

Measurements of the Υ meson production in Au+Au collisions by the STAR experiment



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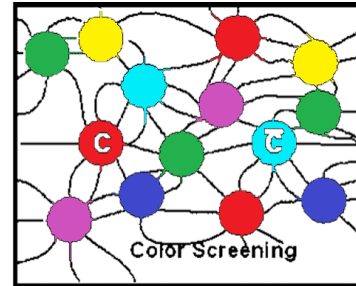
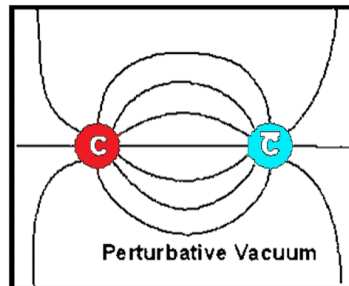


Outline

- 1 **Heavy quarkonia in heavy-ion collisions**
- 2 **The STAR experiment**
- 3 **Measurements of the Υ production at STAR**
- 4 **Comparison with LHC results and model calculations**
- 5 **Summary**

Heavy quarkonia in QGP

- $J/\psi, \Upsilon$ etc. are good candidates to probe QGP
 - $m_c \sim 1.3 \text{ GeV}/c^2, m_b \sim 4.2 \text{ GeV}/c^2 \rightarrow$ early production
 - Typically : $t_{\text{creation}}^{Q\bar{Q}} < t_{\text{creation}}^{QGP} < t_{\text{lifetime}}^{QGP} \ll t_{\text{lifetime}}^{HQ}$



- Suppression of production by colour screening

- Quarkonium expected to dissociate if its radius is greater than the Debye radius:

$$r_{\text{Debye}} \propto 1/T$$

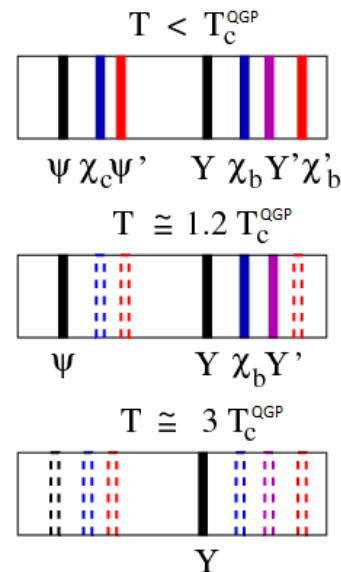
- Sequential melting

- Dissociation depends on the quarkonium binding energy
- Different states expected to melt at different temperatures
- QGP “thermometer”

T. Matsui, H. Satz, PLB 178 (1986) 416

A. Mocsy, EPJ C61 (2009) 705

$$R_{AA} \sim \frac{\#HQ \text{ in } AA}{\#HQ \text{ in } pp \times \# \text{ binary coll.}} = ?$$

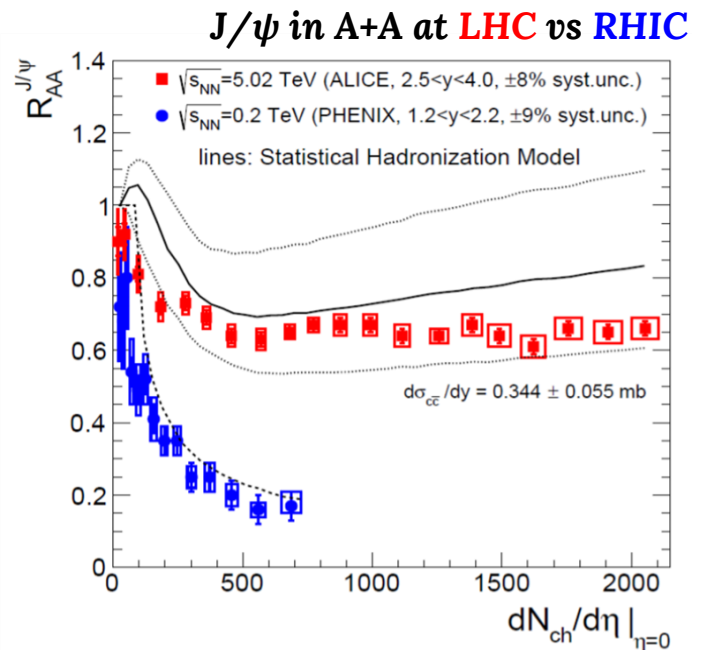
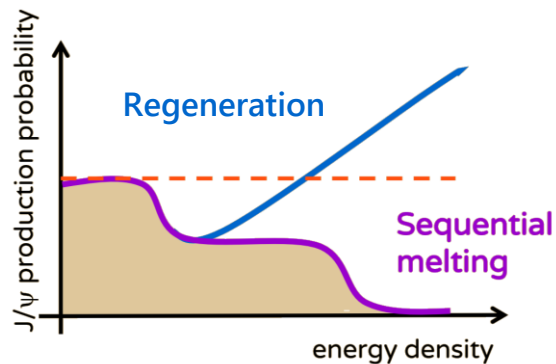
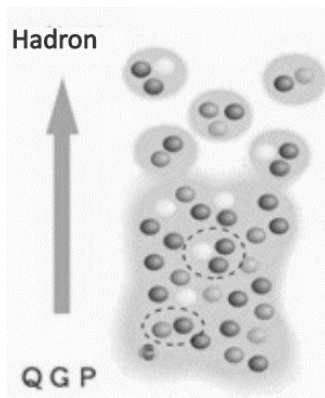


