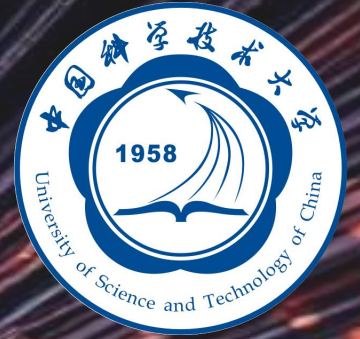




**UPC2025: The second international workshop on
the physics of Ultra Peripheral Collisions**



Coulomb Dissociation Measurement in Isobar UPC at $\sqrt{s_{NN}} = 200$ GeV with The STAR Experiment

Huda Nasrulloh (for STAR Collaboration)
University of Science and Technology of China
- hudanasrulloh@mail.ustc.edu.cn -

Supported in part by the
U.S. DEPARTMENT OF

ENERGY

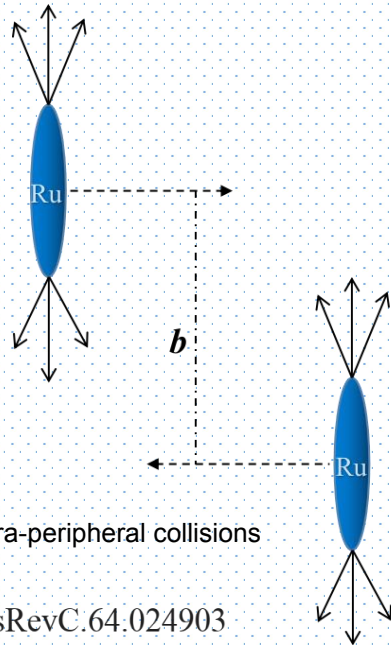
Office of
Science

UPC 2025, Saariselkä, Finland, June 9-13, 2025

Outline

- Introduction
- Detector and data processing
- Approaching method
- Coulomb dissociation *preliminary* result
- Summary and outlook

Electromagnetic Field in Heavy Ion Collisions



a). Ultra-peripheral collisions

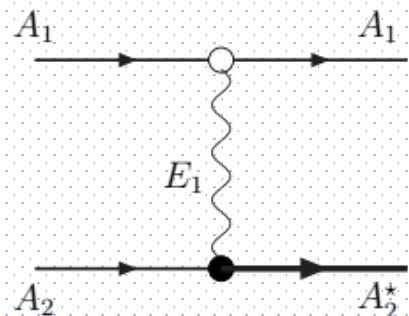
□ Ultra-Peripheral Collisions (UPC):

- Two ions pass each other ($b > R_1 + R_2$)
- Interaction dominated by relativistic electromagnetic field

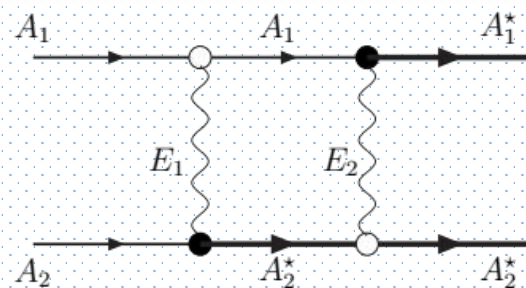
□ Photonuclear Interactions:

- One or both ions emit virtual photons absorbed by the other, causing nuclear excitation and breakup
- Neutron emission occurs through a process known as Coulomb Dissociation (CD)

DOI: 10.1103/PhysRevC.64.024903



b). Single Coulomb dissociation (**SCD**)



c). Mutual Coulomb dissociation (**MCD**)

Isobaric Nuclei (Ru+Ru & Zr+Zr) Collisions

Ru	Zr	Ru/Zr
A = 96	A = 96	1.00
Z = 44	Z = 40	1.10
N / Z ≈ 1.18	N / Z ≈ 1.40	0.84

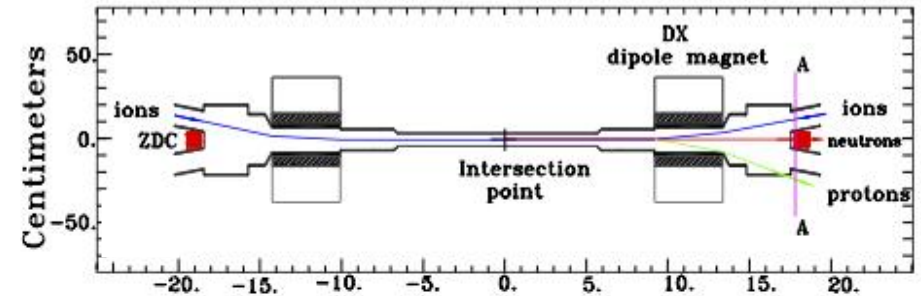
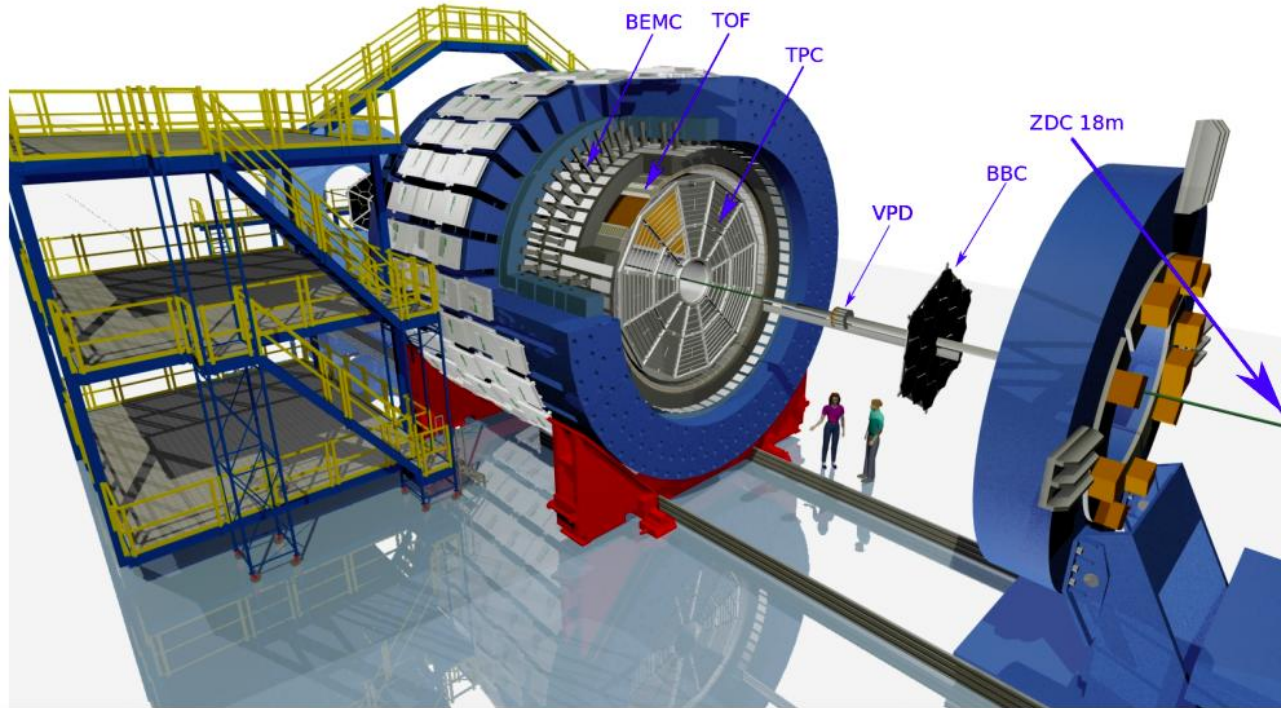
Key points:

- $\gamma \propto Z^2$
- $(Z^2_{\text{Ru}} / Z^2_{\text{Zr}}) \approx 1.21$

□ Isobaric Nuclei: Ruthenium (A=96, Z=44) and Zirconium (A=96, Z=40)

