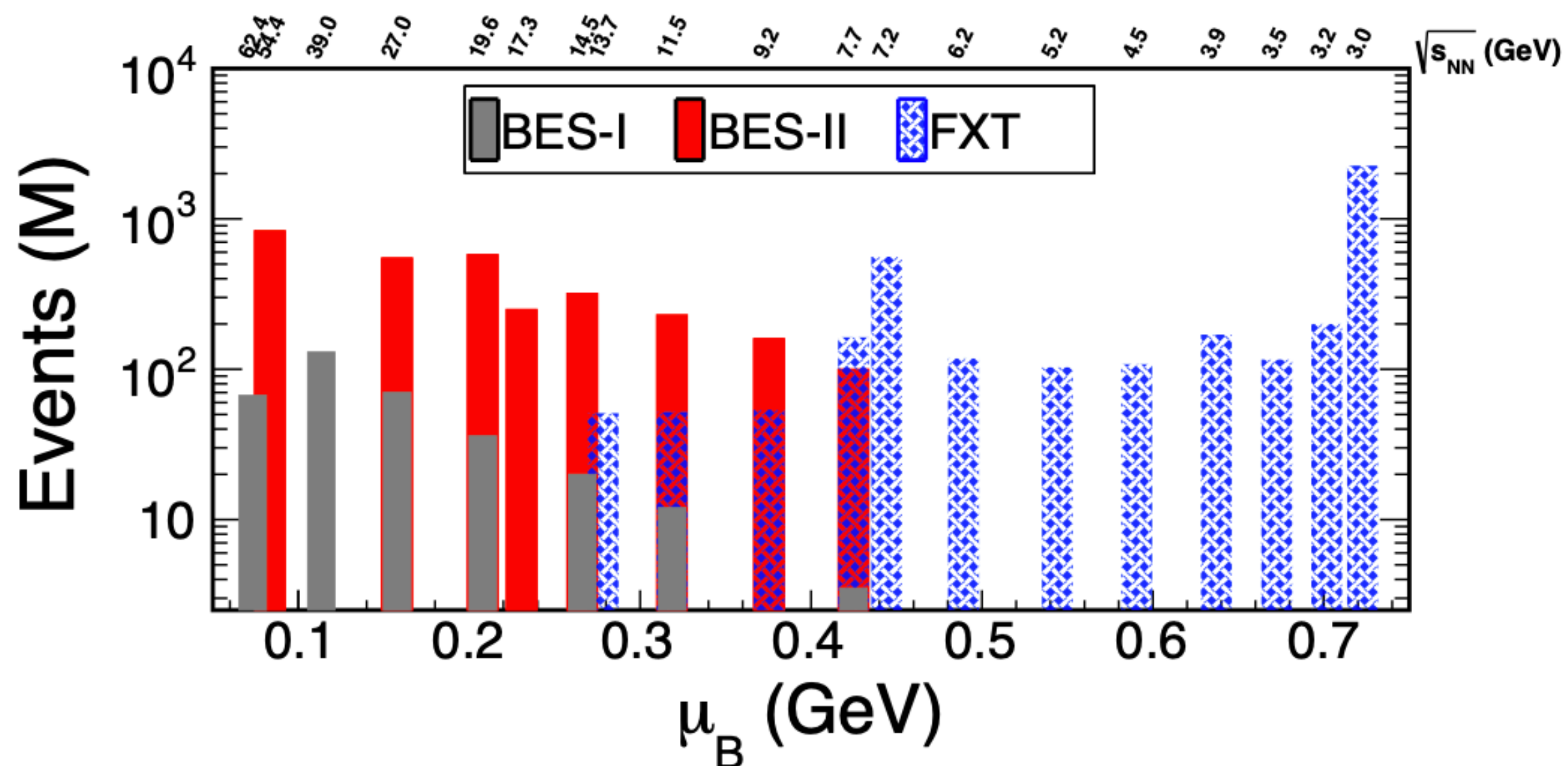
Fixed target mode

Solenoidal Tracker At RHIC originally designed to search for Quark Gluon Plasma.

BES program started at 2010.

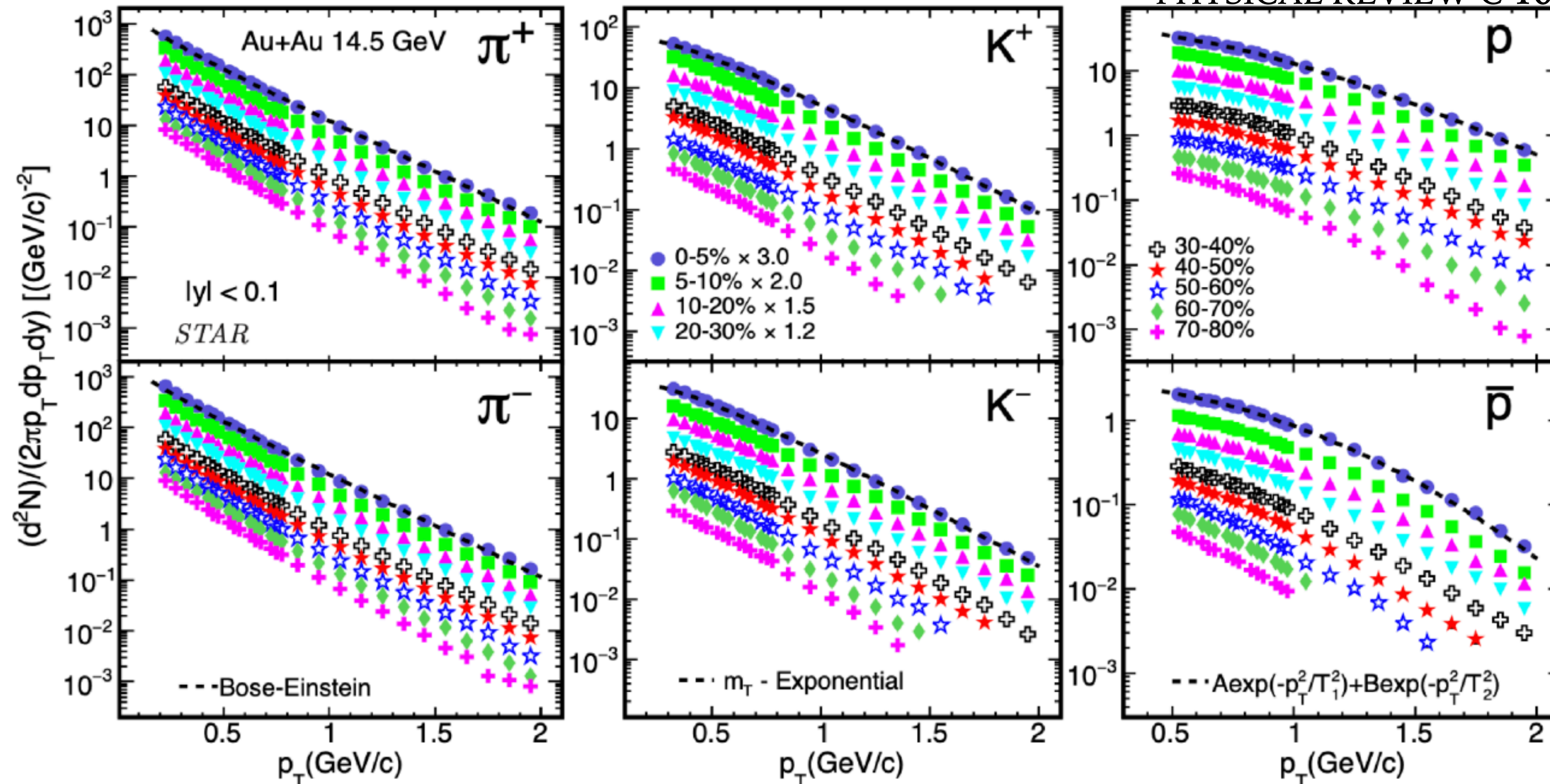
Luminosity of the RHIC collider-mode is unusable for $\sqrt{s_{NN}} < 7.7$ GeV.

Fixed-target (FXT) program extends the collision energy and μ_B coverage.



Particle spectra

PHYSICAL REVIEW C 101, 024905 (2020)

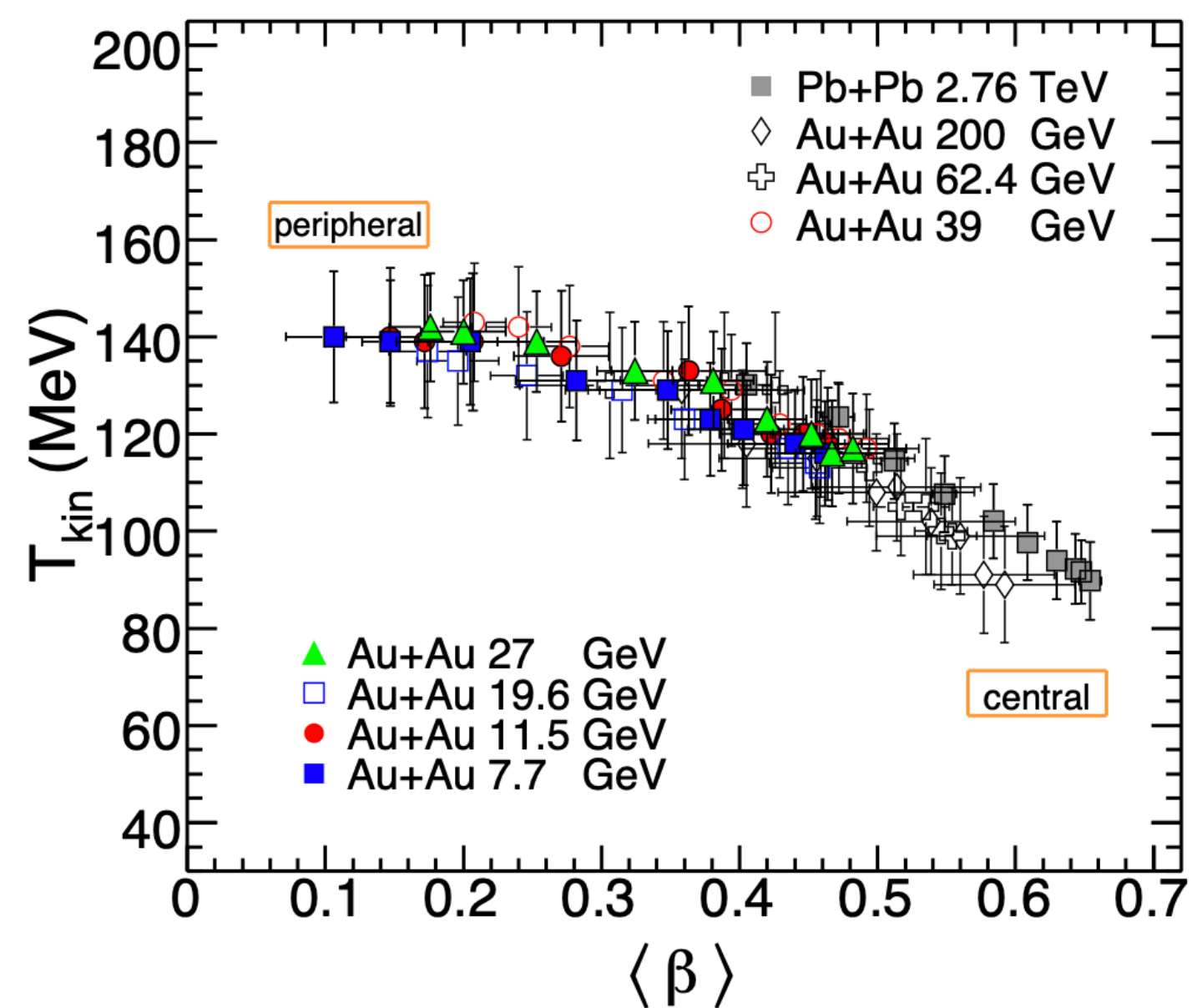


Inverse slopes of the identified hadron spectra follow the order $\pi < K < p$

$$\frac{dN}{m_T dm_T dy} = f(y) \exp\left(-\frac{m_T}{T}\right); \quad m_T = \sqrt{m^2 + p_T^2}$$

Chemical and kinetic freeze-out parameters

PRC 96 (2017) 4, 044904

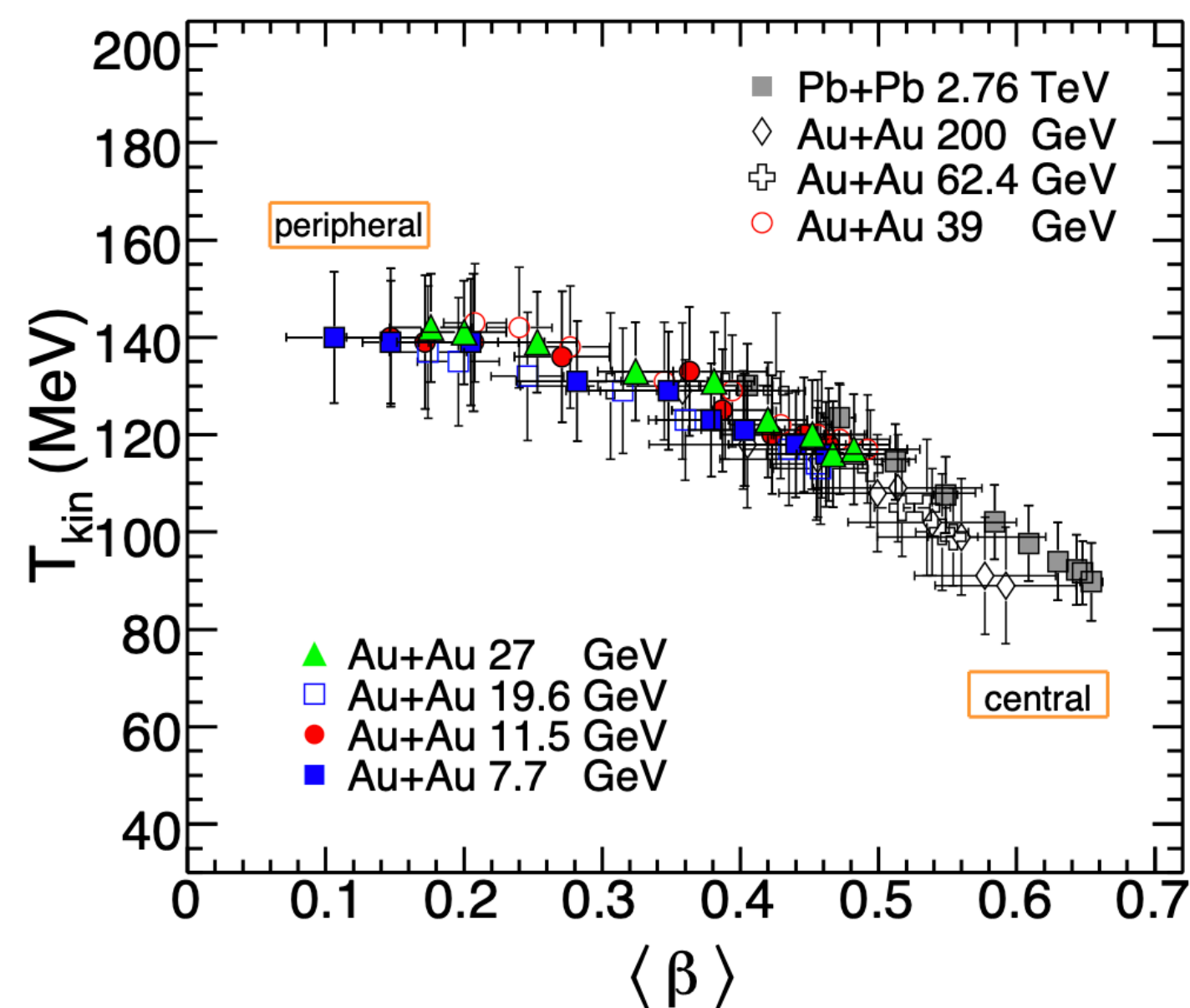


Extracted from spectra:

$$m_T - m \text{ of } \pi, K, p$$

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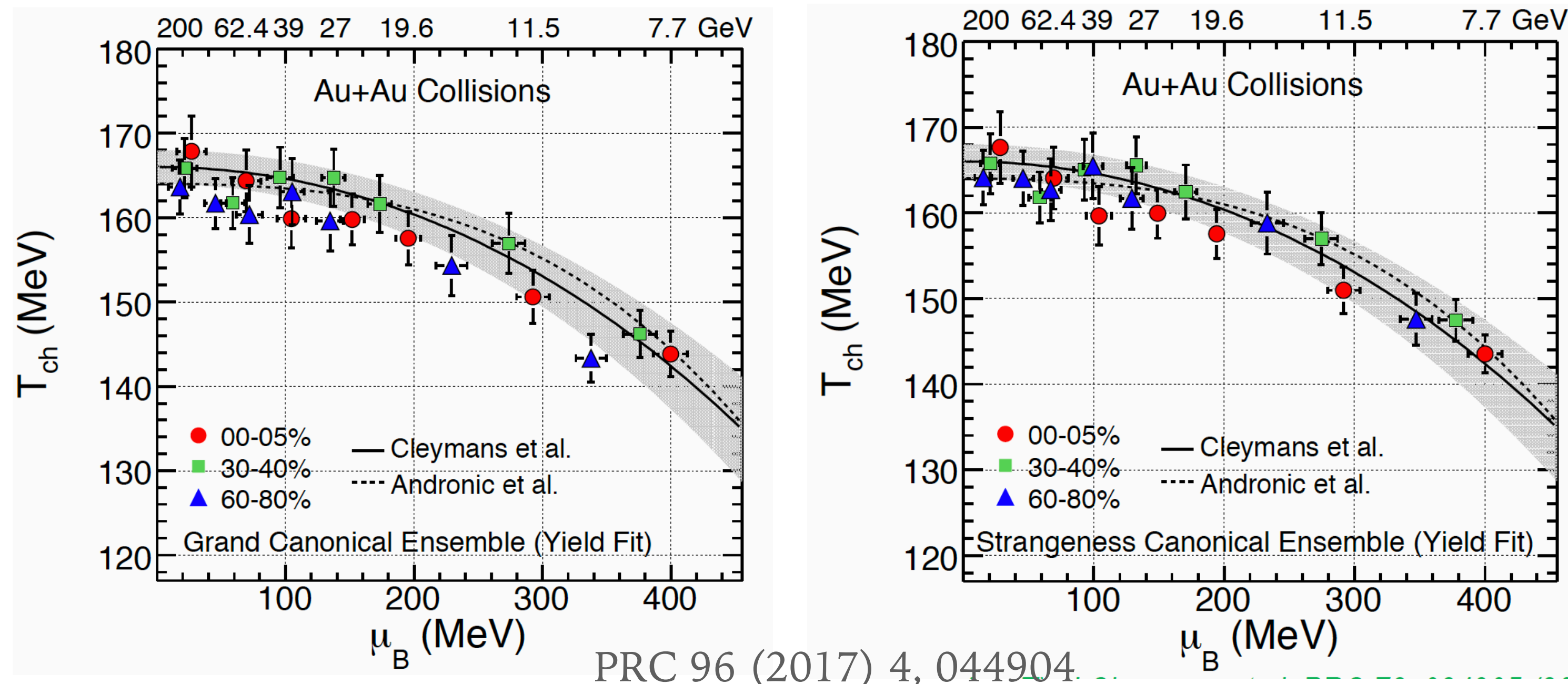
PRC 96 (2017) 4, 044904



Extracted from spectra
(from Blast Wave model):

$$m_T - m \text{ of } \pi, K, p$$

Extracted from particle yields with THERMUS model



Extracted from particle yields with THERMUS model
assuming Grand (Strangeness) Canonical ensemble.

$$\text{BES-I: } \mu_B \sim 20 \text{ MeV} - 420 \text{ MeV}$$

$$\text{BES-II: } \mu_B \sim 205 \text{ MeV} - 720 \text{ MeV}$$

Observables

1. **Onset of QGP** (disappearance of signals of partonic degrees of freedom)

Charge separation w.r.t. EP

NCQ scaling of elliptic flow

2. Observation of **phase transition** (softening of EOS at lower collision energy)

Directed flow v_1

Femtoscopy

3. Existence of **Critical Point** (CP)

Fluctuation analyses

4. **Chiral symmetry restoration**

Low-mass vector mesons, dileptons CME

5. ... and beyond ..

Light nuclei

Lambda's polarization

Hypernuclei production

Nuclear modification factor

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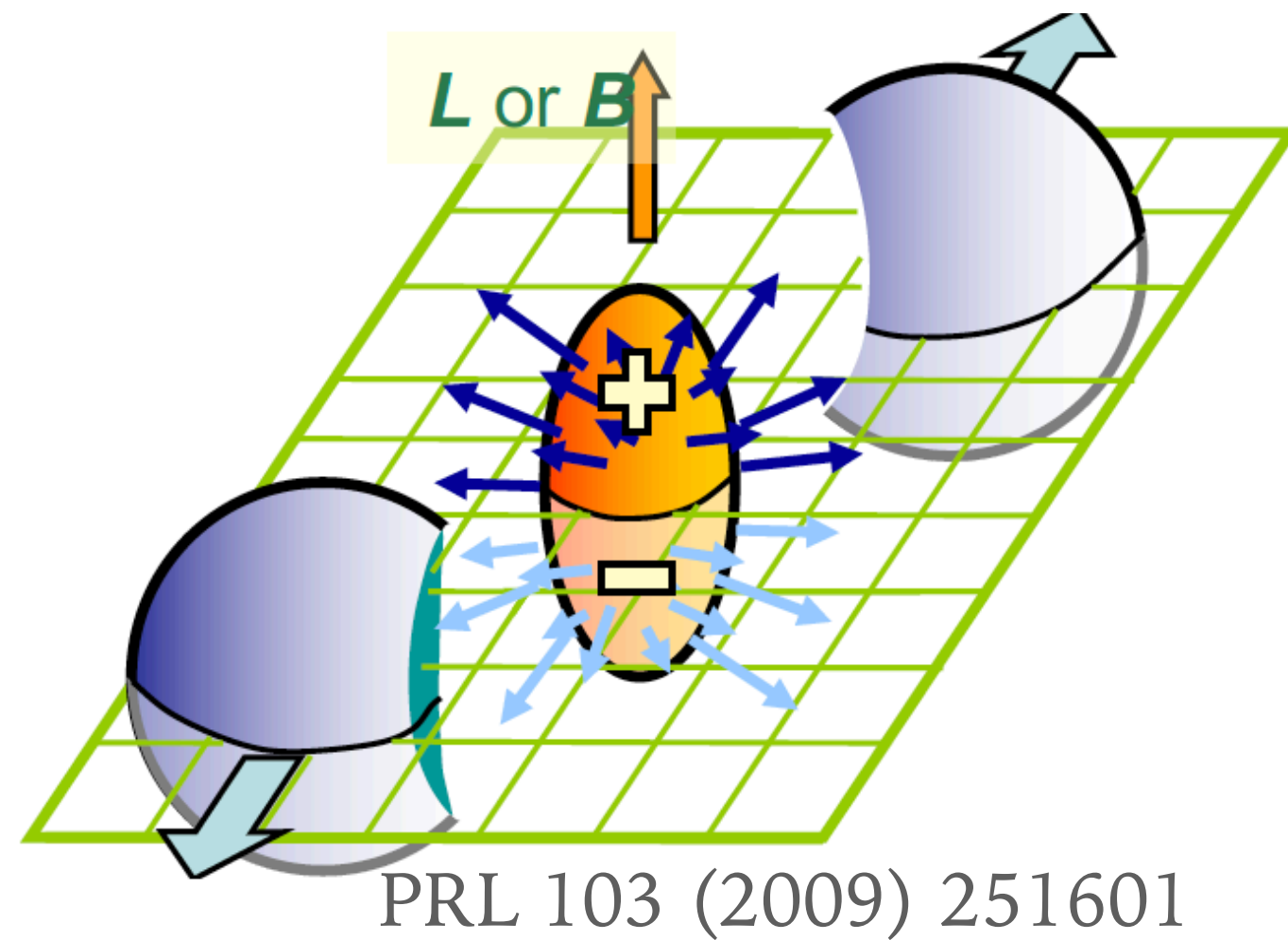
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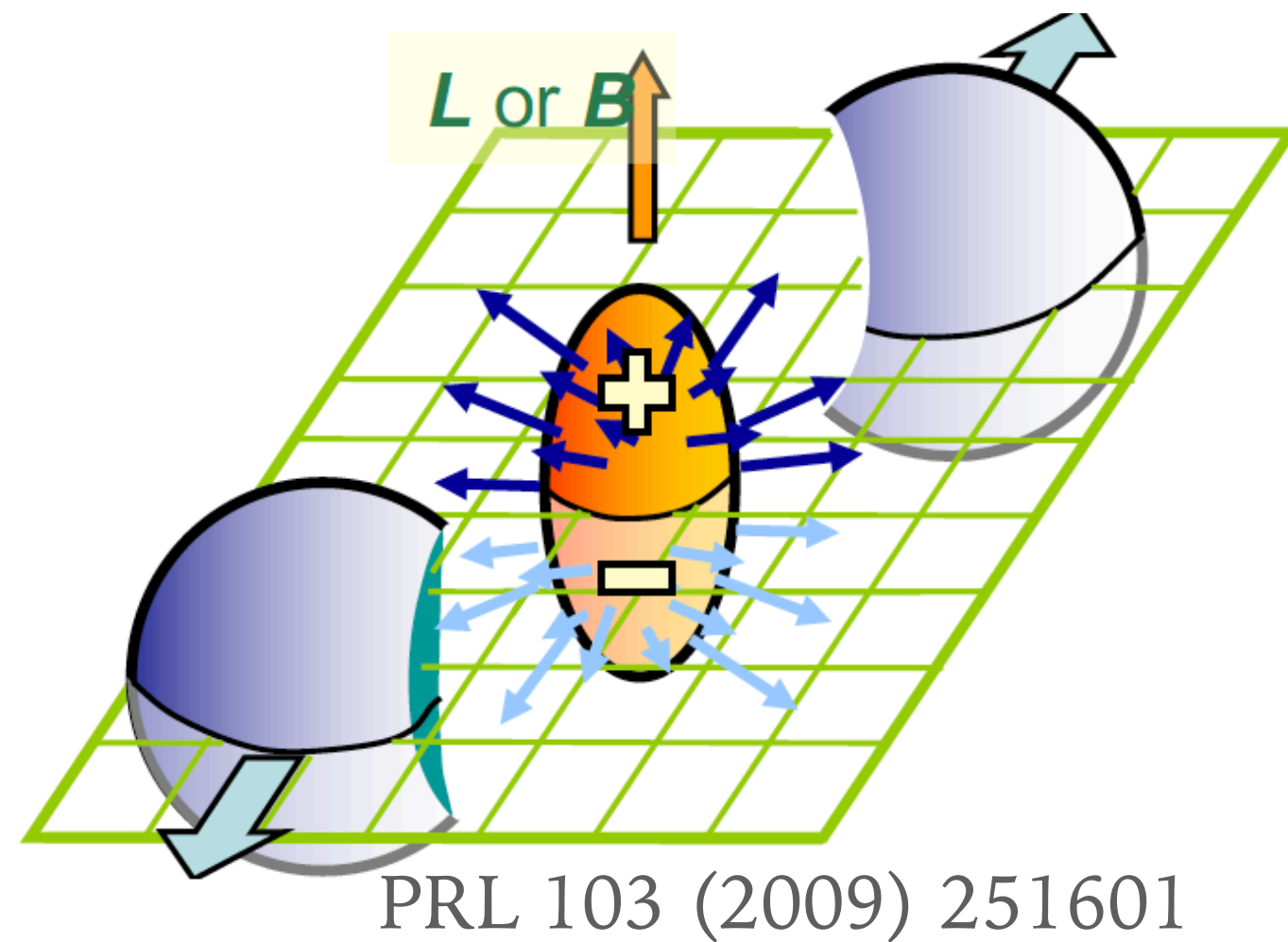
Nuclear modification factor

Charge separation

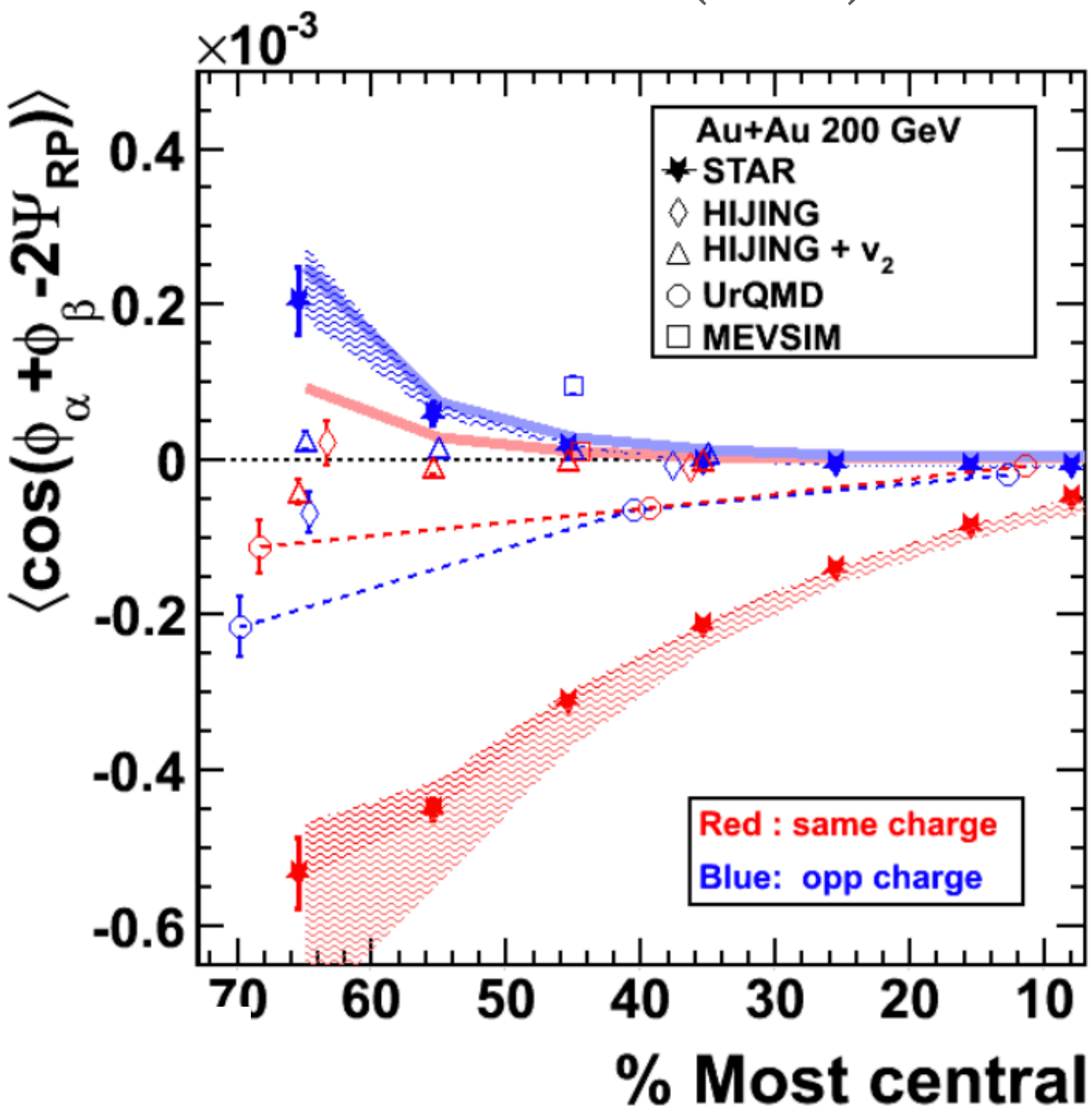


- Strong **B**, system is **deconfinement**, chiral symmetry restoration is reached.
- Chiral symmetry breaking and the origin of hadrons masses related to the existence of gluons field.
- Quarks interactions with gluons fields can change quarks chirality, and may lead to **Local Parity Violation**.
- **Chiral Magnetic Effect**: separation of the charges along the **B** axis (or **L**).