H. Satz, IJMPA 28 (2013) 1330043



Other effects also play a role

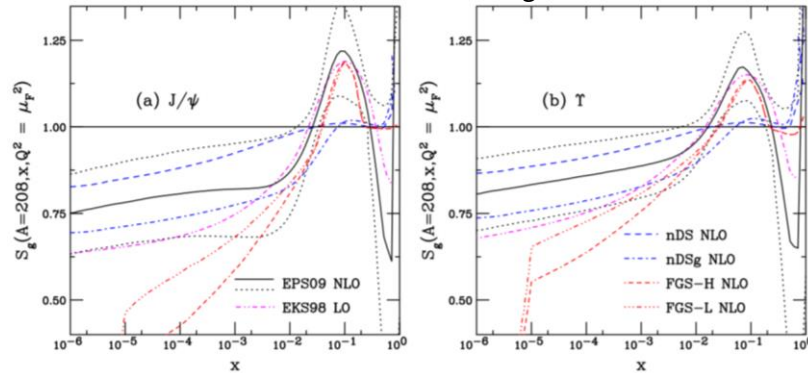
- Other phenomena complicate interpretation of the measured quarkonium suppression
- **Recombination**
 - Coalescence of deconfined quarks



Other effects also play a role

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- Recombination
 - Coalescence of deconfined quarks
- Cold nuclear matter (CNM) effects
 - Initial state: nuclear shadowing, energy loss

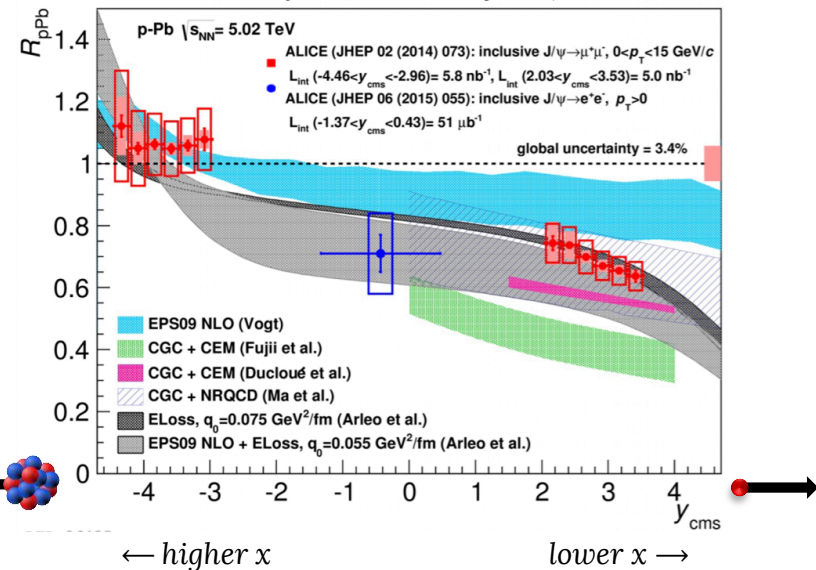
nuclear shadowing ratios



R. Vogt, PRC 92 (2015) 034909

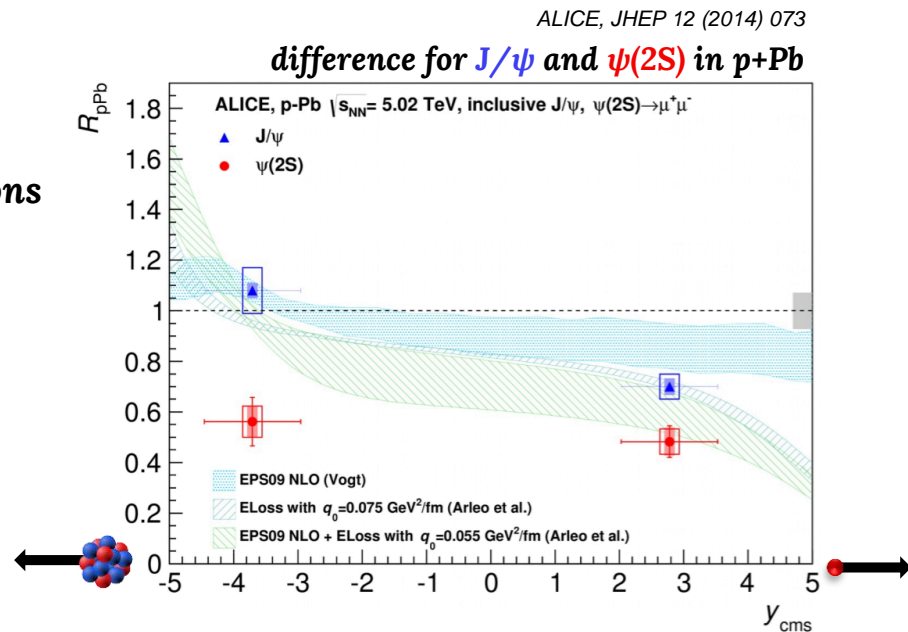
ALICE, JHEP 02 (2014) 073

rapidity dependence for J/ψ in p+Pb



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- Other phenomena complicate interpretation of the measured quarkonium suppression
- Recombination
 - Coalescence of deconfined quarks
- **Cold nuclear matter (CNM) effects**
 - Initial state: nuclear shadowing, energy loss
 - **Final state: inelastic interactions with hadrons**
 - nuclear break-up
 - co-mover absorption



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- **Feed-down**

$\Upsilon(1S) \ p_T > 8 \text{ GeV}/c \text{ at } \sqrt{s} = 1.8 \text{ TeV}$

CDF, PRL 84 (2000) 2094

Prompt $\Upsilon(1s)$	$\sim 51\%$
$\Upsilon(1s)$ from $\chi_b(1P)$ decays	$\sim 27\%$
$\Upsilon(1s)$ from $\chi_b(2P)$ decays	$\sim 10\%$
$\Upsilon(1s)$ from $\Upsilon(2S)$ decays	$\sim 11\%$
$\Upsilon(1s)$ from $\Upsilon(3S)$ decays	$\sim 1\%$



Why measure Υ at STAR?