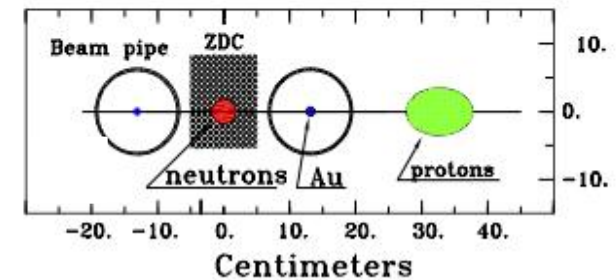
- Same **mass number**, different **proton** number
- Direct comparison of nuclear structures
- Probing the influence of proton number on nuclear properties ($\gamma \propto Z^2$)

STAR-ZDC Detector



arXiv:nucl-ex/0008005v1

A-A



□ STAR Experiment and ZDC Detectors:

- The STAR experiment at RHIC is equipped with the ZDC detectors
- The ZDC detects neutrons emitted in CD events

Data Processing

Data set:

Run18 Ru+Ru / Zr+Zr collisions at

$$\sqrt{s_{NN}} = 200 \text{ GeV}$$

Trigger :

ZDC-Mon

Zerobias

Neutron channels:

ZDC-Mon -> MCD {1n1n, 1nXn, XnXn}

ZB -> SCD {0n1n, 0nXn}

ZDC-Mon	
Collisions Species	Number of Events
Ru+Ru	5.113 M
Zr+Zr	4.988 M

Zerobias	
Collisions Species	Number of Events
Ru+Ru	4.935 M
Zr+Zr	4.843 M

MCD

zdc-mon (history)	37	600017	1.81K	8.81K
zdc-mon (history) --> Requirements:		+ZDC-TAC +ZDC-E +ZDC-W -Laser-protection		

SCD

ZEROBIAS (history)	60	9300	844.48K	4.18M
ZEROBIAS (history) --> Requirements:		-Laser-protection		

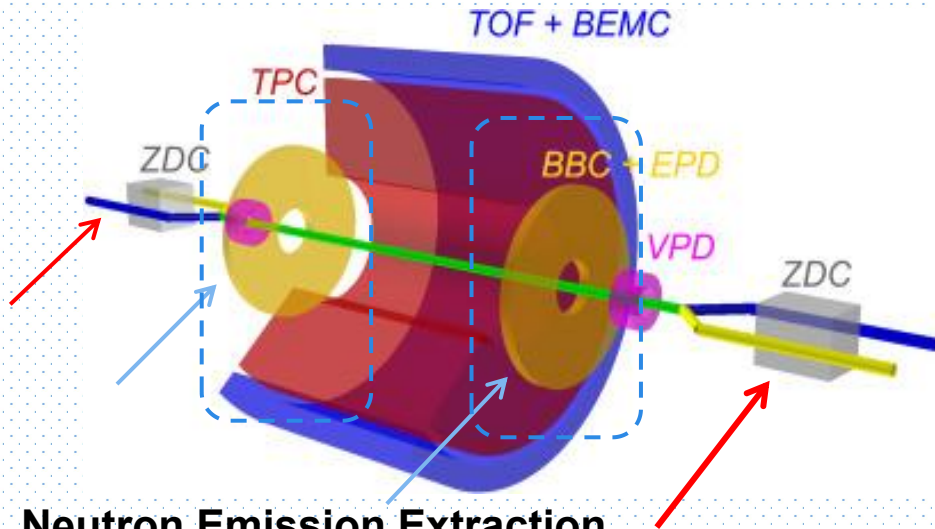
Approaching method

1. Event Filtering:

The BBC detector vetoing **hadronic** events:

- at least one side $BBC_n \leq 20$ (ZDC-Mon)
- at least one side $BBC_n \leq 1$ (Zero-Bias)

BBC_n = BBC coincidence tagging of hadrons, allowing neutron emissions to pass through and be detected by the ZDC.

