

The 8th Asian Triangle Heavy-Ion Conference

# ATHIC2021

5-9 November 2021

Inha University, Incheon, South Korea

## Search for CME with STAR experiment

Fuqiang Wang (Purdue University)

For the STAR Collaboration

Supported in part by



# OUTLINE

- Physics motivation and observables
- Brief historical review of STAR (and other) measurements
- Recent CME measurements from STAR
  - Invariant mass
  - EPD measurements
  - Other observables/approaches
  - **Spectator/participant planes**
  - **Isobar collisions**
- Summary and outlook

# CHIRAL MAGNETIC EFFECT (CME)

## The strong interaction

$$\mathcal{L}_{QCD} = \sum_q \left( \bar{\psi}_{qi} i\gamma^\mu \left[ \delta_{ij} \partial_\mu + ig \left( G_\mu^\alpha t_\alpha \right)_{ij} \right] \psi_{qj} - m_q \bar{\psi}_{qi} \psi_{qi} \right) - \frac{1}{4} G_{\mu\nu}^\alpha G_\alpha^{\mu\nu}$$

quarks
quark-gluon interactions
quarks
gluons

't Hooft vacuum

$$+ \theta \frac{\alpha_s}{8\pi} G_{\mu\nu}^\alpha \tilde{G}_\alpha^{\mu\nu}$$

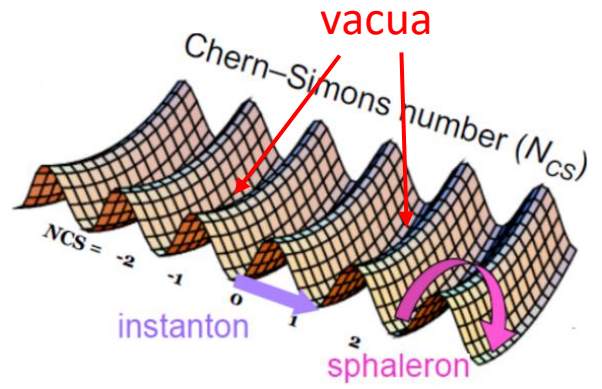
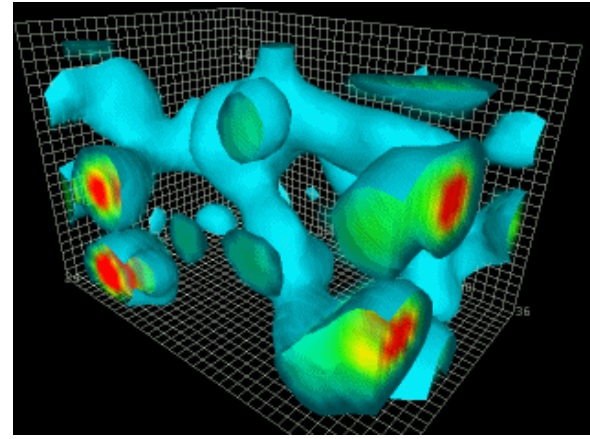
$$= -\theta \frac{\alpha_s}{2\pi} \vec{E}_\alpha \cdot \vec{B}_\alpha$$

to solve the  $U(1)_A$  problem (1976)

E: C-odd, P-odd, T-even  
B: C-odd, P-even, T-odd

Explicitly breaks CP

Early universe ultraviolet  $\theta \approx 1$  ??  $\gg$  current infrared  $\theta \approx 0$



Kharzeev, Pisarski, Tytgat, PRL81(1998)512

QCD vacuum fluctuation, chiral anomaly, topological gluon field

Reaction plane ( $\Psi_R$ )

$\vec{B}$

$B \sim 10^{15} \text{ T}$

X (defines  $\Psi_R$ )

Kharzeev, et al. NPA 803 (2008) 227

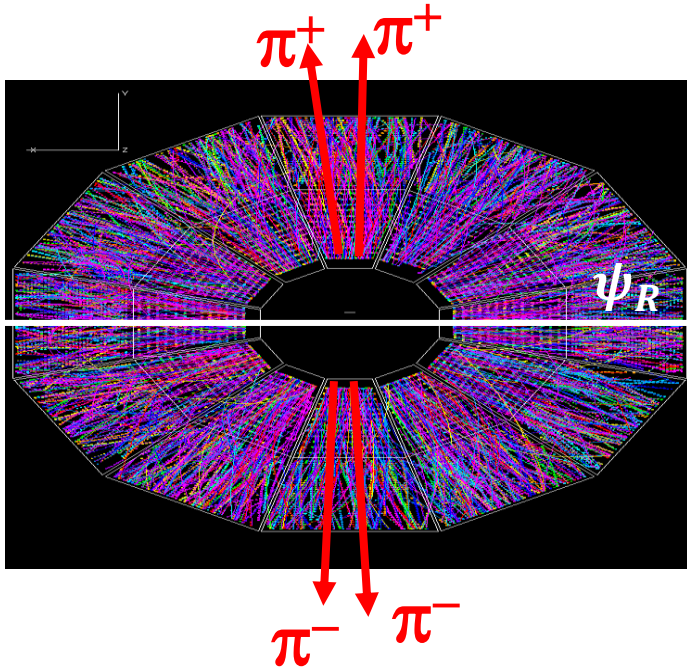
1      2      3

Chiral Magnetic Effect (CME)

Discovery of the CME would imply: Chiral symmetry restoration (current-quark DOF & deconfinement);  
Local P/CP violation that may solve the strong CP problem (matter-antimatter asymmetry)

# THE COMMON $\gamma$ VARIABLE

Voloshin, PRC 70 (2004) 057901

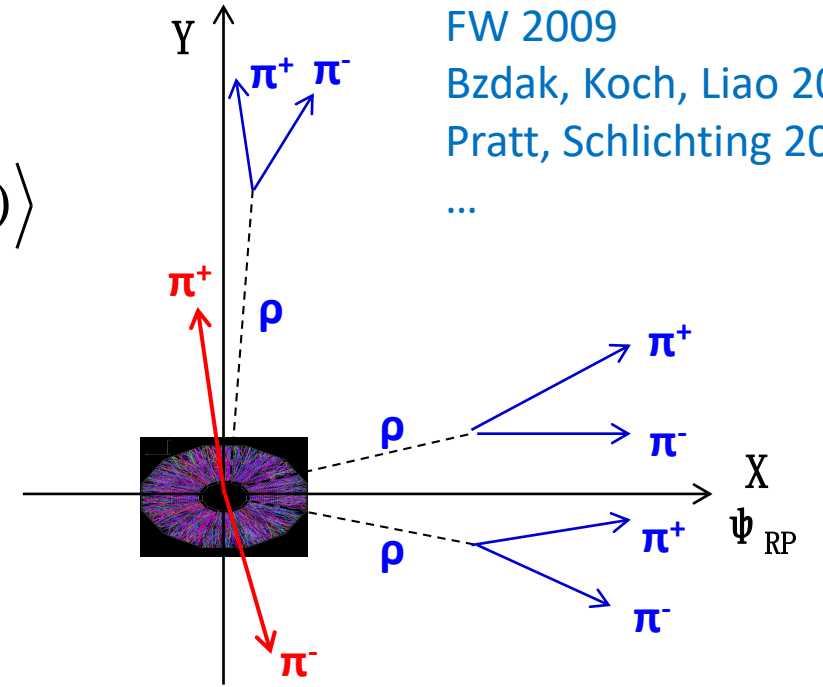


$$\gamma_{\alpha\beta} = \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_{RP}) \rangle$$

$$\gamma_{+,-,+} > 0, \quad \gamma_{+,-,-} < 0$$

$$\Delta\gamma = \gamma_{OS} - \gamma_{SS}$$

$$\Delta\gamma > 0$$



Voloshin 2004  
FW 2009  
Bzdak, Koch, Liao 2010  
Pratt, Schlichting 2010  
...

$$\gamma_{\alpha\beta} = \left[ \langle \cos(\varphi_\alpha - \psi_{RP}) \cos(\varphi_\beta - \psi_{RP}) \rangle - \langle \sin(\varphi_\alpha - \psi_{RP}) \sin(\varphi_\beta - \psi_{RP}) \rangle \right] + \left[ \frac{N_{\text{cluster}}}{N_\alpha N_\beta} \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_{\text{cluster}}) \cos(2\varphi_{\text{cluster}} - 2\varphi_{RP}) \rangle \right]$$

$$= \left[ \langle v_{1,\alpha} v_{1,\beta} \rangle - \langle a_\alpha a_\beta \rangle \right] + [\text{charge-independent Bkg (e.g. mom. conservation)}] + \frac{N_{\text{cluster}}}{N_\alpha N_\beta} \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_{\text{cluster}}) \rangle v_{2,\text{cluster}}$$

$$\Delta\gamma = 2 \langle a_1^2 \rangle + \frac{N_{\text{cluster}}}{N_\alpha N_\beta} \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_{\text{cluster}}) \rangle v_{2,\text{cluster}}$$

# THE R VARIABLE

Ajitanand et al., PRC 83 (2011) 011901  
Magdy et al., PRC 97 (2018) 061901(R)

$$\Delta S = \frac{\sum_1^p \sin\left(\frac{m}{2} \Delta\varphi_m\right)}{p} - \frac{\sum_1^n \sin\left(\frac{m}{2} \Delta\varphi_m\right)}{n}$$

$$R(\Delta S_m) \equiv \frac{N(\Delta S_{m,\text{real}})}{N(\Delta S_{m,\text{shuffled}})} / \frac{N(\Delta S_{m,\text{real}}^\perp)}{N(\Delta S_{m,\text{shuffled}}^\perp)}, \quad m = 2, 3, \dots,$$

Yufu Lin's talk this afternoon

Choudhury et al. arXiv:2105.06044 [nucl-ex],  
CPC in print.

Width of  $R(\Delta S)$  distribution reduces to variance  
 $\sin^* \sin, \cos^* \cos \rightarrow$  equivalently the  $\Delta\gamma$  variable

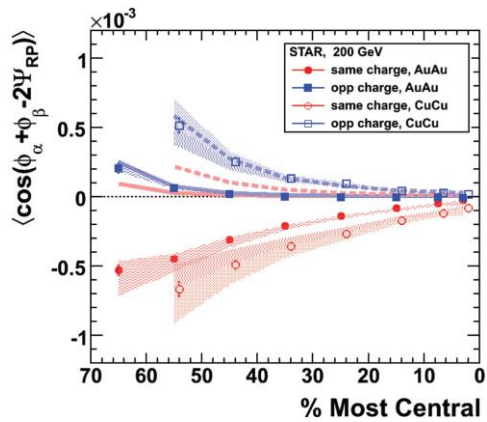
$$\frac{S_{\text{concavity}}}{\sigma_{R2}^2} \approx -\frac{M}{2}(M-1)\Delta\gamma_{112}$$

$$\frac{S_{\text{concavity}}}{\sigma_{R2'}^2} = \frac{S_{\text{concavity}}}{\sigma_{R2}^2} \langle (\Delta S_{2,\text{shuffled}})^2 \rangle \approx -\frac{M}{2}(M-1)\Delta\gamma_{112} \times \frac{2}{M} \approx -M\Delta\gamma_{112}$$

- Established analytical relationship between  $\Delta\gamma$  and  $R_{\Psi_2}(\Delta S)$
- “Equivalence” verified by MC simulations and the EBE-AVFD model
- $\Delta\gamma$  and  $R_{\Psi_2}(\Delta S)$  have similar sensitivities to CME signal and background

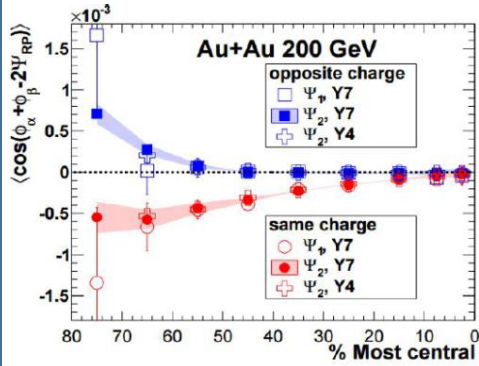
# STAR (and ALICE, CMS) MEASUREMENTS

STAR, PRL 103 (2019) 251601;  
PRC 81, 054908 (2010)



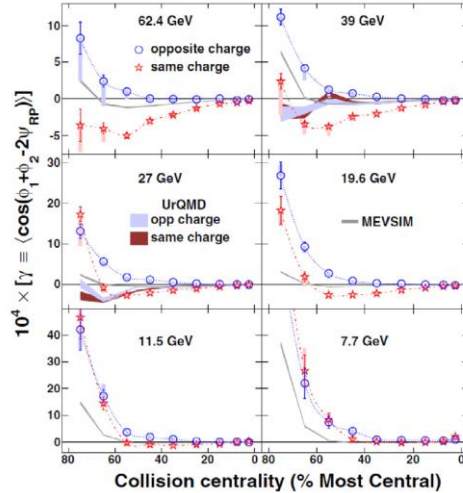
First measurement;  
Large signal

STAR, PRC 88 (2013) 064911



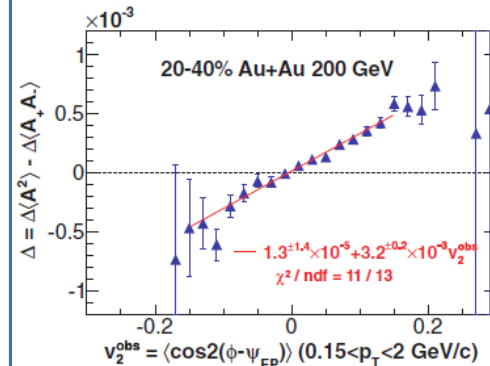
Measurement wrt ZDC  $\Psi_1$ ;  
Similar result wrt TPC  $\Psi_2$

STAR, PRL 113 (2014) 052302



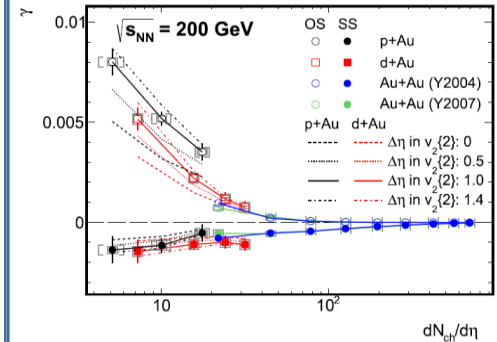
BES; signal disappears  
at low energy

STAR, PRC 89 (2014) 044908



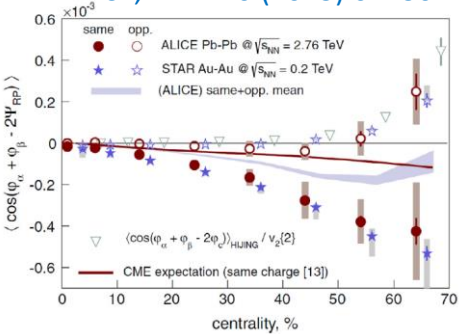
ESE projection to  $v_2=0$ ;  
bkg significantly reduced,  
but not eliminated

STAR, PLB 798 (2019) 134975



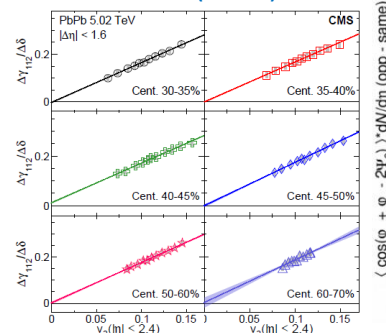
Small system; signal as  
large as heavy ion; large  
bkg contributions