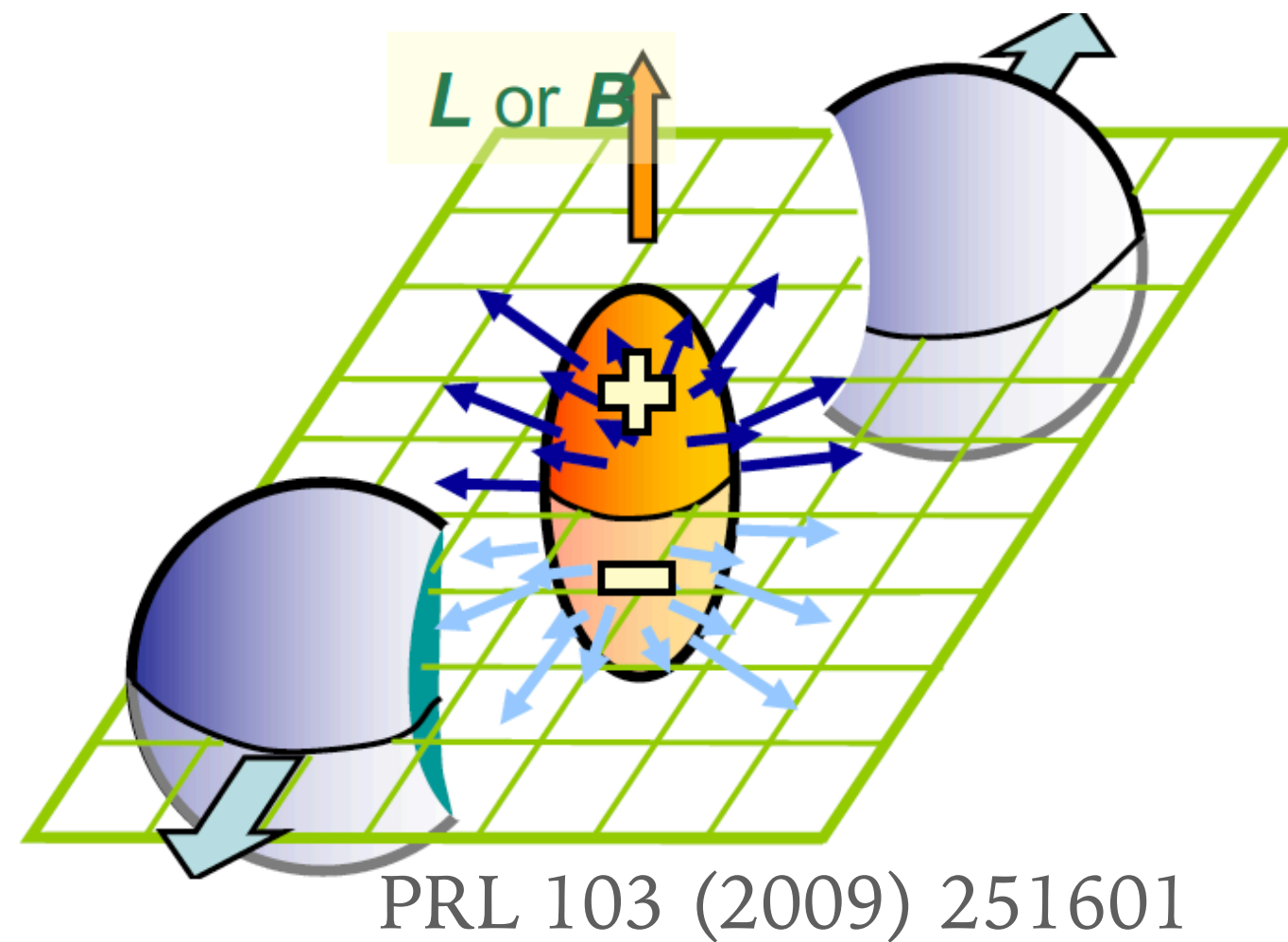
Charge separation



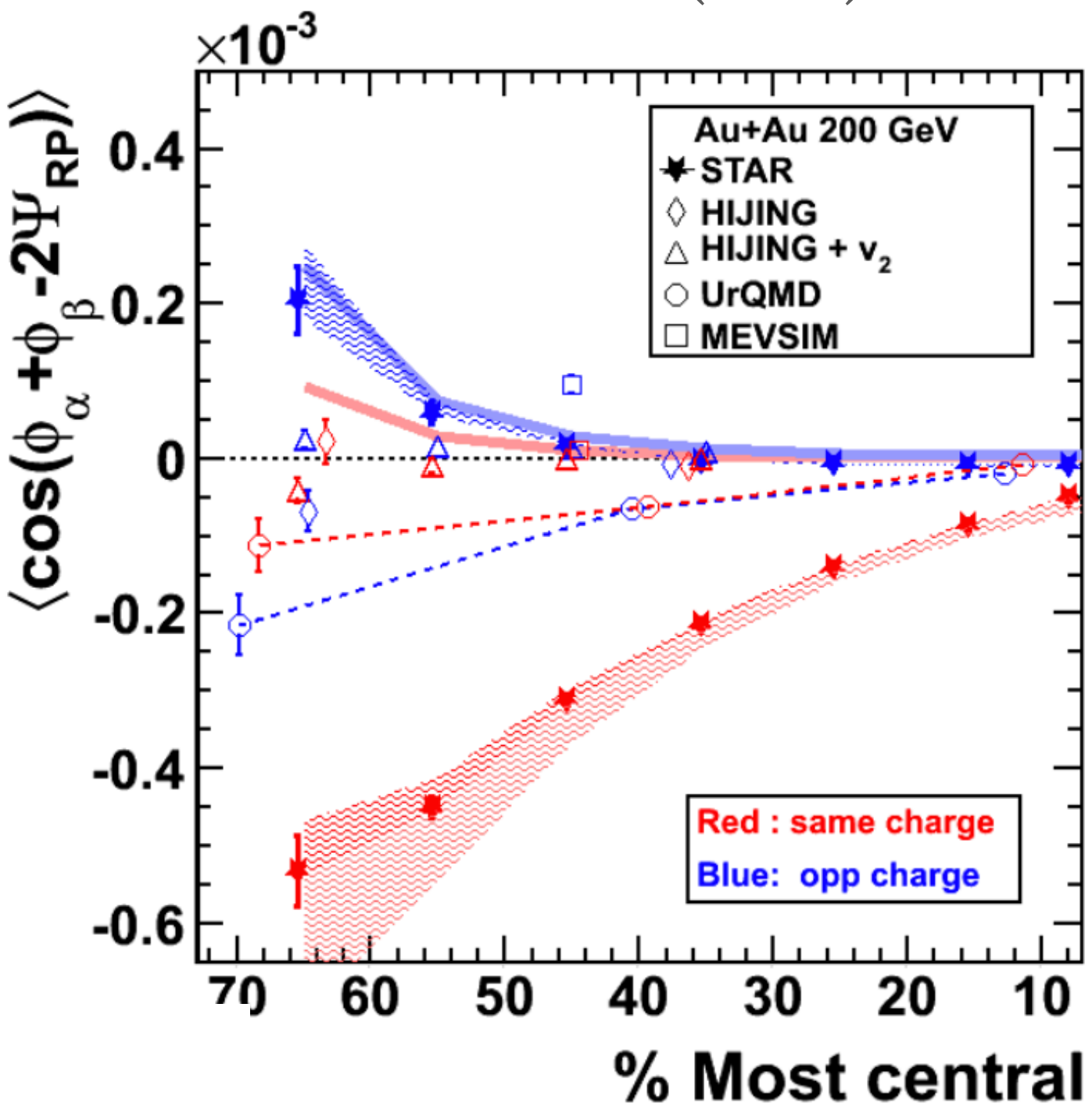
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- Au+Au, U+U and Cu+Cu at top RHIC energies show charge separation measures as $\gamma = \langle \cos(\phi_\alpha + \phi_\beta - 2\Psi_{RP}) \rangle$



Charge separation

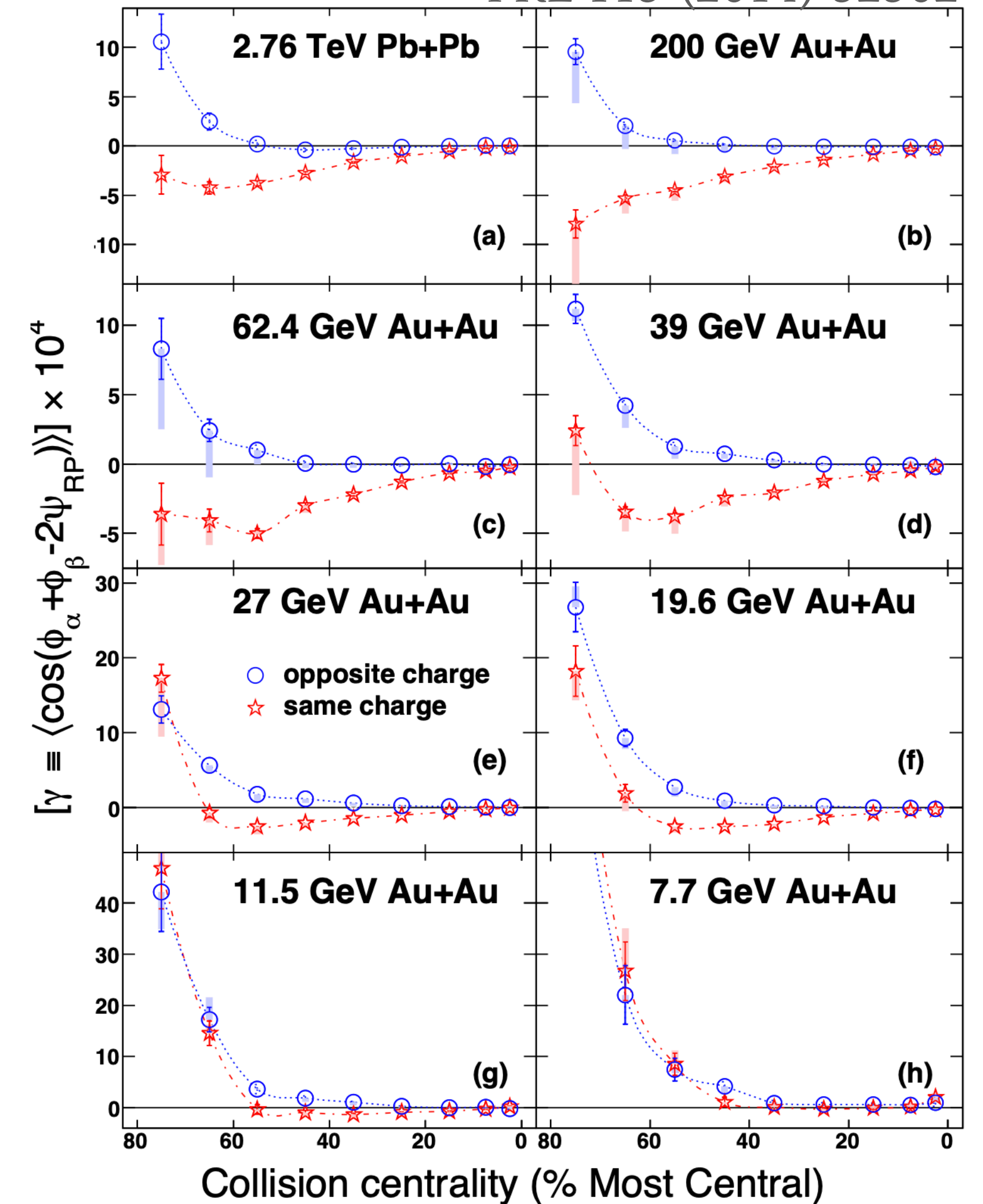


PRL 103 (2009) 251601



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- Au+Au, U+U and Cu+Cu at top RHIC energies show charge separation measures as $\gamma = \langle \cos(\phi_\alpha + \phi_\beta - 2\Psi_{RP}) \rangle$
- Is reduction of signal with decreasing collision energy the signal of turn-off of deconfinement?

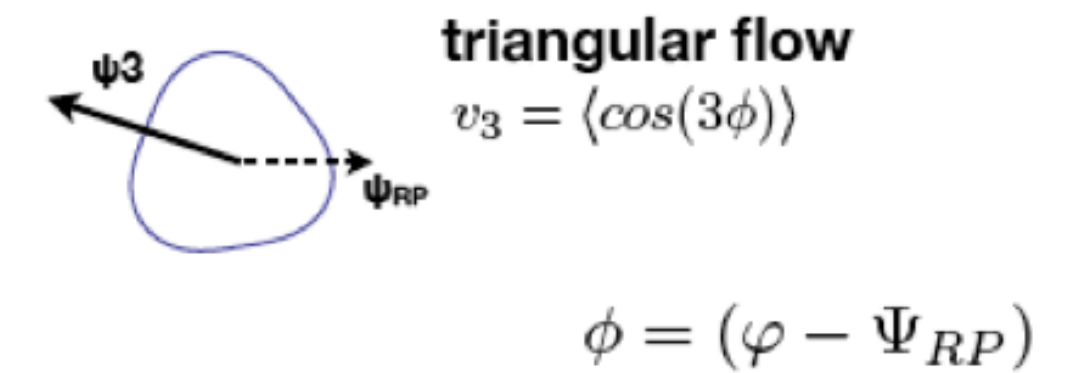
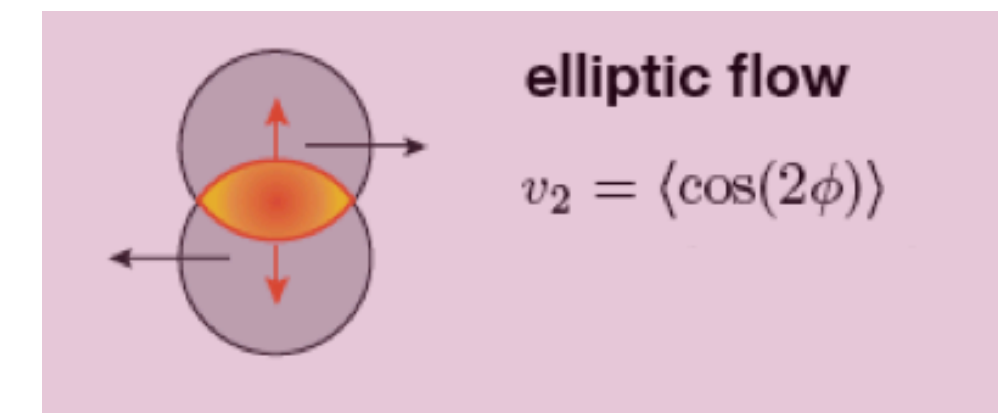
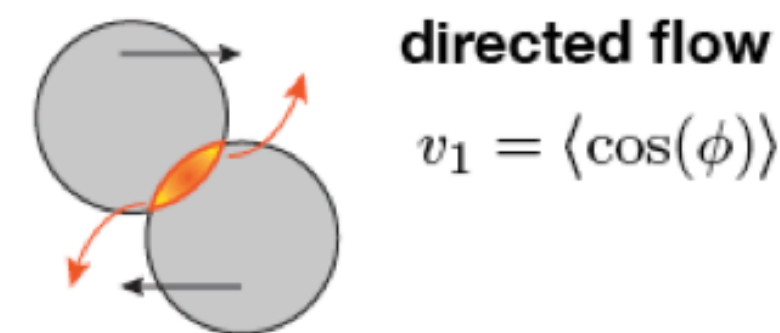
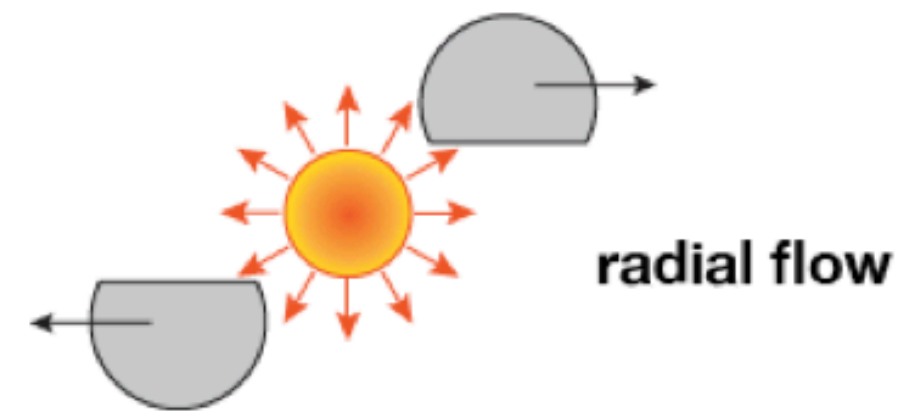
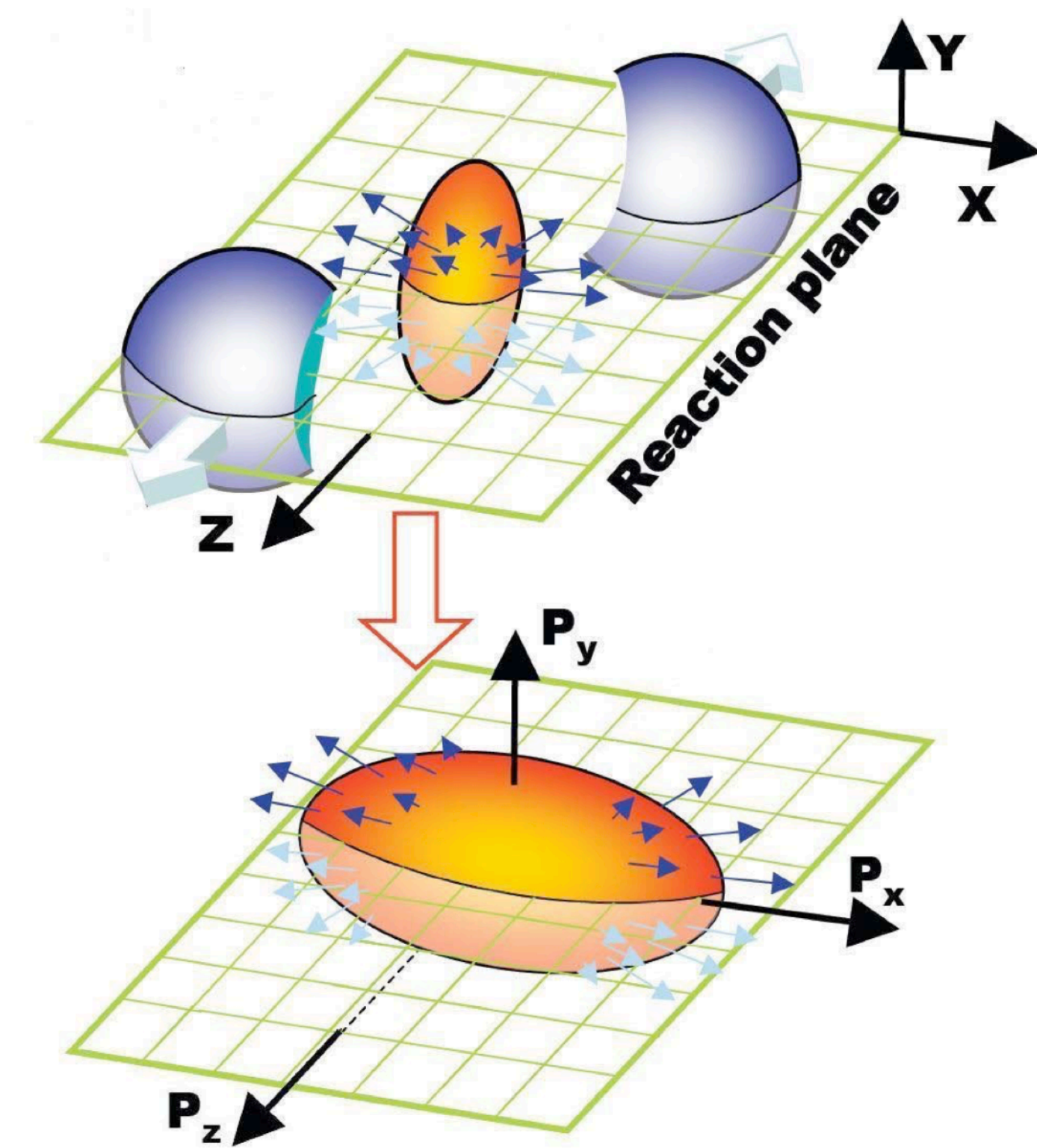
PRL 113 (2014) 52302



Splitting between same- and opposite-sign charges decreases with decreasing collision energy and disappears below $\sqrt{s_{NN}} = 11.5$ GeV

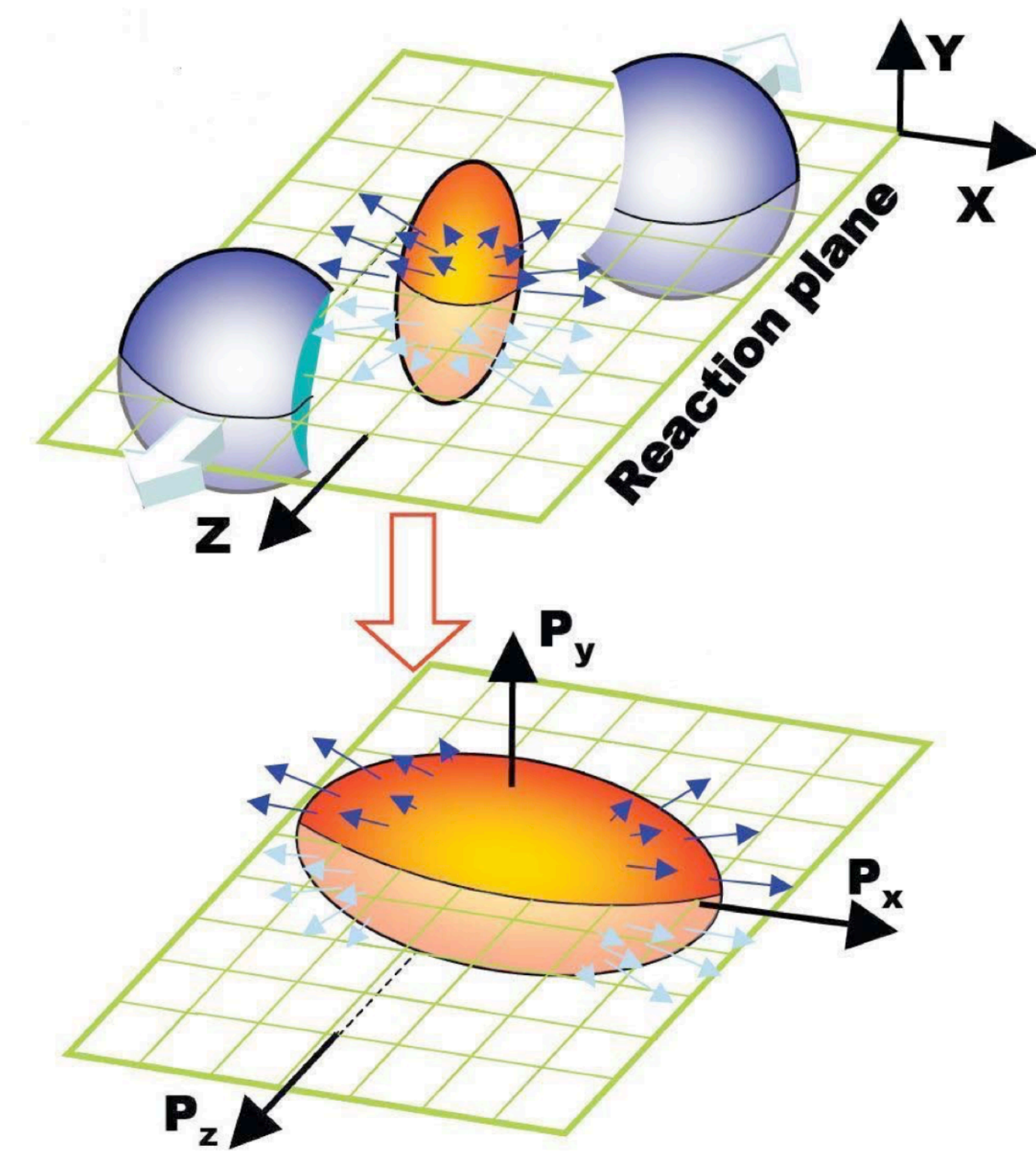
Elliptic flow

Initial spatial anisotropy leads to the final momentum anisotropy

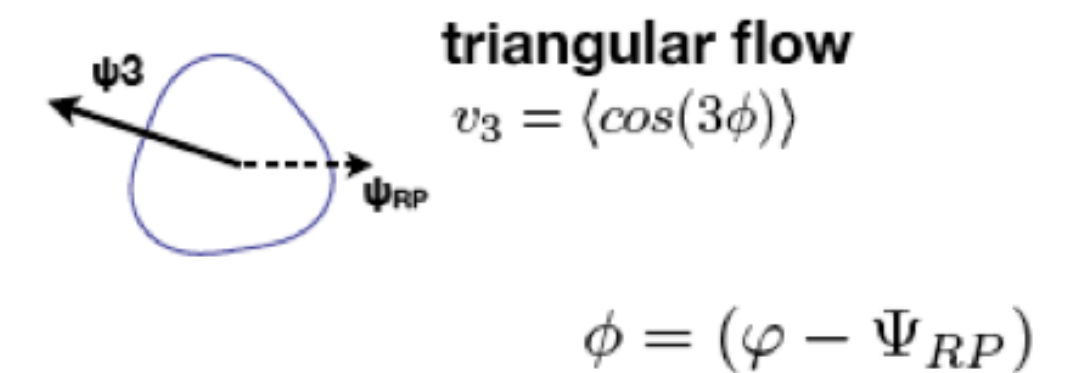
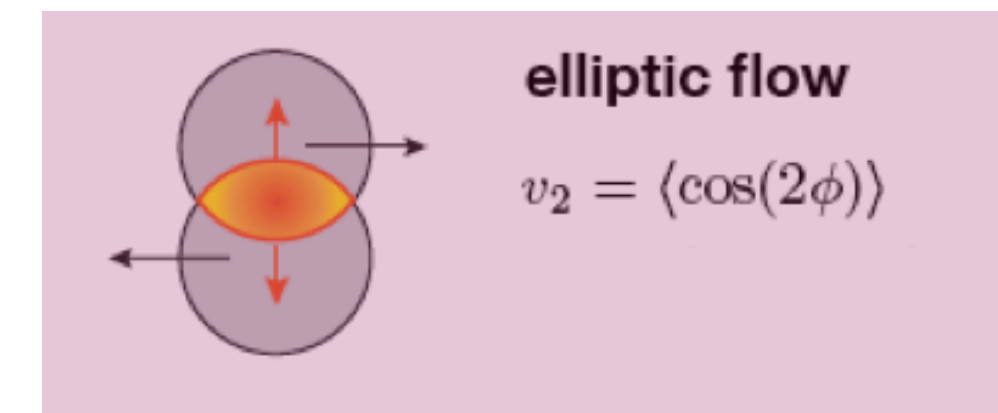
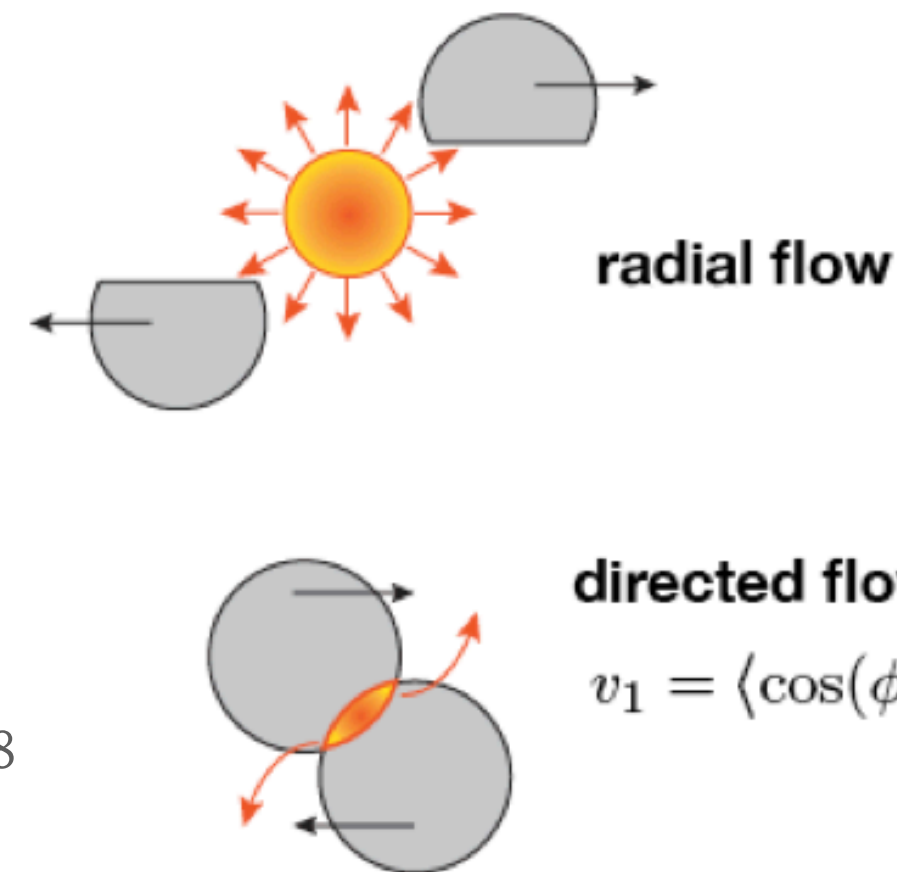


$$E \frac{d^3 N}{d^3 \mathbf{p}} = \frac{1}{2\pi} \frac{d^2 N}{p_t dp_t dy} \left(1 + 2 \sum_{n=1}^{\infty} v_n \cos[n(\varphi - \Psi_{RP})] \right)$$