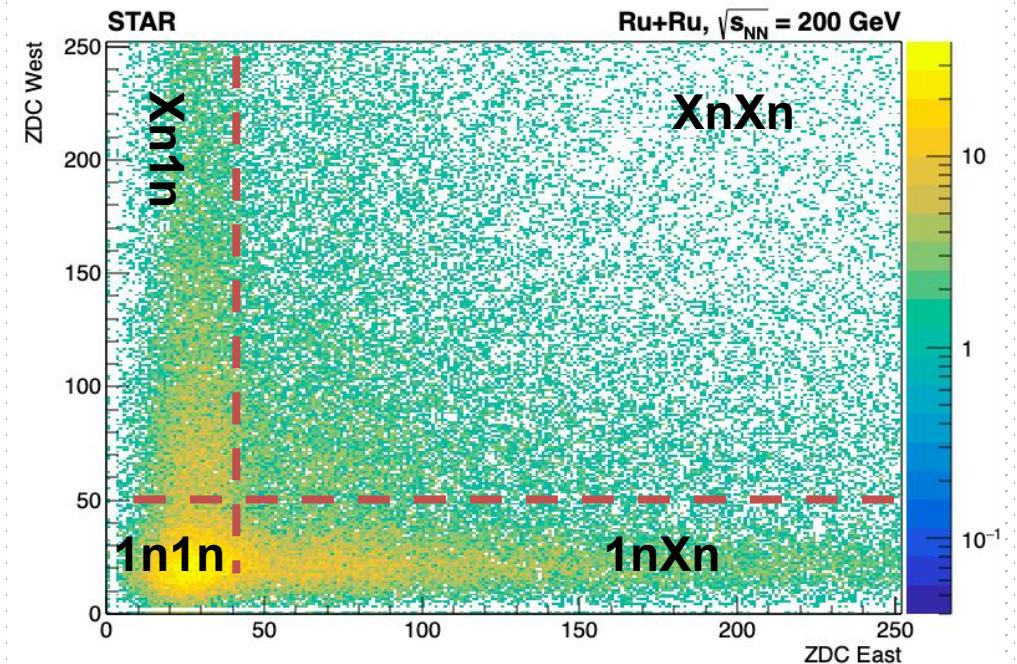


3. Neutron Emission Extraction

Perform multi-gaussian fit to extract neutron emission distribution

2. Neutron tagging of UPC:

Neutron-tagged events are selected using the ZDC



4. Cross-Section Calculation:

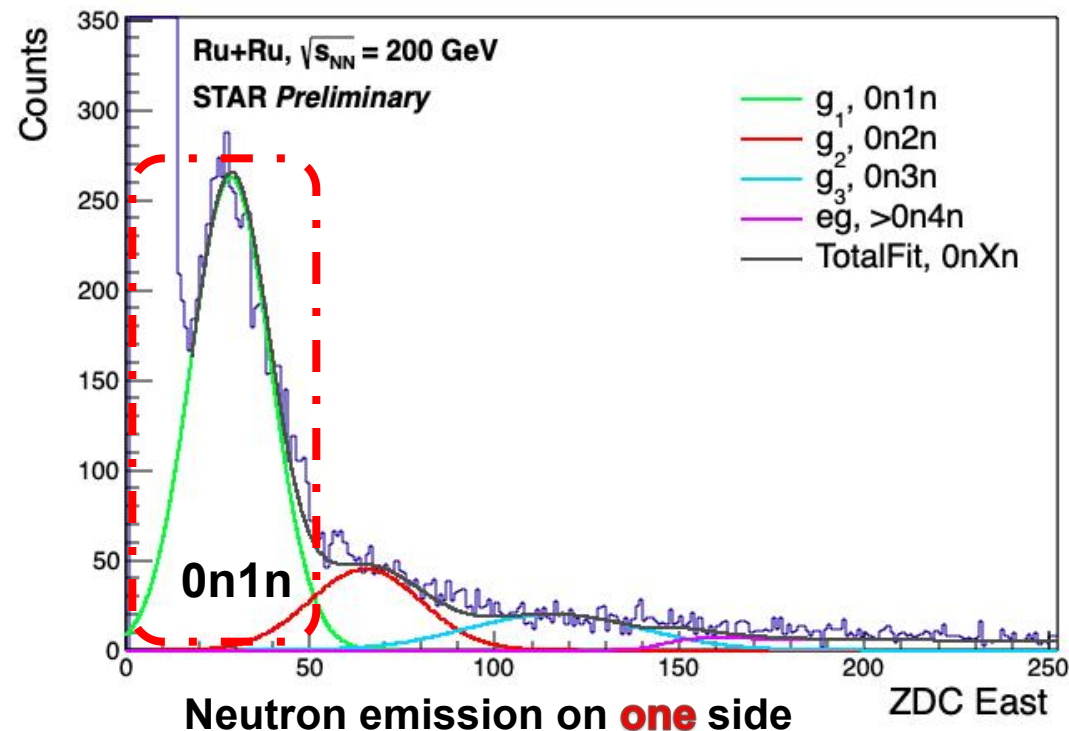
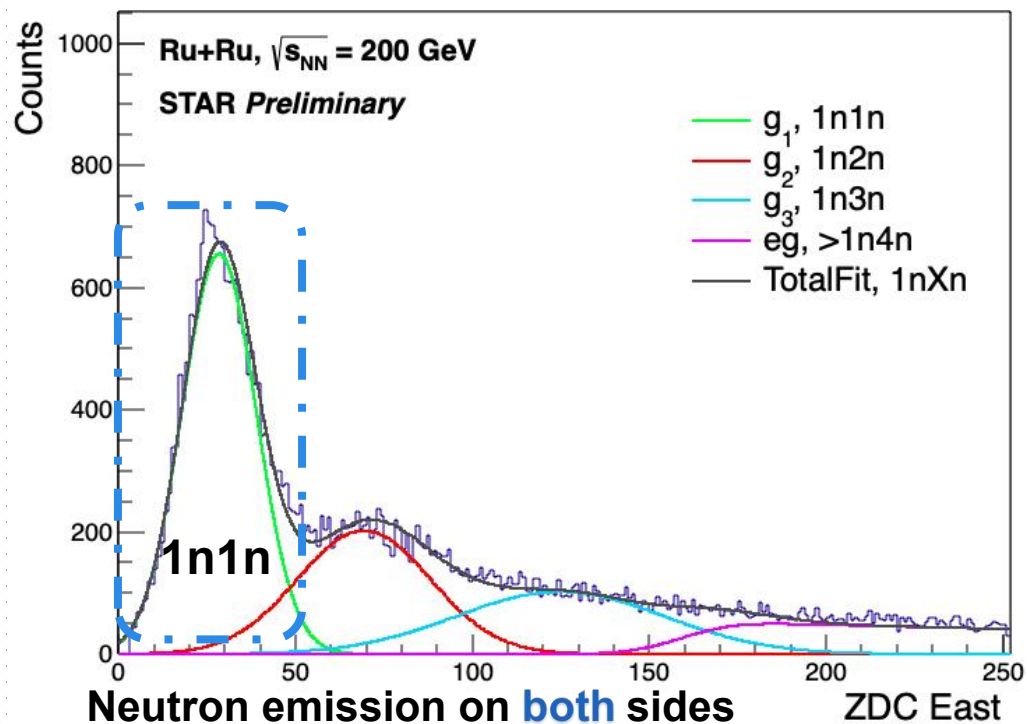
Calculate σ_{cd} for each neutron emission channel

Neutron emission extraction

Ru (A=96, Z=44)

We extracted neutron emission by using the multi-gaussian fitting function

$$\underbrace{\text{TotalFit}}_{Xn} = \underbrace{g_1(x)}_{1n} + \underbrace{g_2(x)}_{2n} + \underbrace{g_3(x)}_{3n} + \underbrace{eg(x)}_{>4n}$$



Coulomb dissociation cross-section

We calculate Coulomb dissociation cross-section using the formula :

$$\sigma_{cd} = \frac{\sum N_{n-channel}}{L_{trigger}},$$

where $\sum N_{n-channel}$ corresponds to each neutron emission channel, for **SCD** ($\sum N_{[0n1n]}$, $\sum N_{[0nXn]}$) and **MCD** ($\sum N_{[1n1n]}$, $\sum N_{[1nXn]}$, & $\sum N_{[XnXn]}$), with the luminosity of each trigger (L_{zb} and $L_{zdc-mon}$) taken into account.

ZDC-Mon Luminosity

The STAR recorded data was analyzed to obtain *luminosity of the ZDC-mon triggered events* (L_{zdcmon})

	Name	Species	Tid	Lum [ub ⁻¹]	Nev [M]
	zdc-mon	RuRu		1.499	6.956
zdc-mon		RuRu	600007	0.019	0.086
zdc-mon		RuRu	600017	1.480	6.870
	zdc-mon	ZrZr		1.530	7.098
zdc-mon		ZrZr	600007	0.001	0.004
zdc-mon		ZrZr	600017	1.529	7.094

ZDC_Mon	
Collision Species	Number of Events
Ru+Ru	5.113M
Zr+Zr	4.988M

- $$L_{\text{zdcmon}}[\text{Ru}](Xn, Xn) = \frac{\sum N(Xn, Xn)_{\text{Ru}} \cdot L_{\text{zdc_coinc}}(\text{Ru})}{\sum N_{\text{zdc_coinc}}(Xn, Xn)_{\text{Ru}}} = \frac{5.113\text{M} \cdot 1.499 \mu\text{b}^{-1}}{6.956\text{M}} \approx 1.101 \mu\text{b}^{-1}$$
- $$L_{\text{zdcmon}}[\text{Zr}](Xn, Xn) = \frac{\sum N(Xn, Xn)_{\text{Zr}} \cdot L_{\text{zdc_coinc}}(\text{Zr})}{\sum N_{\text{zdc_coinc}}(Xn, Xn)_{\text{Zr}}} = \frac{4.989\text{M} \cdot 1.530 \mu\text{b}^{-1}}{7.096\text{M}} \approx 1.076 \mu\text{b}^{-1}$$

ZDC-Coincidence

	Name	Species	Tid	Lum [ub ⁻¹]	Nev [M]
	zdc-mon	RuRu		1.499	6.956
zdc-mon		RuRu	600007	0.019	0.086
zdc-mon		RuRu	600017	1.480	6.870
	zdc-mon	ZrZr		1.530	7.098
zdc-mon		ZrZr	600007	0.001	0.004
zdc-mon		ZrZr	600017	1.529	7.094

The STAR recorded data enables to calculate the *ZDC-coincidence* cross-section, σ_{coinc} , which follow this formula:

$$\sigma_{\text{coinc}} = \frac{\sum N_{\text{coinc}}}{L}$$

where $\sum N$ is number of event of coincidence, and L is integrated luminosity

$$\sigma_{\text{Ru}} \approx 4.64 \text{ b}$$

$$\sigma_{\text{Zr}} \approx 4.64 \text{ b}$$

ZB Trigger Luminosity Estimation

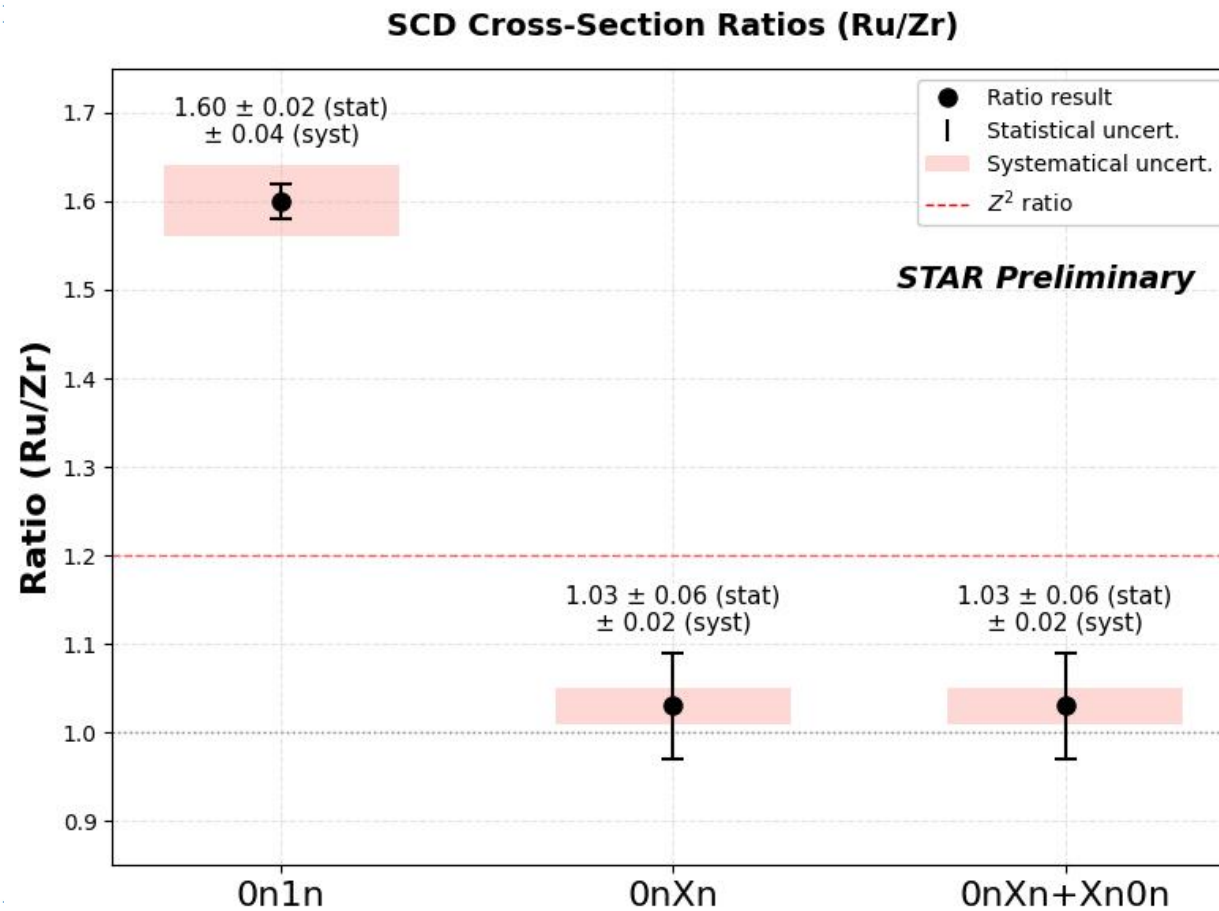
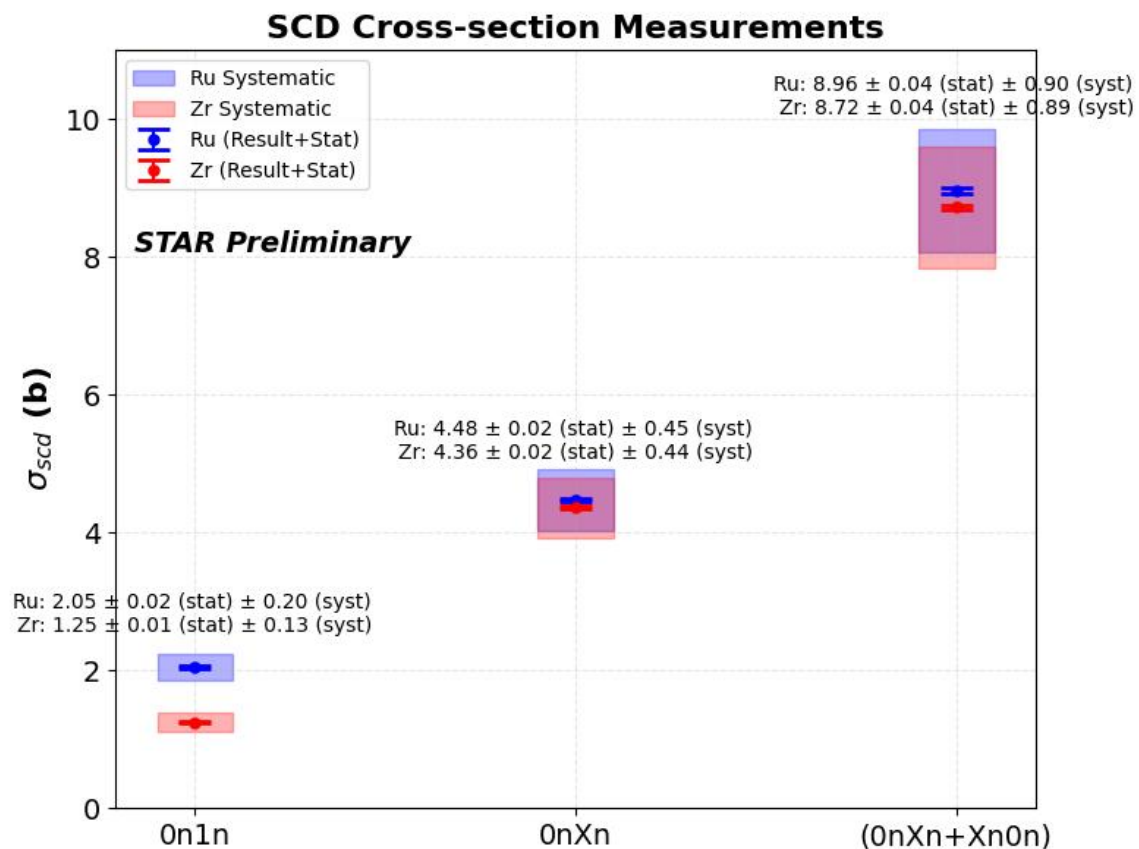
We estimate luminosity of Zerobias trigger (L_{zb}) with the (Xn,Xn) channel :