ALICE, PRL 110 (2013) 012301

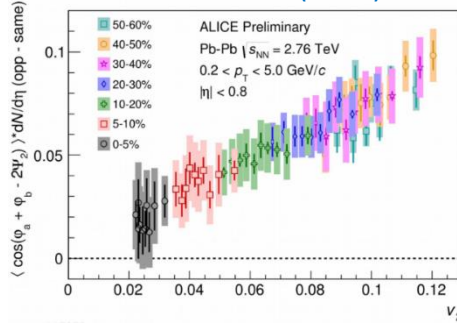


Fuqiang Wang

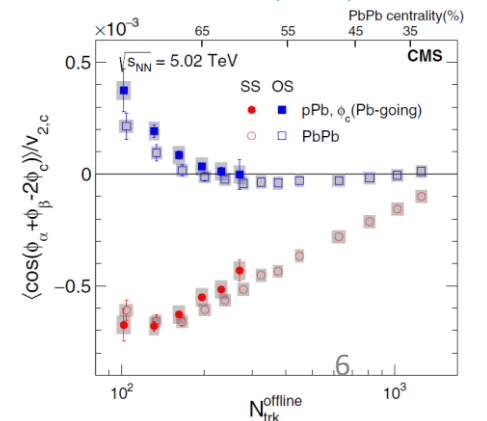
CMS PRC97 (2018) 044912



ALICE PLB777(2018)151

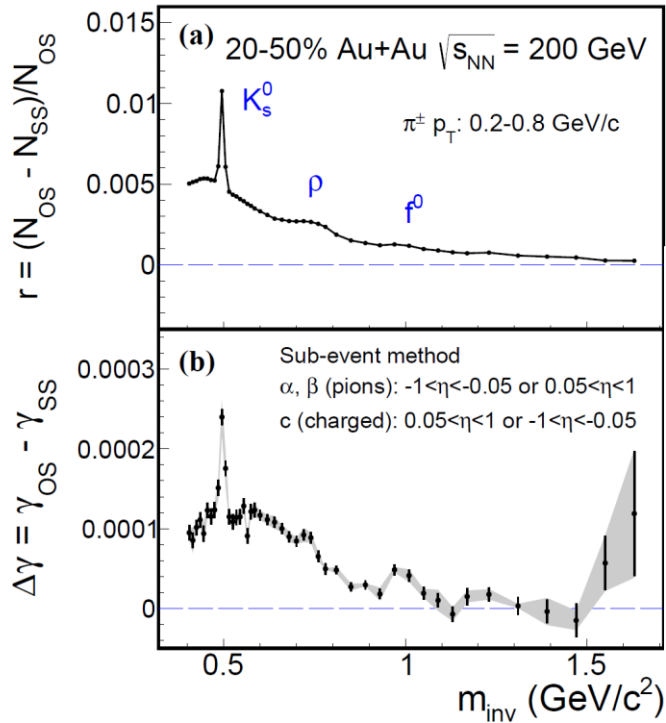


CMS, PRL 118 (2017) 122301

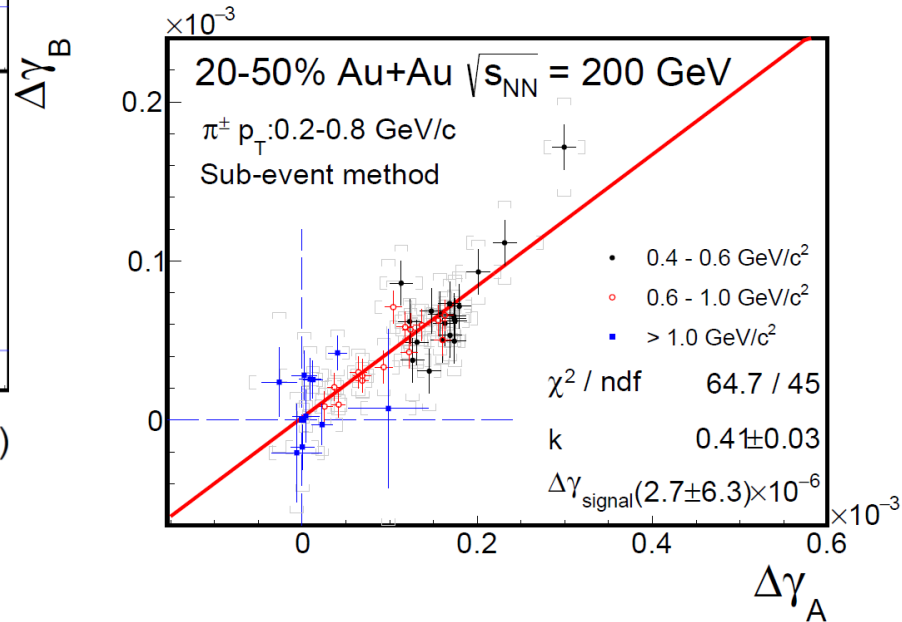


# MEASUREMENT IN INVARIANT MASS

Jie Zhao, Hanlin Li, FW, Eur.Phys.J.C 79 (2019) 168  
 STAR, arXiv:2006.05035



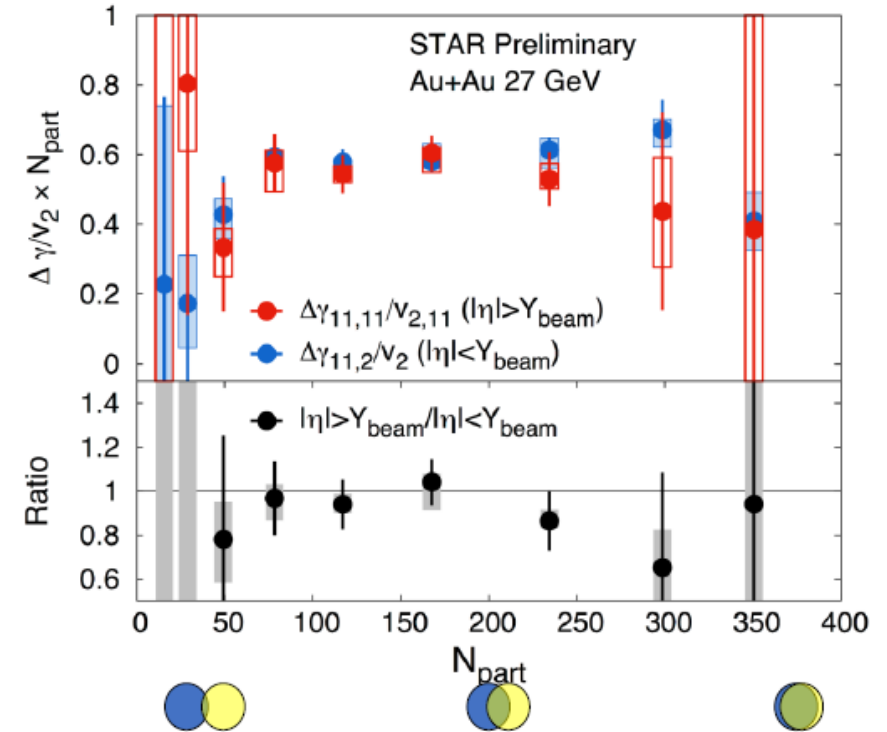
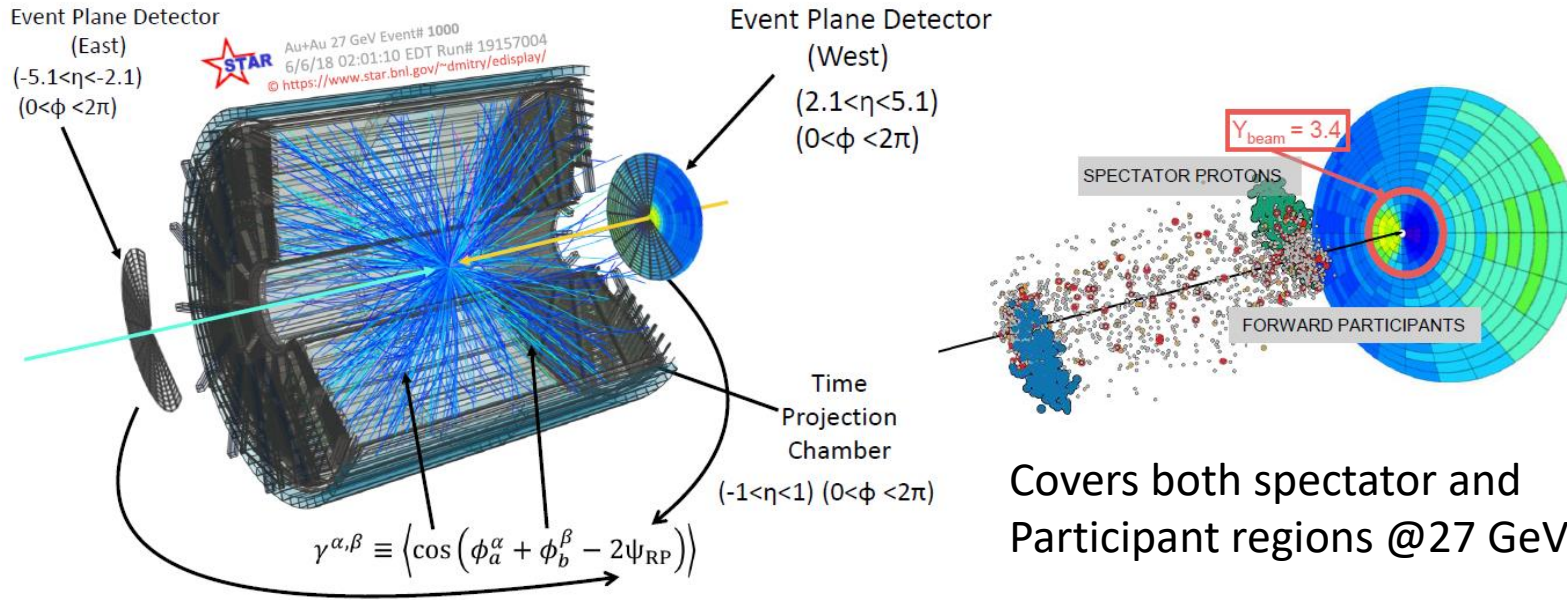
$$\frac{N_\rho}{N_\alpha N_\beta} \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_{clus}) \rangle \times v_{2,clus}$$



- Explicit demonstration of “resonance” background
- Exploit “ESE” to extract CME, assuming CME is mass independent
- Upper limit 15% at 95% CL

# MORE RECENT LOW ENERGY (27 GeV) DATA

Yu Hu (STAR), arXiv:2110.15937, SQM 2021



- Higher statistics, new detector (EPD)
- New approach: inner EPD -> first-order harmonic plane; Outer EPD -> second-order harmonic plane.
- Current data consistent with background contributions

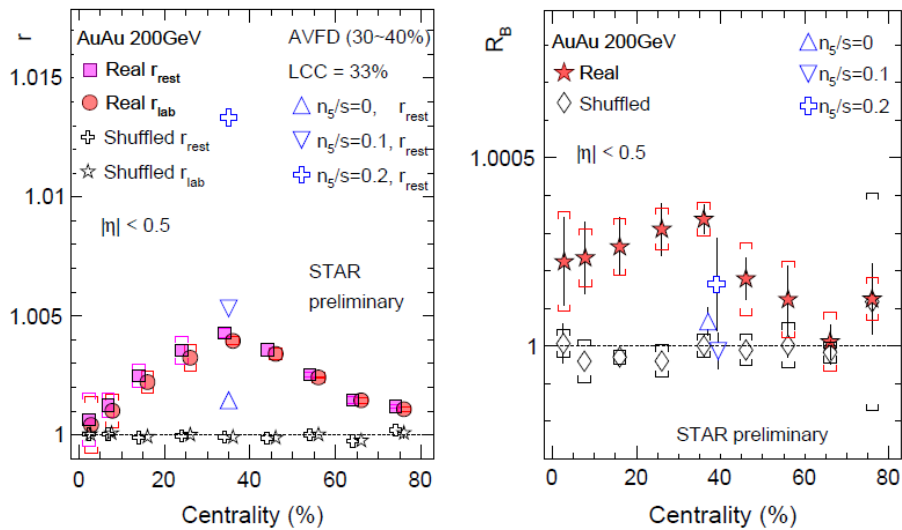


# NEW OBSERVABLES/APPROACHES

## Signed balance function (SPF)

Tang, CPC 44 (2020) 054101

Yufu Lin (STAR), NPA 1005 (2021) 121828, QM 2019



Yufu Lin's talk this afternoon

- $r$  is out-of-plane to in-plane ratio of the SPF momentum-ordering difference
- Both  $r_{rest}$  and  $R_B = r_{rest}/r_{lab}$  are larger than unity, above model calculations without CME.

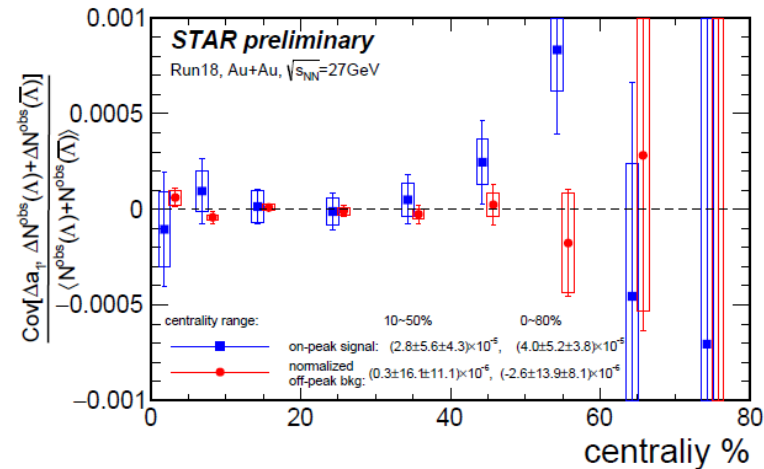
## CME-helicity correlation

Du, Finch, Sandweiss, PRC 78 (2008) 044908

044908

Finch, Murray, PRC 96 (2017) 044911

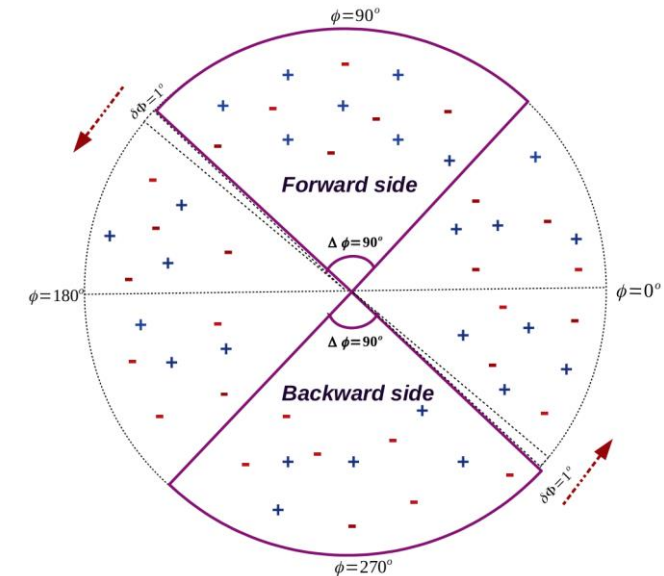
Yicheng Feng (STAR), DNP 2020



- Positive correlation btw CME  $\Delta a_1$  and  $\Lambda$  net-helicity from chirality anomaly
- Current signal consistent with zero within uncertainties