Elliptic flow

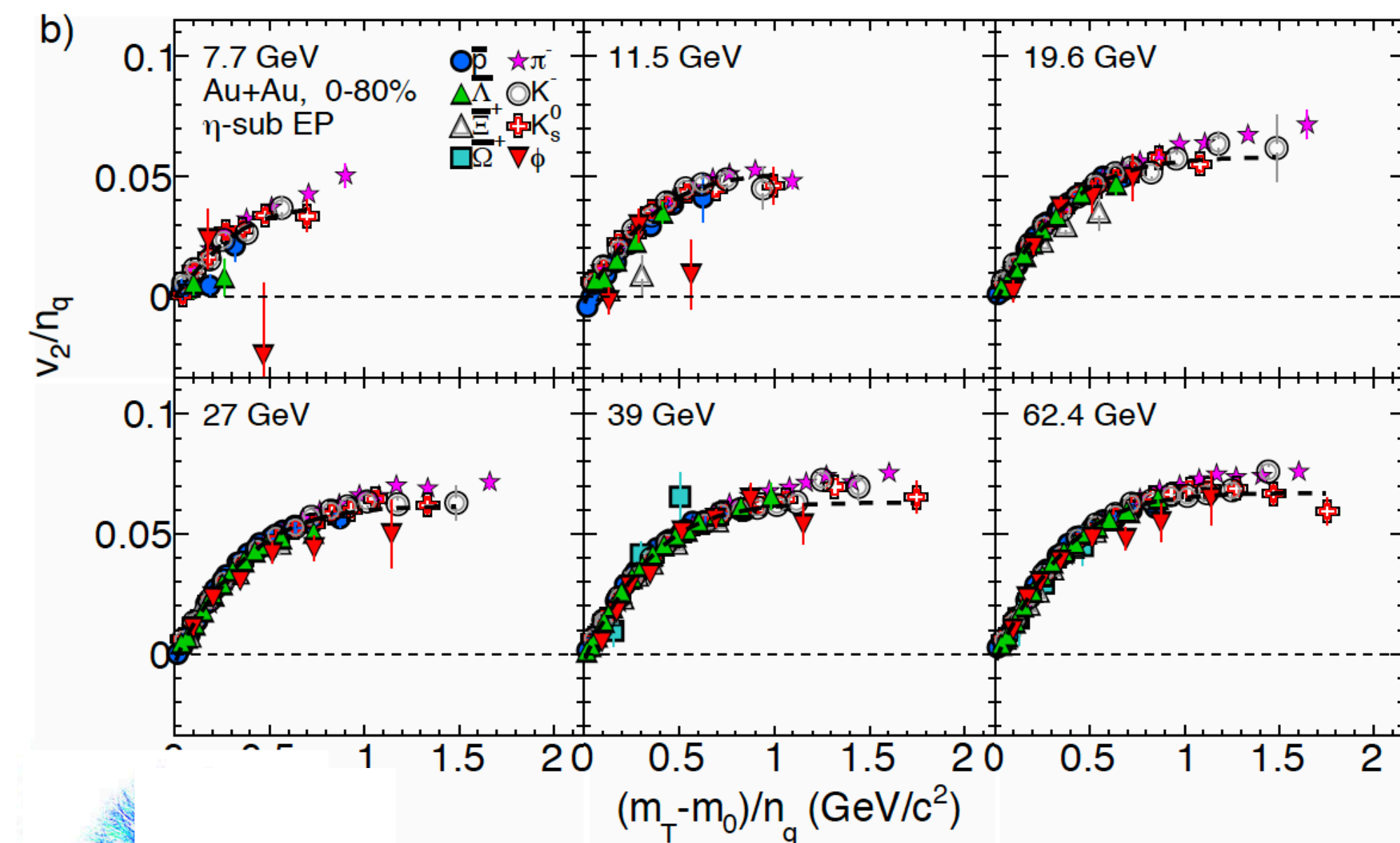
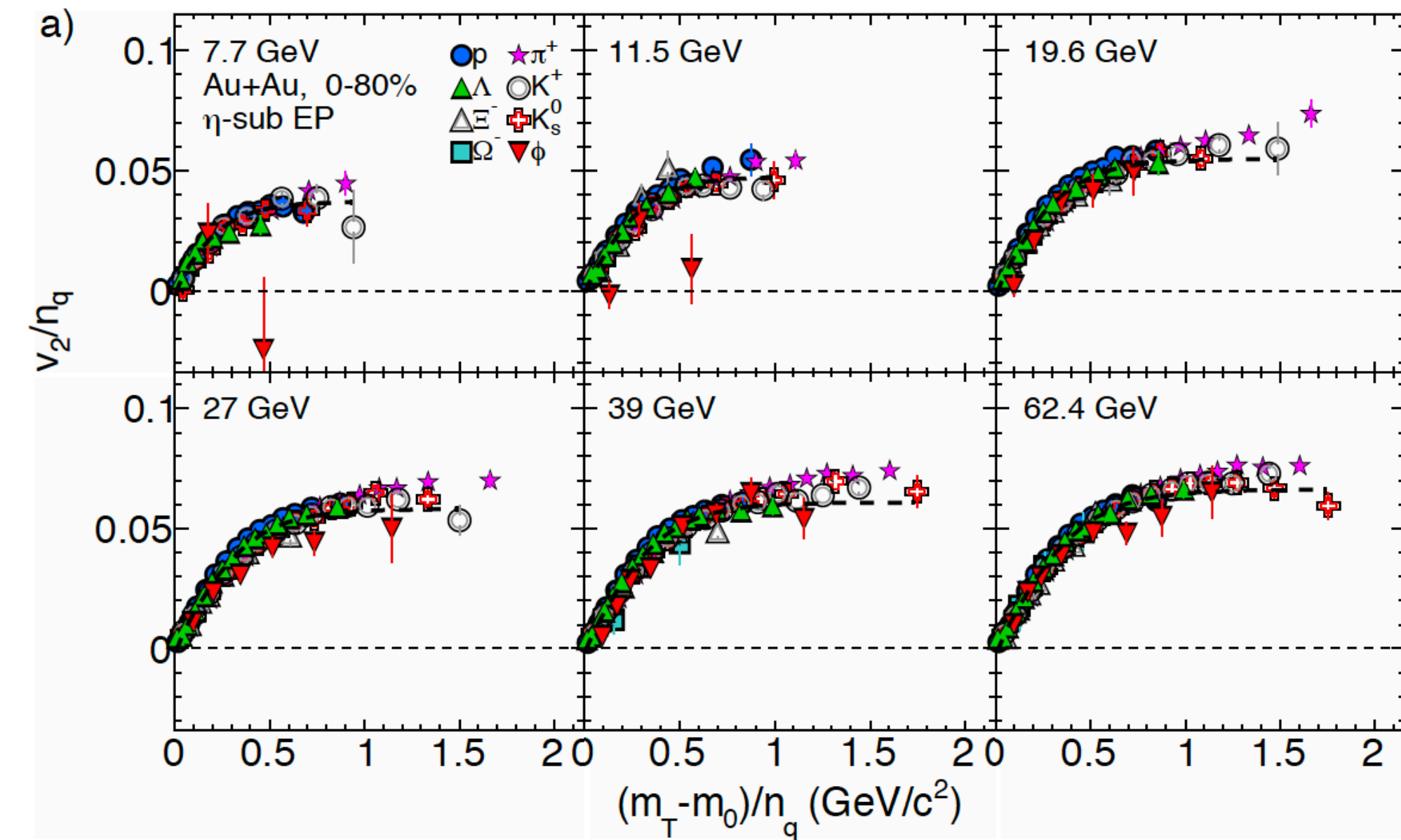


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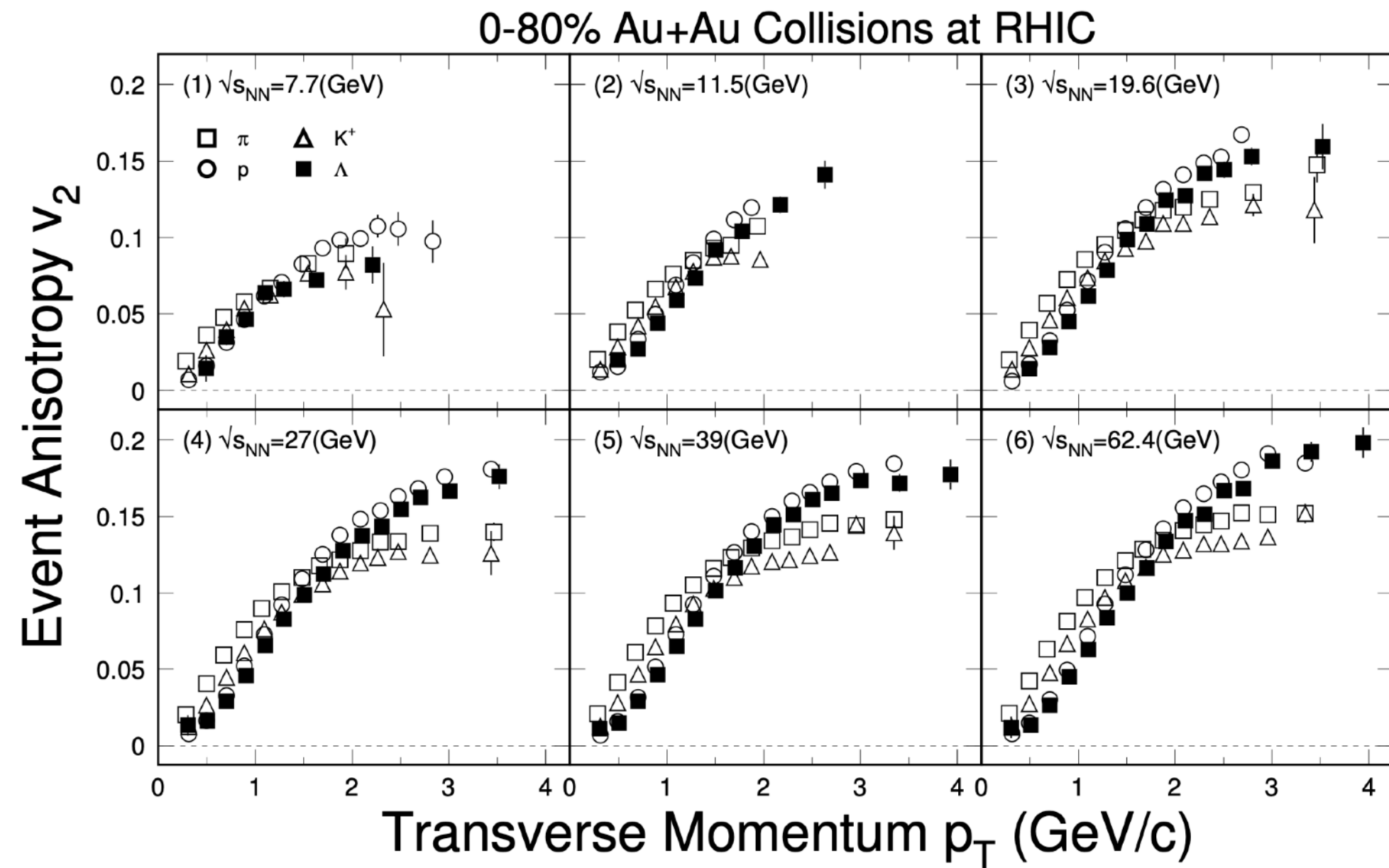


STAR: PRC 88
(2013) 14902
Phys. Rev. C 93,
014907 (2016)
Phys. Rev. Lett.
116, 062301
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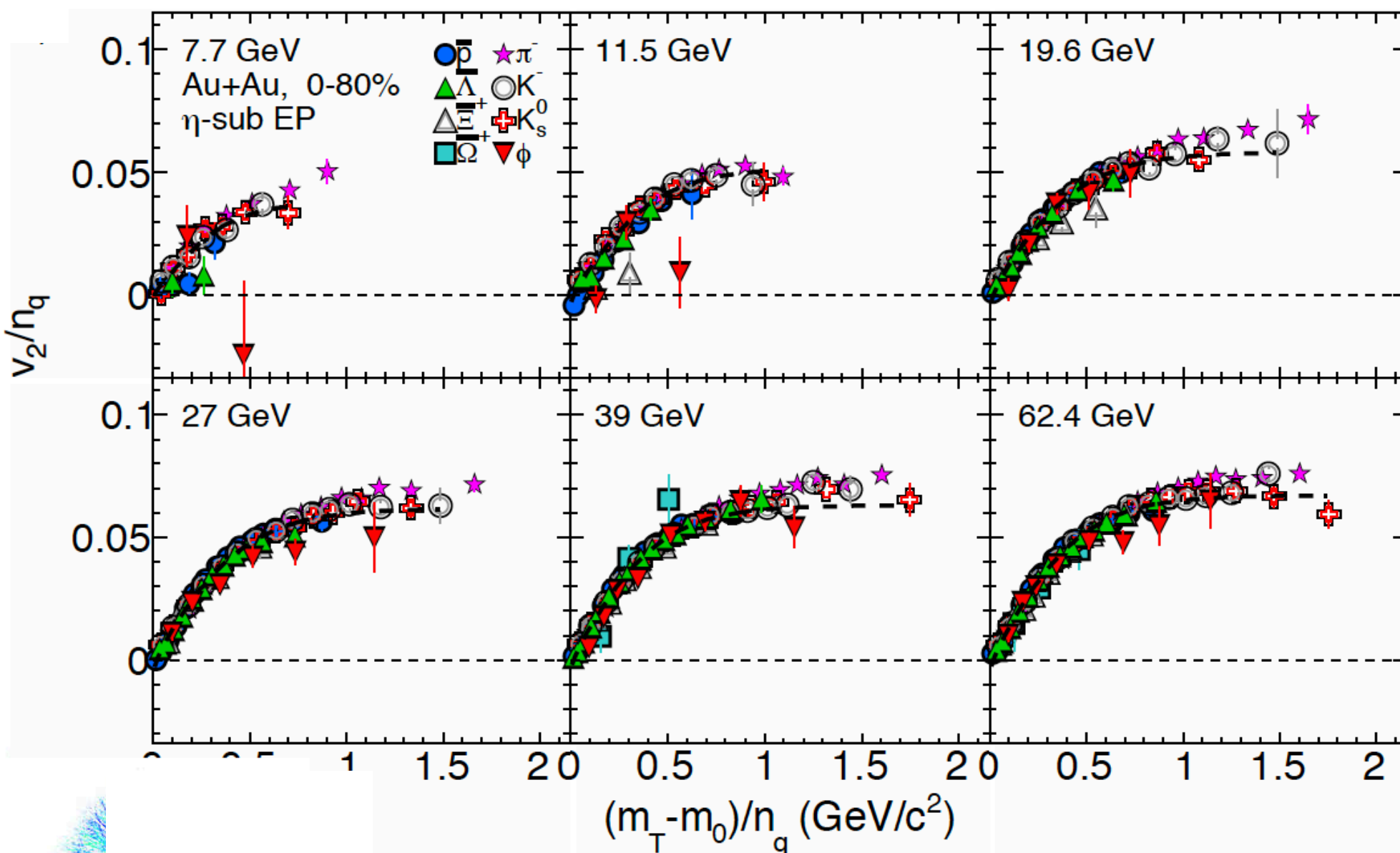


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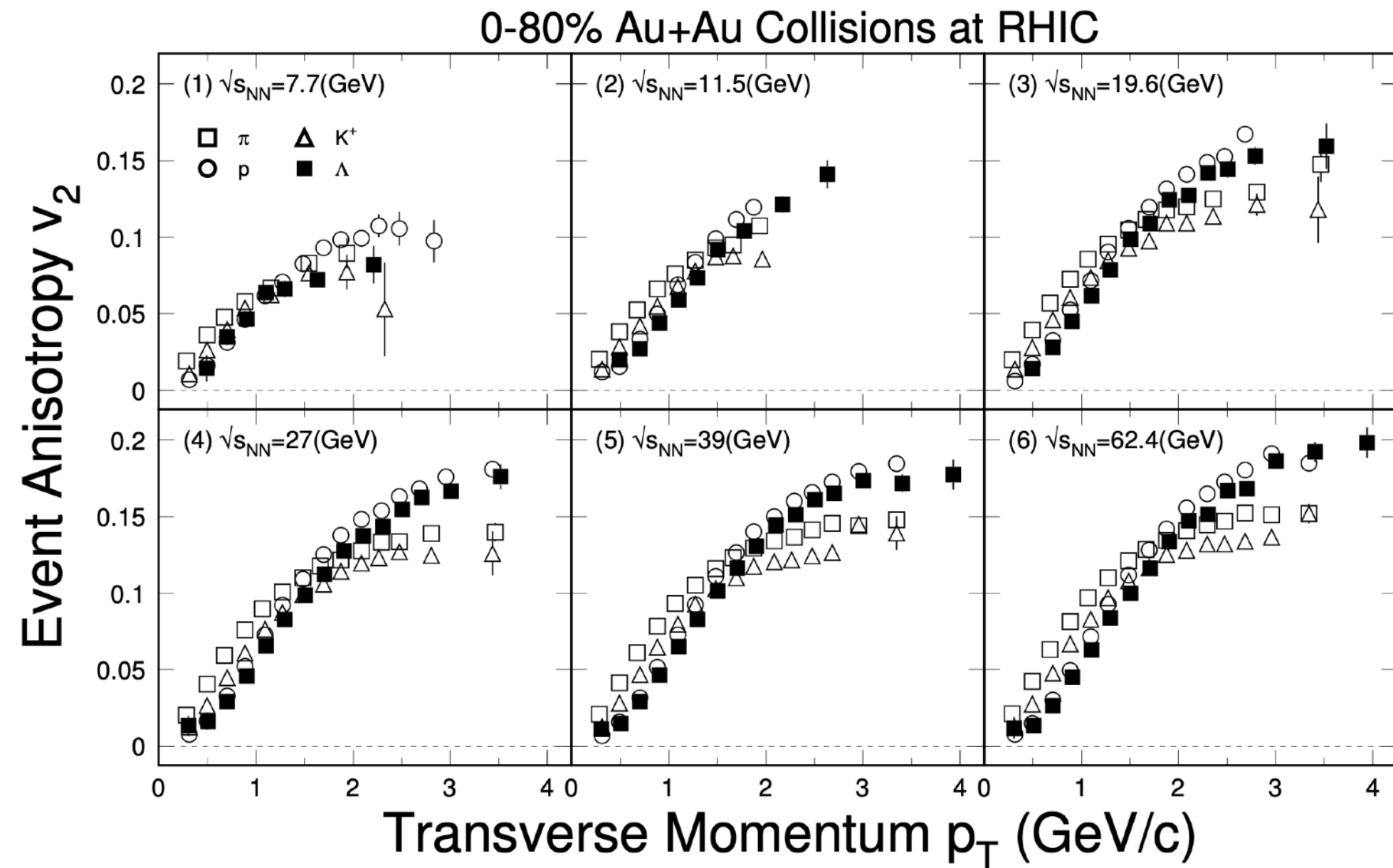
$v_2(p_T)$ are mass ordered

- ϕ meson v_2 fails the trend from other hadrons at $\sqrt{s_{NN}} = 11.5$ GeV, (low statistics)
- NCQ scaling between particles is broken ($\sqrt{s_{NN}} < 19.6$ GeV) and consistent with hadronic interactions becoming dominant



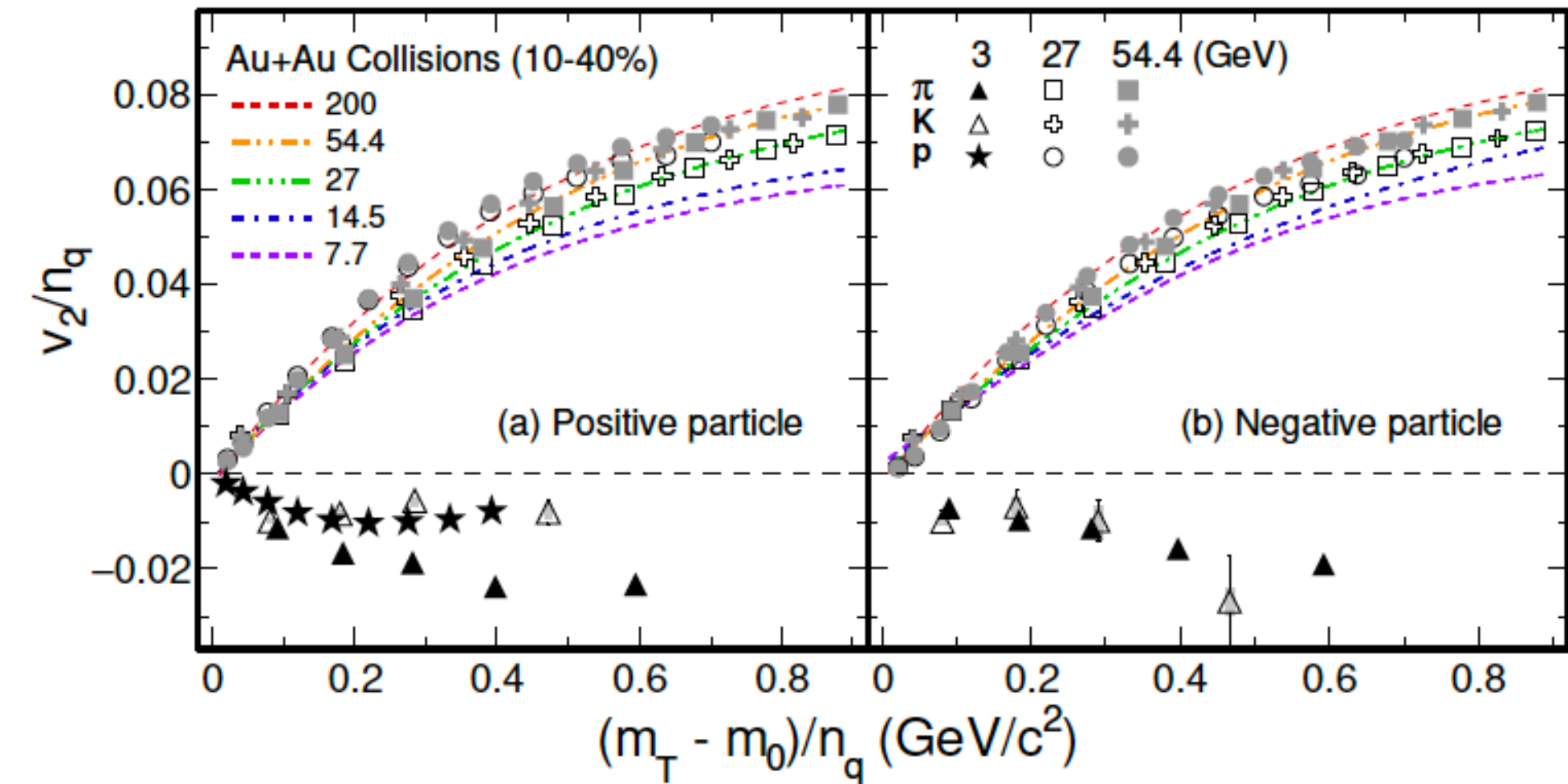
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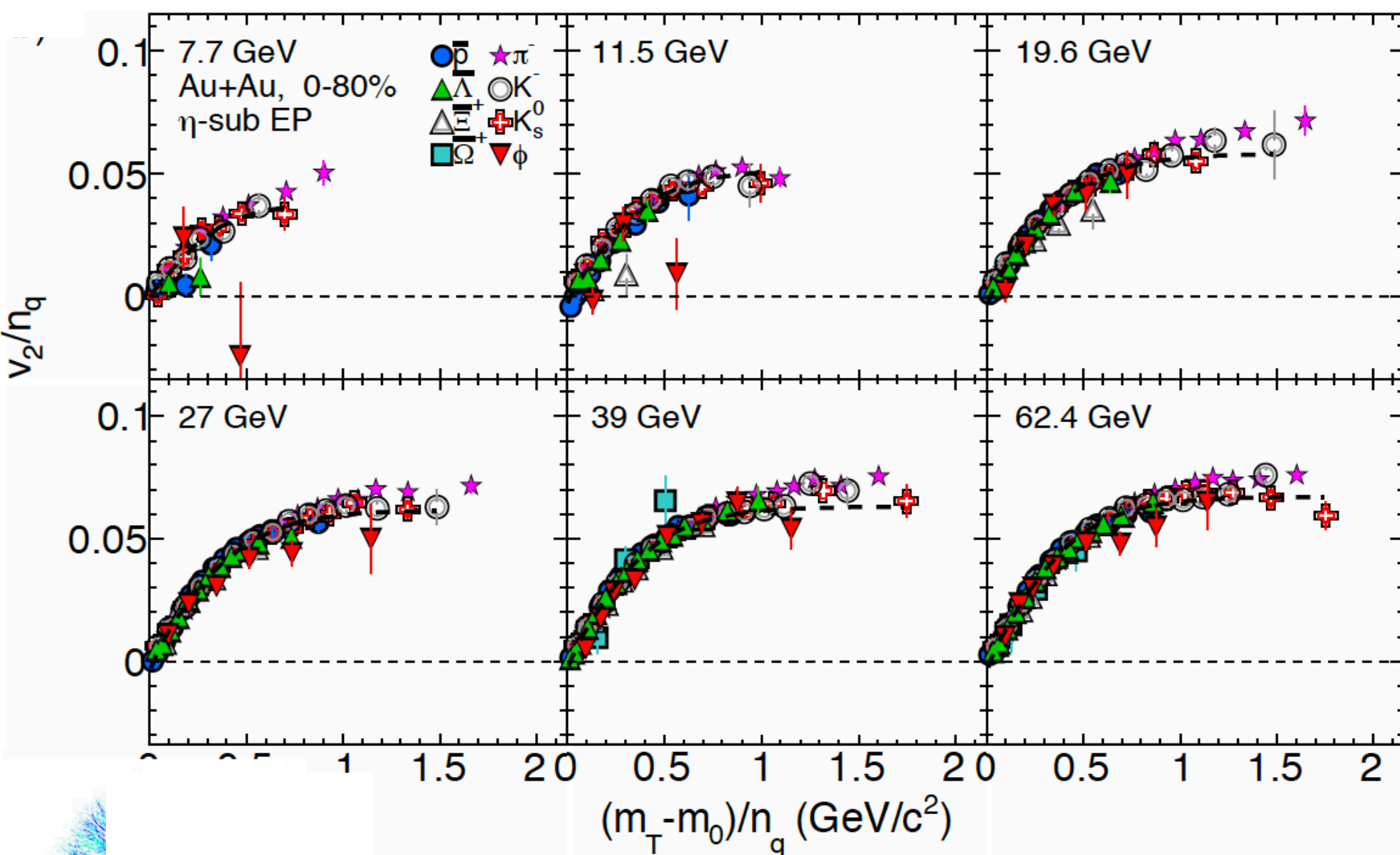


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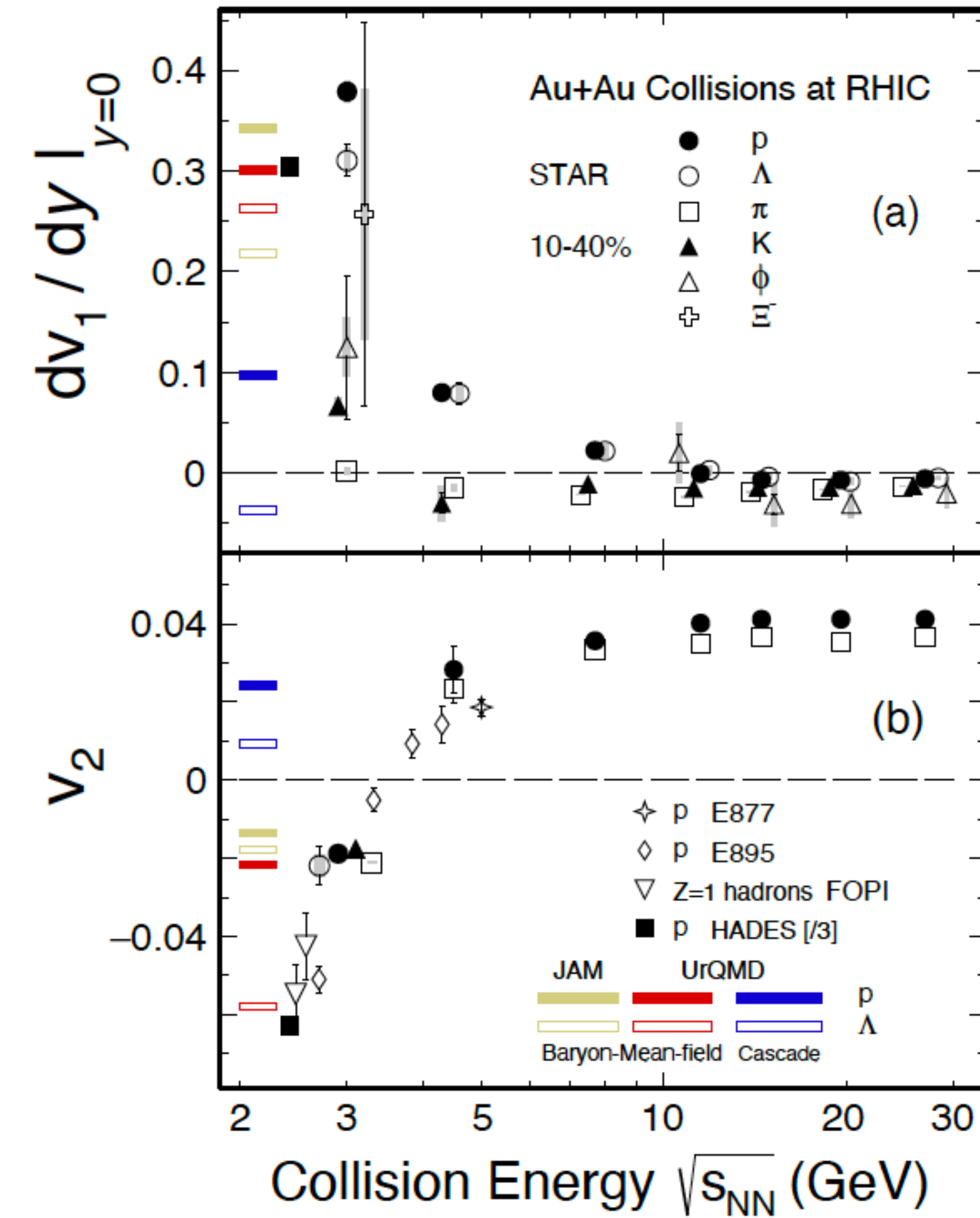