Zerobias	Ru	Zr
	$\sum N$	$\sum N$
(Xn,Xn)	13381	13304

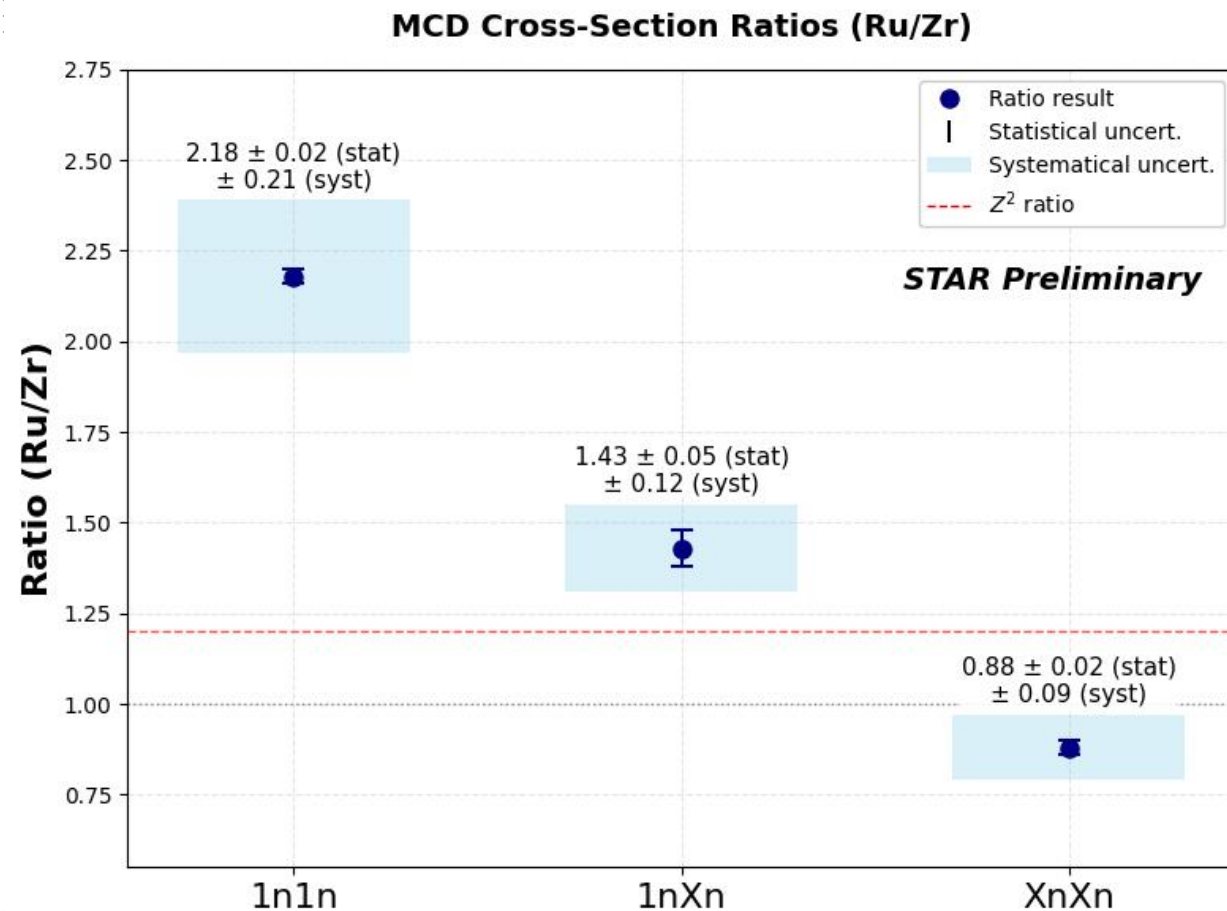
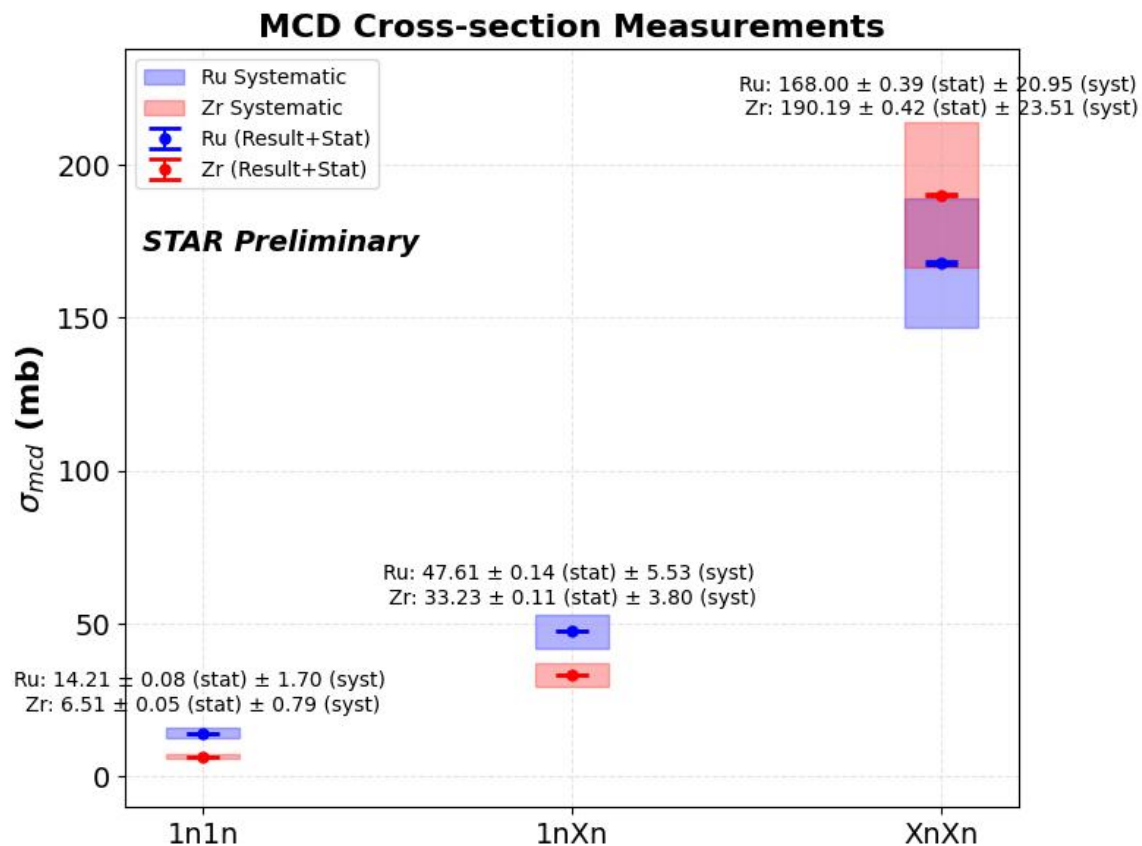
followed by *ZDC coincidence* cross-section, σ_{coinc} , thus :

- $L_{zb[\text{Ru}]}(\text{Xn,Xn}) = \frac{\sum N(\text{Xn,Xn})}{\sigma_{\text{coinc}}(\text{Ru})} = \frac{13381}{4.64 \text{ b}} = 2883.83 \text{ b}^{-1}$
- $L_{zb[\text{Zr}]}(\text{Xn,Xn}) = \frac{\sum N(\text{Xn,Xn})}{\sigma_{\text{coinc}}(\text{Zr})} = \frac{13304}{4.64 \text{ b}} = 2867.85 \text{ b}^{-1}$

Single-Coulomb *dissociation* (SCD) Cross-section



Mutual-Coulomb dissociation (MCD) Cross-section



Summary

□ Conclusion

- The first measurement of Coulomb dissociation in STAR Isobar data including the cross-section calculation of both **SCD** and **MCD** events
- The (Ru/Zr) ratios (1.60 ± 0.02 for 0n1n **SCD**) and (2.18 ± 0.02 for 1n1n **MCD**) demonstrate:
 - Clear Z-dependence, exceeding Z^2 scaling ($Z^2_{\text{Ru}} / Z^2_{\text{Zr}} \approx 1.21$)
 - At Low energies, 1n emission results may be explained by Giant Dipole Resonance (GDR) excitations, where cross-section is enhanced for higher proton of Ru-96.
 - Hints at nuclear structure effect of the Isobar
- The (Ru/Zr) ratios (1.03 ± 0.06 for 0nXn **SCD**) and (0.88 ± 0.02 for XnXn **MCD**) are near unity, implying:
 - Both Isobars are geometrically identical ($A=96$)
 - No Z-dependence which may indicate no significant EM effects

□ Outlook

- To validate the Coulomb dissociation measurements, we will model neutron emission (Giant dipole resonance + Quasi-deuteron + Nucleon resonance + Pomeron) with photon energy up to STAR.

Thank You!

Backup

Multi-gaussian fit

$$\underbrace{\text{TotalFit}}_{\text{Xn}} = \underbrace{g_1(x)}_{1\text{n}} + \underbrace{g_2(x)}_{2\text{n}} + \underbrace{g_3(x)}_{3\text{n}} + \underbrace{eg(x)}_{>4\text{n}}$$

- **Gaussian Peaks (n-neutron events):**

$$g_i(x, \mu_i, \sigma_i) = \frac{A_i}{\sigma_i \sqrt{2\pi}} \exp\left(-\frac{(x - \mu_i)^2}{2\sigma_i^2}\right) \quad \text{for } i = 1, 2, 3$$
$$\begin{aligned} \mu_2 &= 2\mu_1, & \sigma_2 &= \sqrt{2}\sigma_1 & (2\text{n peak}) \\ \mu_3 &= 3\mu_1, & \sigma_3 &= \sqrt{3}\sigma_1 & (3\text{n peak}) \end{aligned}$$

- **Exponential Modified Gaussian (>4n):**

$$eg(x, \mu_4, \sigma_4) = A_4 \left(\frac{\lambda}{2}\right) \exp\left(\frac{\lambda}{2}(2\mu_4 + \lambda\sigma_4^2 - 2x)\right) \left[1 + \operatorname{erf}\left(\frac{x - \mu_4 - \lambda\sigma_4^2}{\sigma_4\sqrt{2}}\right)\right]$$