## Sliding Dumbbell

Jagbir Singh (STAR) QM 2019

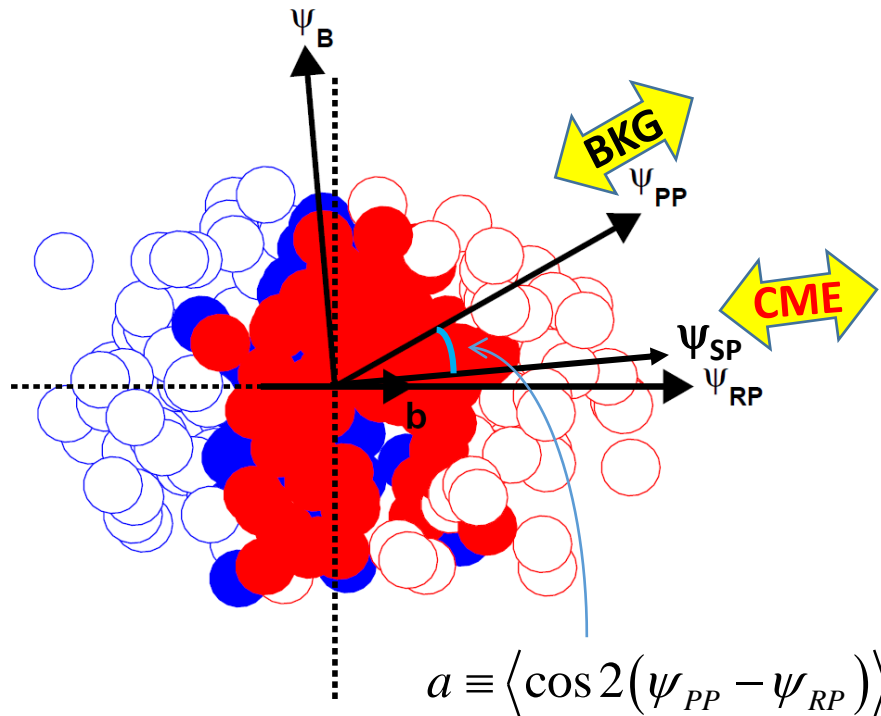


- Select CME enriched sample
- Perform  $\Delta\gamma$  measurement with background subtraction in separate event classes

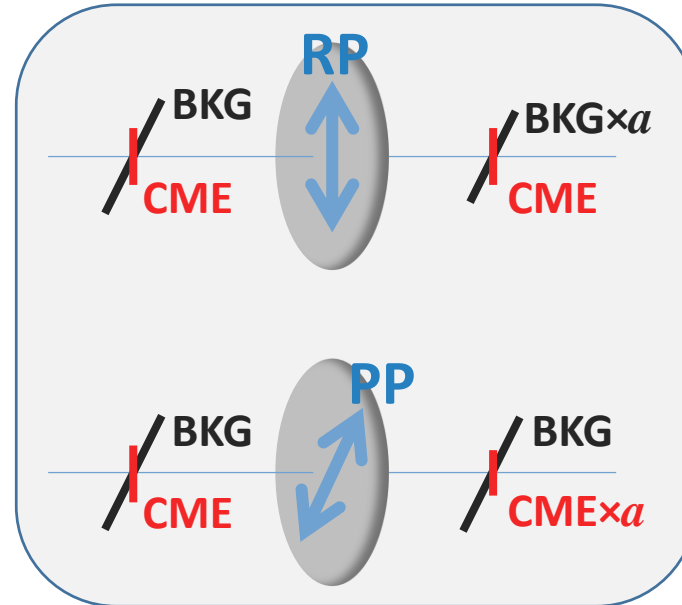
# W.R.T. SPECTATOR & PARTICIPANT PLANES, 2021

Haojie Xu et al., CPC 42 (2018) 084103, arXiv:1710.07265

S.A. Voloshin, PRC 98 (2018) 054911, arXiv:1805.05300



## INTRA-EVENT "CME- $v_2$ FILTER"



IN THE SAME EVENT

$$\Delta\gamma_{\{SP\}} = a\Delta\gamma_{Bkg\{PP\}} + \Delta\gamma_{CME\{PP\}} / a$$

$$\Delta\gamma_{\{PP\}} = \Delta\gamma_{Bkg\{PP\}} + \Delta\gamma_{CME\{PP\}}$$

$$A = \Delta\gamma_{\{SP\}} / \Delta\gamma_{\{PP\}}$$

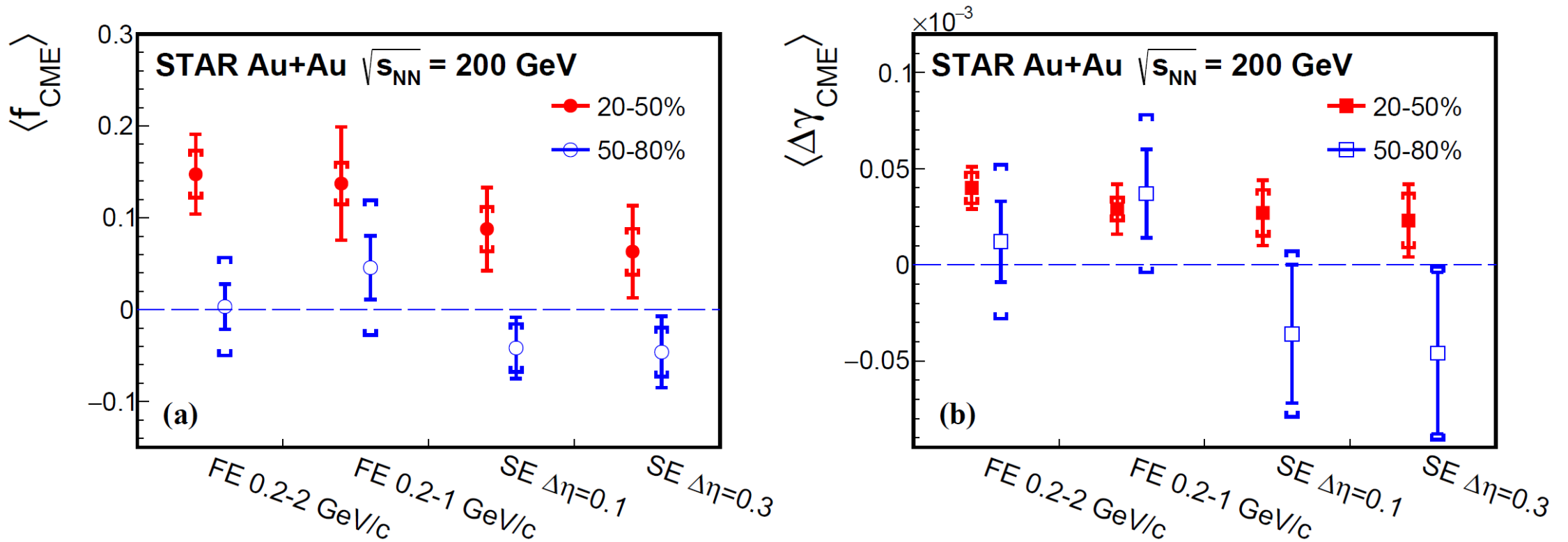
$$a = v_2\{SP\} / v_2\{PP\}$$

$$\Delta\gamma_{\{SP\}} / a - \Delta\gamma_{\{PP\}} = (1/a^2 - 1)\Delta\gamma_{CME\{PP\}}$$

$$f_{CME} = \frac{\Delta\gamma_{CME\{PP\}}}{\Delta\gamma_{\{PP\}}} = \frac{A/a - 1}{1/a^2 - 1}$$

# Au+Au Collisions at 200 GeV (2.4B MB)

STAR, arXiv:2106.09243



- Consistent-with-zero signal in peripheral 50-80% collisions with relatively large errors
- Indications of finite signal in mid-central 20-50% collisions, with 1-3 $\sigma$  significance
- Possible remaining nonflow effects

# REMAINING NONFLOW EFFECTS

Feng et al., arXiv:2106.15595

$$f_{\text{CME}} = \frac{\Delta\gamma_{\text{CME}}\{\text{PP}\}}{\Delta\gamma\{\text{PP}\}} = \frac{A/a - 1}{1/a^2 - 1}$$

$$\frac{A}{a} = \frac{\Delta\gamma\{\text{SP}\} / v_2\{\text{SP}\}}{\Delta\gamma\{\text{PP}\}^* / v_2\{\text{PP}\}^*} = \frac{C_3\{\text{SP}\} / v_2^2\{\text{SP}\}}{C_3\{\text{PP}\}^* / v_2^2\{\text{PP}\}^*} = \frac{1 + \epsilon_{\text{nf}}}{1 + \frac{\epsilon_3 / \epsilon_2}{Nv_2^2\{\text{PP}\}}}$$

$$C_3\{\text{SP}\} = \frac{C_{2\text{p}}N_{2\text{p}}}{N^2} v_{2,2\text{p}}\{\text{SP}\}v_2\{\text{SP}\},$$

Nonflow in  $\Delta\gamma$   
→ negative  $f_{\text{CME}}$

$$C_3^*\{\text{EP}\} = \frac{C_{2\text{p}}N_{2\text{p}}}{N^2} v_{2,2\text{p}}\{\text{EP}\}v_2\{\text{EP}\} + \frac{C_{3\text{p}}N_{3\text{p}}}{2N^3}.$$

$$\epsilon_2 \equiv \frac{C_{2\text{p}}N_{2\text{p}}v_{2,2\text{p}}}{Nv_2}$$

$$\epsilon_3 \equiv \frac{C_{3\text{p}}N_{3\text{p}}}{2N}$$

$$\Delta\gamma_{\text{bkgd}} = \frac{N_{2\text{p}}}{N^2} \langle \cos(\phi_\alpha + \phi_\beta - 2\phi_{2\text{p}}) \rangle v_{2,2\text{p}}$$

$$C_{2\text{p}} = \langle \cos(\phi_\alpha + \phi_\beta - 2\phi_{2\text{p}}) \rangle$$

$$C_{3\text{p}} = \langle \cos(\phi_\alpha + \phi_\beta - 2\phi_c) \rangle_{3\text{p}}$$

$$v_2^*\{\text{EP}\} = \sqrt{v_2^2\{\text{EP}\} + v_{2,\text{nf}}^2}$$

$$\epsilon_{\text{nf}} \equiv v_{2,\text{nf}}^2 / v_2^2$$

Nonflow in  $v_2$   
→ positive  $f_{\text{CME}}$

$$f_{\text{CME}}^* \approx \left( \epsilon_{\text{nf}} - \frac{\epsilon_3 / \epsilon_2}{Nv_2^2\{\text{EP}\}} \right) / \left( \frac{1 + \epsilon_{\text{nf}}}{a^2} - 1 \right)$$

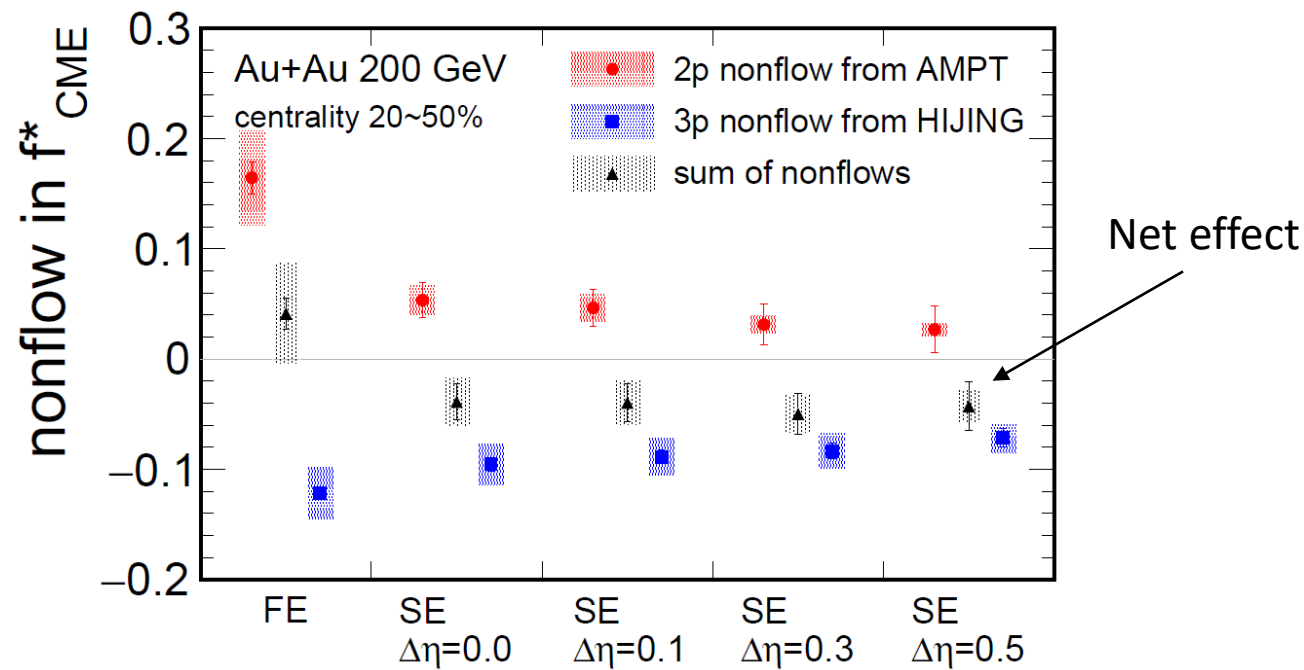
$$f_{\text{CME}}^* = \left( \frac{1 + \epsilon_{\text{nf}}}{1 + \frac{\epsilon_3 / \epsilon_2}{Nv_2^2\{\text{EP}\}}} - 1 \right) / \left( \frac{1 + \epsilon_{\text{nf}}}{a^2} - 1 \right)$$

$$= \left( \frac{1 + \epsilon_{\text{nf}}}{1 + \frac{(1 + \epsilon_{\text{nf}})\epsilon_3 / \epsilon_2}{Nv_2^{*2}\{\text{EP}\}}} - 1 \right) / \left( \frac{1}{a^{*2}} - 1 \right)$$