- $v_2 > 0 \rightarrow$ formation of the QGP, scaling of NCQ
- $v_2 < 0$, slope of the $v_1 < 0$ ($\sqrt{s_{NN}} = 3$ GeV) \rightarrow NCQ scaling absent

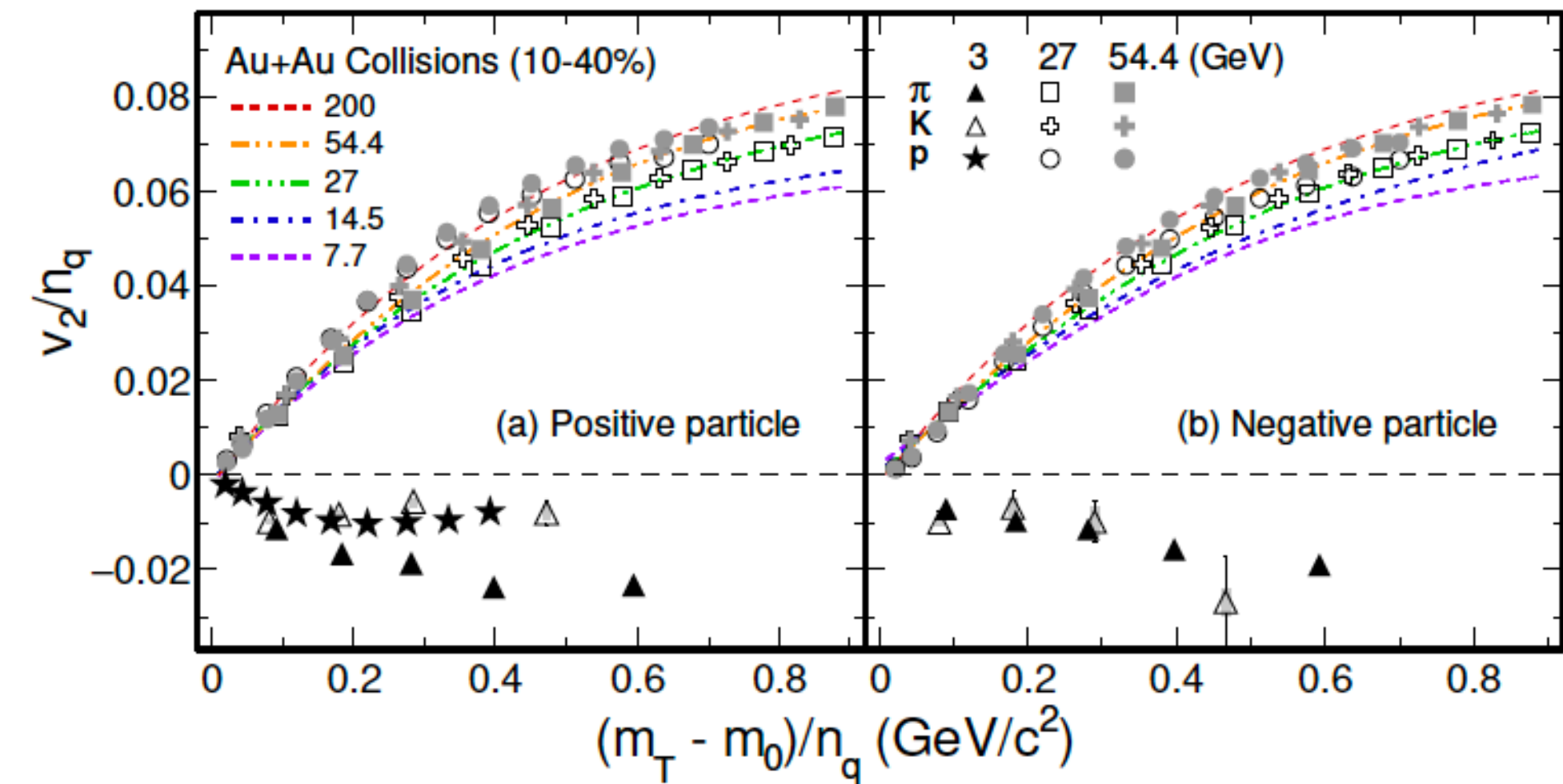


STAR: PRC 88 (2013) 14902
 Phys. Rev. C 93, 014907 (2016)
 Phys. Rev. Lett. 116, 062301 (2016)

Elliptic flow



- $v_2 > 0 \rightarrow$ partonic collectivity \rightarrow formation of QGP \rightarrow scaling of NCQ
- $v_2 < 0$, slope of the $v_1 < 0$ ($\sqrt{s_{NN}} = 3$ GeV) \rightarrow NCQ scaling absent



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- JAM, UrQMD mean field reproduced results.
- **Vanishing of partonic collectivity and a new EOS, dominated by baryonic interactions in the high baryon density region.**

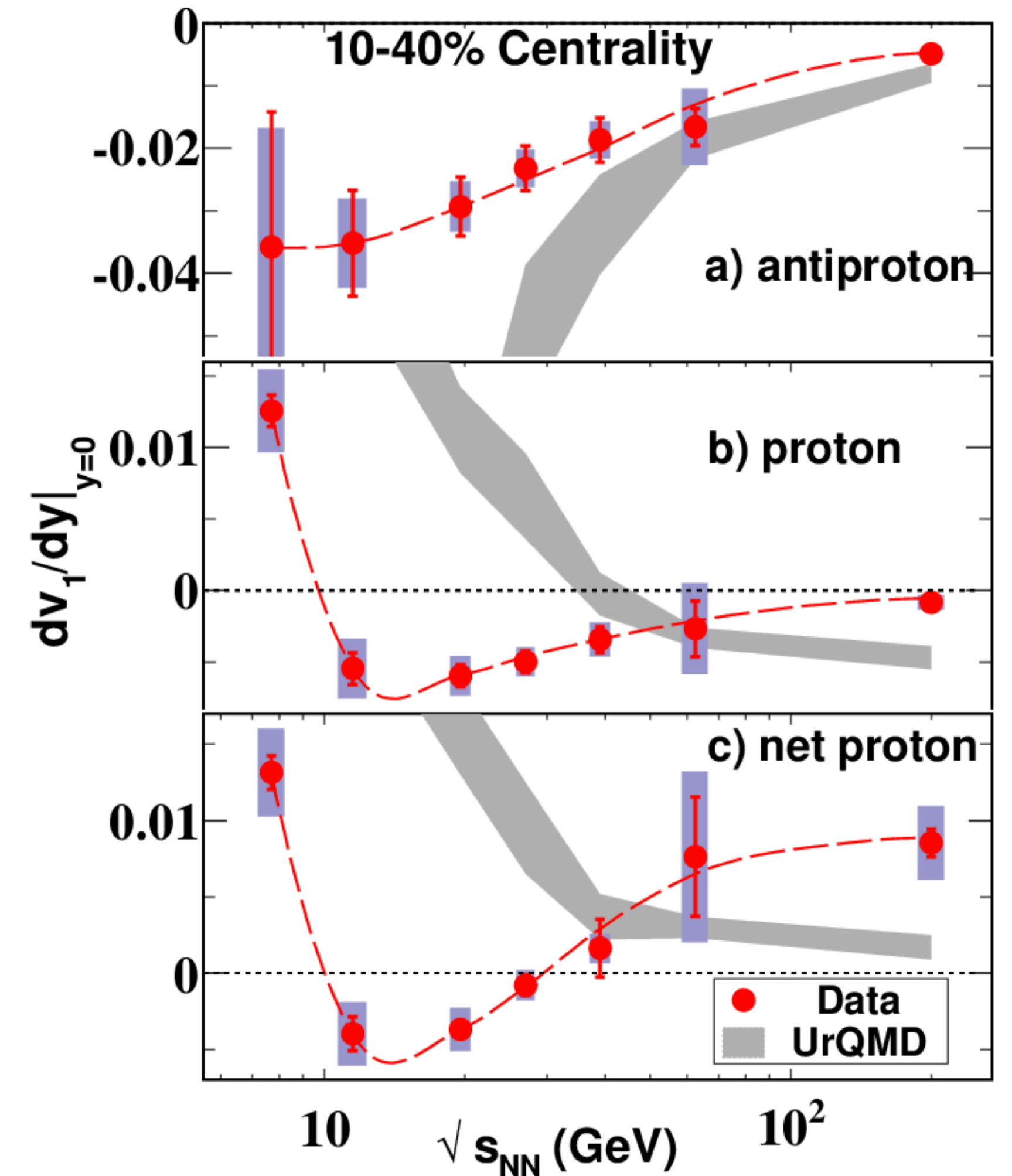
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Directed flow and $\langle m_T \rangle - m$ dependence

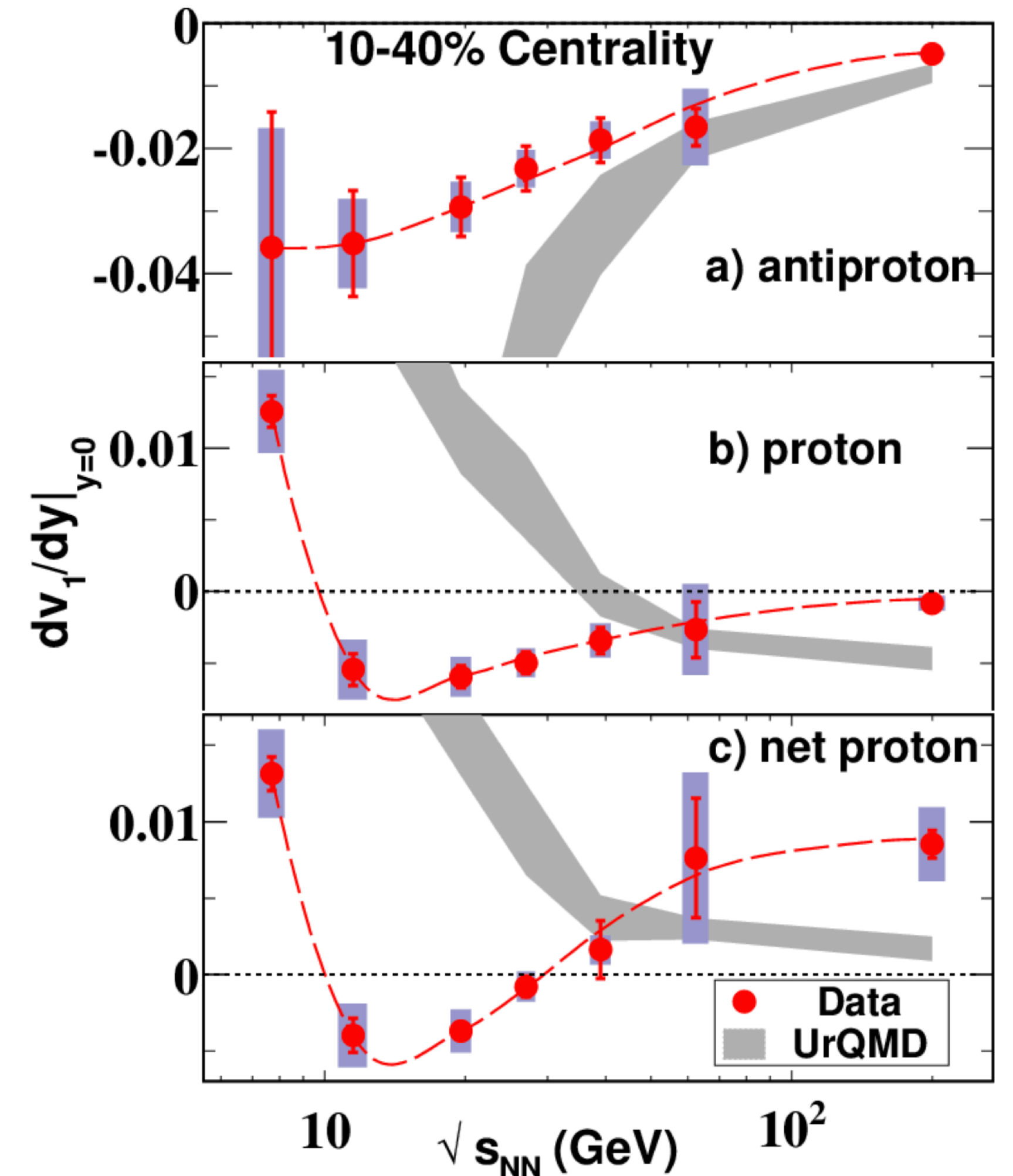
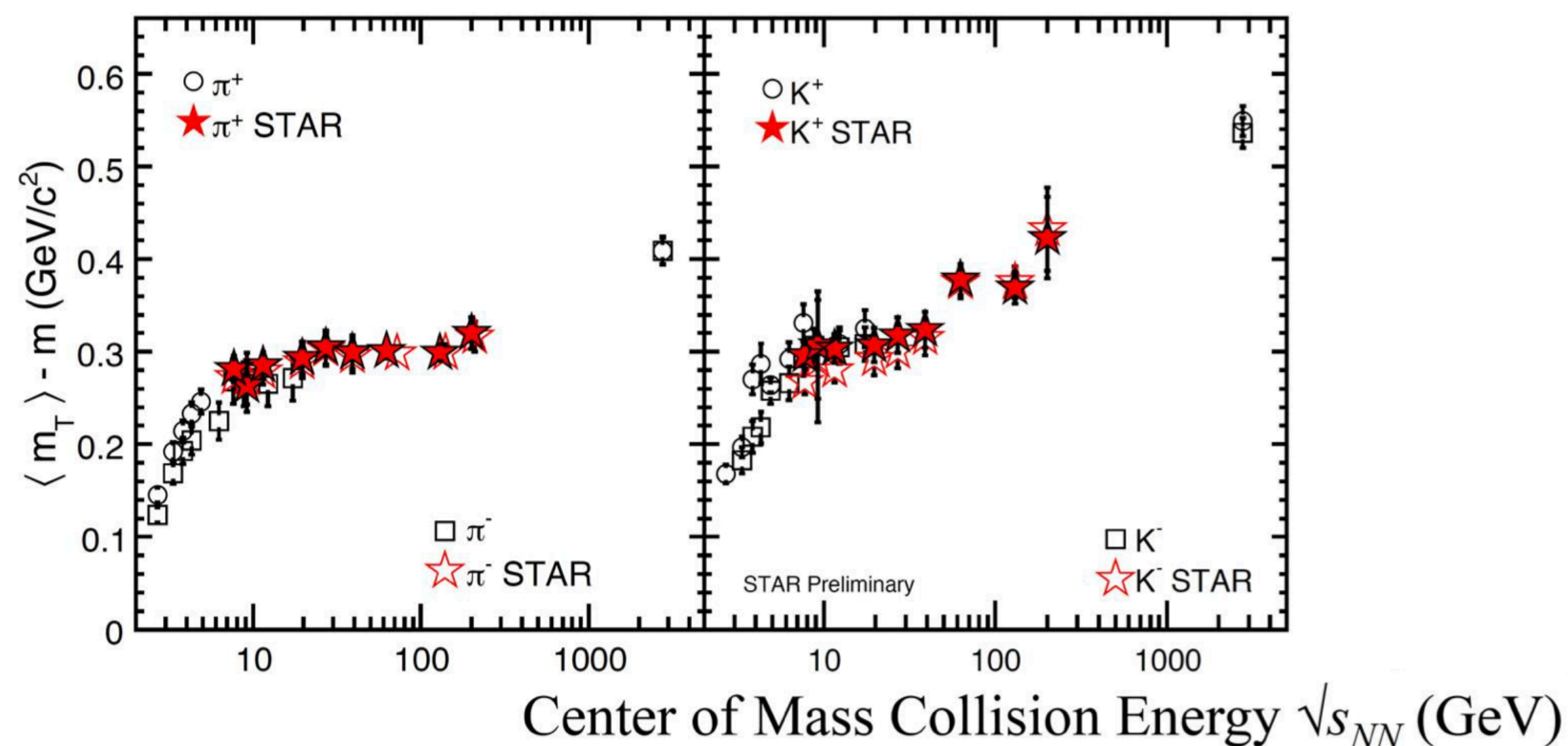
- v_1 probes early stage of collision;
- v_1 sensitive to compression;
- v_1 should be sensitive to the first-order phase transition;
- change of sign in the slope of $\frac{dv_1}{dy}$ (for baryons, or net-baryons) as a probe to the softening of EOS and/or the first-order phase transition;
- If a system undergoes a first-order phase transition, due to formation of mixed phase, pressure gradient is small (minimum in the v_1 slope parameter);



STAR, PRL 112, 162301 (2014)

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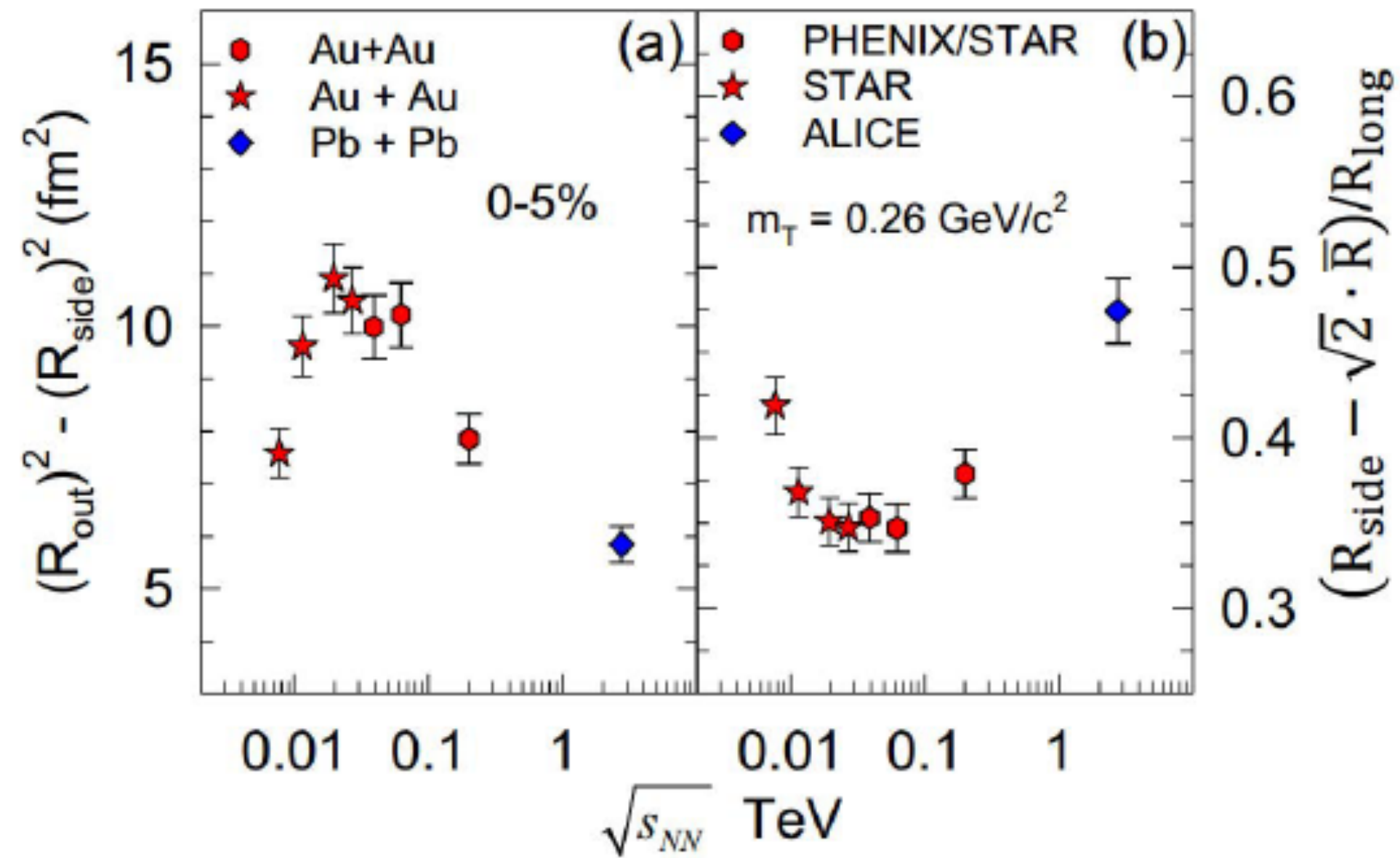
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- If a system undergoes a first-order phase transition, due to formation of mixed phase, pressure gradient is small (minimum in the v_1 slope parameter);
- $\langle m_T \rangle - m$ measures thermal excitation in the transverse direction



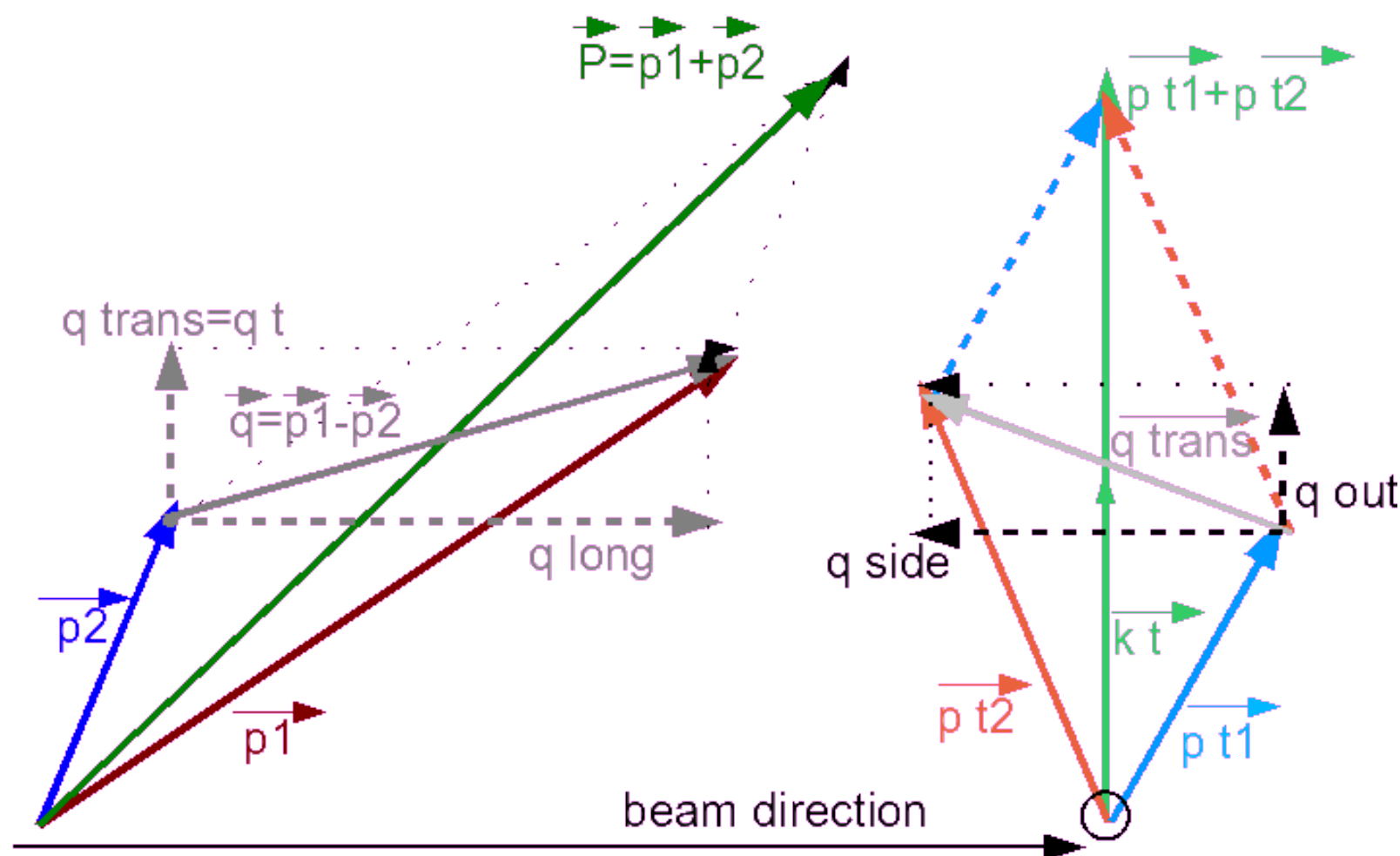
STAR, PRL 112, 162301 (2014)

Femtoscscopy

PHENIX Collaboration, arXiv:1410.2559

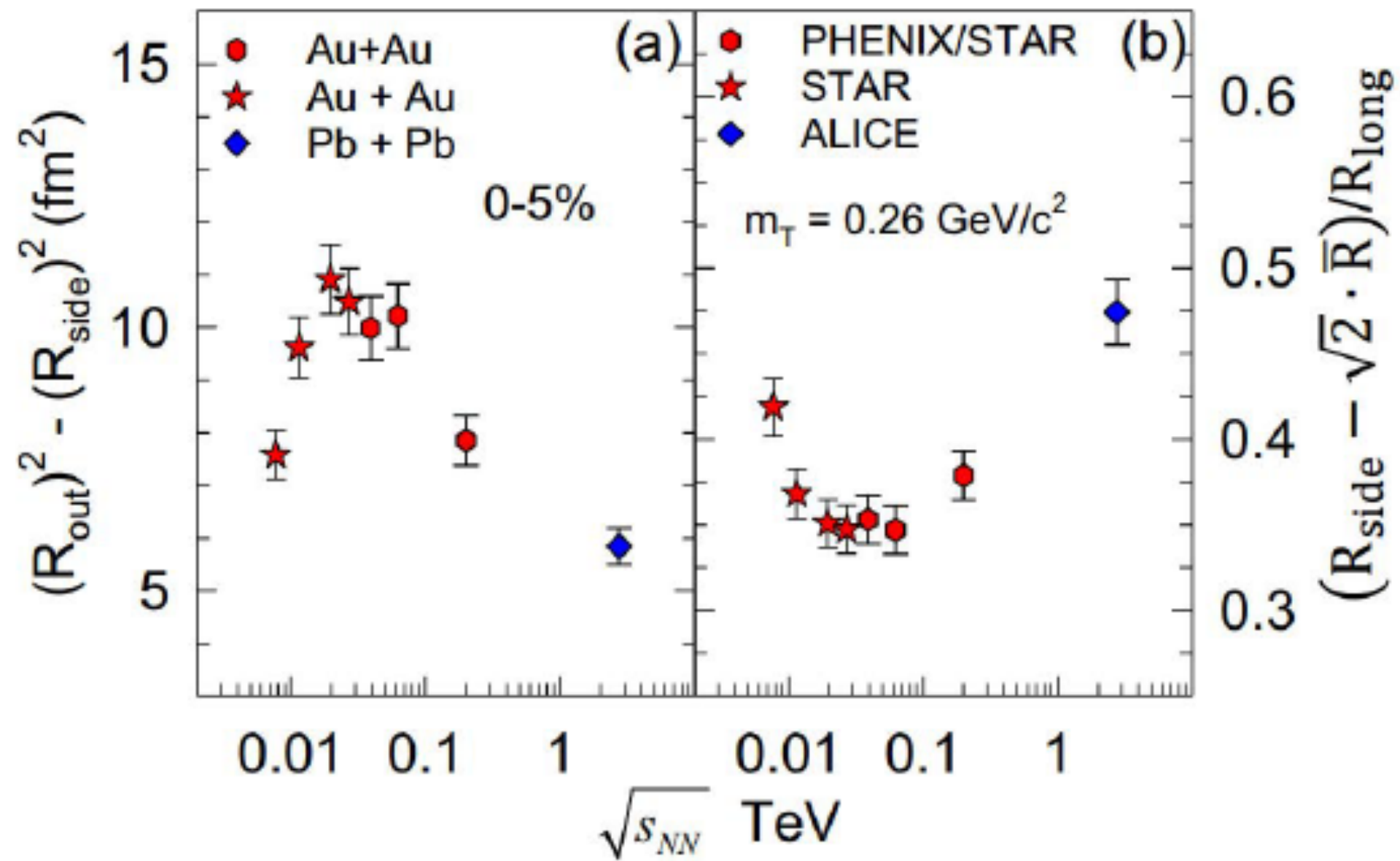


- $R_{out}^2 - R_{side}^2 = \beta_t^2 \Delta\tau^2$: related to emission duration
- $(R_{side} - \sqrt{2}\bar{R})/R_{long}$: related to expansion velocity, \bar{R} : initial transverse size
- Indication of the critical behavior?