I. Das, QM2015,
<https://indico.cern.ch/event/355454/contributions/838966/>

- Very **small production by recombination**

A. Emerick, X. Zhao, R. Rapp, EPJ A48 (2012) 72

	RHIC 200 GeV	LHC 2.76 TeV
# $c\bar{c}$ / event	13	115
# $b\bar{b}$ / event	0.1	3

- Precursor $b\bar{b}$ pair is more probable to survive propagation through a nuclear medium than $c\bar{c}$ since:

$$\sigma_{\text{eff}}^{\Upsilon} \sim \left(\frac{m_c}{m_b} \right)^2 \sigma_{\text{eff}}^{J/\psi} \simeq 0.1 \sigma_{\text{eff}}^{J/\psi} \quad \text{E. Ferreira, et al., PoS 157 (2012) 159}$$

- For co-movers (pions): $\sigma^{co-\Upsilon(1S)} \ll \sigma^{co-J/\psi} \rightarrow$ **break-up by co-movers insignificant**

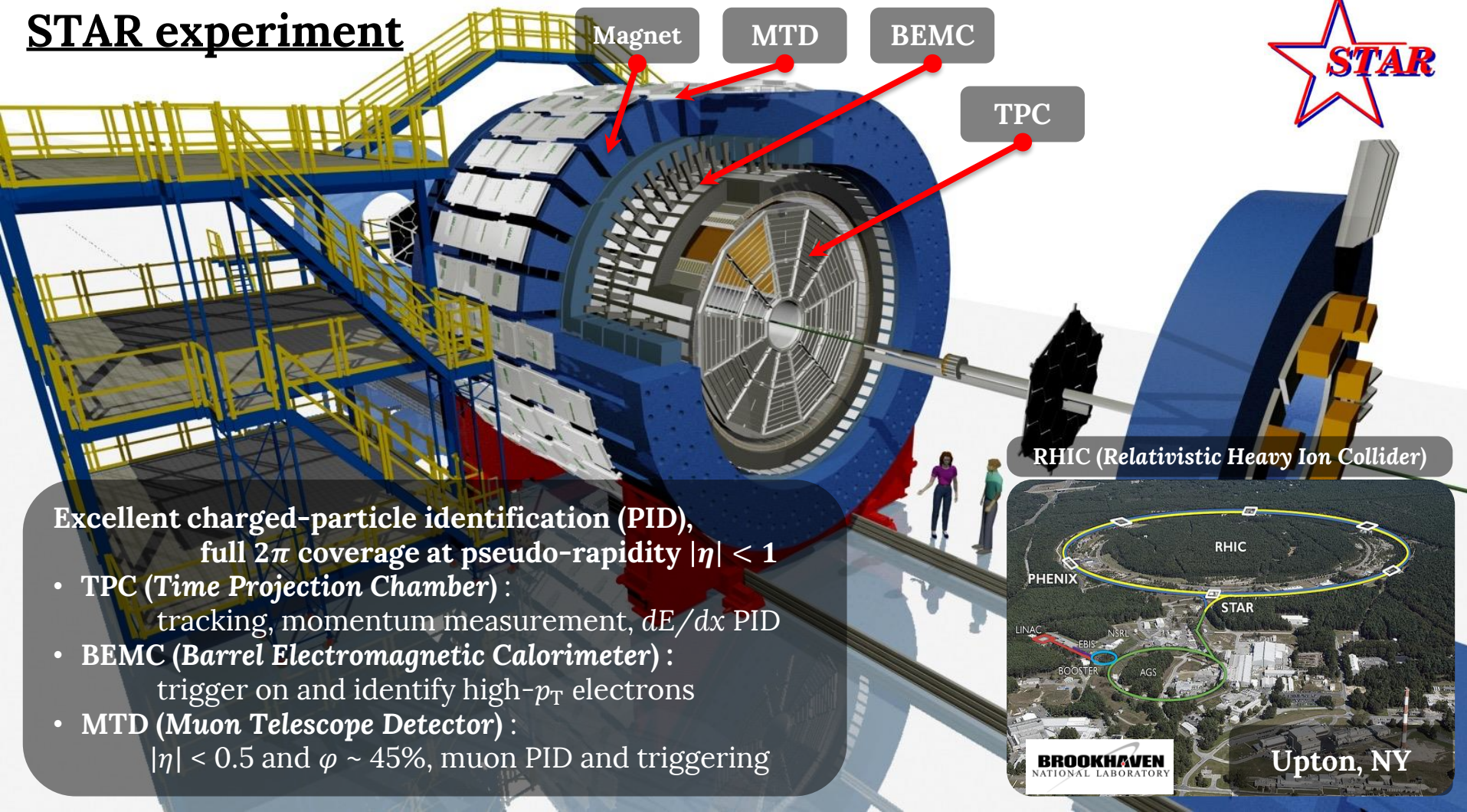
Z. Lin, C. Ko, PLB 503 (2001) 104

**$\rightarrow \Upsilon$'s at RHIC $\sqrt{s_{\text{NN}}} = 200$ GeV are a cleaner probe
of the screening effect!**