Preliminary Result

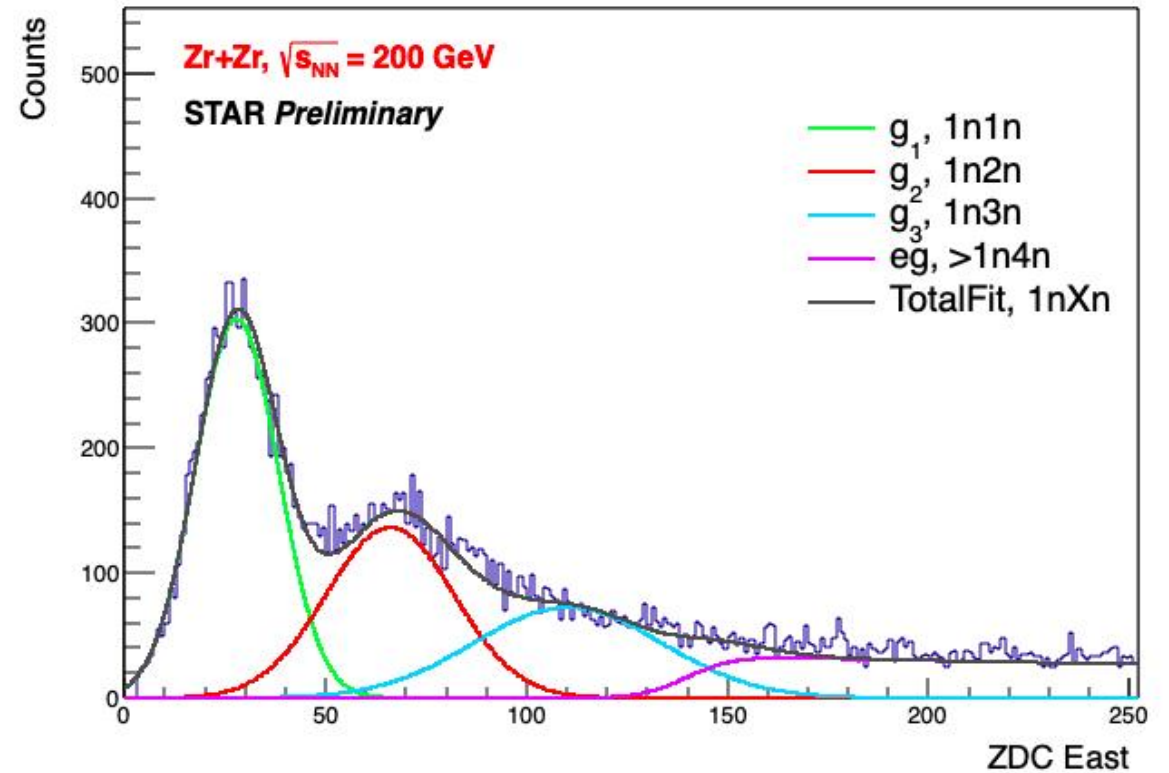
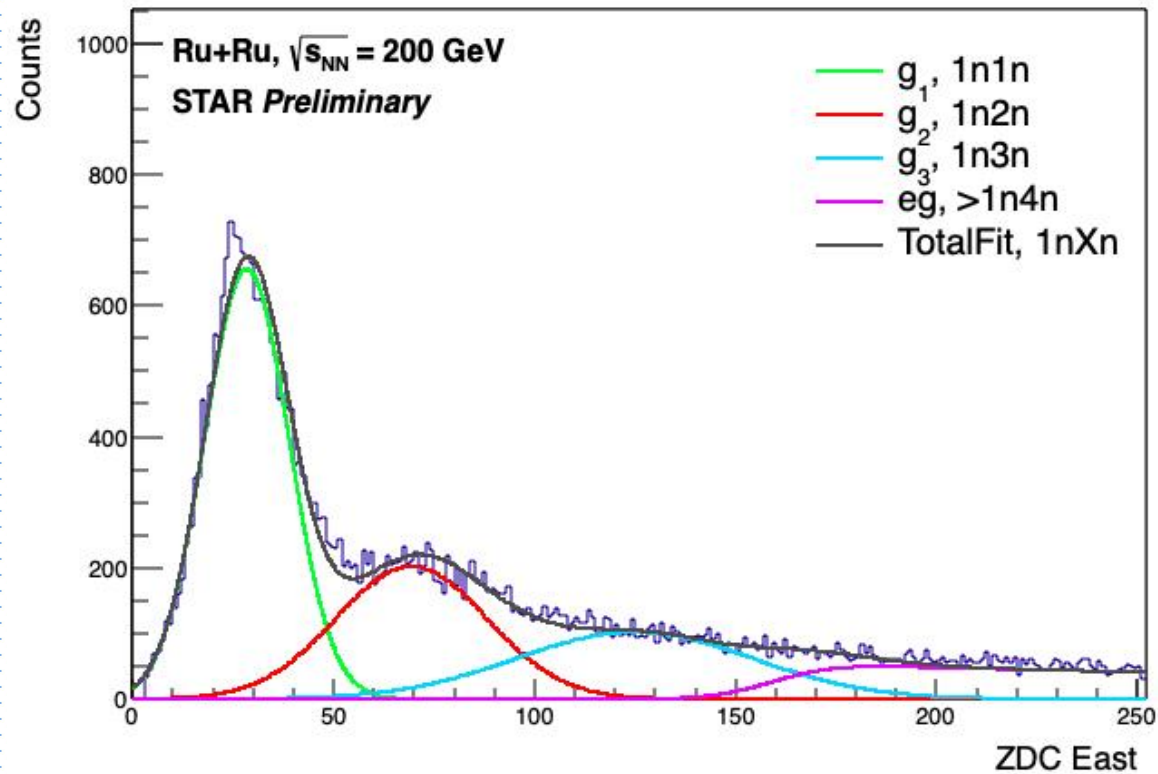
Table 1: Mutual Coulomb dissociation (MCD) cross-section of Isobar data

Neutron	Ru (mb)			Zr (mb)			Ratio (Ru/Zr)		
1n1n	14.21	± 0.08	± 1.70	6.51	± 0.05	± 0.79	2.18	± 0.02	± 0.21
1nXn	47.61	± 0.14	± 5.53	33.23	± 0.11	± 3.80	1.43	± 0.05	± 0.12
XnXn	168.00	± 0.39	± 20.95	190.19	± 0.42	± 23.51	0.88	± 0.02	± 0.09

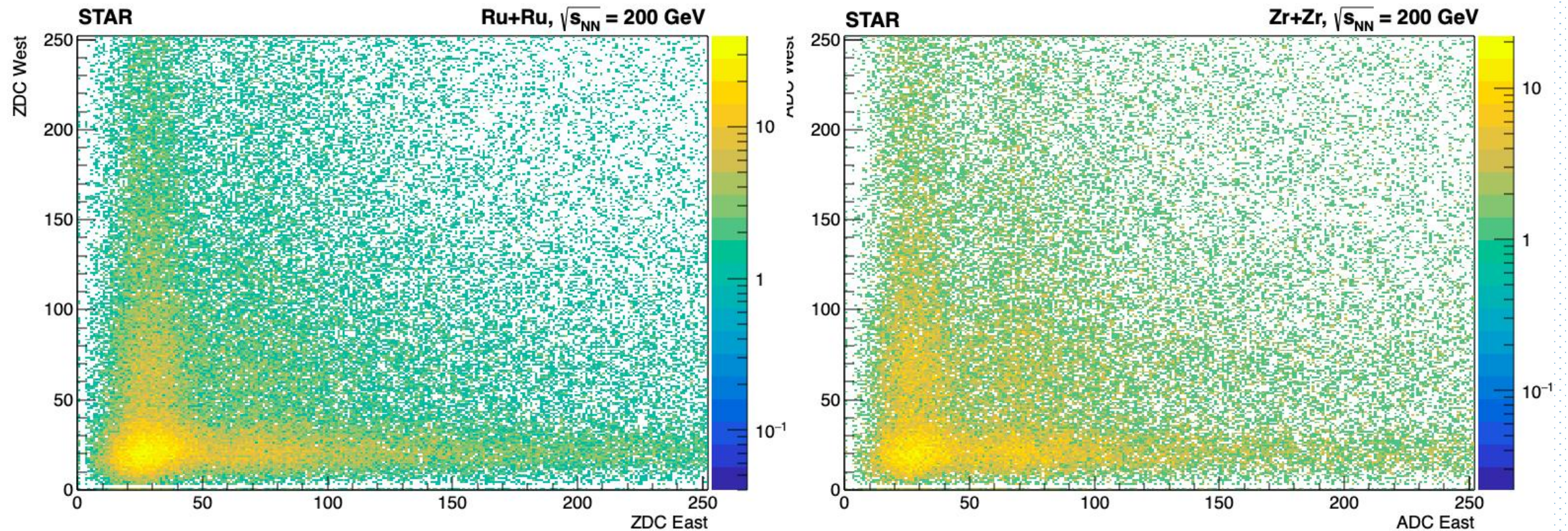
Table 2: Single Coulomb dissociation (SCD) cross-section of Isobar data