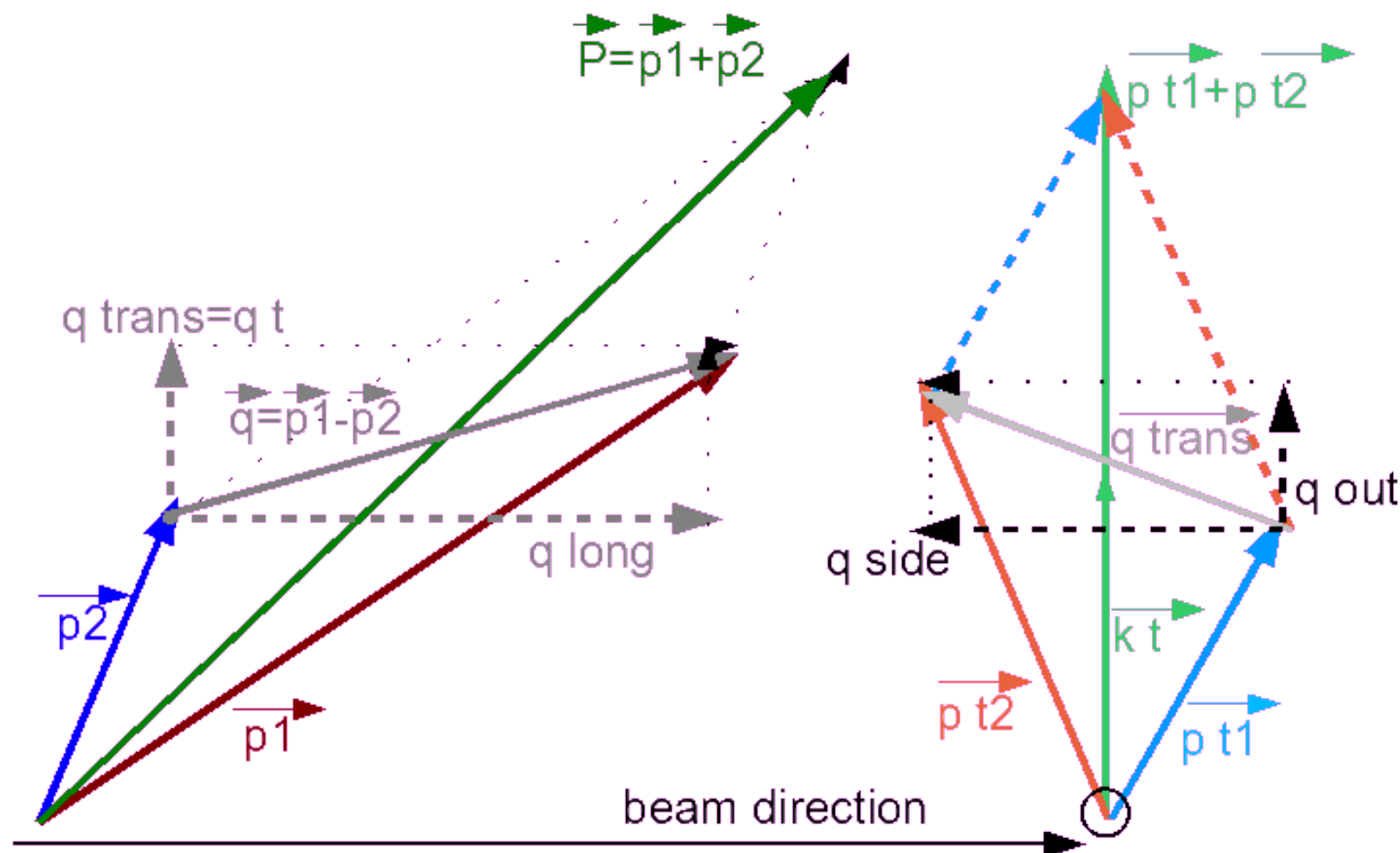


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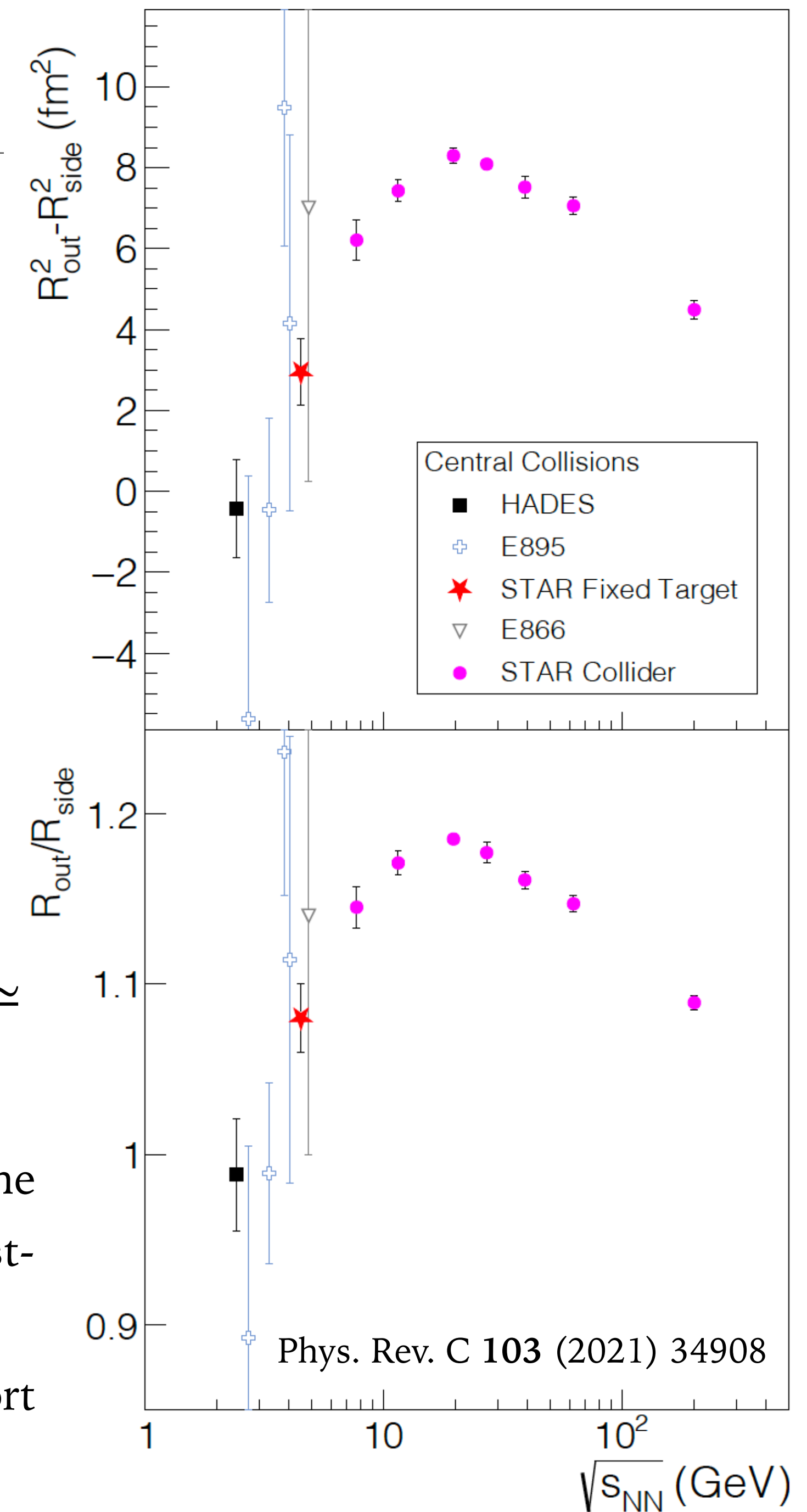
PHENIX Collaboration, arXiv:1410.2559



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- Visible **peak** in $\frac{R_{out}}{R_{side}}(\sqrt{s_{NN}})$ near the $\sqrt{s_{NN}} \approx 20$ GeV
- QCD calculations predict a peak near to the QGP transition threshold - signature of first-order phase transition?
- Theoretical attention from hydro and transport models needed



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Fluctuations and correlations

$$\delta N = N - \langle N \rangle$$

$$C_1 = \langle N \rangle$$

$$C_2 = \langle (\delta N)^2 \rangle = \sigma^2$$

$$C_3 = \langle (\delta N)^3 \rangle$$

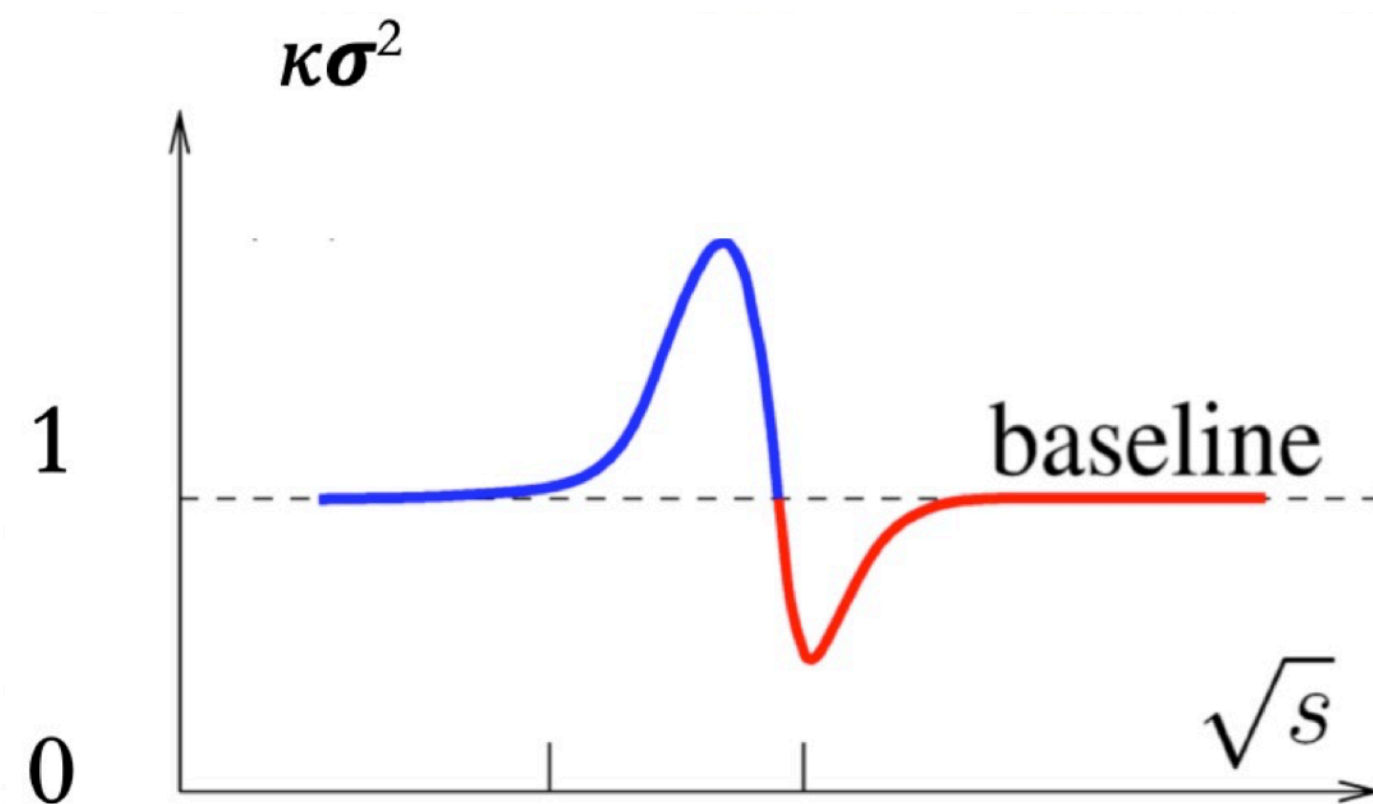
$$C_4 = \langle (\delta N)^4 \rangle - 3\langle (\delta N)^2 \rangle^2$$

$$S\sigma = \frac{C_3}{C_2} \quad \kappa\sigma^2 = \frac{C_4}{C_2}$$

$$C_2 \sim \xi^2 \quad C_3 \sim \xi^{4.5} \quad C_4 \sim \xi^7$$

- Near the QCD CP the divergence of the correlation length expected
- Non-monotonic correlations and fluctuations related to conserved quantities (B, Q, S) could indicate CP
- Higher moments of conserved quantities measure non-Gaussian nature of fluctuations, and are more sensitive (than e.g. variance) to CP fluctuations (leads to correlation length)

The higher cumulant order, the more sensitive to the correlation length



4th order: predicts a non-monotonic energy dependence due to contribution from QCD critical point

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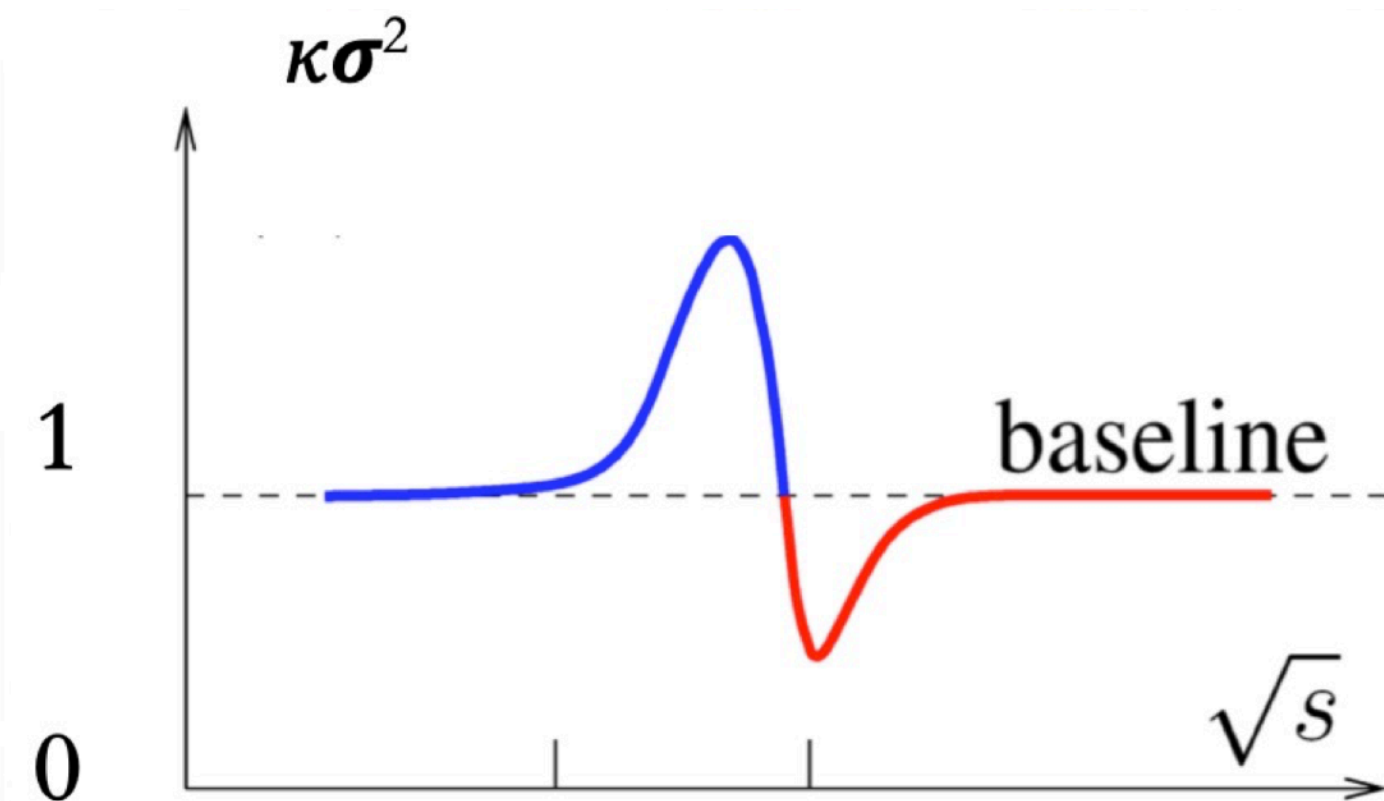
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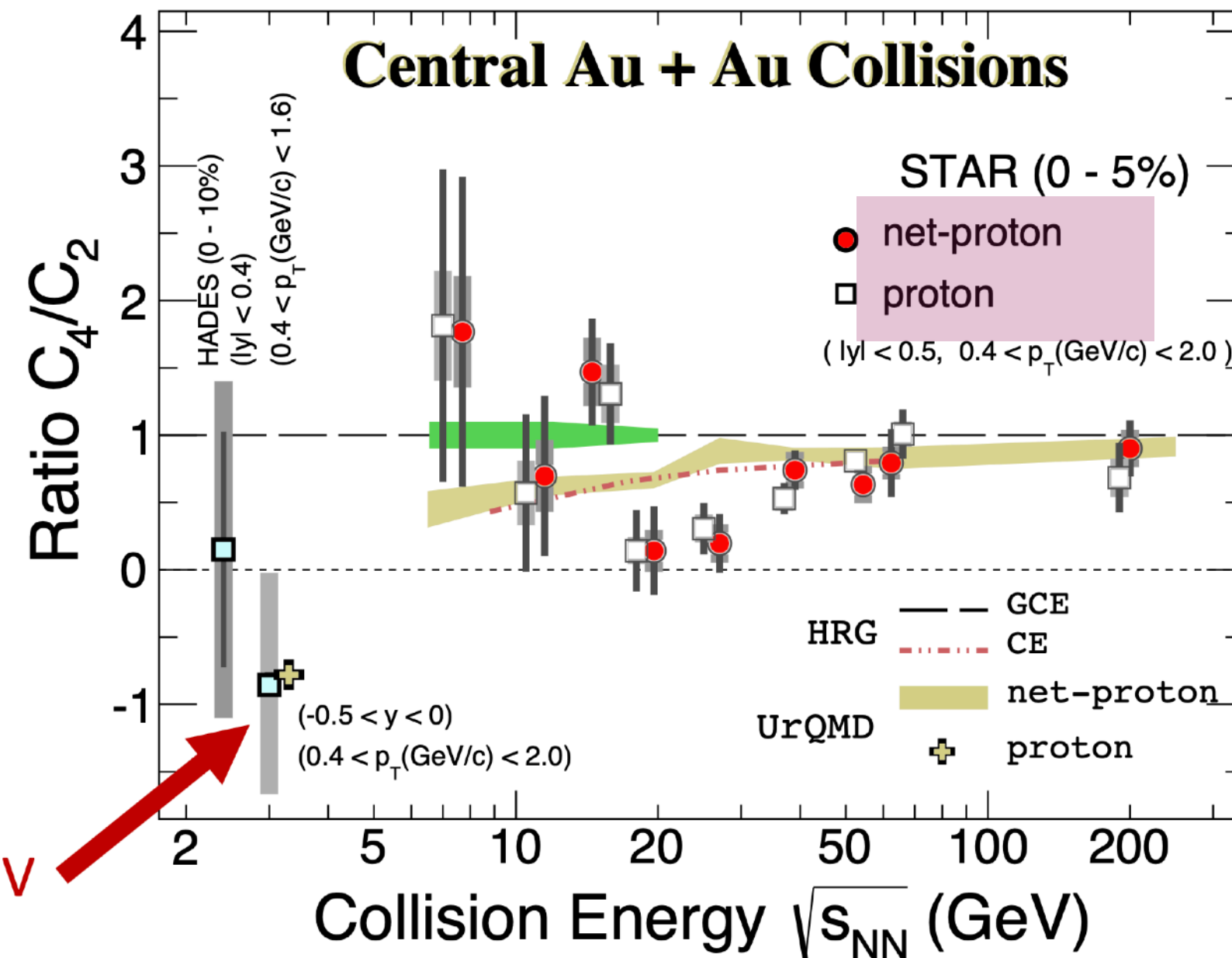
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- Non-monotonic correlations and fluctuations related to conserved quantities (B, Q, S) could indicate CP
- Higher moments of conserved quantities measure non-Gaussian nature of fluctuations, and are more sensitive (than e.g. variance) to CP fluctuations (leads to correlation length)
- Non-monotonic energy dependence of net-proton consistent with theoretical expectation with a CP.
- The suppression of C_4/C_2 consistent with fluctuations driven by baryon number conservation indicating a hadronic interaction dominated region at $\sqrt{s_{NN}} = 3 \text{ GeV}$
- The QCD critical point, (if exists in heavy ion collisions), could be located at $\sqrt{s_{NN}} > 3 \text{ GeV}$; STAR, PRL 126, 092301 (2021), PRC 104.024902 (2021), PRL 128.202303 (2022)

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Thermal dileptons at BES

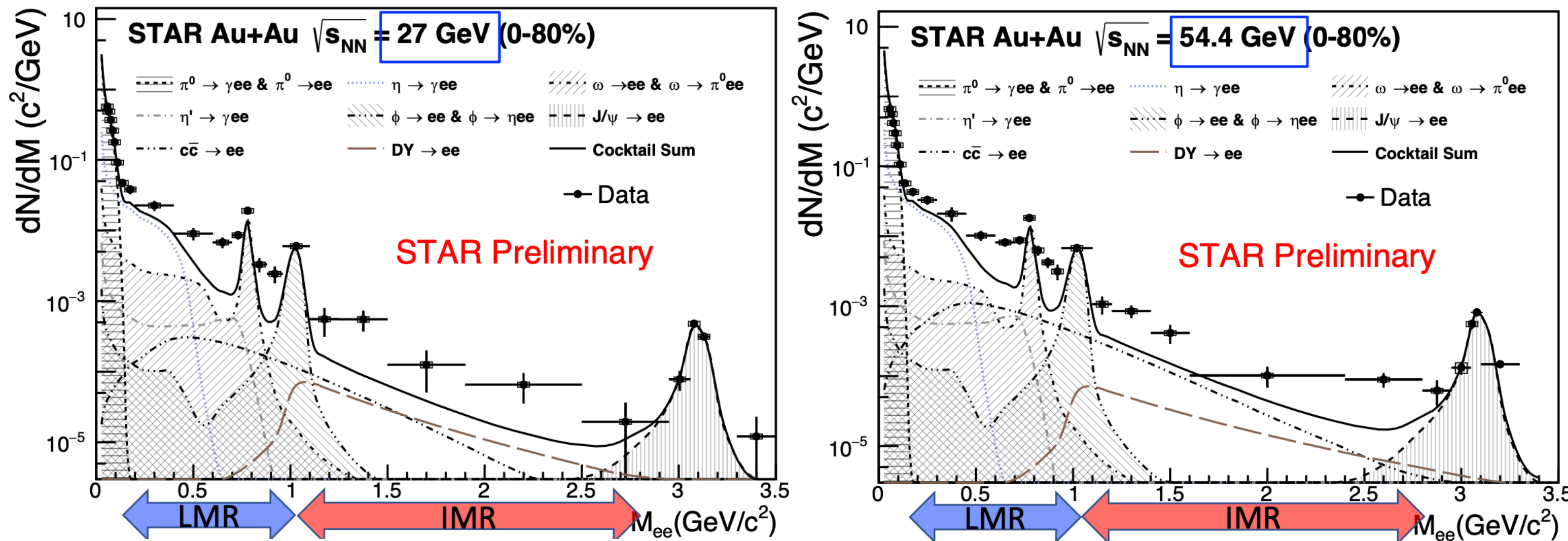
Hadrons (yields, p_T) \rightarrow Statistical thermal models $\rightarrow T_{ch}$
 \rightarrow Hydro-inspired models $\rightarrow T_k$

- EM probes (photons, leptons): emitted from early to final stages of HIC, carry original information of emission source, probe earlier and hotter phases of medium
- Why dileptons?
Mass spectral of thermal dileptons could reveal temperature of the hot medium at both QGP phase and hadronic phase
- How to measure thermal dileptons?
Interested signal (in medium ρ decays, QGP radiation) + **Physics background** (cocktails) = **Inclusive signal** (space-time integral)

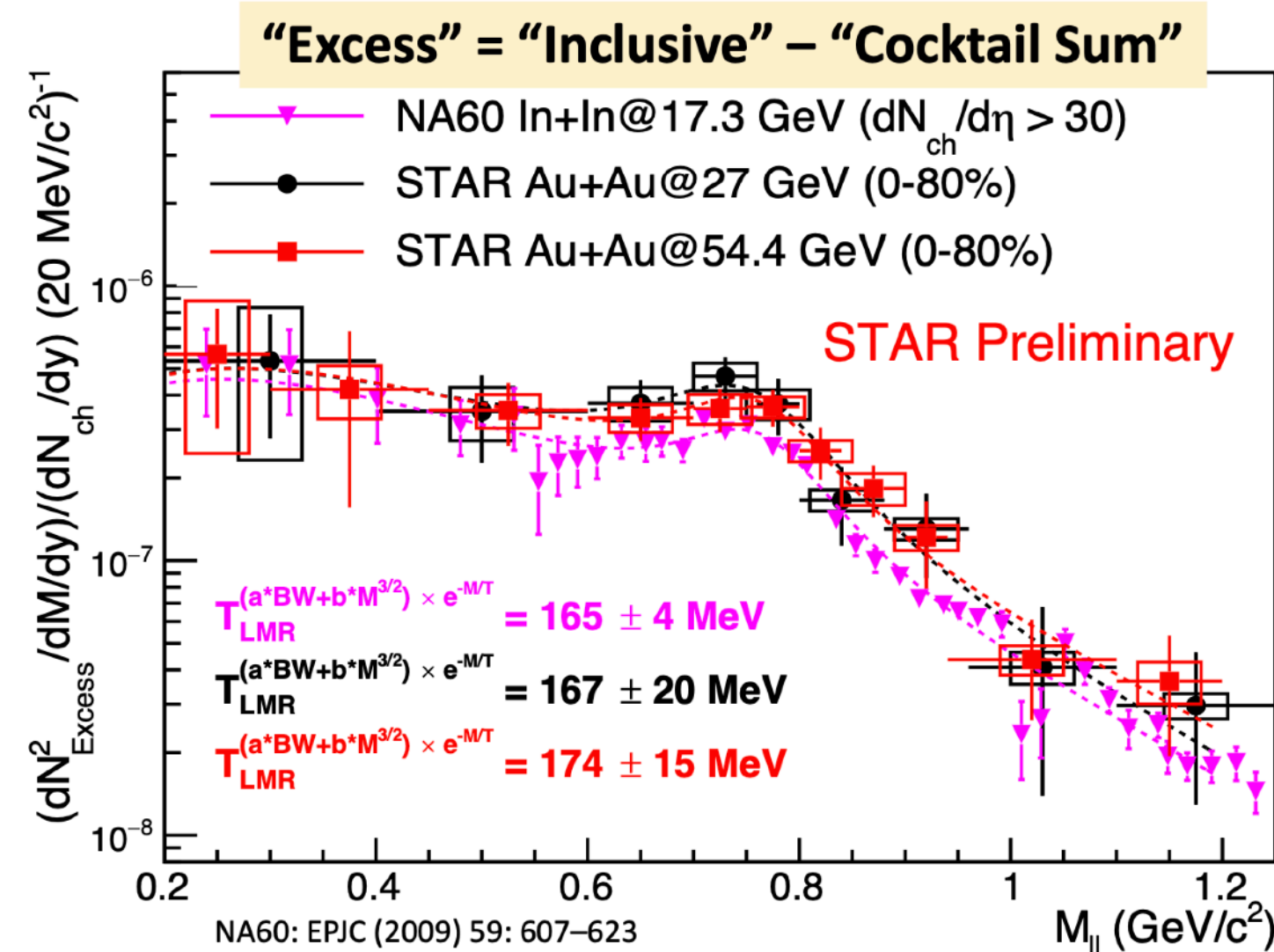
Thermal dileptons at BES

Hadrons (yields, p_T) \rightarrow Statistical thermal models $\rightarrow T_{ch}$
 \rightarrow Hydro-inspired models $\rightarrow T_k$

- EM probes (photons, leptons): emitted from early to final stages of HIC, carry original information of emission source, probe earlier and hotter phases of medium



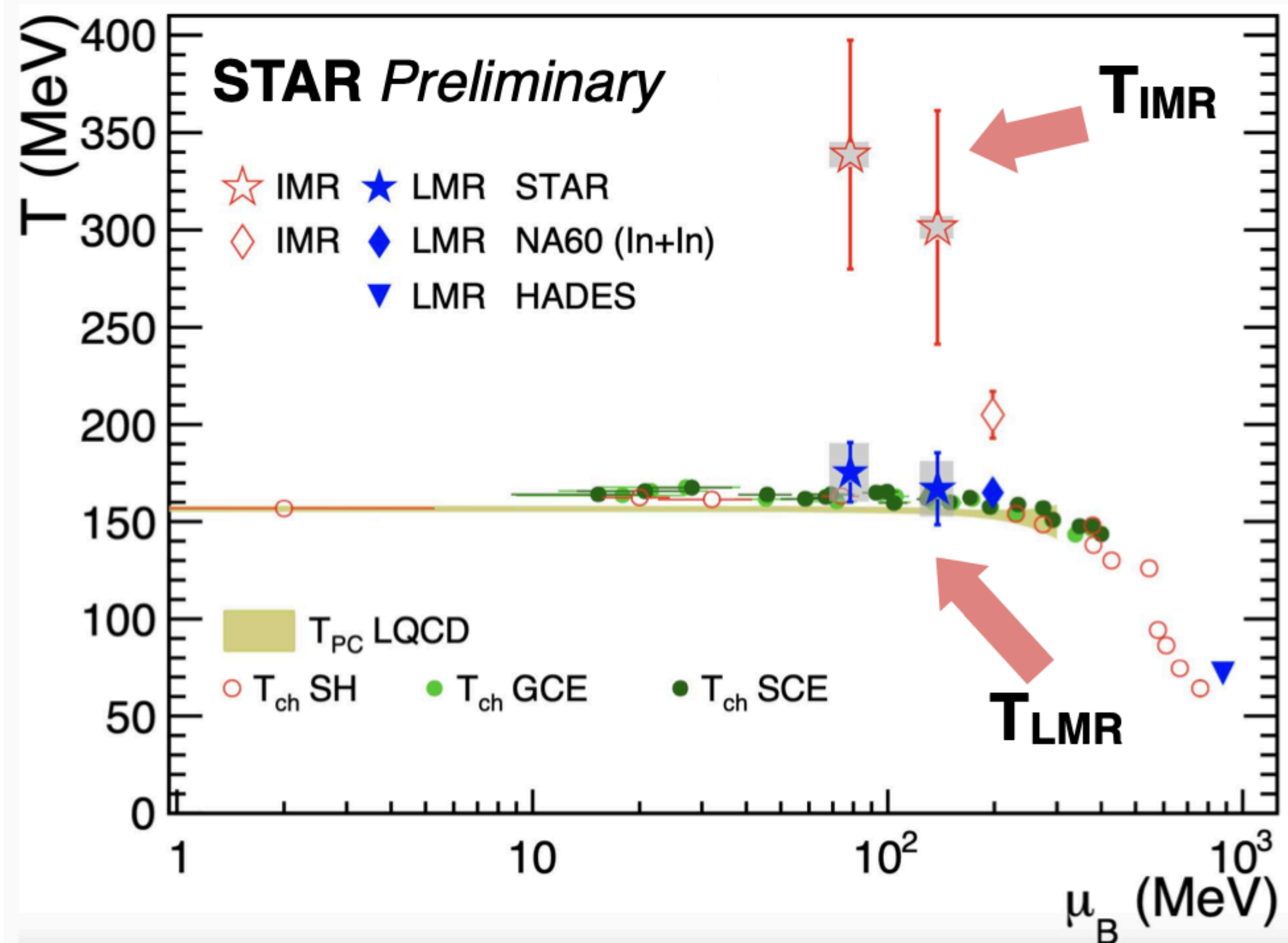
- Why dileptons? Mass spectral of thermal dileptons could reveal temperature of the hot medium at both QGP phase and hadronic phase