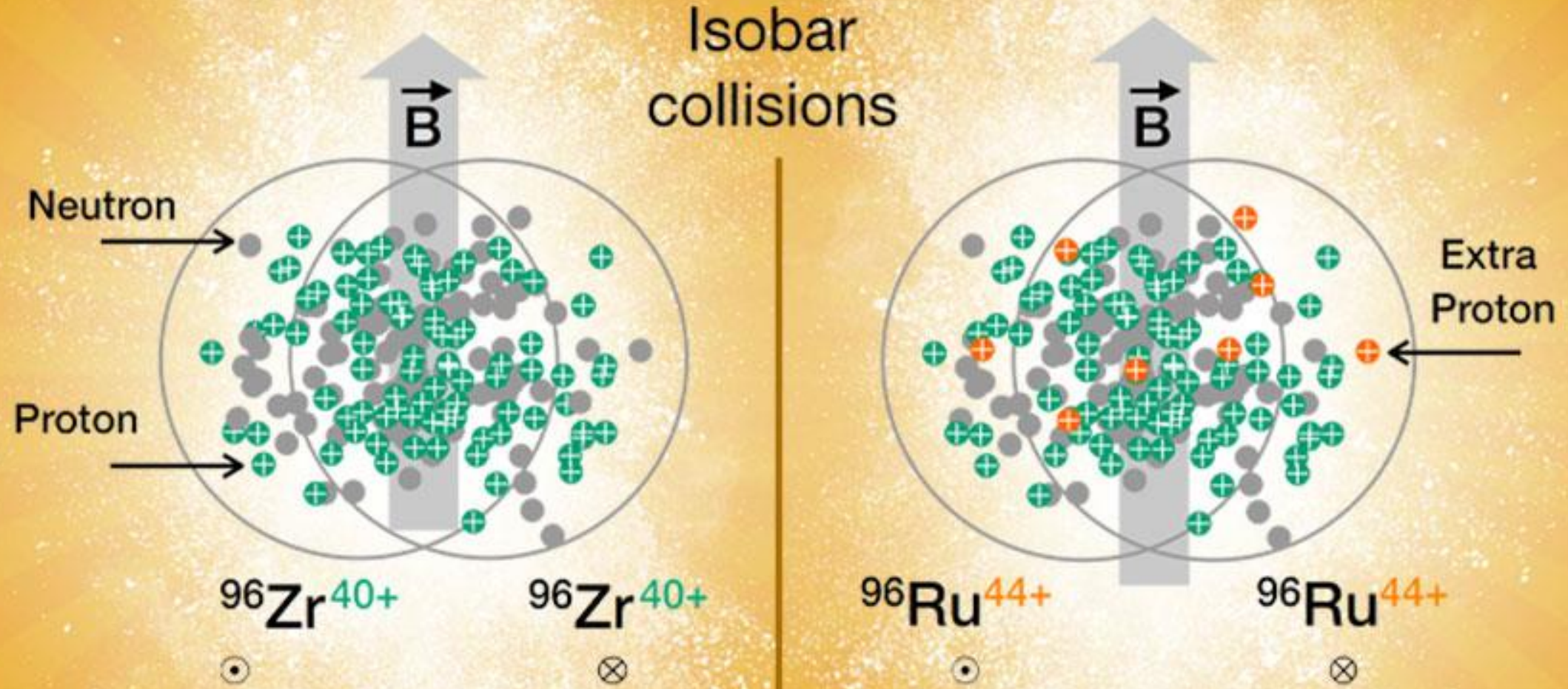
# MODEL ESTIMATES OF NONFLOW

Feng et al., arXiv:2106.15595

$$f_{\text{CME}}^* \approx \left( \epsilon_{\text{nf}} - \frac{\epsilon_3/\epsilon_2}{Nv_2^2\{\text{EP}\}} \right) / \left( \frac{1 + \epsilon_{\text{nf}}}{a^2} - 1 \right)$$

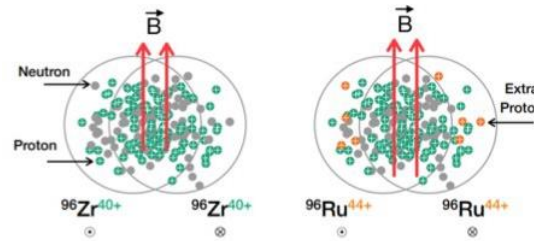


- 2-particle nonflow estimates from AMPT
- 3-particle nonflow estimates from HIJING
- Net effect on  $f_{\text{CME}}^*$  can possibly be negative (model dependent)
- Further, additional model studies



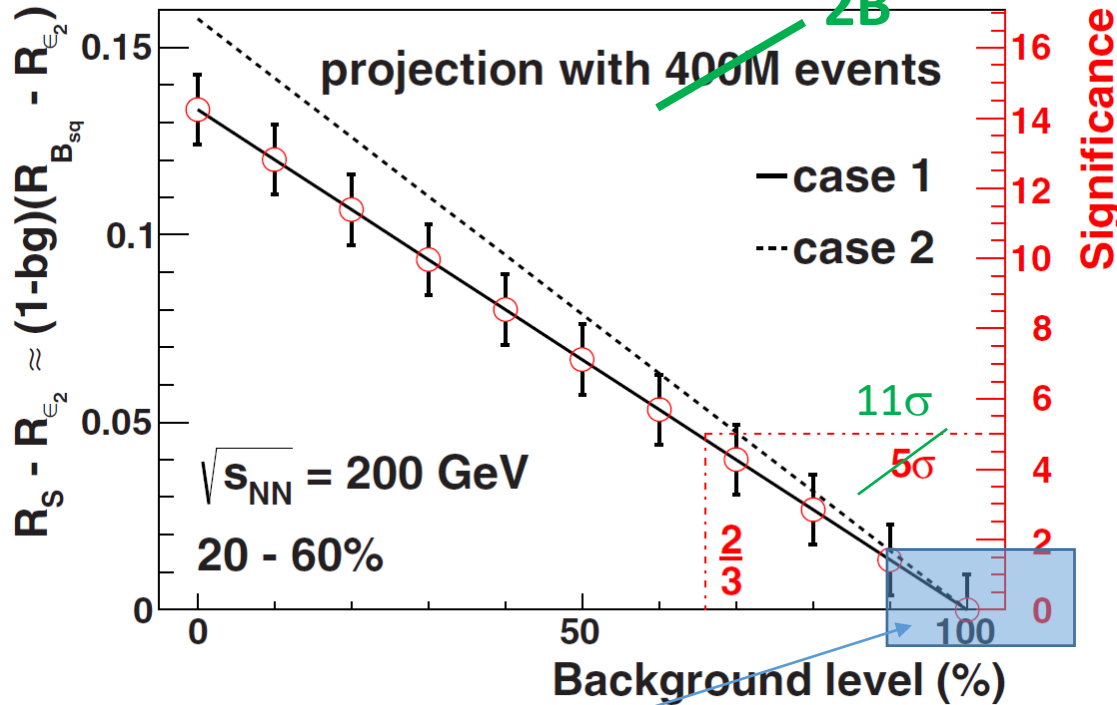
# ISOBAR COLLISIONS

Voloshin, PRL 105 (2010) 172301



Same A  $\rightarrow$  same background  
Different Z  $\rightarrow$  different signal

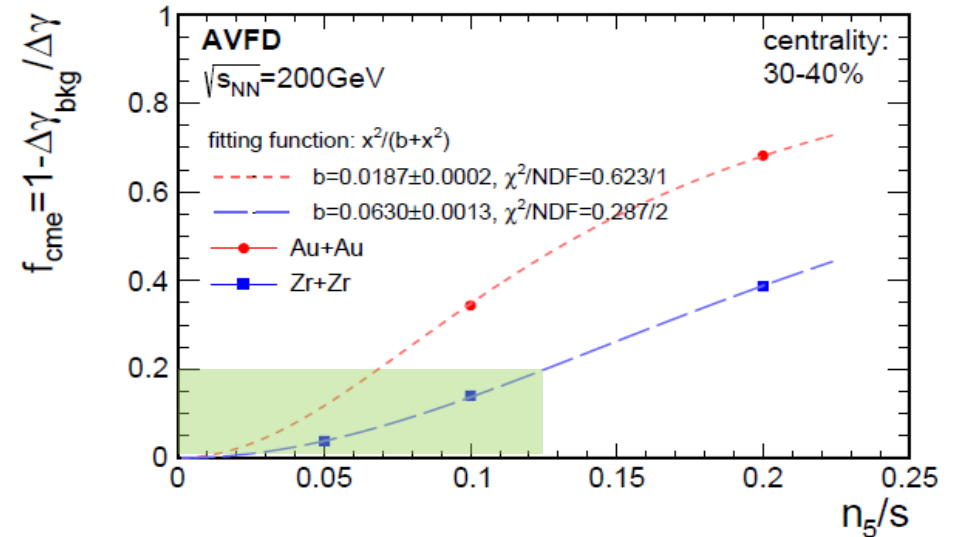
Deng et al. PRC 94 (2016) 041901(R)



3.3 $\sigma$  effect if isobar  $\approx$  AuAu ( $f_{cme} = 10\%$ )

$$Ru/Zr = 1 + 15\% * 10\% = 1.015 \pm 0.004$$

Yicheng Feng, Yufu Lin, Jie Zhao, FW, arXiv:2103.10378



Background  $\propto 1/N \rightarrow$  isobar/AuAu  $\sim 2$

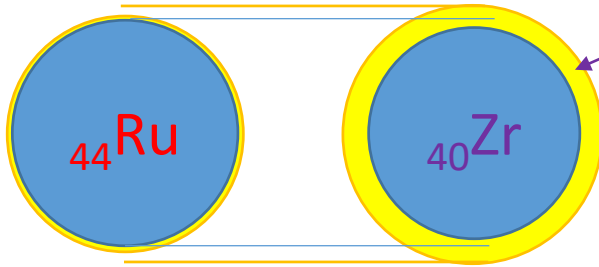
Mag. field  $B \sim A^{1/3} \rightarrow$  Signal: AuAu/isobar  $\sim 1.5$

Could be  $\times 3$  reduction in  $f_{CME}$  at the same  $n_5/s$

If AuAu  $f_{CME} = 10\%$ , then isobar 3% (1 $\sigma$  effect)

Caveats: Axial charge densities and sphaleron transition probabilities could be different between Au+Au and isobar, e.g. AVFD-glasma  $\mu_5/s$ : isobar/AuAu  $\sim 1.5$

# ISOBAR SYSTEMS ARE NOT IDENTICAL: MULTIPLICITY



Halo-type

Larger charge radius  
Little neutron skin

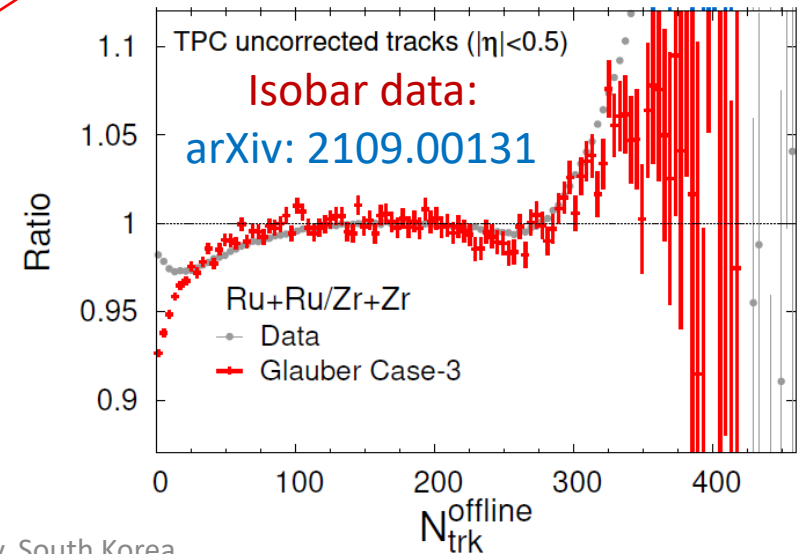
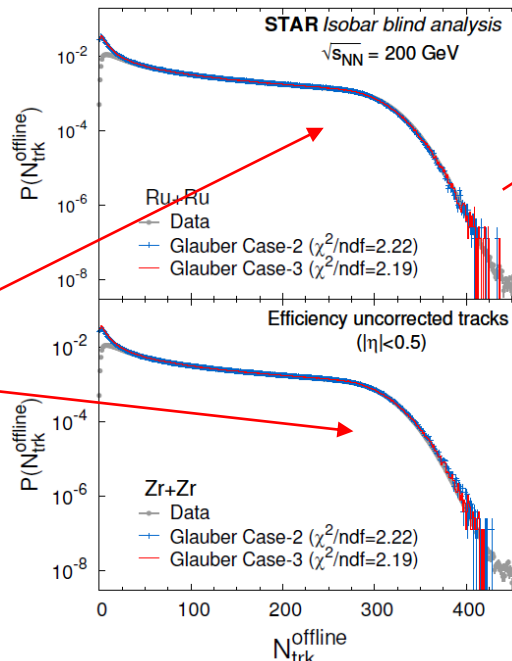
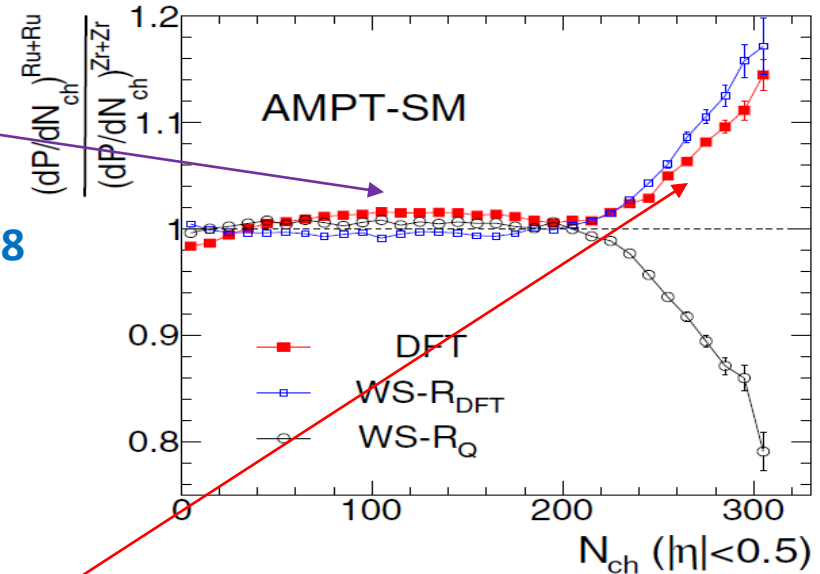
Smaller charge radius  
Much thicker neutron skin

Overall size: Ru < Zr

→ Larger energy density in Ru > Zr

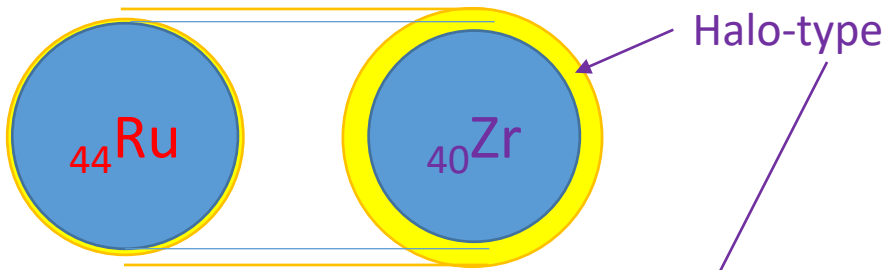
→ Larger multiplicity in Ru > Zr

Predicted by DFT:  
Hanlin Li et al. PRC 98  
(2018) 054907





# ISOBAR SYSTEMS ARE NOT IDENTICAL: $V_2$

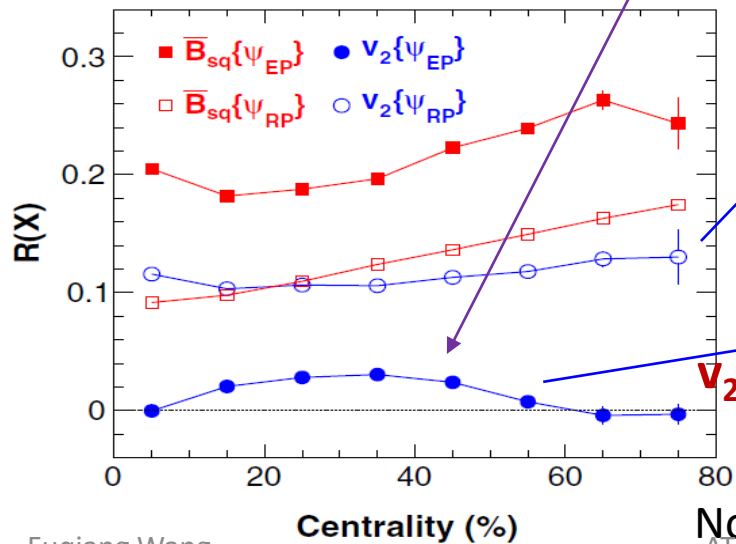


Larger charge radius  
Little neutron skin

Smaller charge radius  
Much thicker neutron skin

Redicted by DFT:

Haojie Xu et al. PRL **121** (2018) 022301

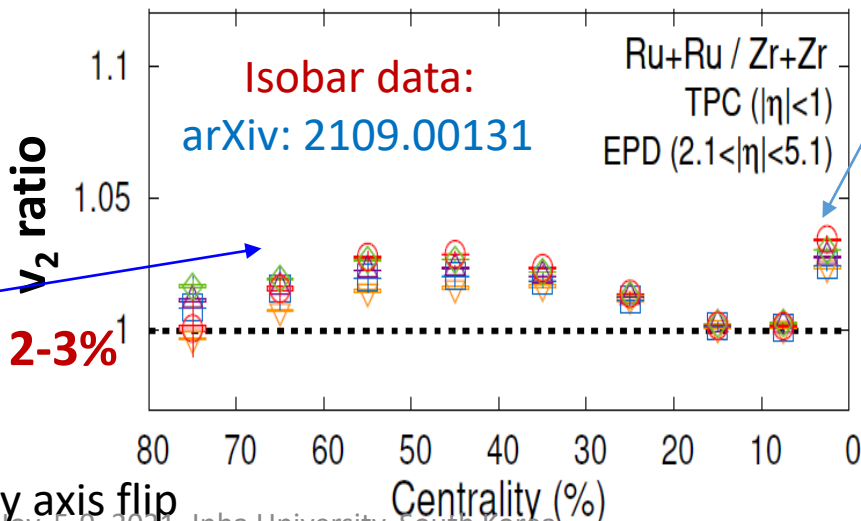
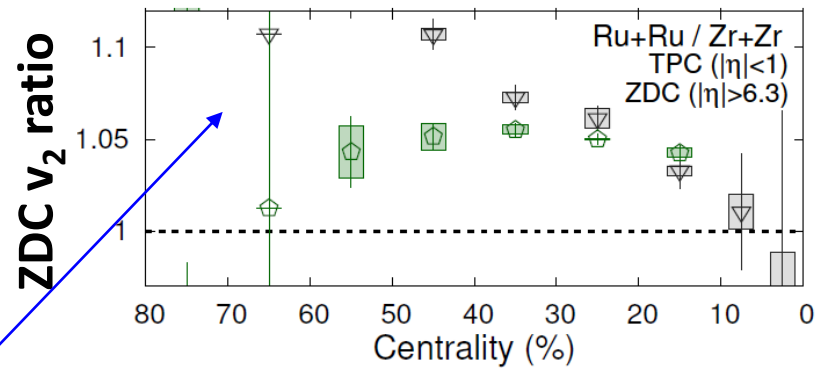


$v_2$  differs by 2-3%

Note centrality axis flip

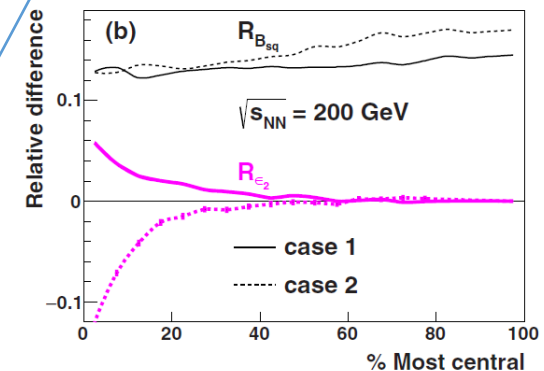
$$\Delta\gamma_{\text{bkgd}} = \frac{4N_{2p}}{N^2} \langle \cos(\phi_\alpha + \phi_\beta - 2\phi_{2p}) \rangle v_{2,2p}$$

Normalize by  $v_2$  and  $N \rightarrow N\Delta\gamma/v_2$



J. Jia, C. Zhang (Mon.)

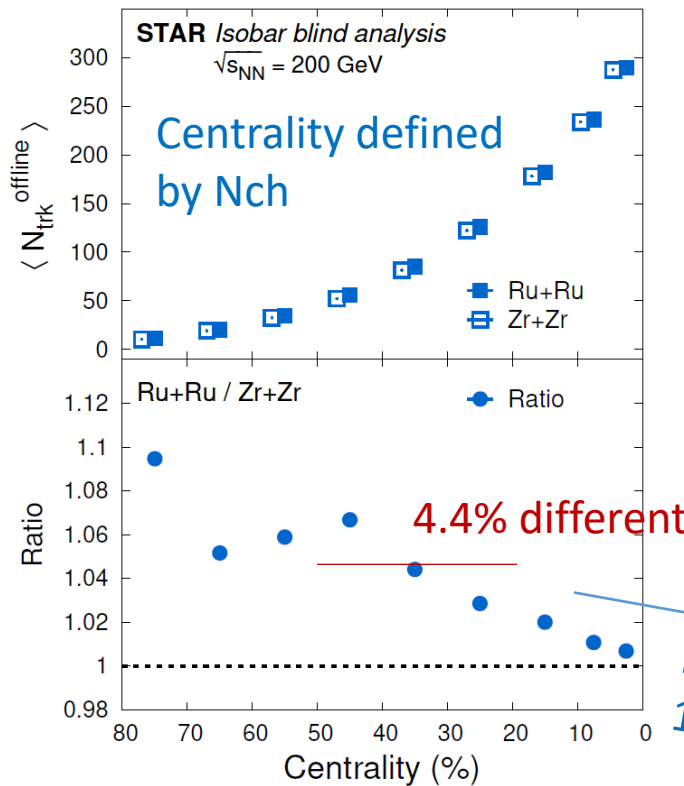
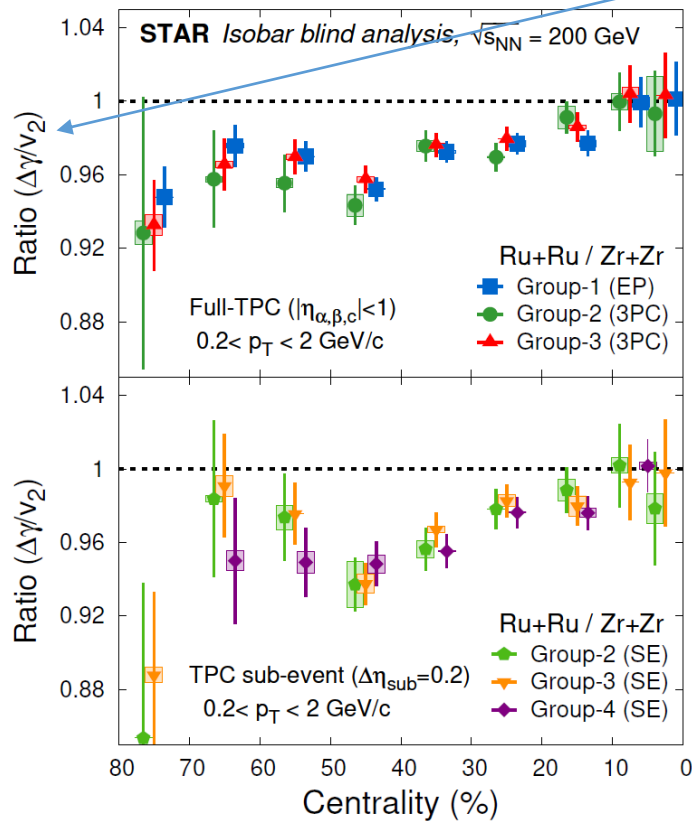
Nuclear deformity:  
Deng et al. PRC 94  
(2016) 041901(R)



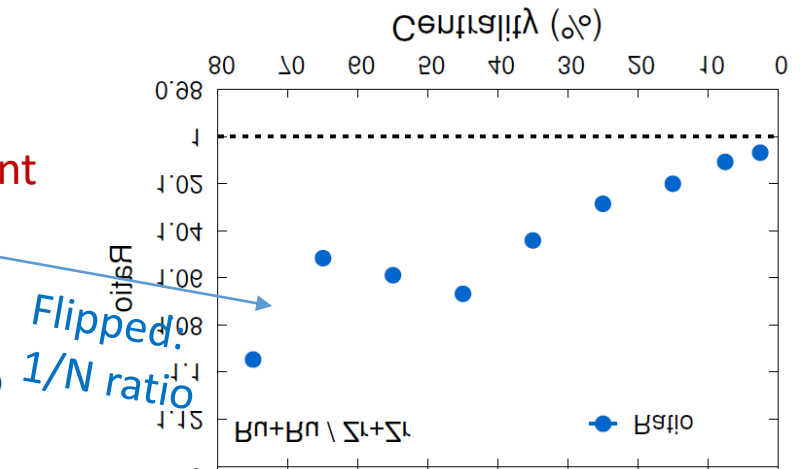
# $\Delta\gamma/v_2$ RESULTS FROM MULTIPLE GROUPS

$$\Delta\gamma_{\text{bkgd}} = \frac{4N_{2p}}{N^2} \langle \cos(\phi_\alpha + \phi_\beta - 2\phi_{2p}) \rangle v_{2,2p}$$

Under the assumption of flowing clusters, scales with overall multiplicity, then  $\Delta\gamma$  is diluted by  $1/N$

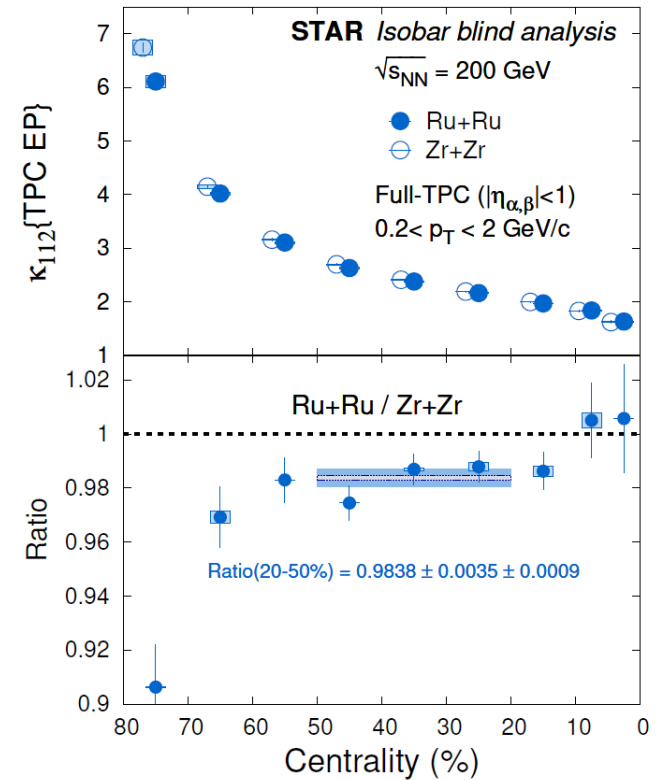
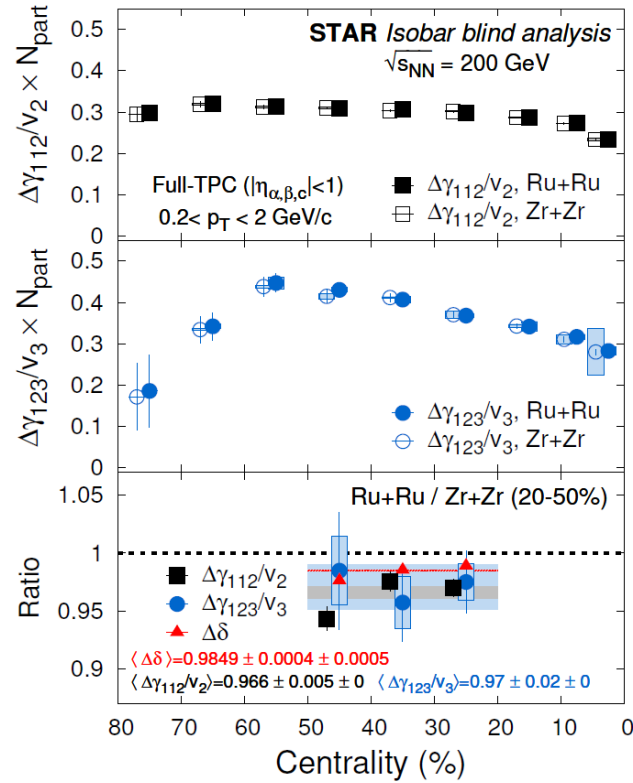
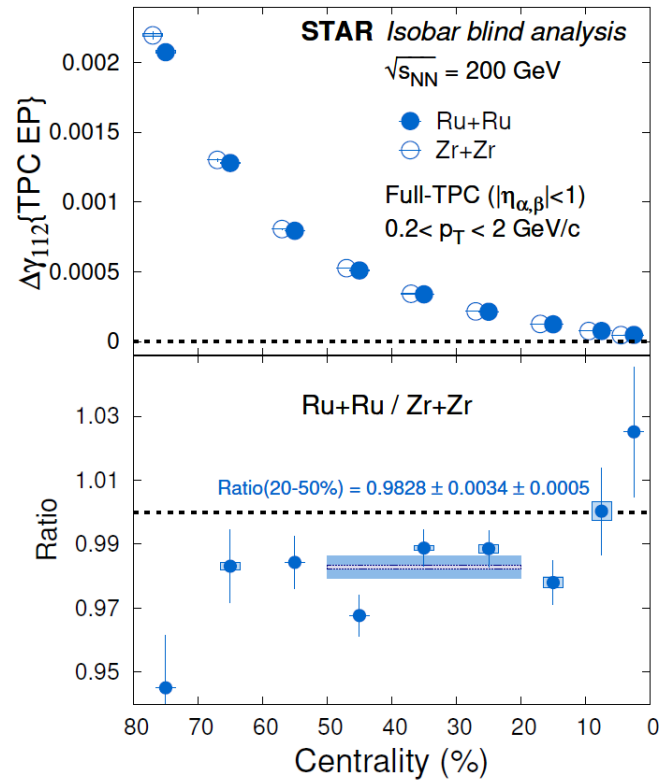


- Trivial multiplicity dilution effect
- Not included in the predefined observable
- $N\Delta\gamma/v_2$  would be better



All groups are consistent.  $\Delta\gamma/v_2$  follows closely with  $N_{\text{ch}}$