- Challenge: smaller production cross-section compared to the LHC



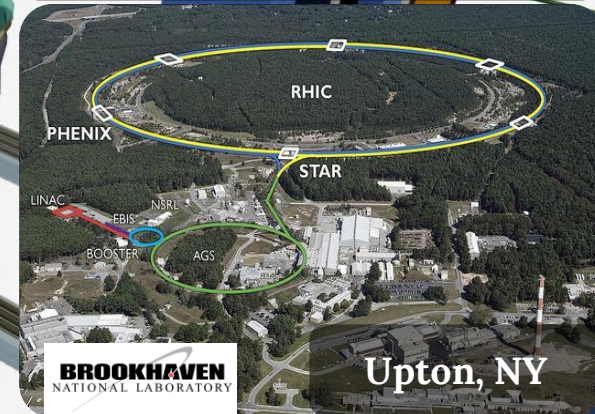
STAR experiment



Excellent charged-particle identification (PID),
full 2π coverage at pseudo-rapidity $|\eta| < 1$

- TPC (Time Projection Chamber) :
tracking, momentum measurement, dE/dx PID
- BEMC (Barrel Electromagnetic Calorimeter) :
trigger on and identify high- p_T electrons
- MTD (Muon Telescope Detector) :
 $|\eta| < 0.5$ and $\varphi \sim 45\%$, muon PID and triggering

RHIC (Relativistic Heavy Ion Collider)

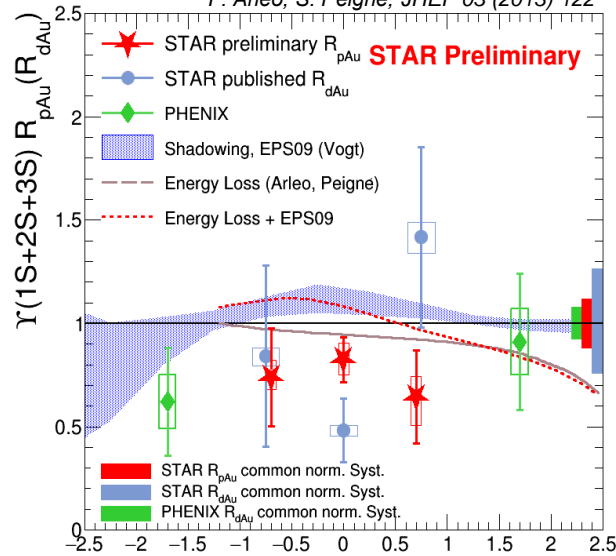
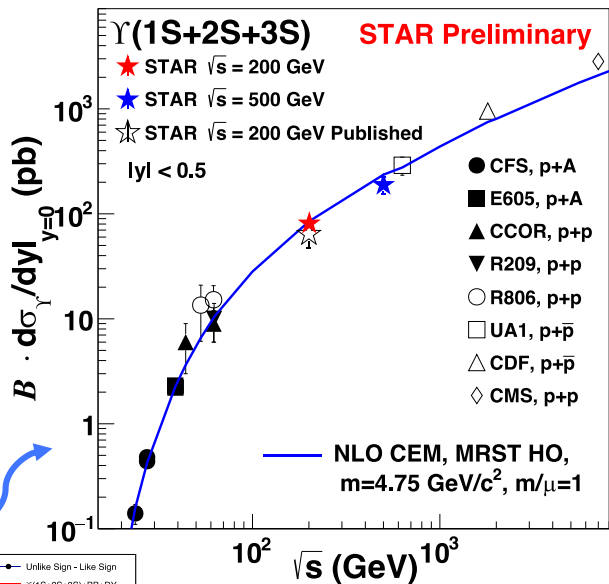


Results from p+p and p+Au collisions

STAR, PLB 735 (2014) 127

PHENIX, PRC 87 (2013) 044909

F. Arleo, S. Peigné, JHEP 03 (2013) 122

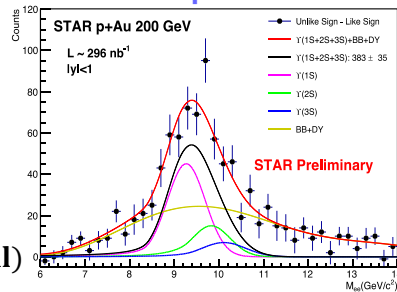
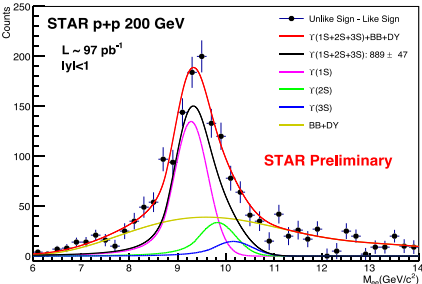