T are similar, despite significant difference in collision species and energies, indicate that ρ is dominantly emitted around phase transition

- How to measure thermal dileptons? **Interested signal** (in medium ρ decays, QGP radiation) + **Physics background** (cocktails) = **Inclusive signal** (space-time integral)

Thermal dileptons at BES



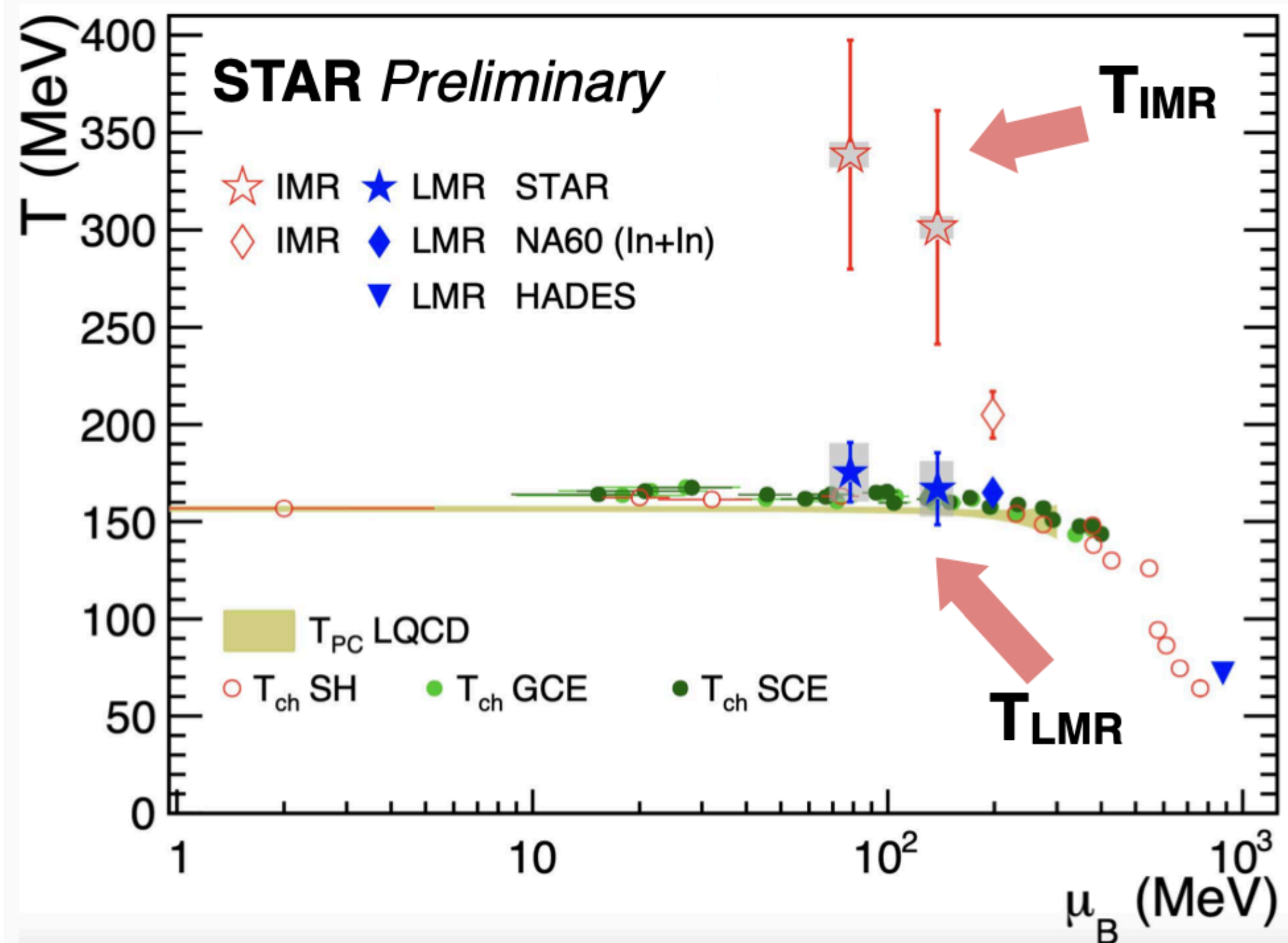
Thermal dileptons in IMR

- T always higher than phase transition T at RHIC and SPS
- Emitted from QGP phase

Thermal dileptons in LMR

- T consistent with T_{ch} from statistical models
- Emitted from hadronic phase, dominantly around phase transition region at RHIC and SPS

Thermal dileptons at BES



Thermal dileptons in IMR

- T always higher than phase transition T at RHIC and SPS
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Thermal dileptons in LMR

- T consistent with T_{ch} from statistical models
- Emitted from hadronic phase, dominantly around phase transition region at RHIC and SPS

- $T_{LMR} \sim 170$ MeV: First experimental evidence that in-medium ρ (hadronic matter) is dominantly produced around temperature higher than T at the phase transition
- Normalized dilepton yield is higher in RHIC Au+Au than that in SPS In+In: indicate a longer medium lifetime for larger collision systems

$T_{IMR} \sim 300$ MeV: First QGP temperature measurement at RHIC

QGP produced at RHIC is hotter than that at SPS

Summary from BES

Continue to look for the **Critical Point** and the **first-order phase transition**.

Data exists over wide range of the collision energy for heavy and light ions, p(d), and pp

High statistics exploration of QCD phase diagram and its key features has already begun

More coming soon (BES-II, SPS, FAIR)

Turn trends and features into definite conclusions

More interesting questions appeared..

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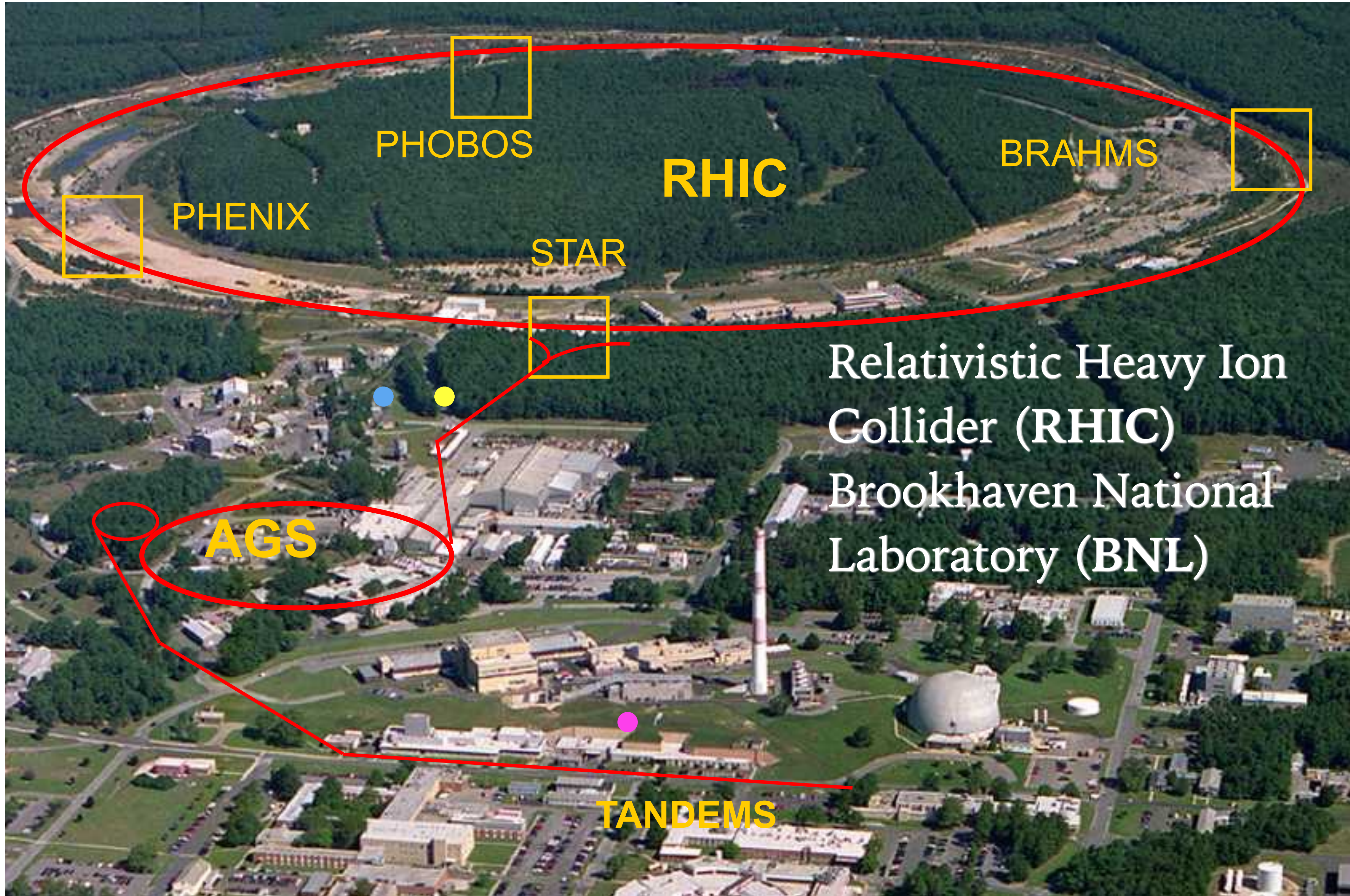
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*Thank You for
your attention!*

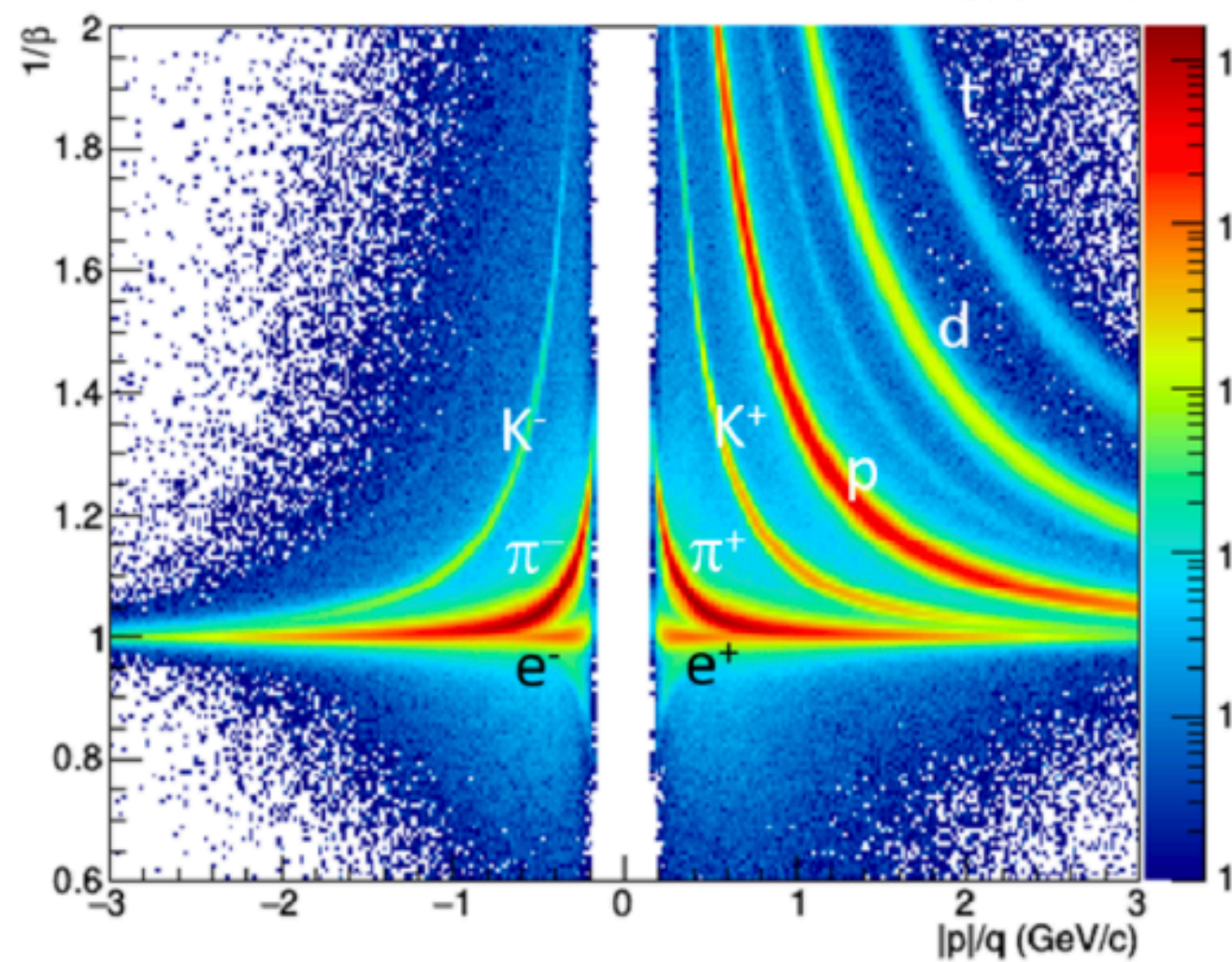
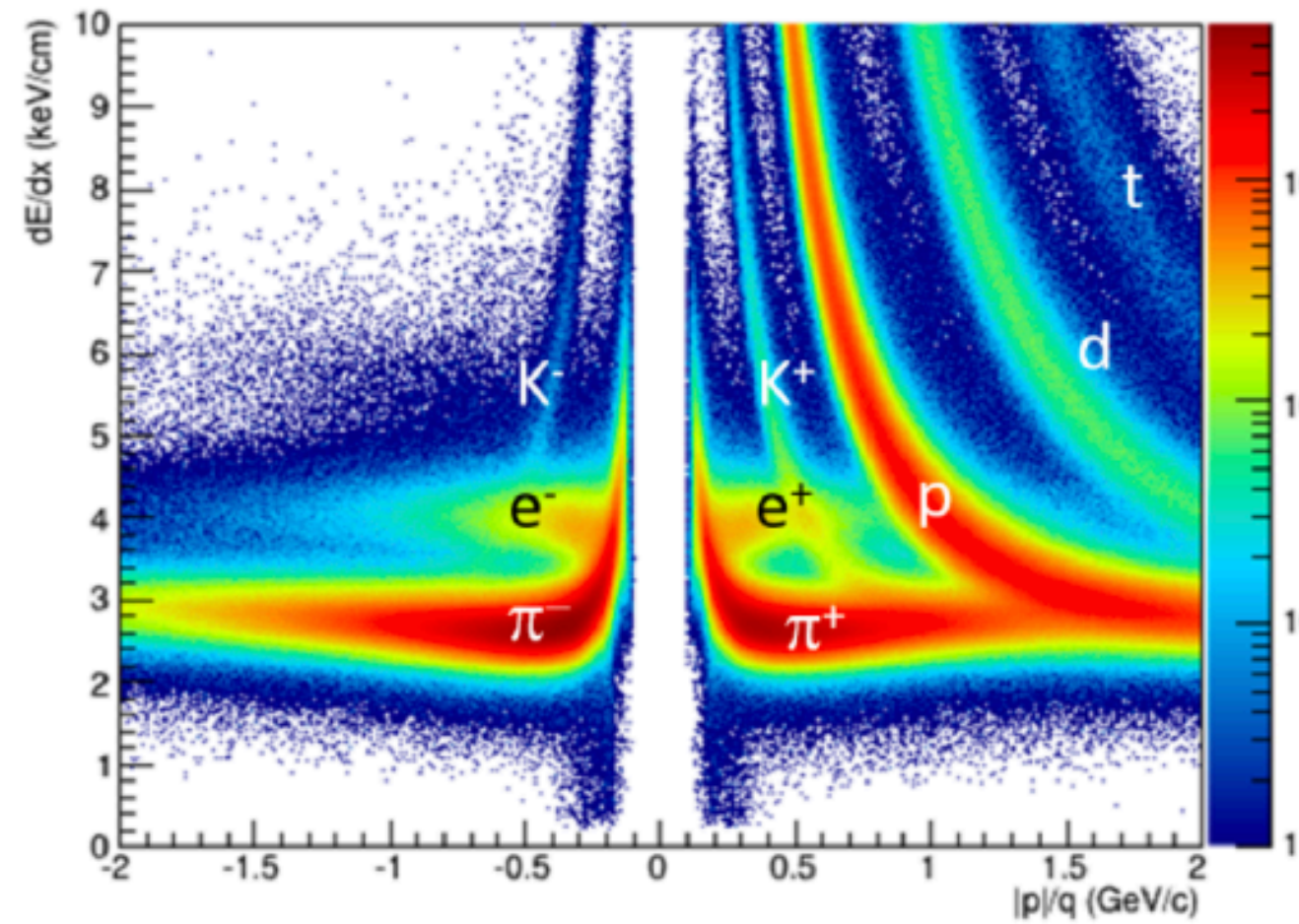
- 3.83 km circumference
- Two independent rings
- Collides so far:
 - Au+Au, p+p, d+Au, Cu+Cu, U+U, Cu+Au, $^3\text{He}+\text{Au}$, p+Au
- Top Center-of-Mass Energy
 - 510 GeV for p-p
 - 200 GeV/nucleon for Au-Au



Particle identification

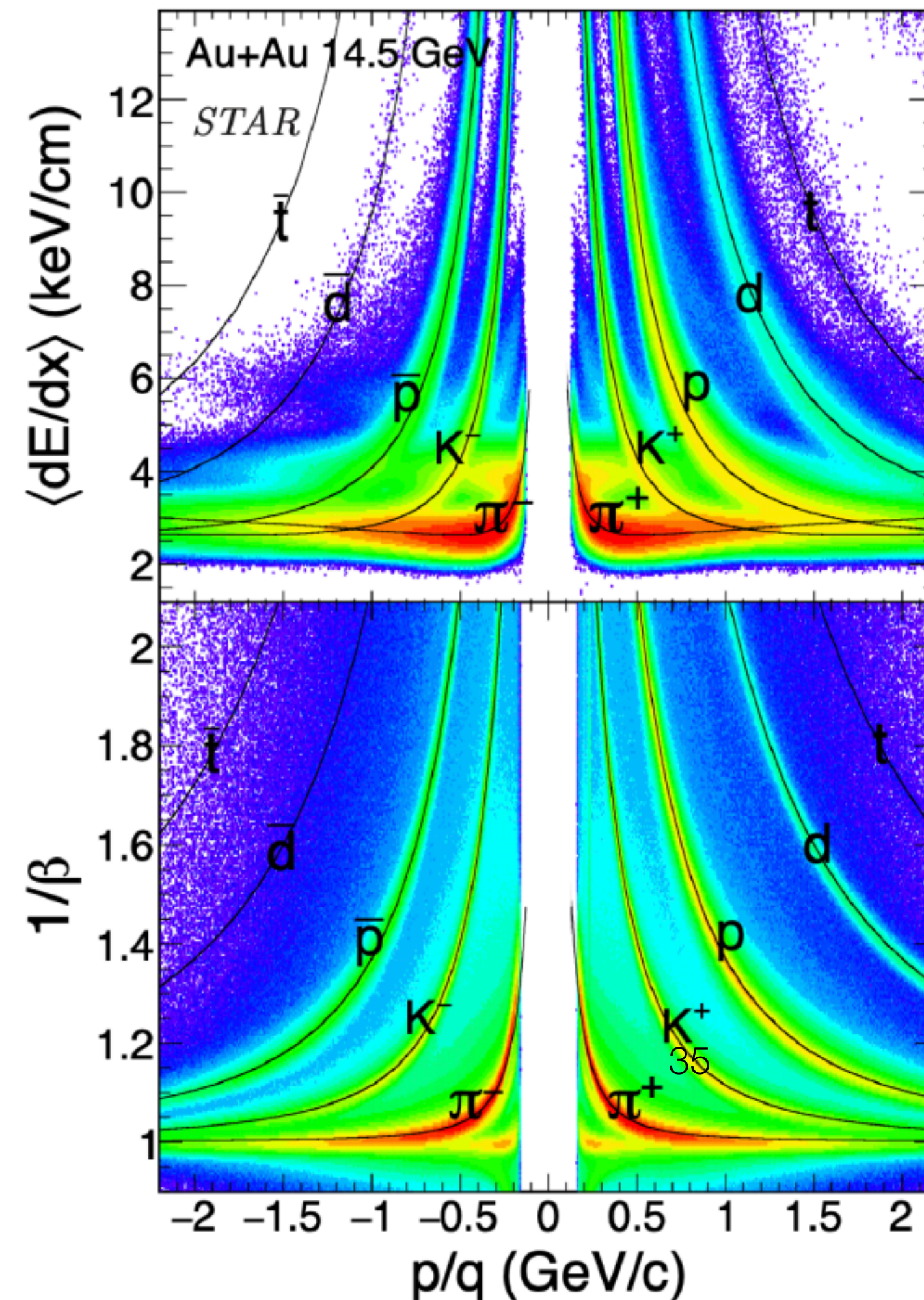
arXiv:2007.14005 (STAR)

$\sqrt{s_{NN}}=4.5\text{GeV}$ Fixed Target



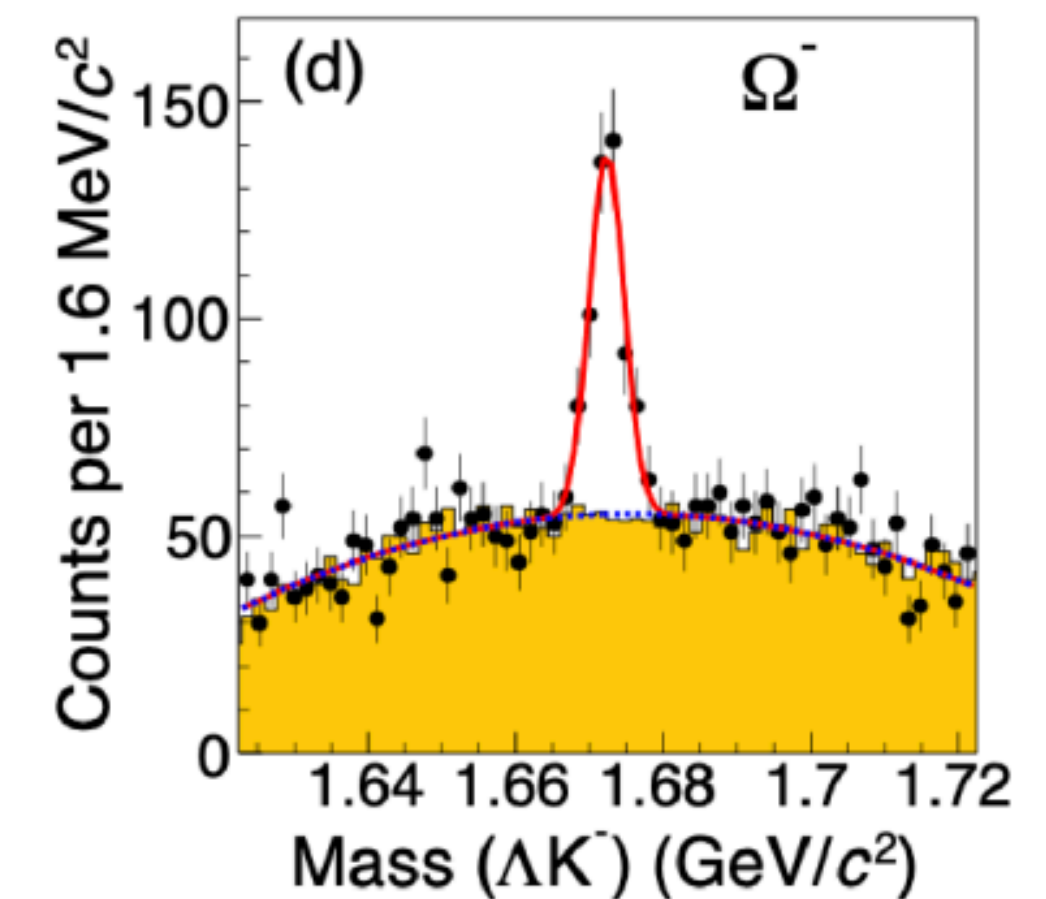
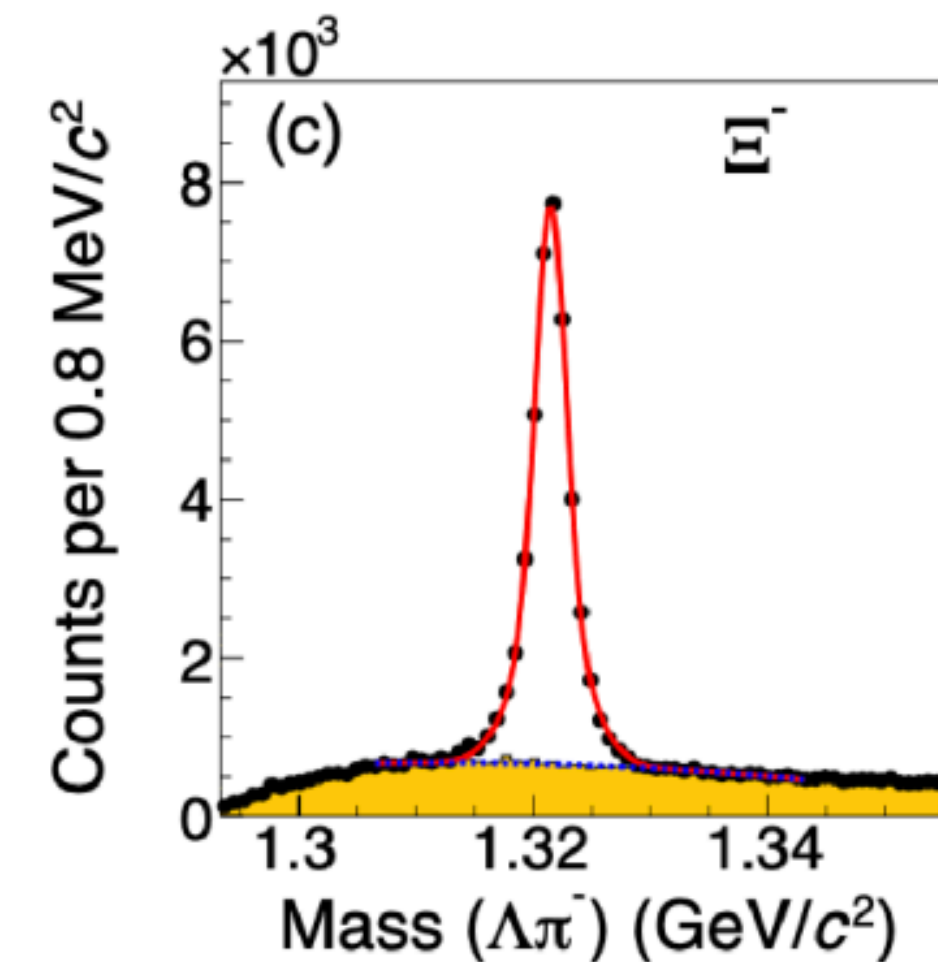
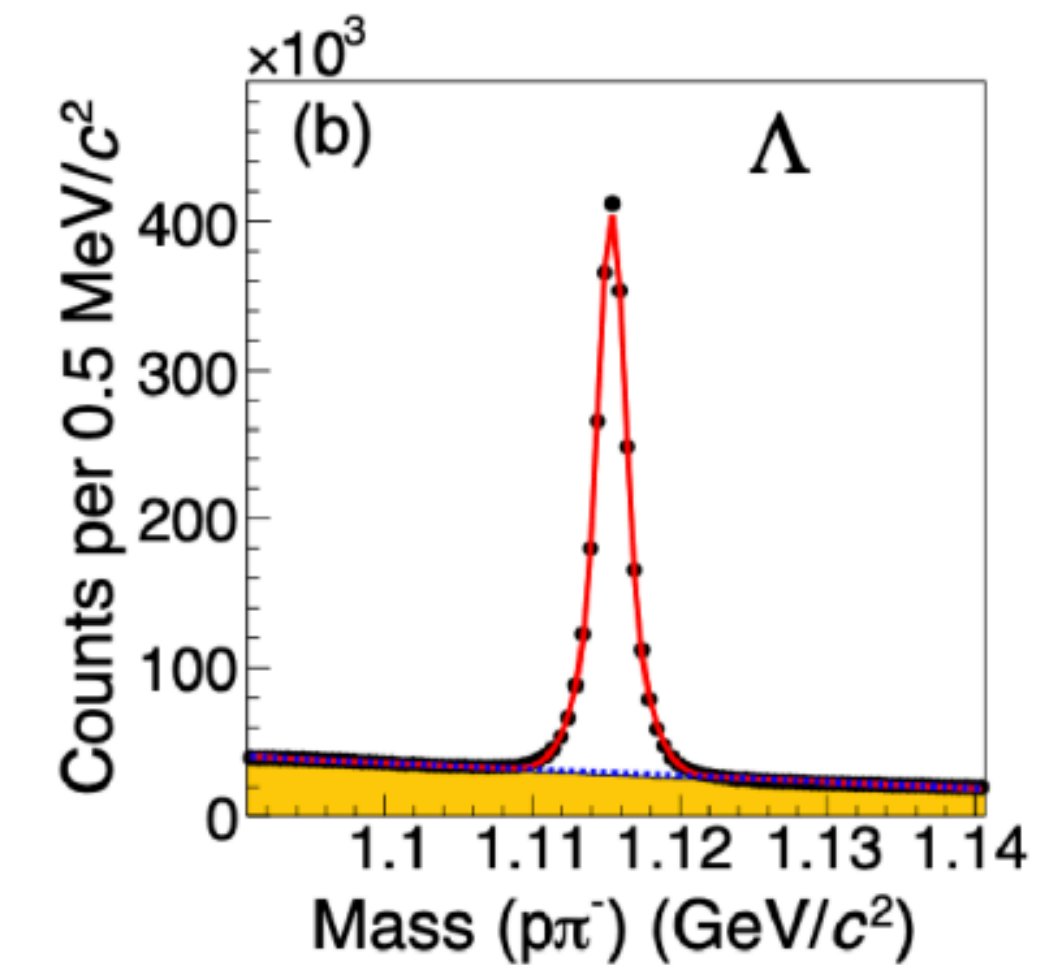
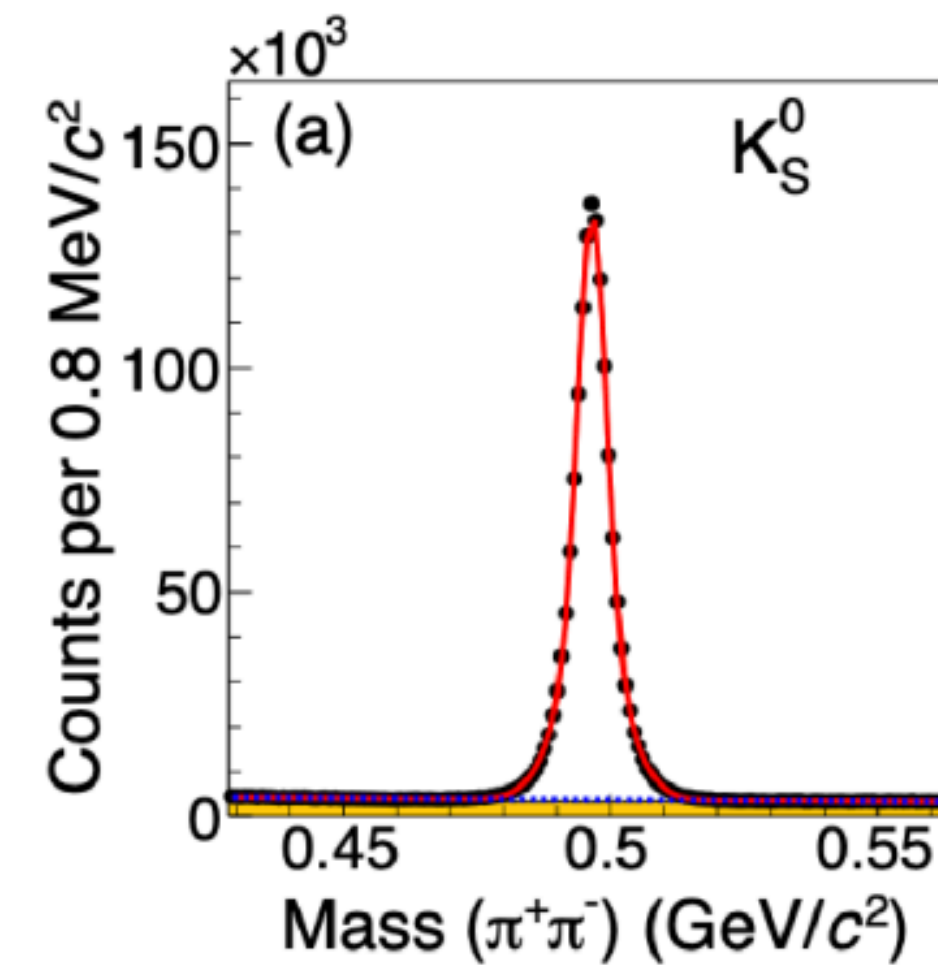
PRC 101 (2020) 24905 (STAR)