Neutron	Ru (b)			Zr (b)			Ratio (Ru/Zr)		
0n1n	2.05	± 0.02	± 0.20	1.25	± 0.01	± 0.13	1.60	± 0.02	± 0.04
0nXn	4.48	± 0.02	± 0.45	4.36	± 0.02	± 0.44	1.03	± 0.06	± 0.02
0nXn+Xn0n	8.96	± 0.04	± 0.90	8.72	± 0.04	± 0.89	1.03	± 0.06	± 0.02

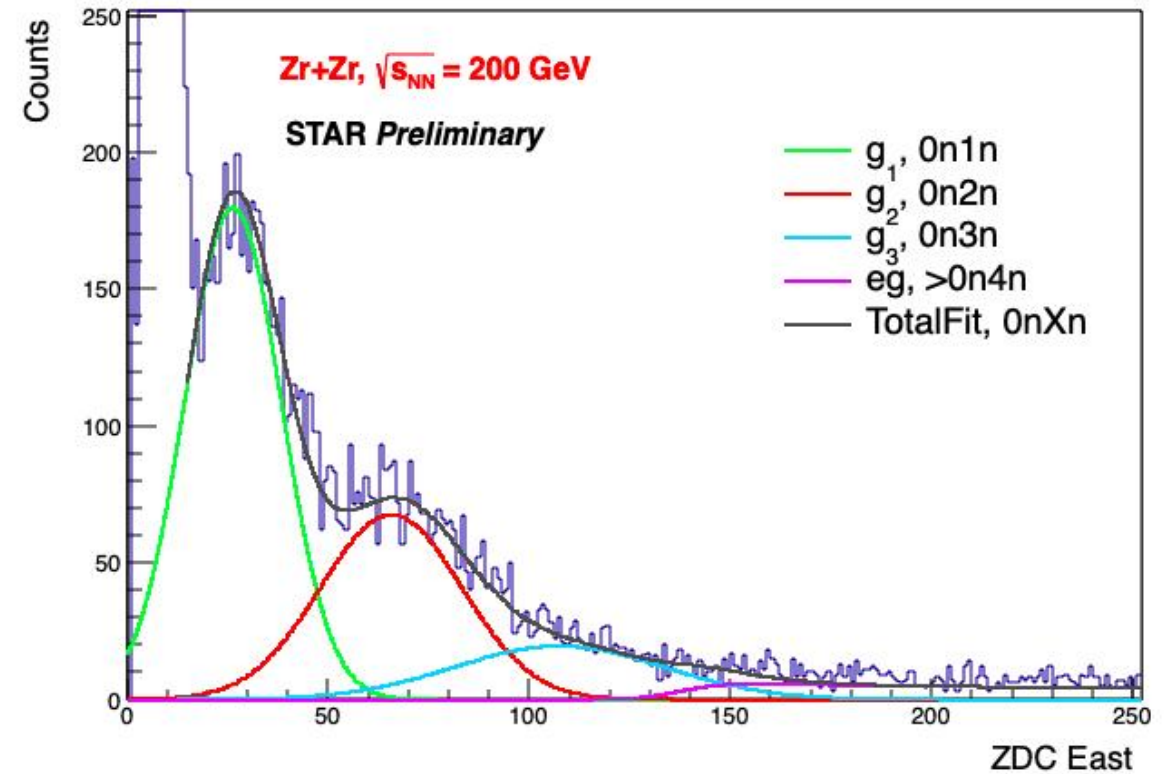
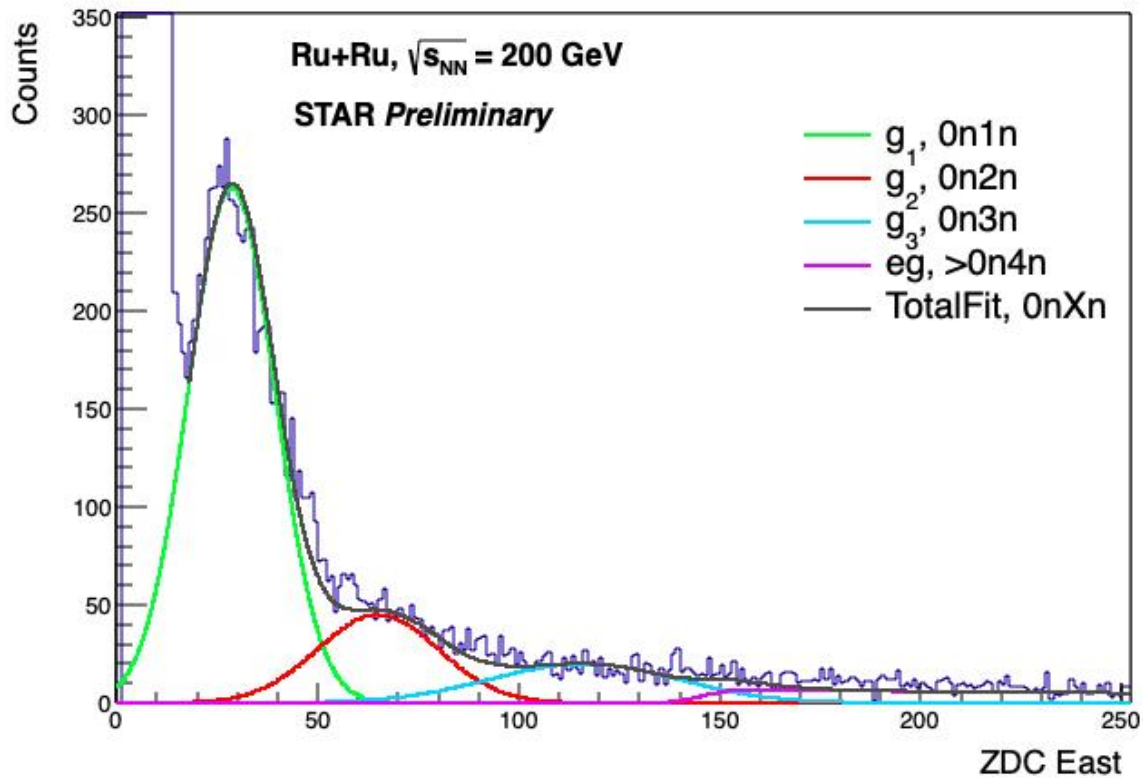
1nXn (Ru vs Zr)



XnXn (Ru vs Zr)