- p+p: precise baseline for comparison with Au+Au collisions

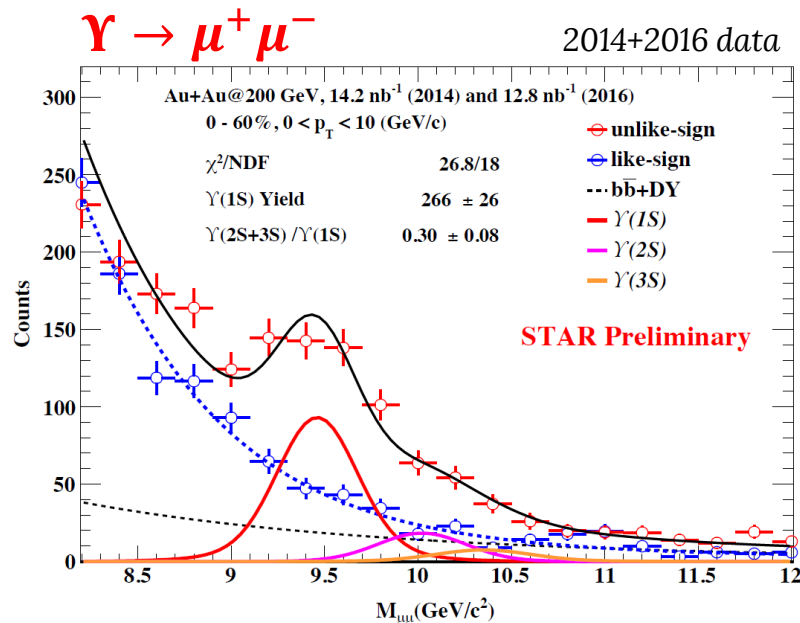
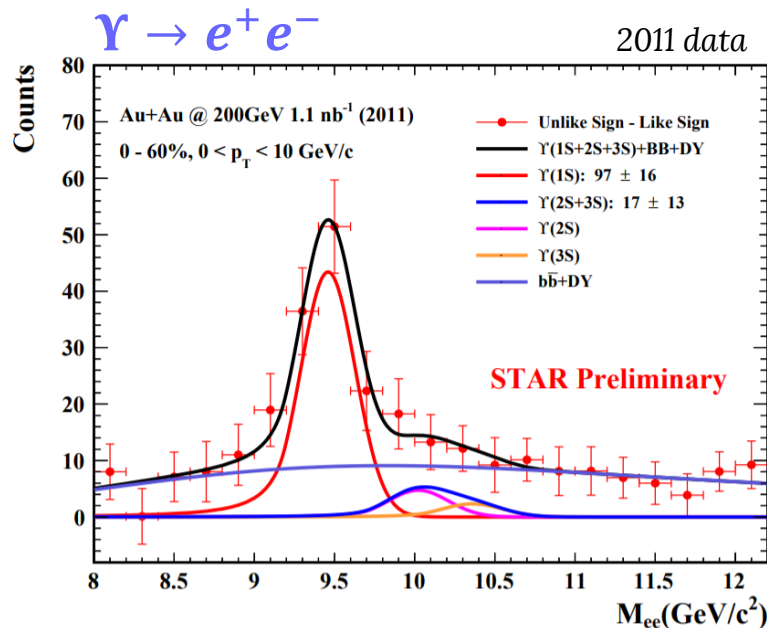
$$\sigma = 81 \pm 5 \text{ (stat.)} \pm 8 \text{ (syst.) pb}$$

- p+Au: quantification of CNM effects

$$R_{pAu}(|y| < 0.5) = 0.82 \pm 0.10 \text{ (stat.)} {}^{+0.07}_{-0.08} \text{ (syst.)} \pm 0.10 \text{ (global)}$$



Signal in Au+Au collisions



- Background sources:

→ combinatorial background (estimated as $N_{l^+l^+} + N_{l^-l^-}$)