$\sqrt{s_{NN}}=14.5\text{GeV}$ Collider



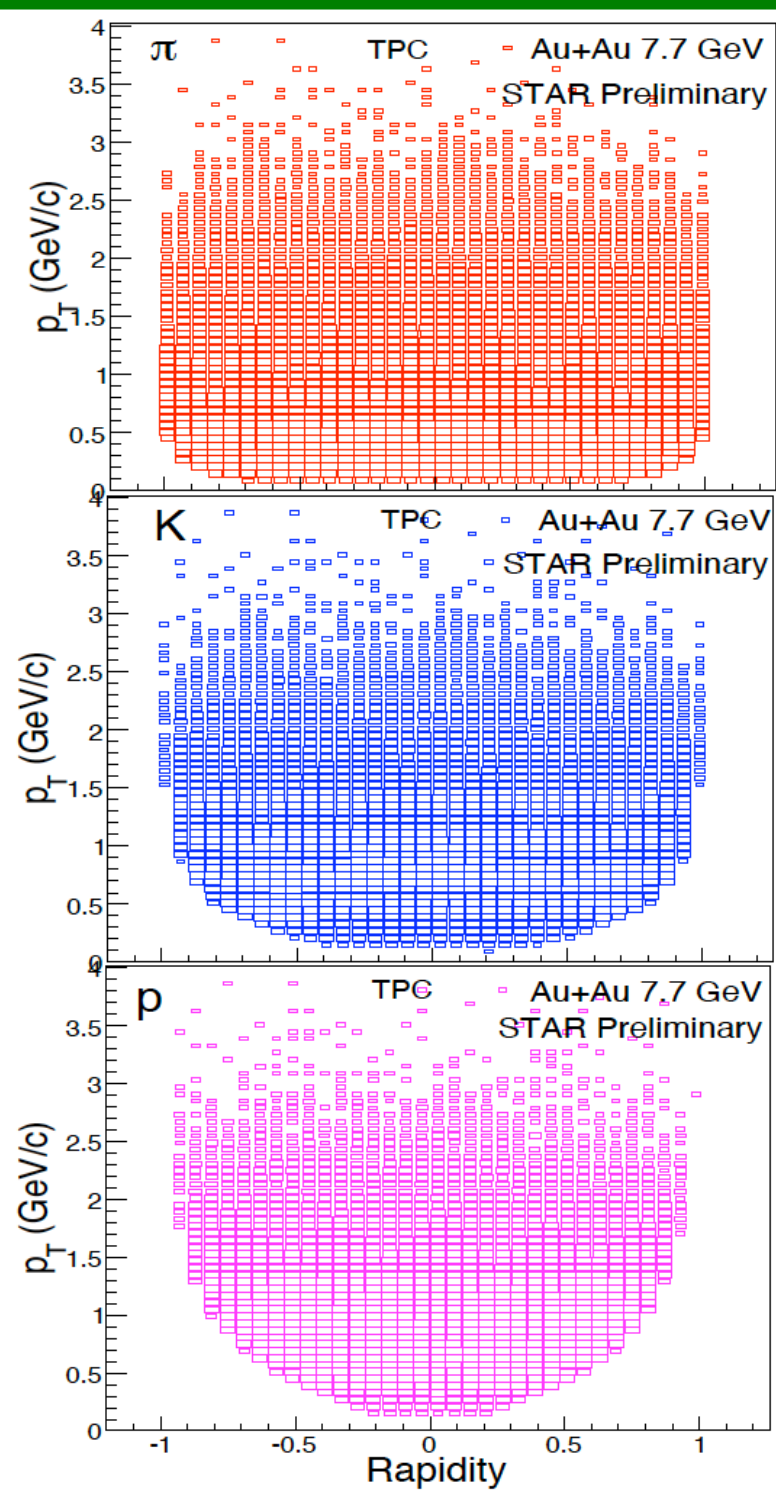
PRC 102 (2020) 34909 (STAR)

$\sqrt{s_{NN}}=7.7\text{GeV}$ Collider

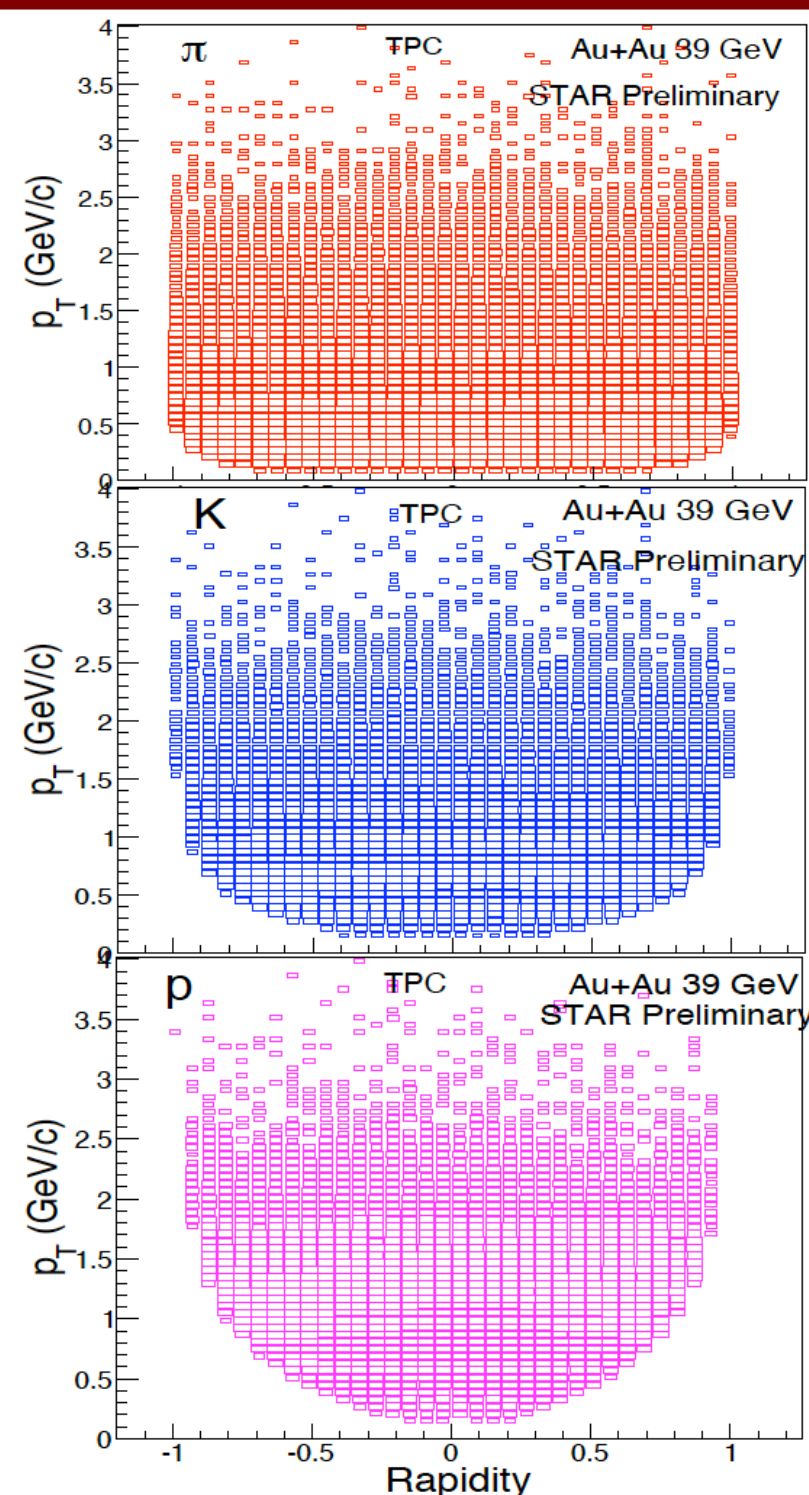


Collider and fixed-target acceptance

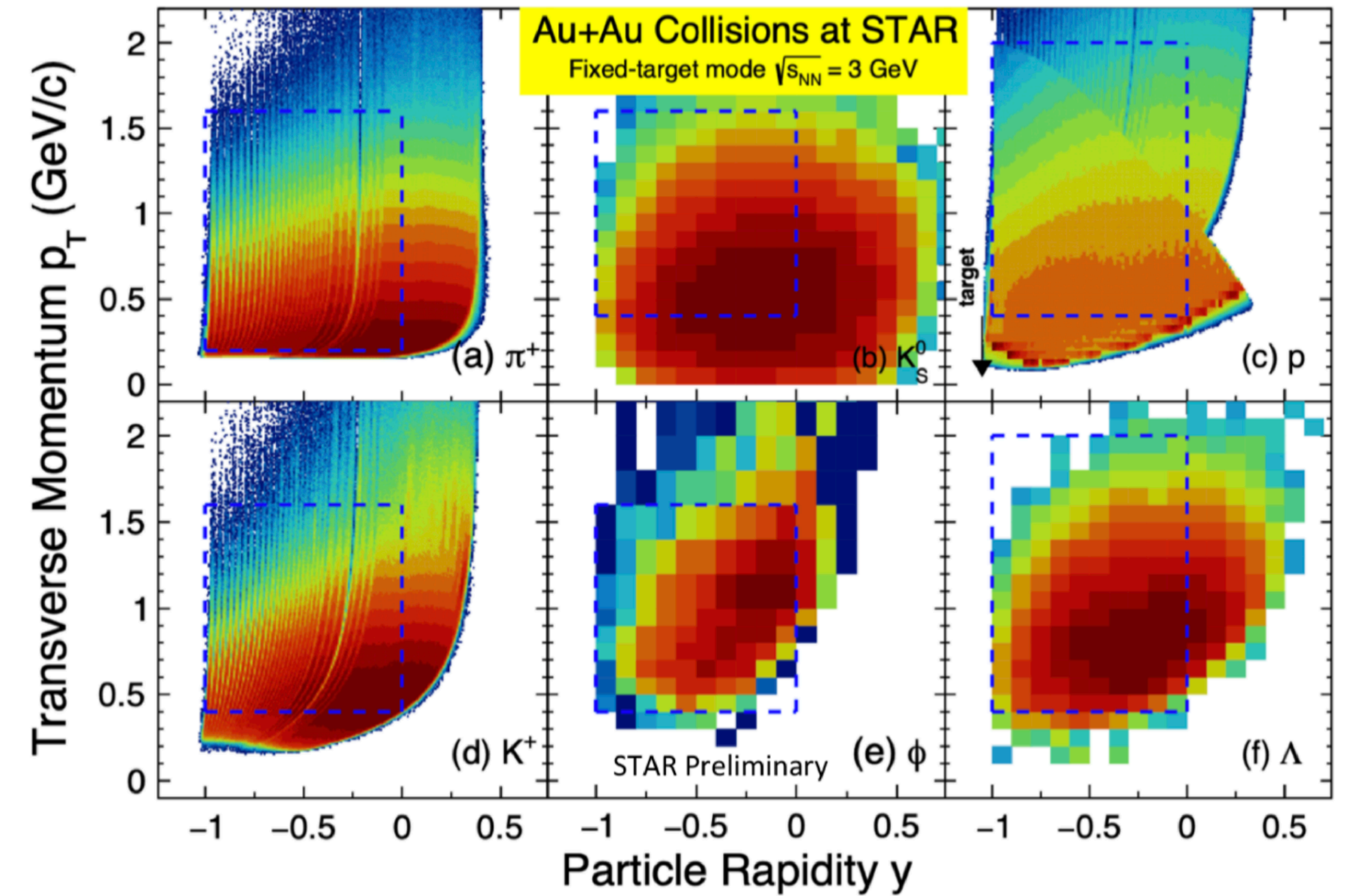
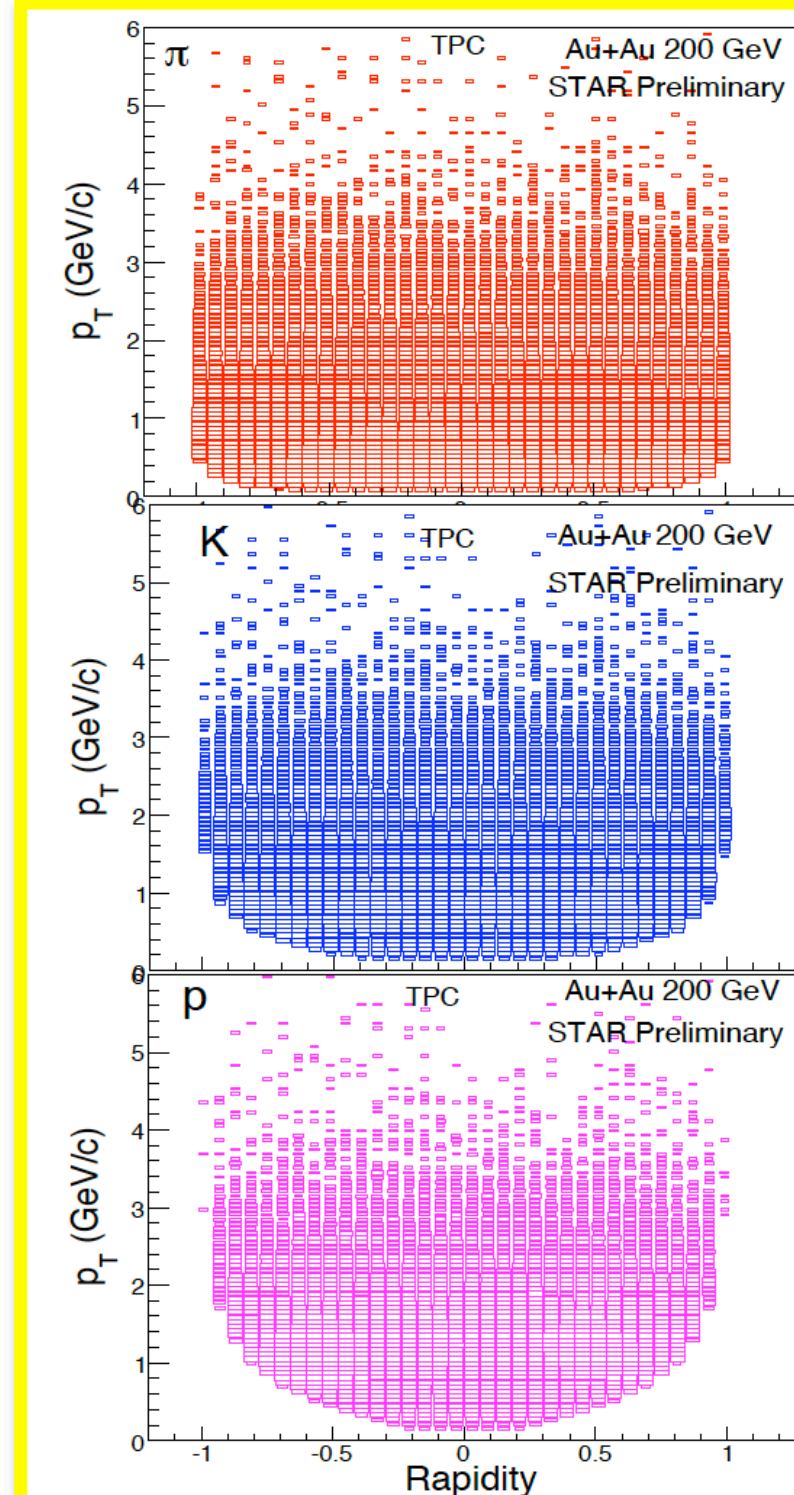
Au+Au at 7.7 GeV



Au+Au at 39 GeV

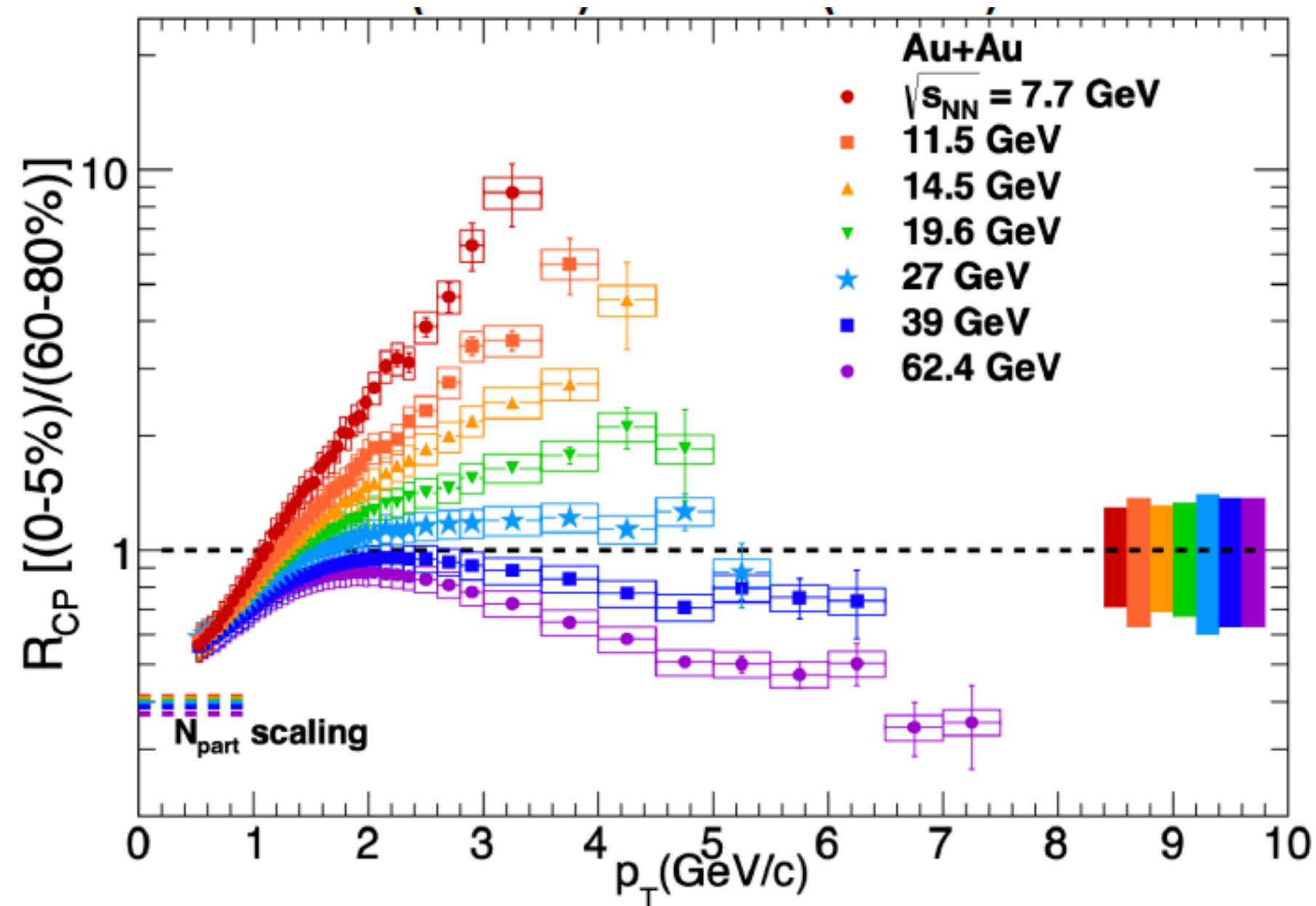


Au+Au at 200 GeV



Similar acceptance for all particles and energies

R_{CP} for charged particles



PRL 121 (2018) 32301 (STAR)

$$R_{CP} = \frac{\langle N_{bin}^{per} \rangle d^3 N_{AA}^{cen} / d\eta d^2 p_T}{\langle N_{bin}^{cen} \rangle d^3 N_{AA}^{per} / d\eta d^2 p_T}$$

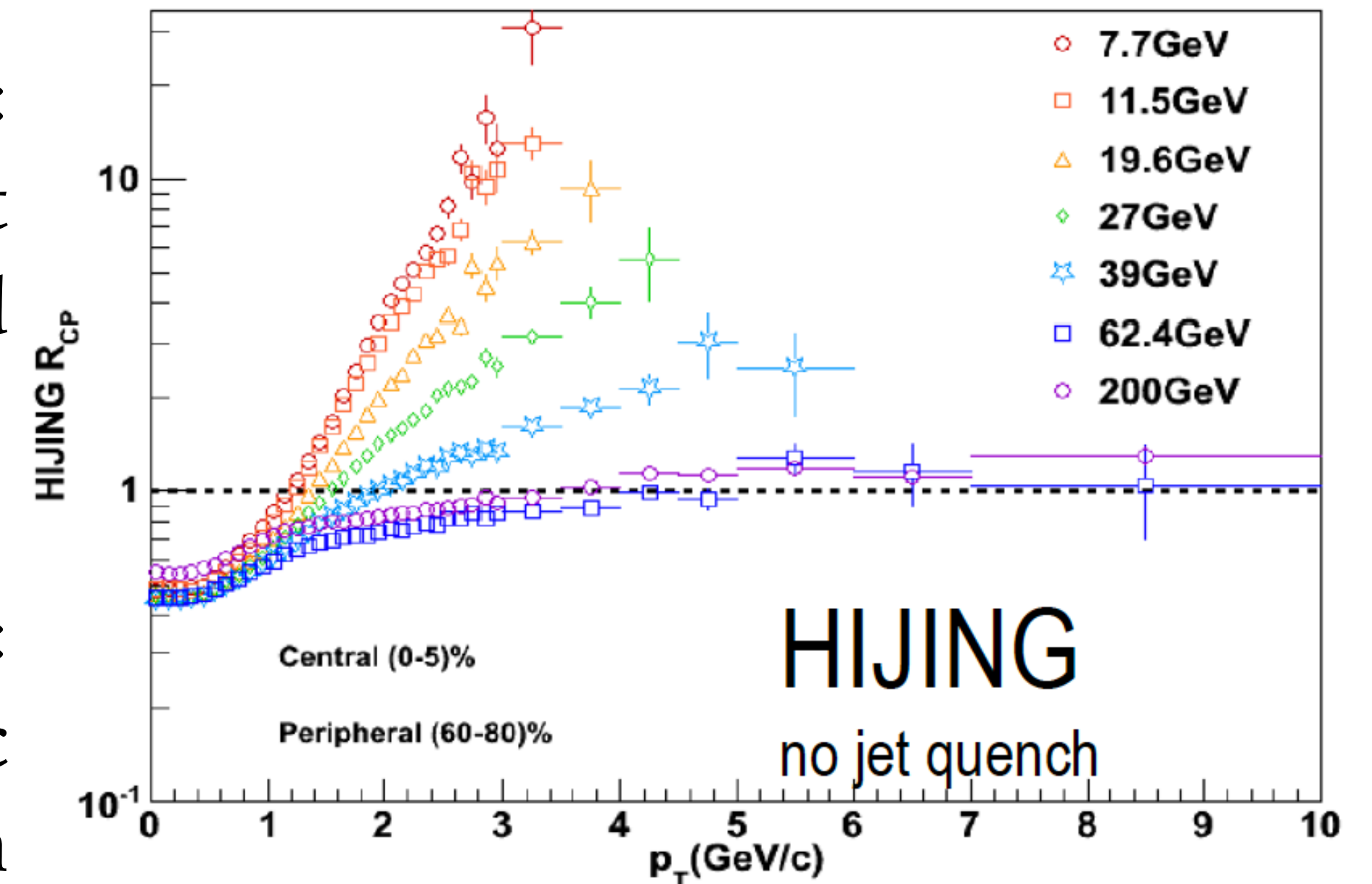
increases from suppression at $\sqrt{s_{NN}} = 62.4$ GeV to enhancement at $\sqrt{s_{NN}} = 7.7$ GeV.

Suppression $R_{CP} < 1$:
opacity of deconfinement medium of quarks and gluons.

Enhancement $R_{CP} > 1$:
dominance of hadronic interactions: Cronin effect, cold matter effect, radial flow, ...

$R_{CP} > 1$ does not mean that QGP is not formed.