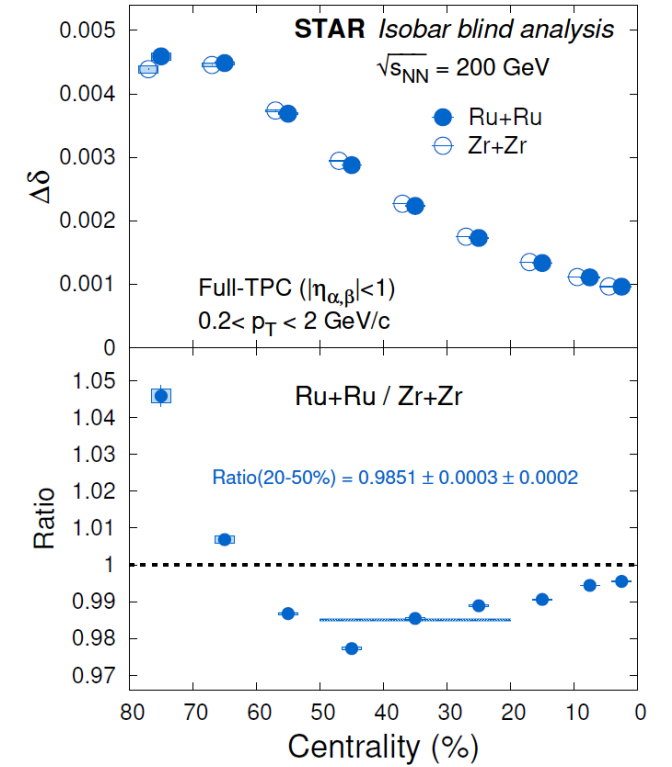
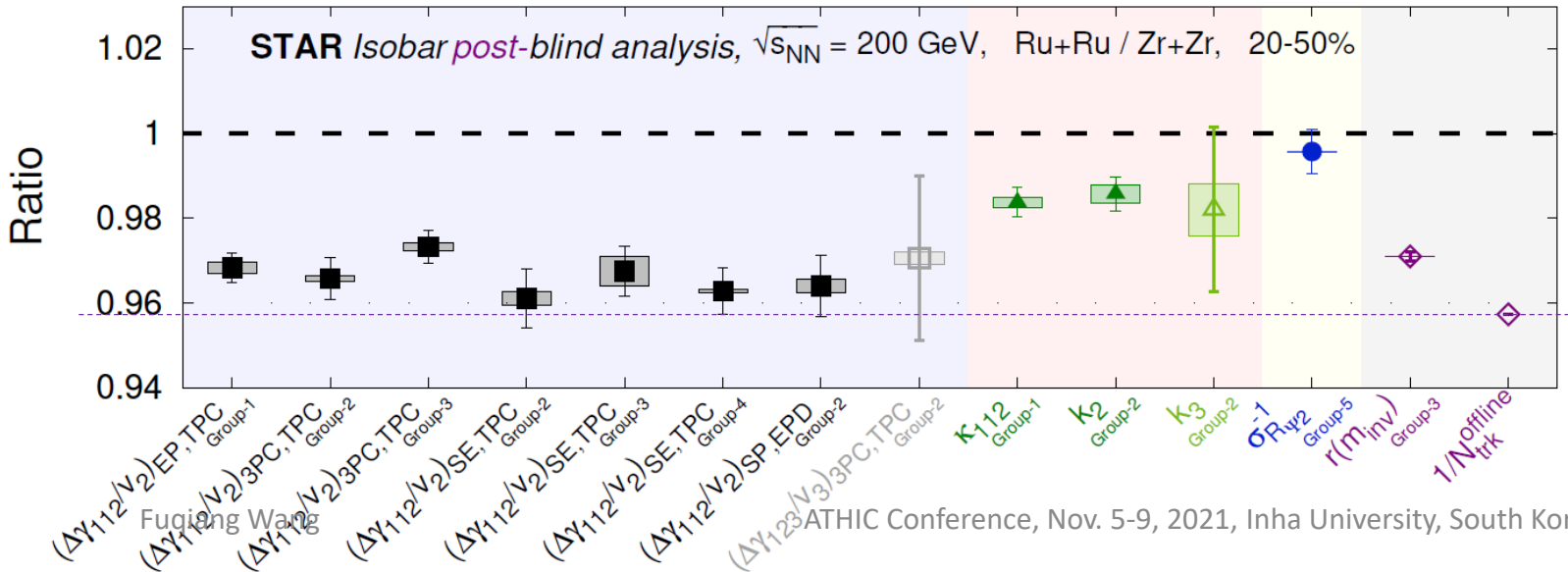
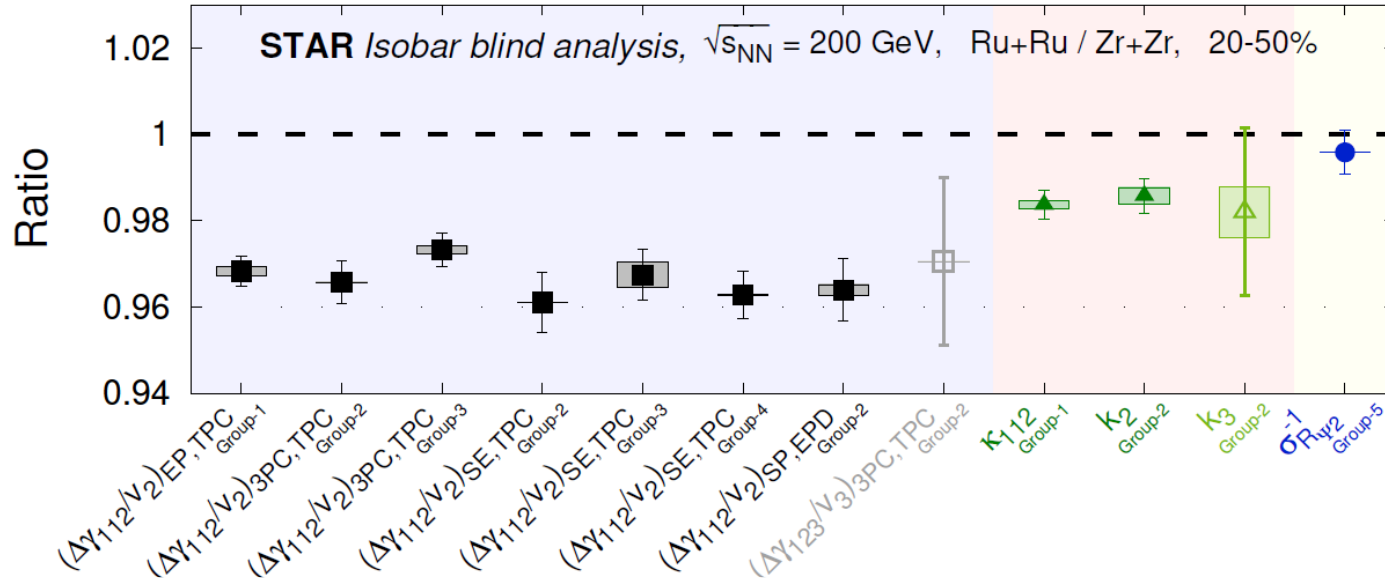
# $\Delta\gamma$ , $\Delta\gamma/v_2$ , $\kappa=\Delta\gamma/(\Delta\delta*v_2)$ MEASUREMENTS



Indeed a precision of 0.4% is achieved!

Ru+Ru/Zr+Zr ratios all below unity, naively unexpected; main reason is the 4.4% Nch difference

# MONEY PLOTS



Nonflow: 
$$\frac{(N\Delta\delta)^{Ru+Ru}}{(N\Delta\delta)^{Zr+Zr}} \approx 1.03$$

**Nonflow difference is important!**

# REMAINING NONFLOW EFFECTS

Feng et al., arXiv:2106.15595

$$f_{\text{CME}} = \frac{\Delta\gamma_{\text{CME}}\{\text{PP}\}}{\Delta\gamma\{\text{PP}\}} = \frac{A/a - 1}{1/a^2 - 1}$$

$$\frac{A}{a} = \frac{\Delta\gamma\{\text{SP}\} / v_2\{\text{SP}\}}{\Delta\gamma\{\text{PP}\}^* / v_2\{\text{PP}\}^*} = \frac{C_3\{\text{SP}\} / v_2^2\{\text{SP}\}}{C_3\{\text{PP}\}^* / v_2^2\{\text{PP}\}^*} = \frac{1 + \varepsilon_{\text{nf}}}{1 + \frac{\varepsilon_3 / \varepsilon_2}{Nv_2^2\{\text{PP}\}}}$$

$$C_3\{\text{SP}\} = \frac{C_{2\text{p}}N_{2\text{p}}}{N^2} v_{2,2\text{p}}\{\text{SP}\}v_2\{\text{SP}\},$$

Nonflow in  $\Delta\gamma$   
 $\rightarrow$  negative  $f_{\text{CME}}$

$$v_2^*\{\text{EP}\} = \sqrt{v_2^2\{\text{EP}\} + v_{2,\text{nf}}^2}$$

$$\frac{\left(N\Delta\gamma / v_2^*\right)^{\text{Ru}}}{\left(N\Delta\gamma / v_2^*\right)^{\text{Zr}}} = \frac{\left(NC_3 / v_2^{*2}\right)^{\text{Ru}}}{\left(NC_3 / v_2^{*2}\right)^{\text{Zr}}} = \frac{\left(C_{2\text{p}} \frac{N_{2\text{p}}}{N} \frac{v_{2,2\text{p}}}{v_2}\right)^{\text{Ru}}}{\left(C_{2\text{p}} \frac{N_{2\text{p}}}{N} \frac{v_{2,2\text{p}}}{v_2}\right)^{\text{Zr}}} \cdot \frac{\left(1 + \varepsilon_{\text{nf}}\right)^{\text{Zr}}}{\left(1 + \varepsilon_{\text{nf}}\right)^{\text{Ru}}} \cdot \frac{\left(1 + \frac{\varepsilon_3 / \varepsilon_2}{Nv_2^2}\right)^{\text{Ru}}}{\left(1 + \frac{\varepsilon_3 / \varepsilon_2}{Nv_2^2}\right)^{\text{Zr}}}$$

$$C_3 = \frac{C_{2\text{p}}N_{2\text{p}}}{N^2} v_{2,2\text{p}}v_2 + \frac{C_{3\text{p}}N_{3\text{p}}}{2N^3}$$

$$\varepsilon_2 = \frac{C_{2\text{p}}N_{2\text{p}}}{N} \cdot \frac{v_{2,2\text{p}}}{v_2}$$

$$\varepsilon_3 = \frac{C_{3\text{p}}N_{3\text{p}}}{2N}$$

- Depending on the relative Ru+Ru/Zr+Zr difference of various nonflow effects, **the baseline can be above, equal, or below unity**
- **Final isobar conclusion will require detailed nonflow studies**

# SUMMARY AND OUTLOOK

- CME is very important physics. Significant efforts in theory and experiments.
- STAR has pioneered and played significant role in the CME search. Primary efforts in understanding and removing backgrounds.
- The possible CME is a small fraction of the measured  $\Delta\gamma$  signal. Most recent STAR data indicate **a finite CME signal with 1-3 $\sigma$  significance**; nonflow effects under investigation.
- Isobar blind analysis is a tour de force. Anticipated **precision down to 0.4%** is achieved. No CME signal is observed in the blind analysis; not inconsistent with Au+Au data. **Further (nonflow) investigations** needed to quantify significance.
- Current data **2.4B MB Au+Au, 3.8B isobar events. Expect 20B Au+Au from 2023+25 runs**, together with large BES-II data samples.

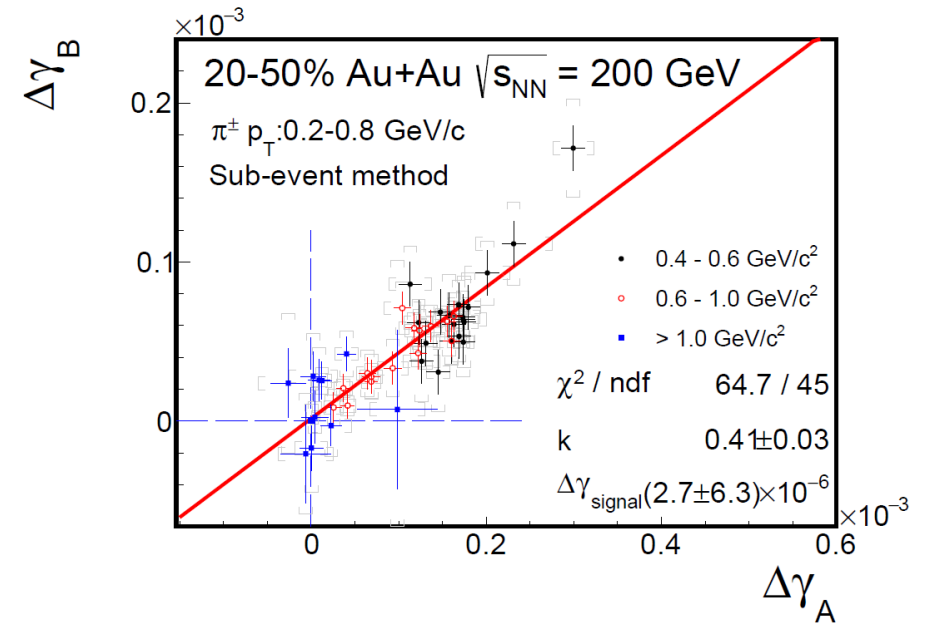
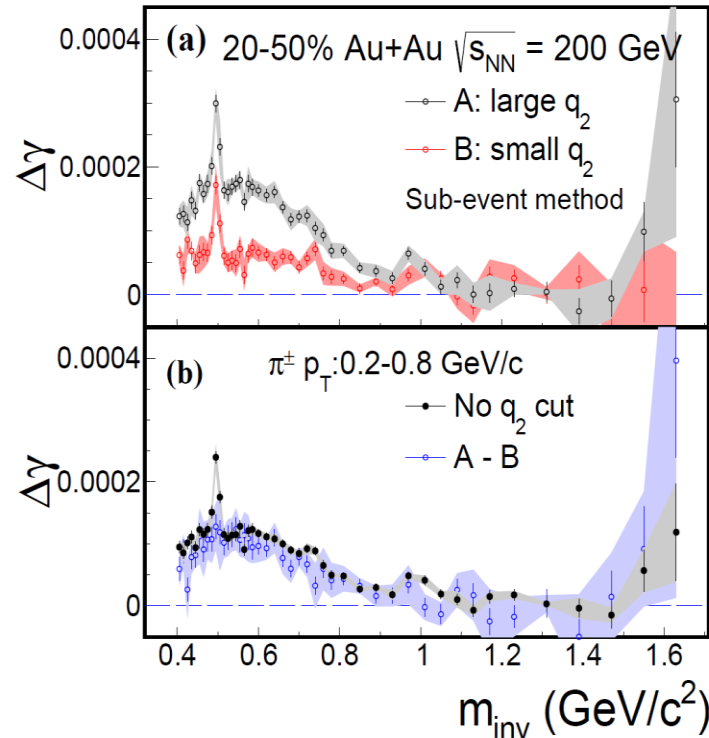
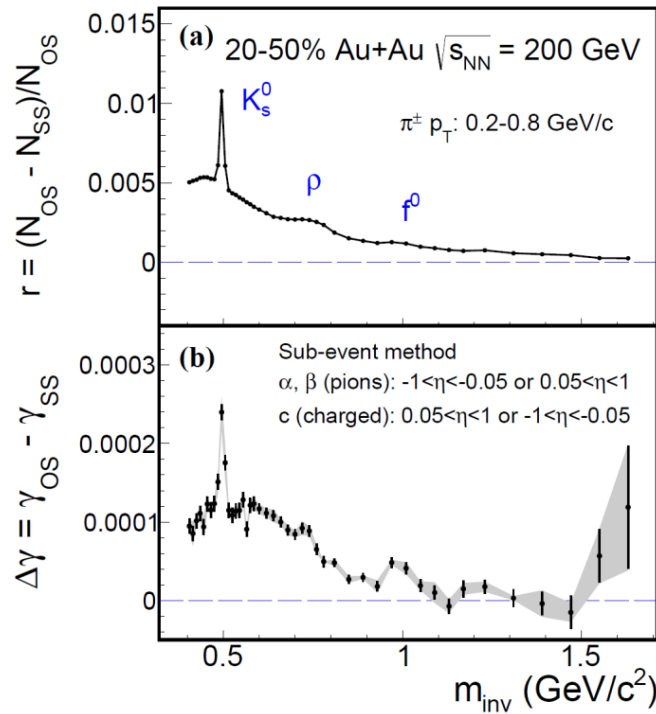
# Backup slides