→ Drell-Yan process, $B\bar{B}$ semi-leptonic decays



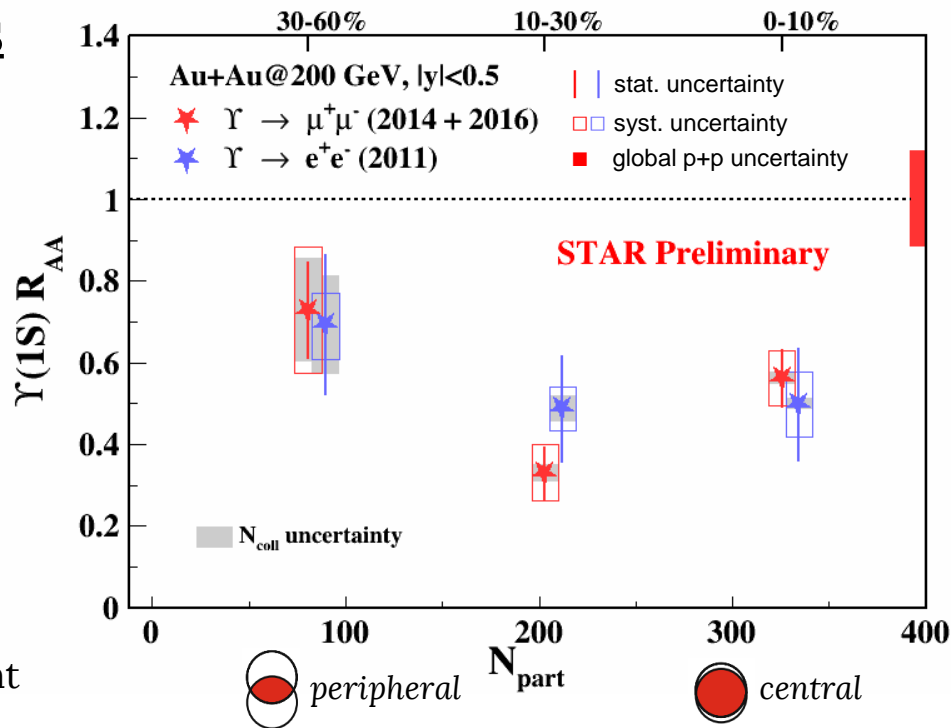
Results from Au+Au collisions

- Nuclear modification factor

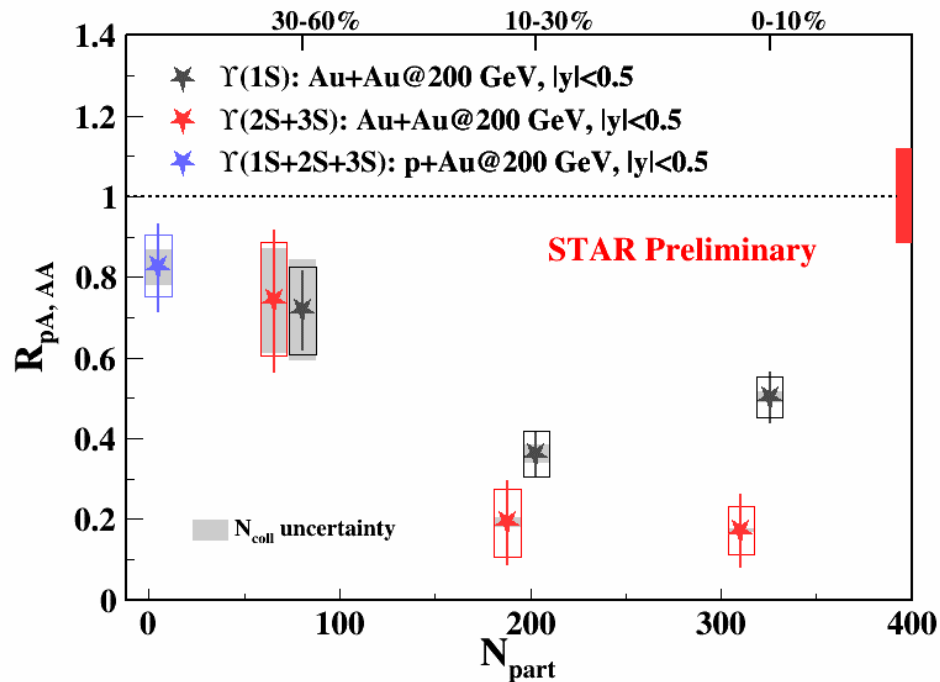
$$R_{AA} = \frac{\sigma_{\text{inel}}^{\text{pp}}}{\langle N_{\text{coll}} \rangle} \frac{d^2 N_{AA}/dp_T dy}{d^2 \sigma_{\text{pp}}/dp_T dy} \quad \text{as a function of mean number of participants } N_{\text{part}}$$

- σ_{pp} baseline taken from the Υ measurements in p+p collisions at STAR

- Di-muon** and **di-electron** results consistent with each other within the uncertainties
→ **results combined for increased statistical precision**