Lack of baseline $p + p$ and $p + A$ measurements



HIJING (including Cronin effect) similar to $\sqrt{s_{NN}}$ dependence at low energies

Higher order cumulants

$$\delta N = N - \langle N \rangle$$

$$C_1 = \langle N \rangle$$

$$C_2 = \langle (\delta N)^2 \rangle = \sigma^2$$

$$C_3 = \langle (\delta N)^3 \rangle$$

$$C_4 = \langle (\delta N)^4 \rangle - 3 \langle (\delta N)^2 \rangle^2$$

$$S\sigma = \frac{C_3}{C_2} \quad \kappa\sigma^2 = \frac{C_4}{C_2}$$

Factorial cumulants:

$$\kappa_1 = C_1$$

$$\kappa_2 = -C_1 + C_2$$

$$\kappa_3 = 2C_1 - 3C_2 + C_3$$

$$\kappa_4 = -6C_1 + 11C_2 - 6C_3 + C_4$$

$$\kappa_5 = 24C_1 - 50C_2 + 35C_3 - 10C_4 + C_5$$

$$\kappa_6 = -120C_1 + 27C_2 - 225C_3 + 85C_4 - 15C_5 + C_6$$

Cross-over: $C_5, C_6 < 0$ - LQCD, FRG

No transition: $C_5, C_6 > 0$ - HRG, UrQMD

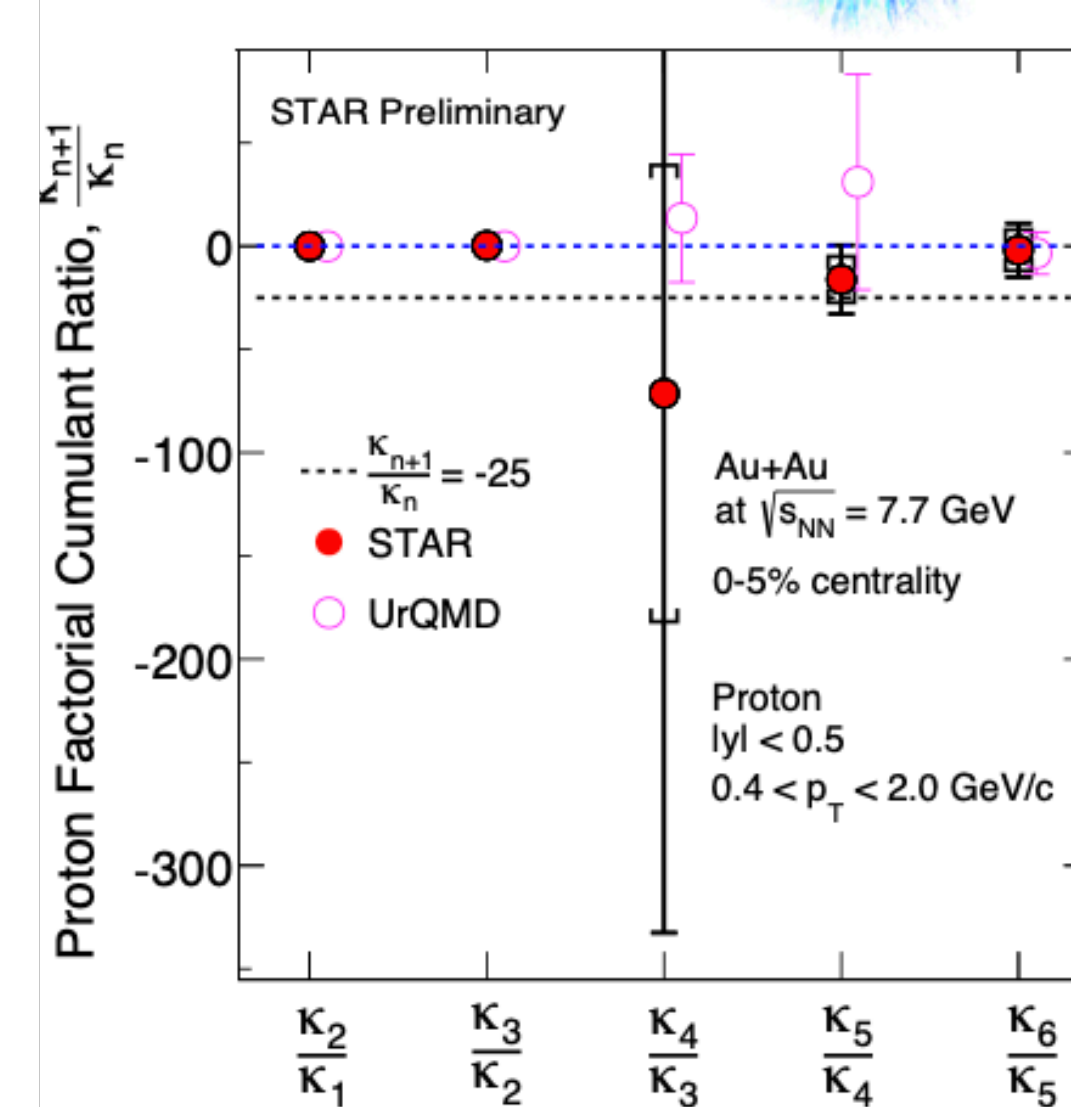
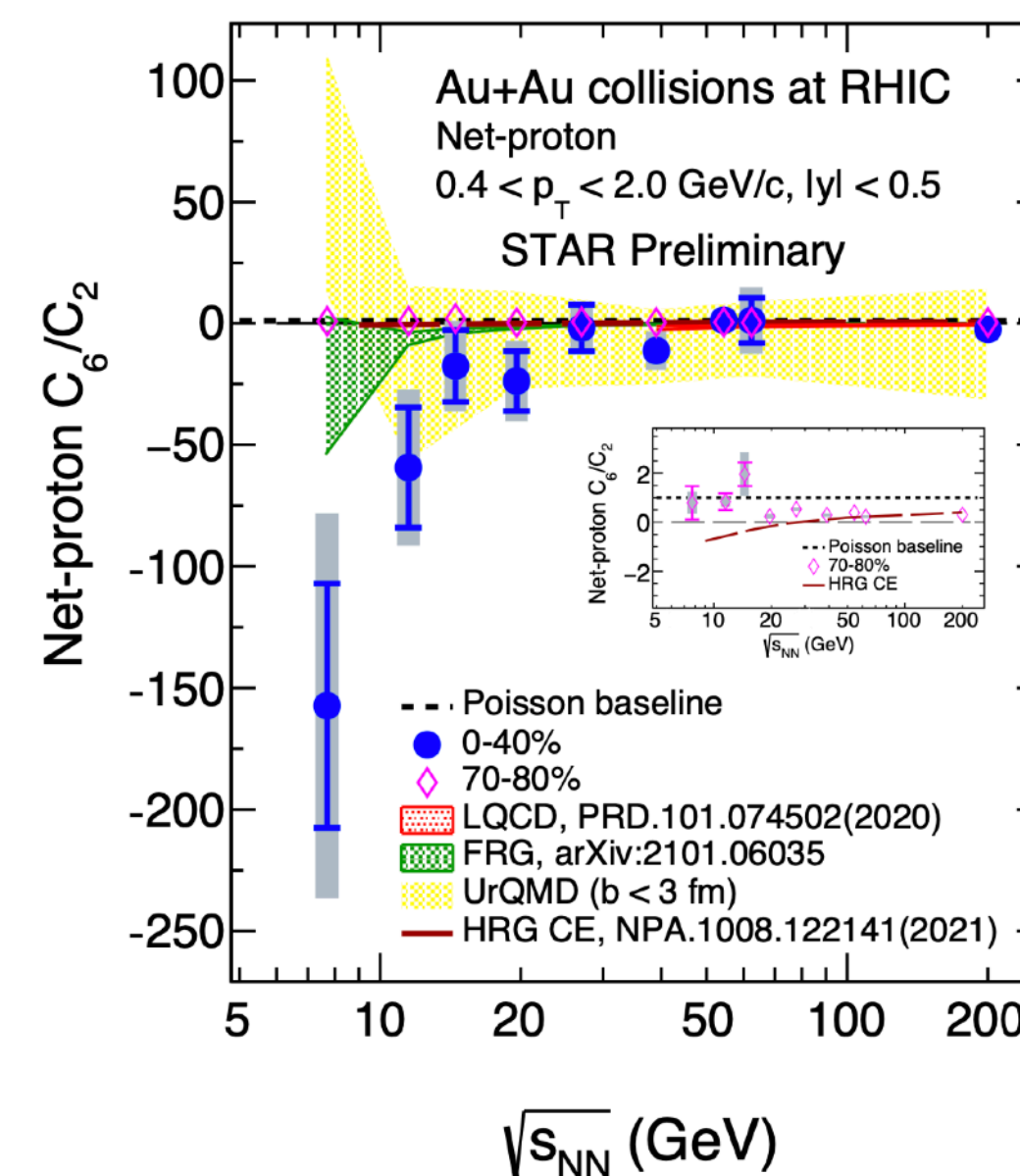
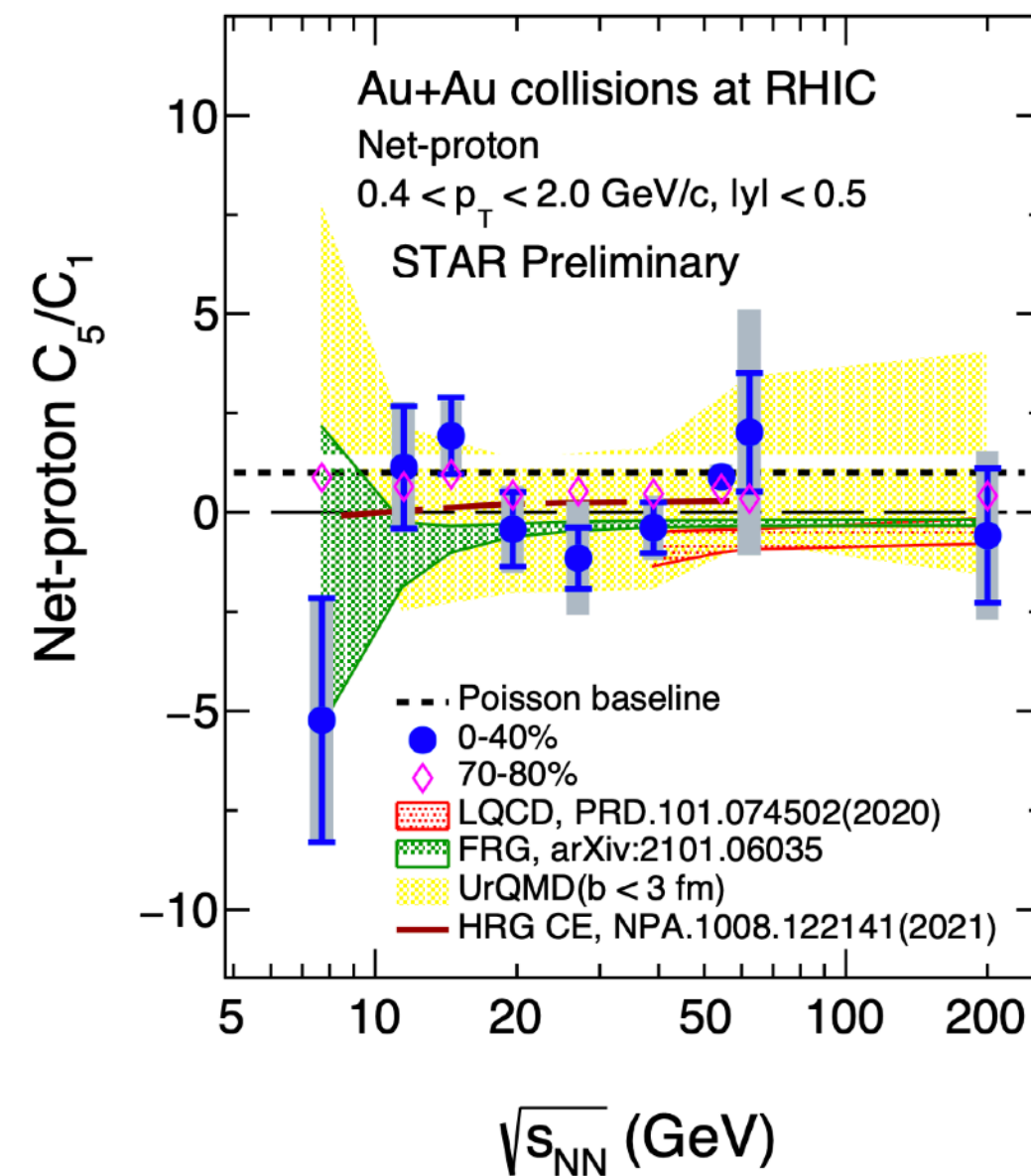
First order phase transition:

- multiplicity distribution bi-modal
(contribution from two phases)

- $\kappa_{n+1}/\kappa_n < 0$ for higher orders

$$C_5 = \langle (\delta N)^5 \rangle - 5 \langle (\delta N)^3 \rangle \langle (\delta N)^2 \rangle$$

$$C_6 = \langle (\delta N)^6 \rangle - 15 \langle (\delta N)^4 \rangle \langle (\delta N)^2 \rangle - 10 \langle (\delta N)^3 \rangle^2 + 30 \langle (\delta N)^2 \rangle^3$$



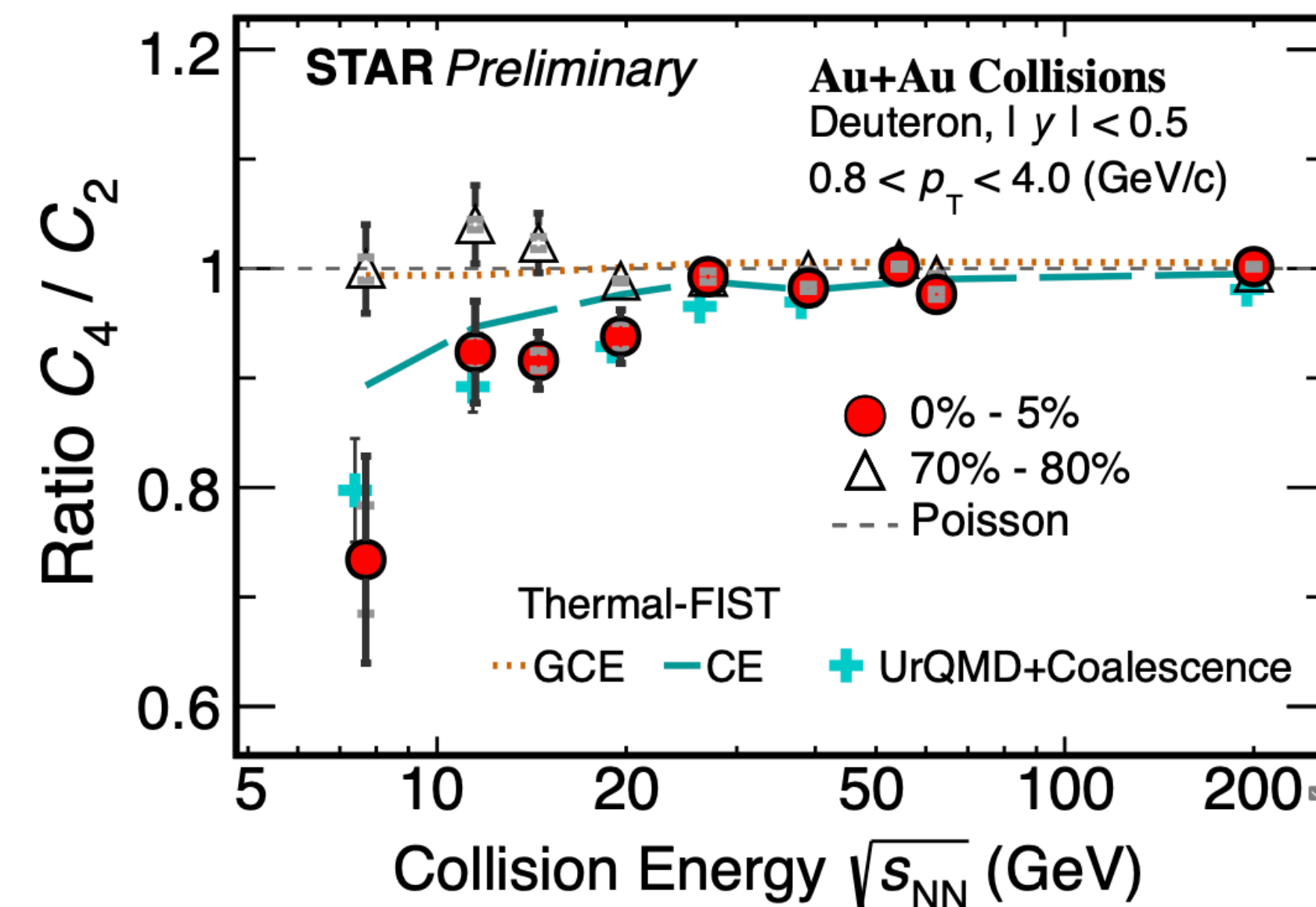
κ_5/κ_4 and κ_6/κ_5 consistent with zero.

High order fluctuations are crucial for determining the QCD phase structure. Precision measurements are necessary in order to confirm the observed trend in the fifth and sixth order cumulants.

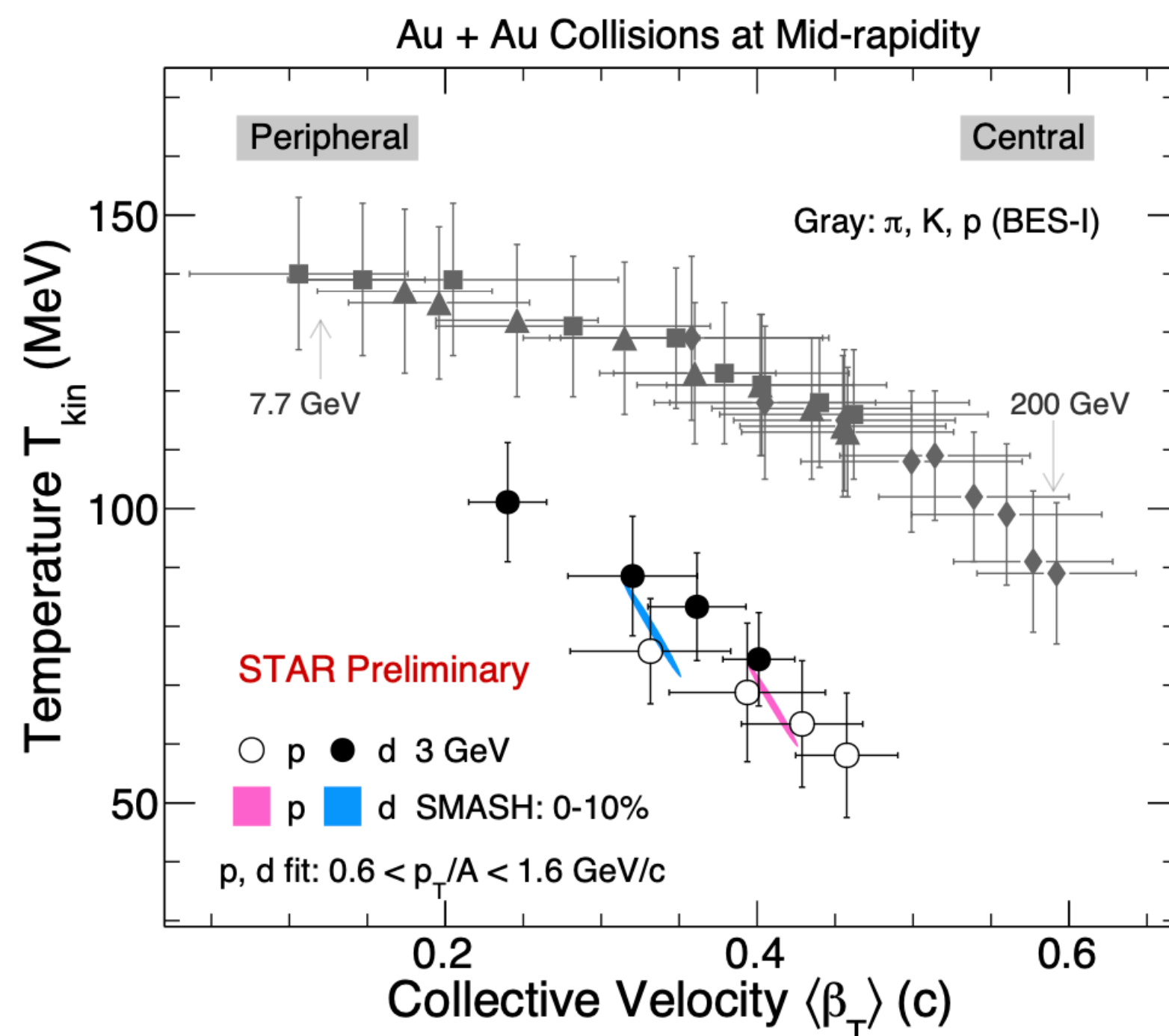
Light nuclei

1. Light nuclei carry information about local baryon density fluctuations
2. It provides a probe to study first-order phase transition and the QCD Critical Point

$\sqrt{s_{NN}} = 3$ GeV collisions: First measurements of yields of p and light nuclei ($d, t, {}^3\text{He}, {}^4\text{He}$)



- Difference with net-proton: role of different freeze-out and smaller yield of deuterons should be investigated,
- With lowering the collision energy, C_4/C_2 is decreasing,
- **UrQMD+Coalescence model reproduces collision energy dependence.**

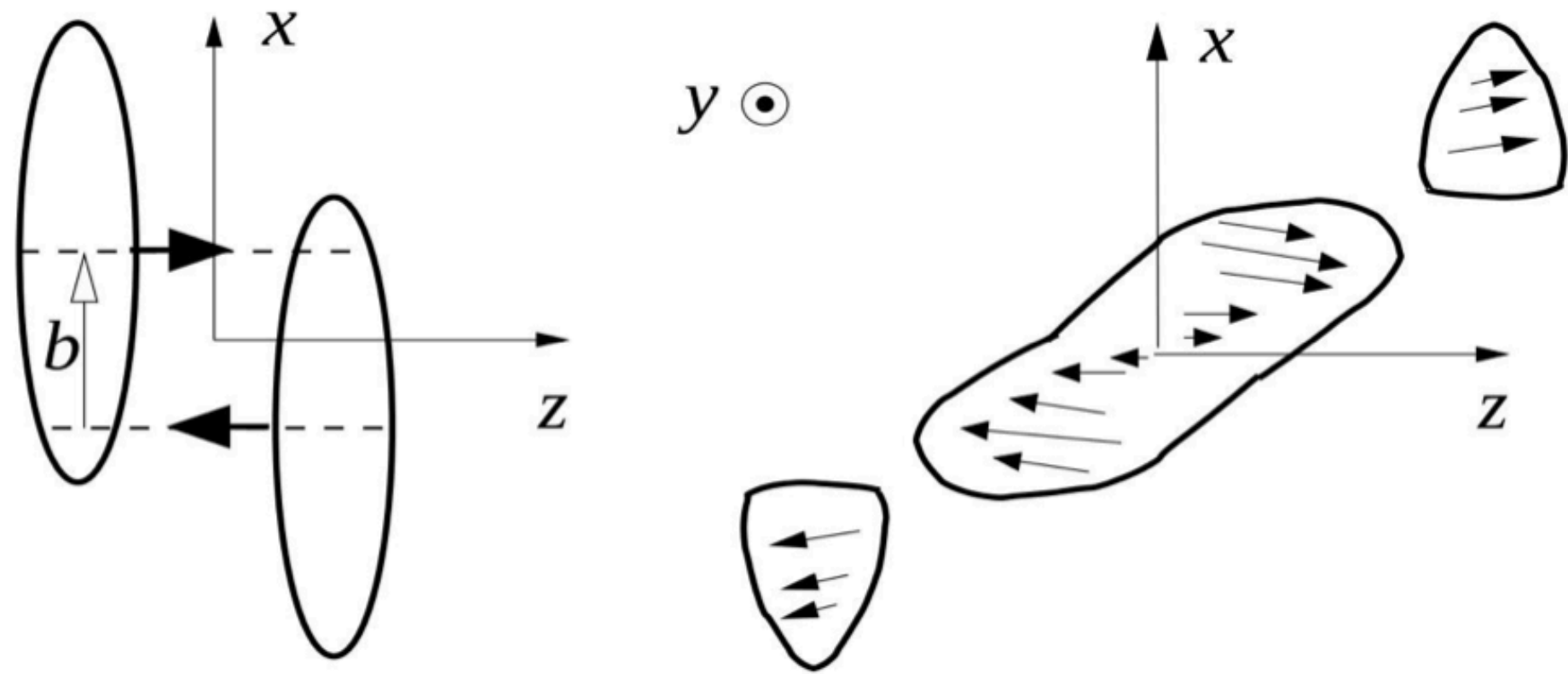


The freeze-out parameters show different trend compared to that of higher energy collisions: $T_{kin}(d) > T_{kin}(p)$,

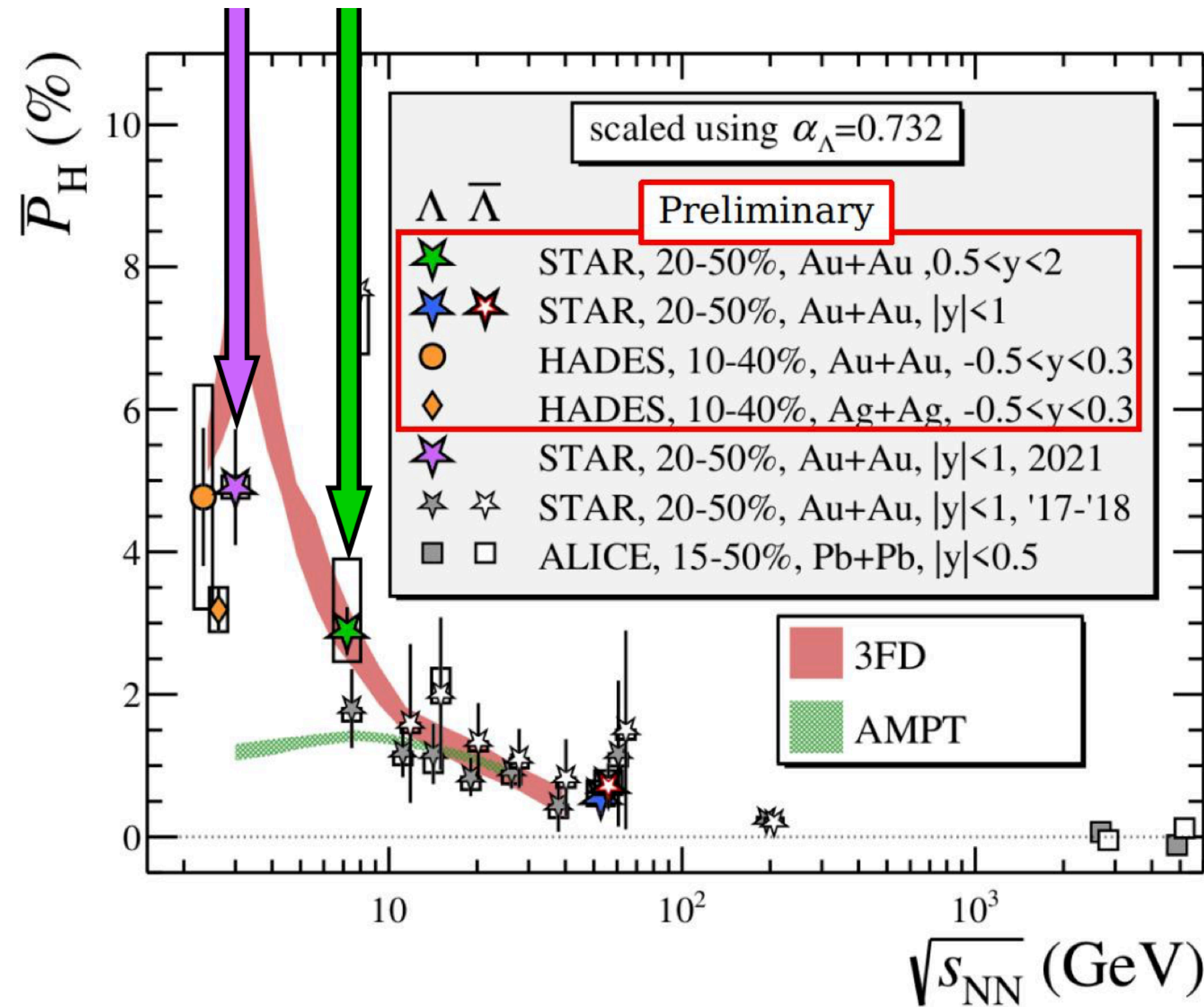
Hot and dense medium created at $\sqrt{s_{NN}} = 3$ GeV collisions seem different from that of high energy collisions

Hyperon polarization at BES

In non-central heavy-ion collisions large orbital momentum: $\vec{L} = \vec{r} \times \vec{p} \sim bA\sqrt{s_{NN}}$ can polarize quarks due to „spin-orbit” interaction leading to non-zero polarization of produced particles along with \vec{L}

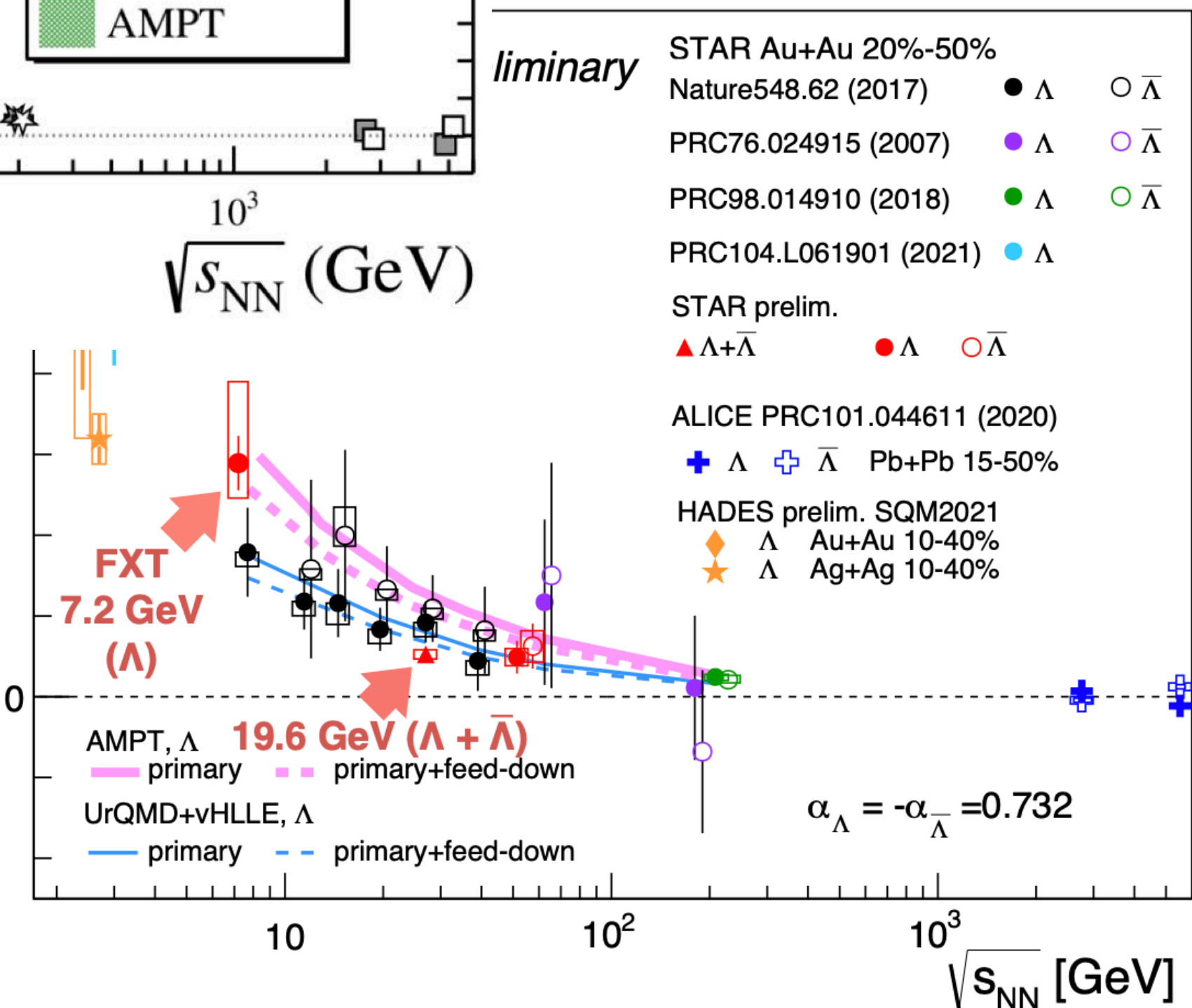


Z.-T. Liang and X.-N. Wang, Phys. Rev. Lett. 94, 102301 (2005),
Erratum:ibid. 96, 039901 (2006).



$\sqrt{s_{NN}} = 7.7$ GeV:
polarization consistent with 3FD

$\sqrt{s_{NN}} = 3.0$ GeV: trend of increasing polarization persist



Hypernuclei production

- Provides access to the hyperon–nucleon (Y-N) interaction, strangeness in high density nuclear matter: EOS of neutron star, hadronic phase of a heavy ion collision.
- Understanding hypernuclear structure may give more constraints on the Y-N interaction. The formation of loosely bound states in heavy-ion collisions are not well understood

Hypernuclei ${}^3_{\Lambda}\text{H}$, (${}^4_{\Lambda}\text{H}$) are reconstructed via
2 and 3 body decays ${}^3\text{He} + \pi$ and $p + d + \pi$, (${}^4\text{He} + \pi$)