# THE INVARIANT MASS METHOD

Zhao, Li, Wang, Eur.Phys.J.C 79 (2019) 2, 168

$$\frac{N_\rho}{N_\alpha N_\beta} \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_{clus}) \rangle \times v_{2,clus}$$

STAR, arXiv:2006.05035

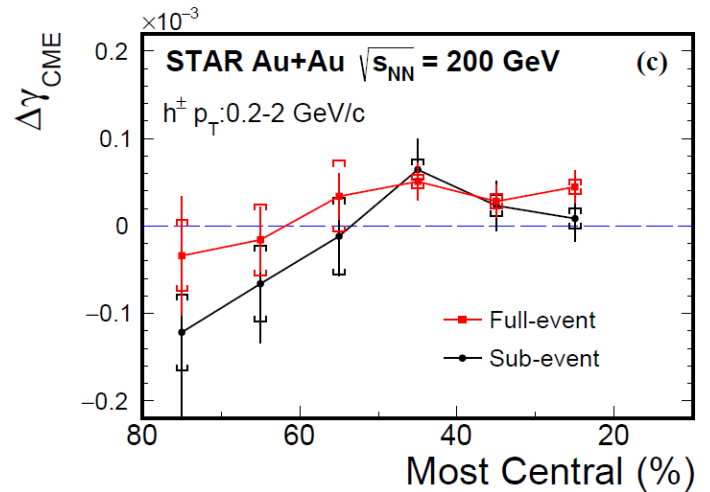
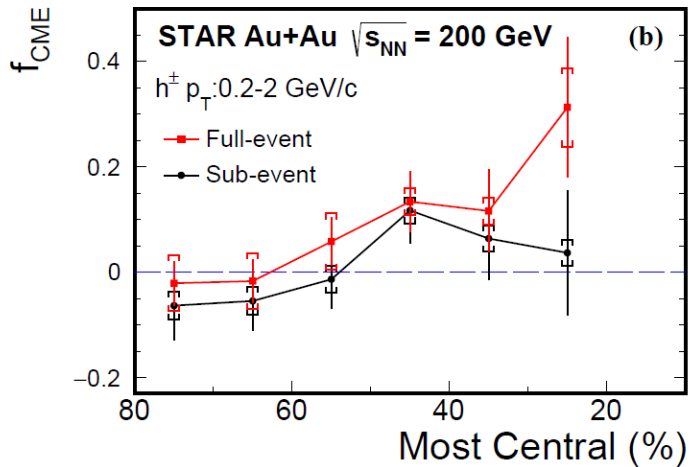
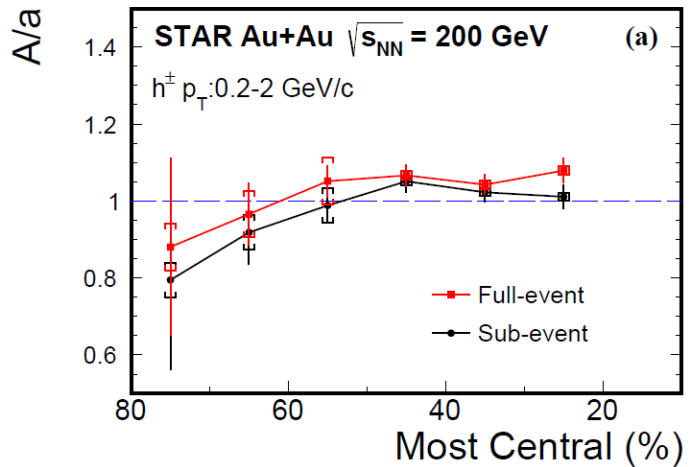
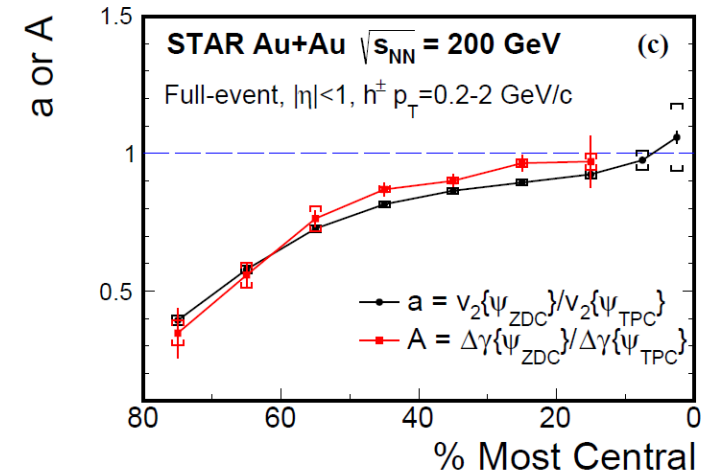
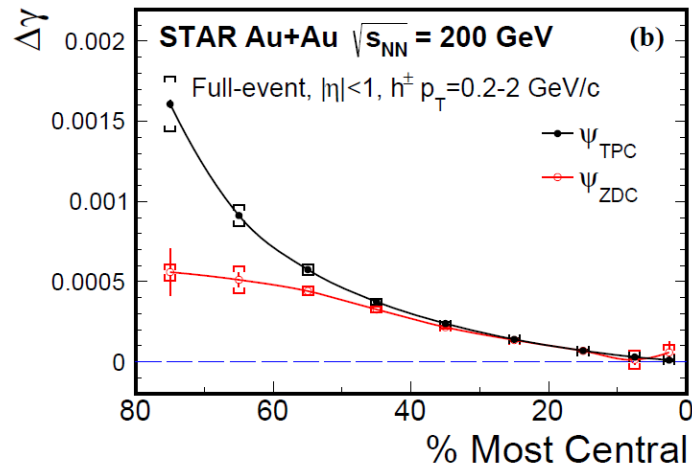
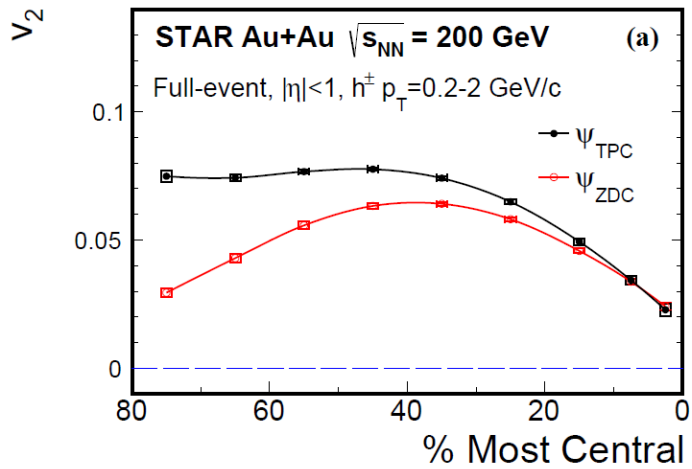


CME fraction =  $(2 \pm 4 \pm 5)\%$   
 CME upper limit 15% at 95% CL



# Au+Au Collisions at 200 GeV (2.4B MB)

STAR, arXiv:2106.09243

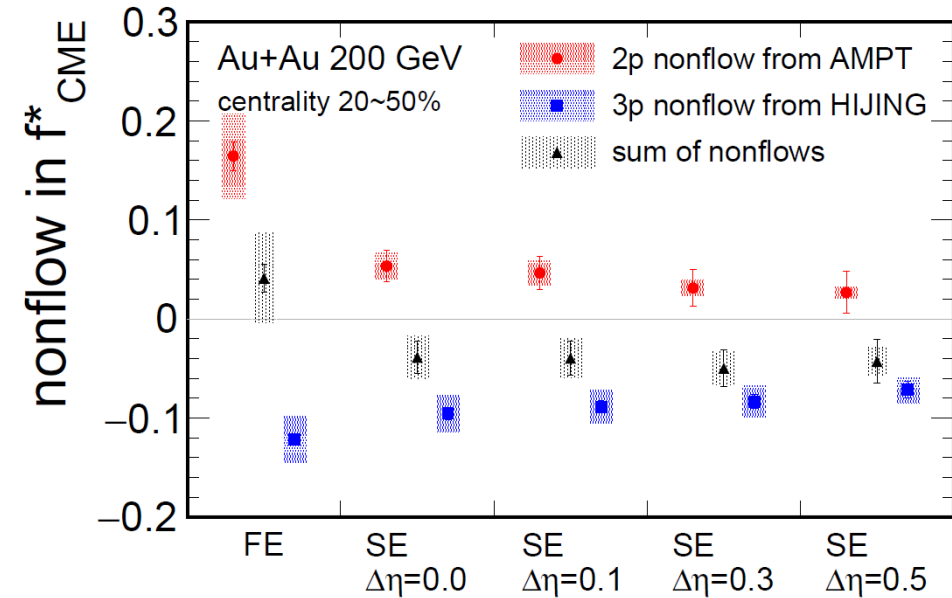
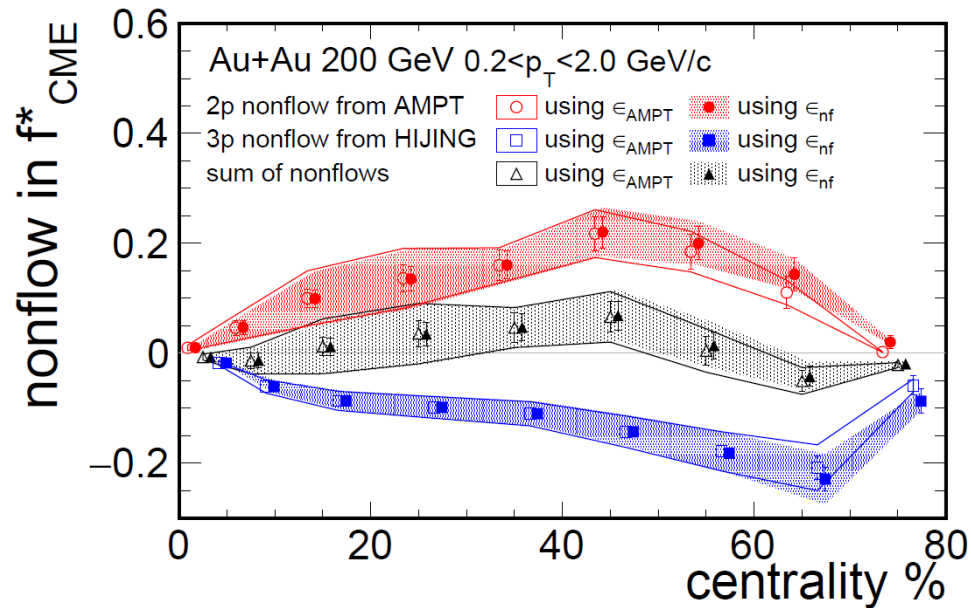


# MODEL ESTIMATES OF NONFLOW

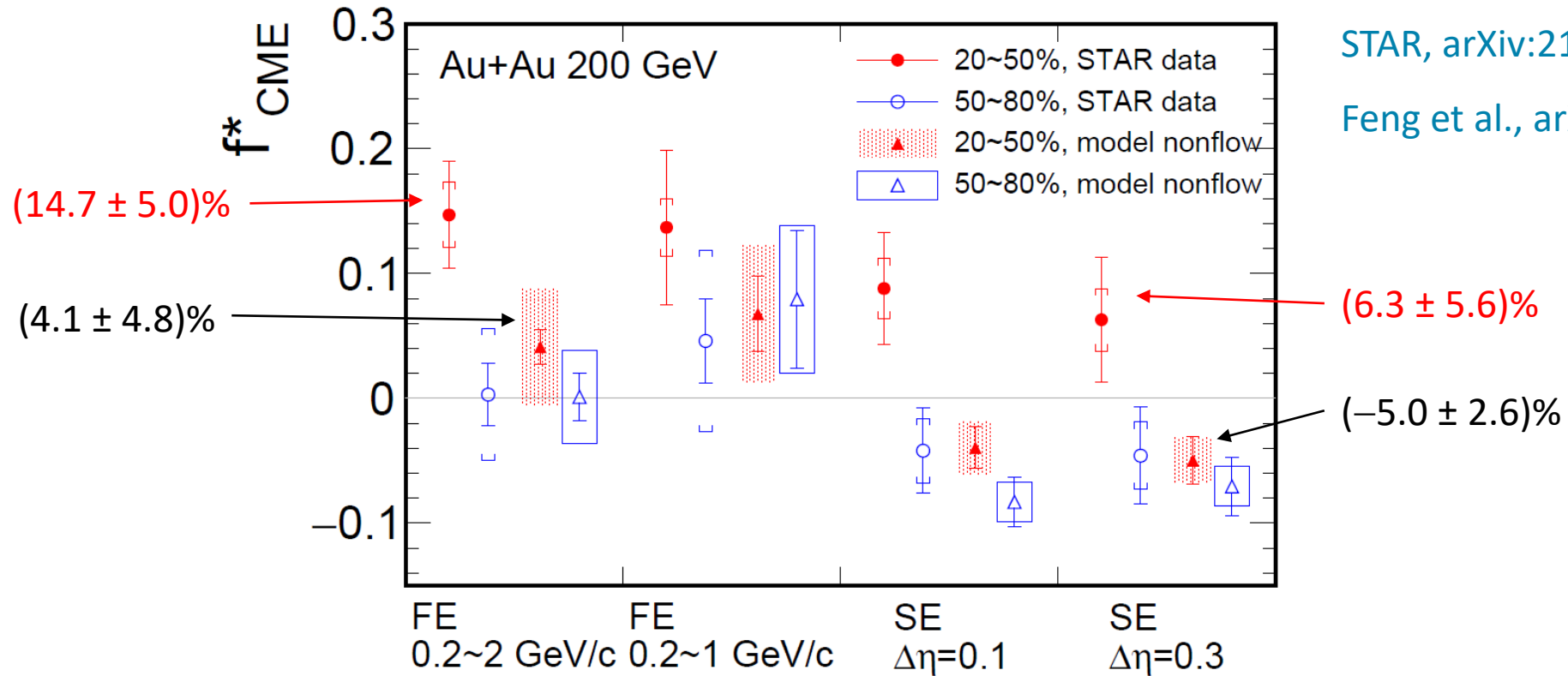
Feng et al., arXiv:2106.15595

$$\frac{A}{a} = \frac{\Delta\gamma\{\text{SP}\} / v_2\{\text{SP}\}}{\Delta\gamma\{\text{PP}\}^* / v_2\{\text{PP}\}^*} = \frac{C_3\{\text{SP}\} / v_2^2\{\text{SP}\}}{C_3\{\text{PP}\}^* / v_2^2\{\text{PP}\}^*} = \frac{1 + \epsilon_{\text{nf}}}{1 + \frac{\epsilon_3 / \epsilon_2}{Nv_2^2\{\text{PP}\}}}$$

$$f_{\text{CME}}^* \approx \left( \epsilon_{\text{nf}} - \frac{\epsilon_3 / \epsilon_2}{Nv_2^2\{\text{EP}\}} \right) / \left( \frac{1 + \epsilon_{\text{nf}}}{a^2} - 1 \right)$$