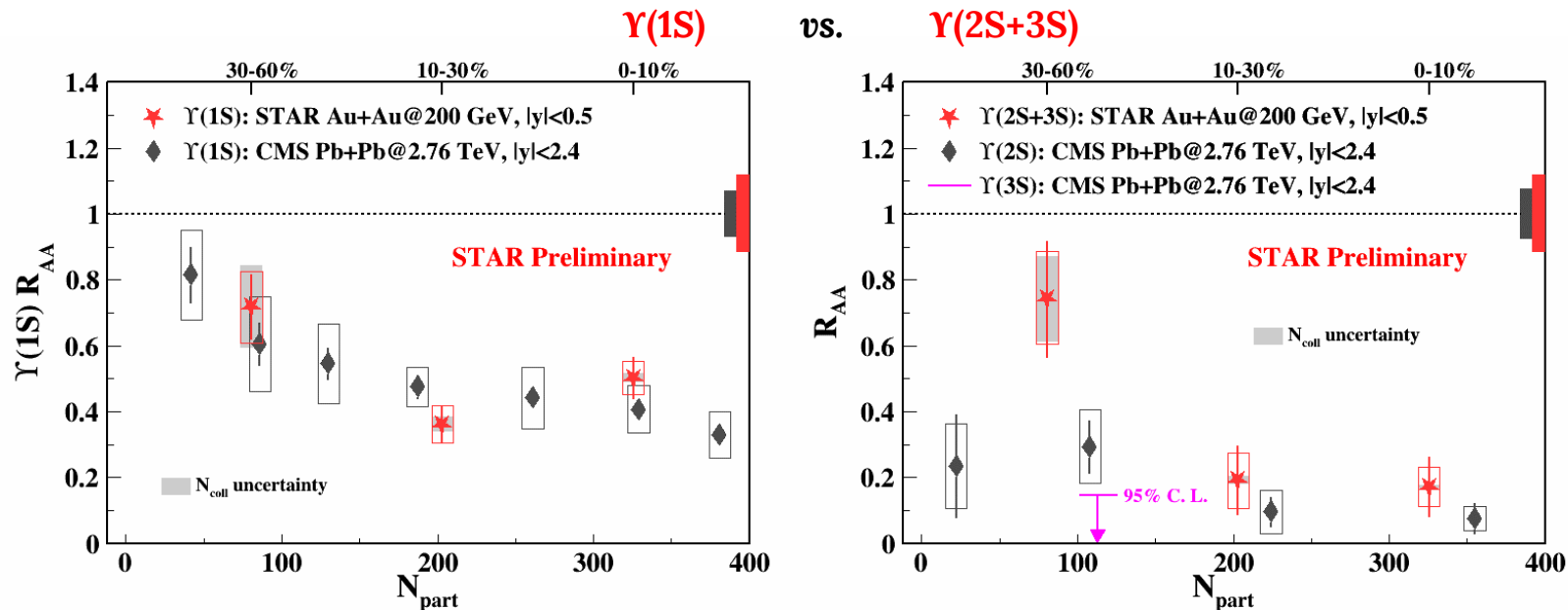
R_{AA} vs. N_{part} at RHIC



- $\Upsilon(2S+3S)$ *more suppressed* than $\Upsilon(1S)$ in central collisions



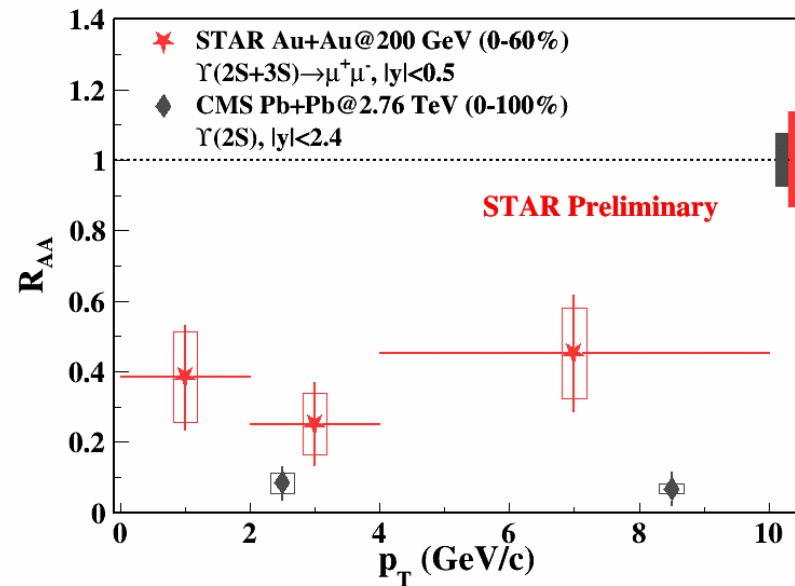
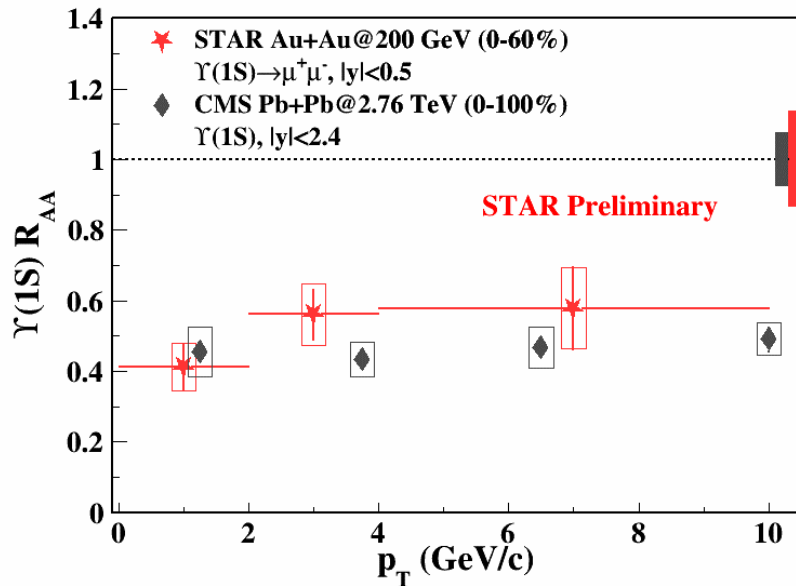
R_{AA} vs. N_{part} at RHIC + comparison with the LHC



- Comparison with the LHC: *CMS, PLB 770 (2017) 357*
 - comparable suppression for inclusive $\Upsilon(1S)$
 - hint of **less suppression** for $\Upsilon(2S+3S)$ at RHIC than at LHC



Suppression vs. p_T



RHIC vs. LHC

CMS, PLB 770 (2017) 357

- Comparable suppression for inclusive $Y(1S)$
- Signs of **less suppression** at high- p_T for $Y(2S+3S)$ at RHIC



Comparison with models

- **Krouppa, Rothkopf, Strickland:**

PRD 97 (2018) 016017

- No CNM effects, no regeneration
- Uses a lattice QCD vetted complex potential embedded in a hydrodynamically evolving medium

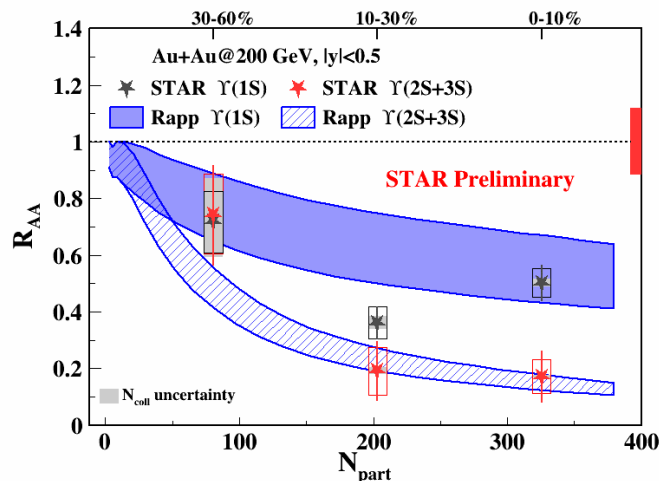
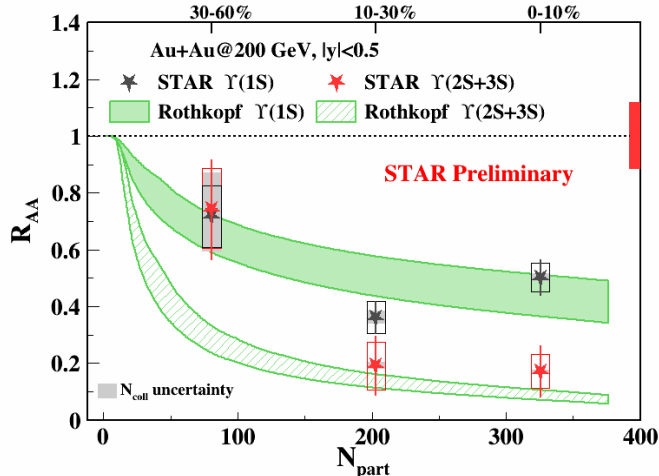
- **Du, He, Rapp:**