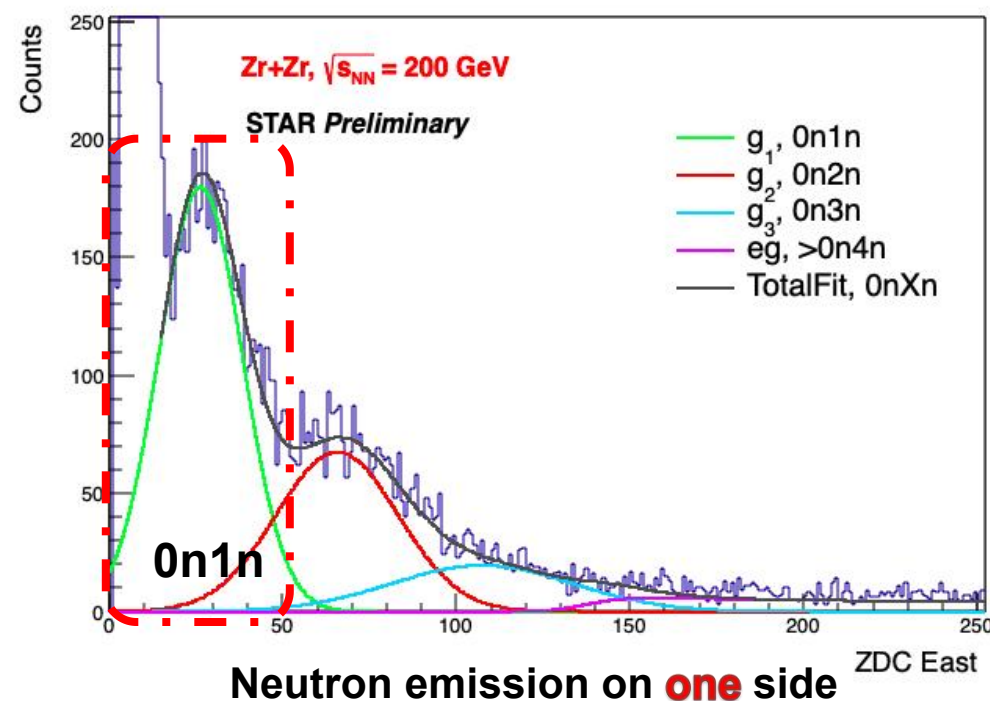
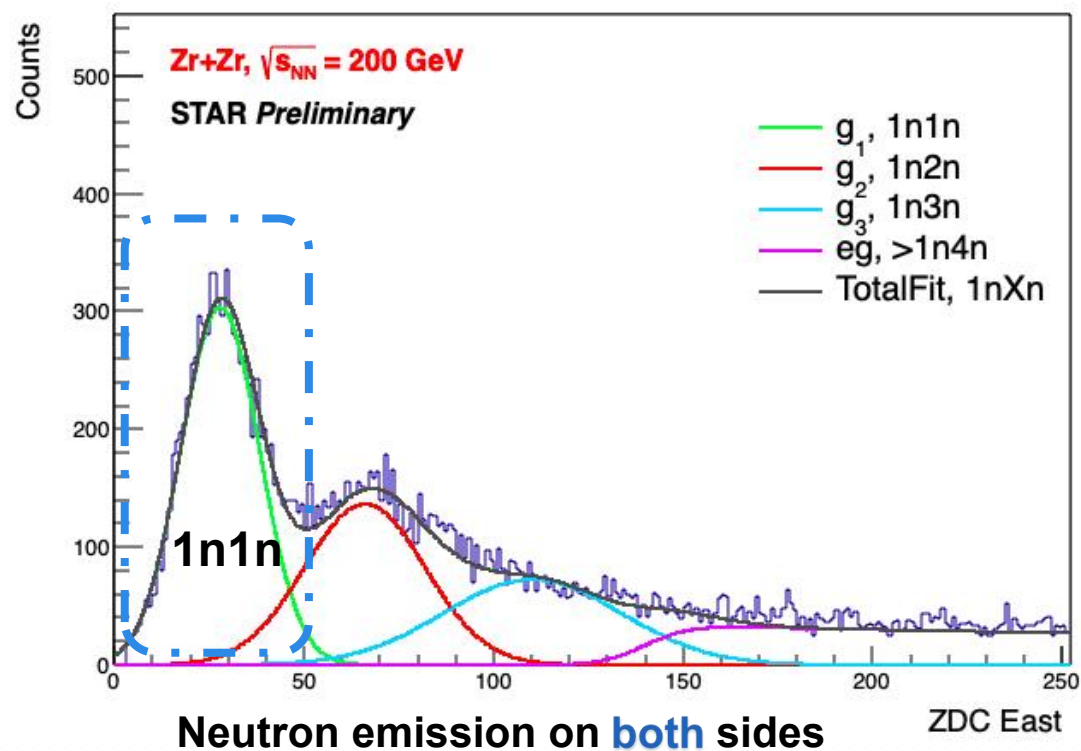
0nXn (Ru vs Zr)



Neutron emission extraction

Zr (A=96, Z=40)

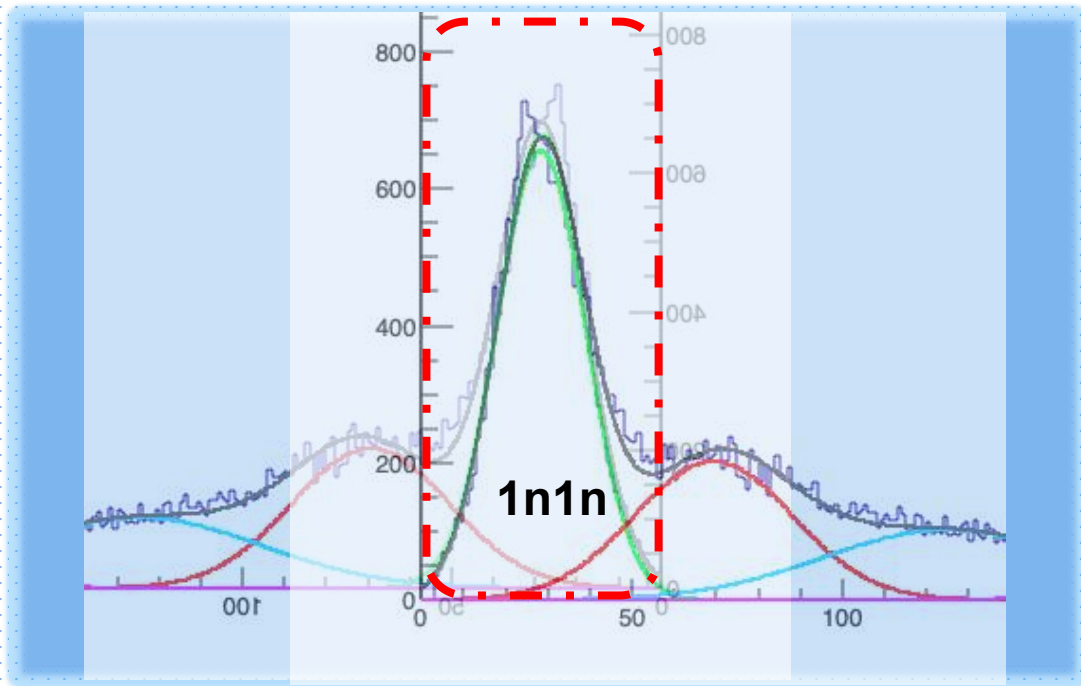
We extracted neutron emission by fitting function

$$\underbrace{\text{TotalFit}}_{Xn} = \underbrace{g_1(x)}_{1n} + \underbrace{g_2(x)}_{2n} + \underbrace{g_3(x)}_{3n} + \underbrace{eg(x)}_{>4n}$$


Efficiency and purity of 1n event

Ru (A=96, Z=44)

The 1n emission probably also contains 2n. Therefore, we calculate efficiency & purity



$$1n_{(corrected)} = \underbrace{g_1(x)}_{1n} \cdot \frac{\epsilon_{1n}}{p_{1n}}$$

with $\epsilon = \frac{N_{1n,selected}}{N_{1n,total}}$ (Efficiency) $p = \frac{N_{1n,selected}}{N_{all,selected}}$ (Purity)