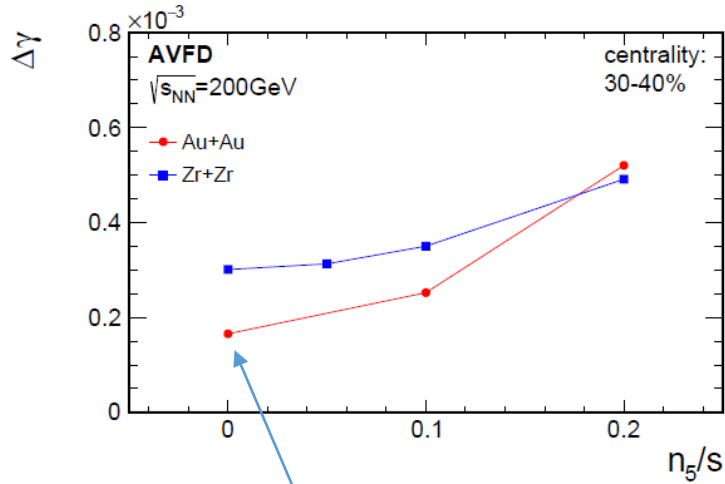
# NONFLOW EFFECTS IN $f_{\text{CME}}^*$



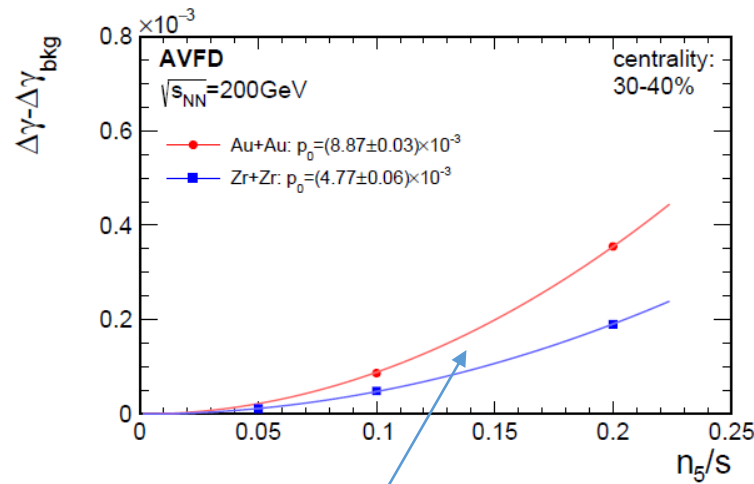
There may indeed be hint of CME in the data,  $\sim 2\sigma$

# Au+Au DATA AND ISOBAR ARE CONSISTENT

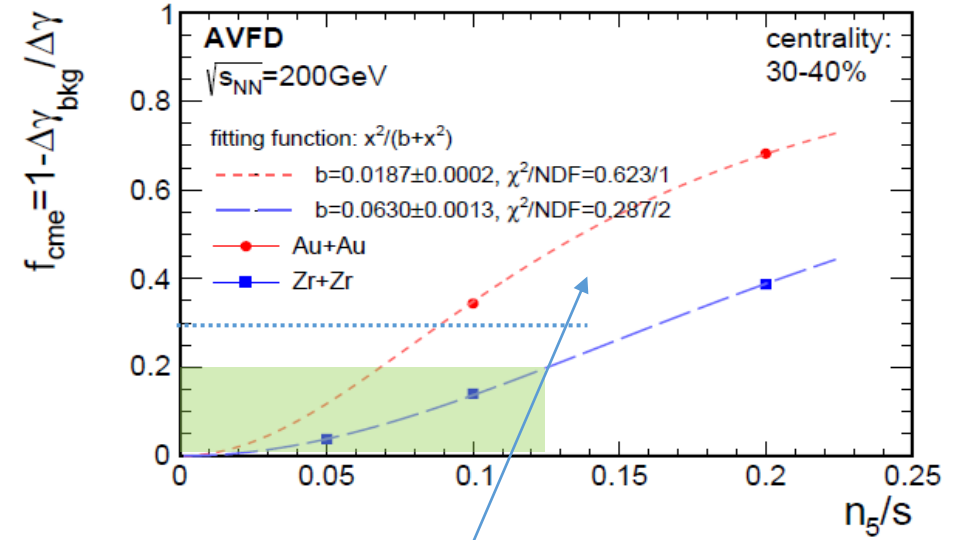
Yicheng Feng, Yufu Lin, Jie Zhao, FW, arXiv:2103.10378



Background  $\propto 1/N$   
 isobar/AuAu  $\sim 2$



Mag. field  $B \sim A/A^{2/3} \sim A^{1/3}$   
 $\Delta\gamma_{\text{CME}} \sim B^2 \sim A^{2/3}$   
 Signal: AuAu/isobar  $\sim 1.5$



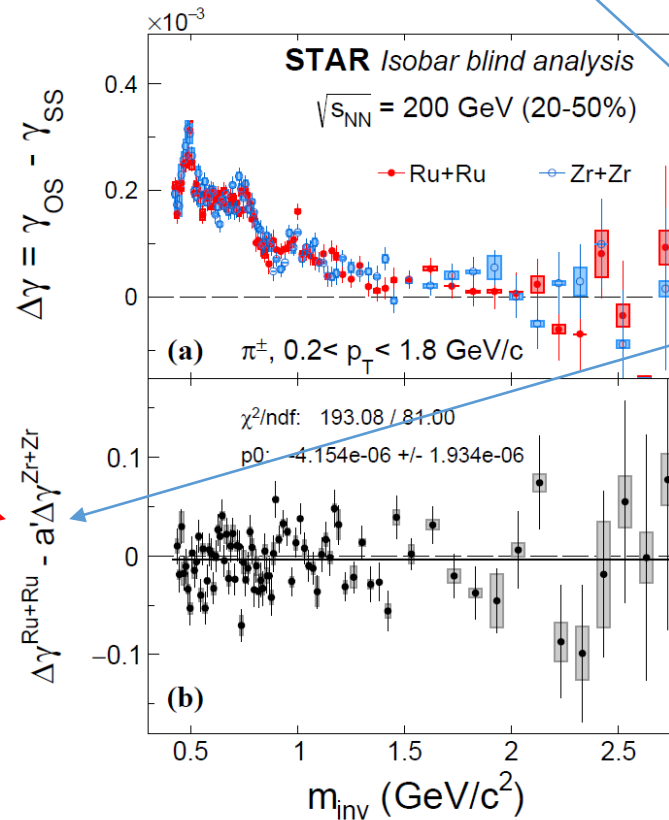
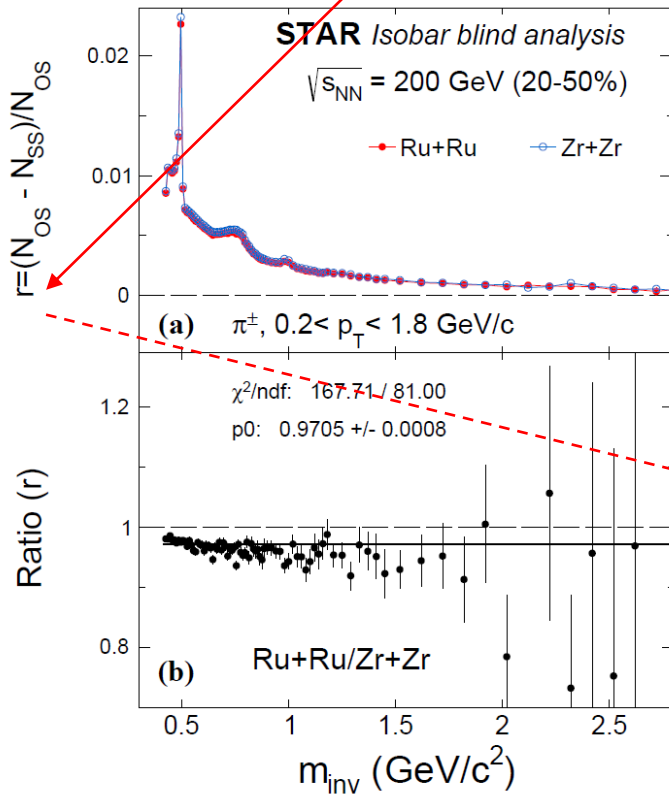
Could be x3 reduction in  $f_{\text{CME}}$  at the same  $n_5/s$   
 If AuAu  $f_{\text{CME}}=10\%$ , then isobar 3% ( $1\sigma$  effect)  
 $\text{Ru/Zr} = 1 + 15\% \cdot 3\% = 1.005 (\pm 0.004)$

Caveats: Axial charge densities and sphaleron transition probabilities could be different between Au+Au and isobar, e.g. AVFD-glasma  $\mu_5/s$ : isobar/AuAu  $\sim 1.5$

# INVARIANT MASS MEASUREMENT

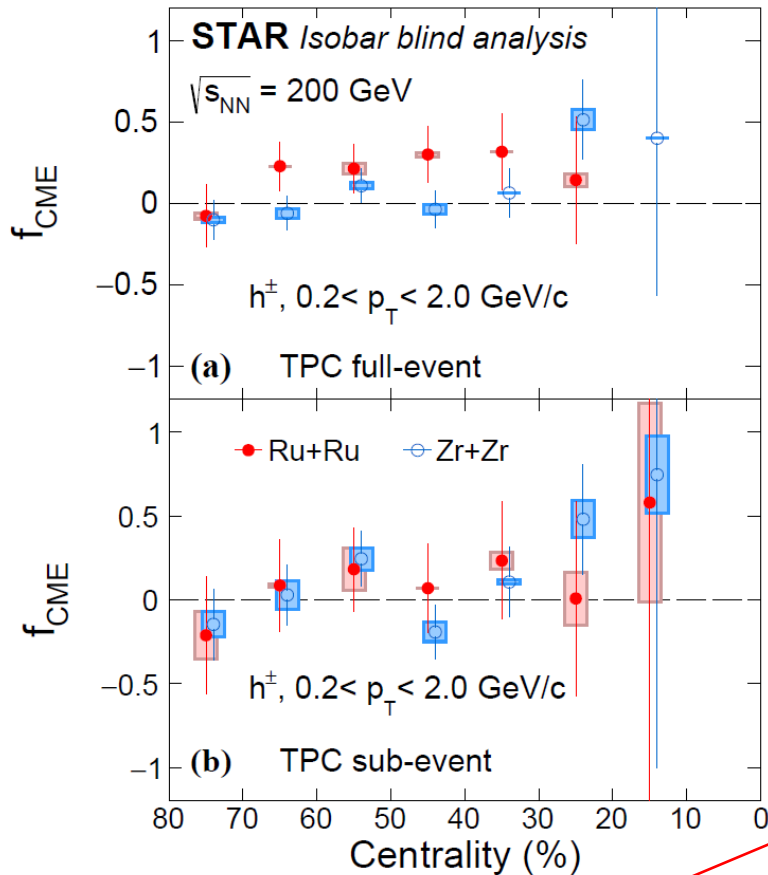
$$\Delta\gamma_{\text{bkgd}} = \frac{4N_{2p}}{N^2} \langle \cos(\phi_\alpha + \phi_\beta - 2\phi_{2p}) \rangle v_{2,2p}$$

$$r = \frac{N_{\text{OS}} - N_{\text{SS}}}{N_{\text{OS}}} \quad \text{Relative pair multiplicity difference}$$



- $r$  deviates from unity, qualitatively consistent with  $1/N$  ratio.
- $a' = v_2^{\text{Ru+Ru}} / v_2^{\text{Zr+Zr}}$
- $r$  not included in the predefined  $a'$
- Including  $r$  into  $a'$ ,  $\Delta\gamma^{\text{Ru+Ru}} - a' \Delta\gamma^{\text{Zr+Zr}}$  becomes numerically positive but within  $1\sigma$  from zero.

# CME FRACTION MEASUREMENTS



$$f_{CME} = \frac{\Delta\gamma_{CME}\{TPC\}}{\Delta\gamma\{TPC\}} = \frac{A/a - 1}{1/a^2 - 1}$$

$$A = \Delta\gamma\{ZDC\}/\Delta\gamma\{TPC\}$$

$$a = v_2\{ZDC\}/v_2\{TPC\}$$

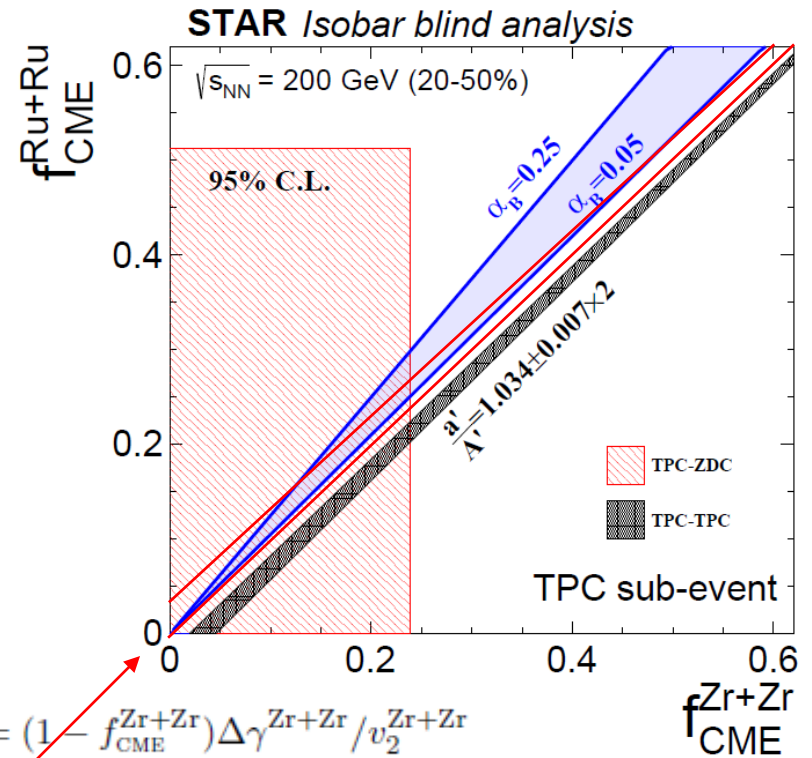
$$(1 - f_{CME}^{Ru+Ru})\Delta\gamma^{Ru+Ru}/v_2^{Ru+Ru} = (1 - f_{CME}^{Zr+Zr})\Delta\gamma^{Zr+Zr}/v_2^{Zr+Zr}$$

$$f_{CME}^{Ru+Ru} = \left(\frac{a'}{A'}\right) f_{CME}^{Zr+Zr} + \left(1 - \frac{a'}{A'}\right)$$

$$A' = \Delta\gamma^{Ru+Ru}/\Delta\gamma^{Zr+Zr}$$

$$a' = v_2^{Ru+Ru}/v_2^{Zr+Zr}$$

N ratio not included in the predefined  $a'$ .  
 Including it,  $a'/A' = 0.990 \pm 0.007$



B-field expectation:

$$f_{CME}^{Ru+Ru} / f_{CME}^{Zr+Zr} = 1 + \alpha_B$$

where  $\alpha_B = 0.15 \pm 0.05$