PRC 96 (2017) 054901

- Incorporates regeneration as well as CNM effects
- Uses in-medium binding energies predicted by thermodynamic T-matrix calculations with internal-energy potentials

→ Both models show agreement with the $\Upsilon(1S)$ data from STAR

→ Rothkopf model seems to underestimate the $\Upsilon(2S+3S)$ R_{AA} in the 30-60% centrality



Summary

- **Physics of heavy quarkonia**
 - Quarkonia can be used to constrain the QGP temperature thanks to the *colour screening effect*
 - Υ 's at RHIC are a cleaner probe than e.g. the J/ψ , due to lesser influence of CNM effects and regeneration
- **Υ measurements at STAR**
 - Υ 's are reconstructed in the *di-electron channel* and *di-muon channel* in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV
 - Υ 's are also measured in p+p and p+Au collisions to provide a baseline and a quantification of the CNM effects
 - Inclusive $\Upsilon(1S)$ at RHIC are strongly suppressed in semi-central and central collisions
 - Excited states $\Upsilon(2S+3S)$ at RHIC are more suppressed than $\Upsilon(1S)$ in central collisions (*sequential melting*)
 - R_{AA} for $\Upsilon(1S)$ at RHIC is similar to the LHC, but $\Upsilon(2S+3S)$ seems to be less suppressed
- These results can be used to test models of different quarkonium behaviour in QGP and to help infer the medium temperature

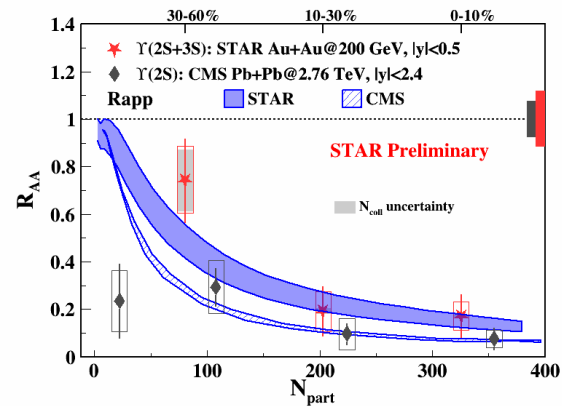
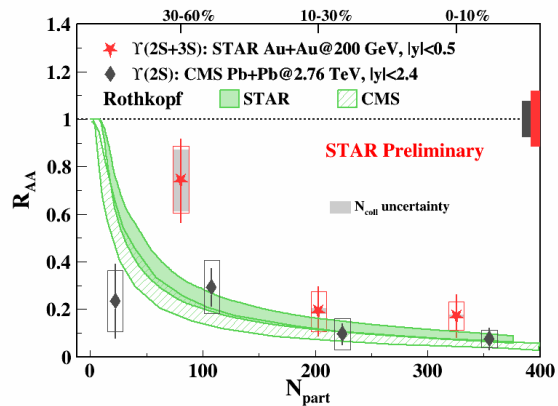
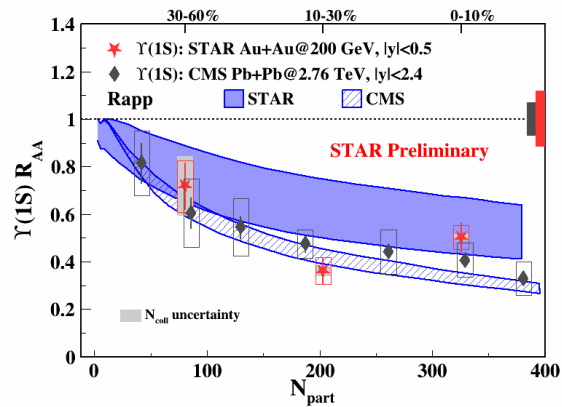
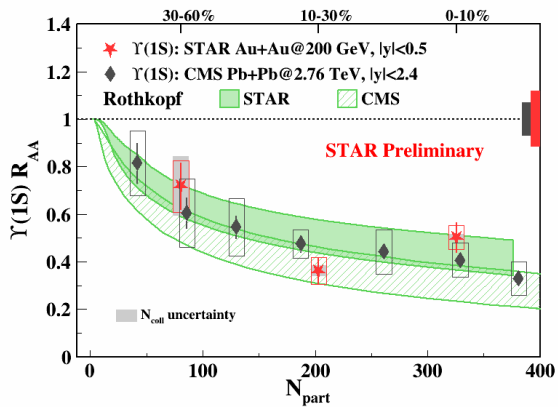
Thanks for your attention!



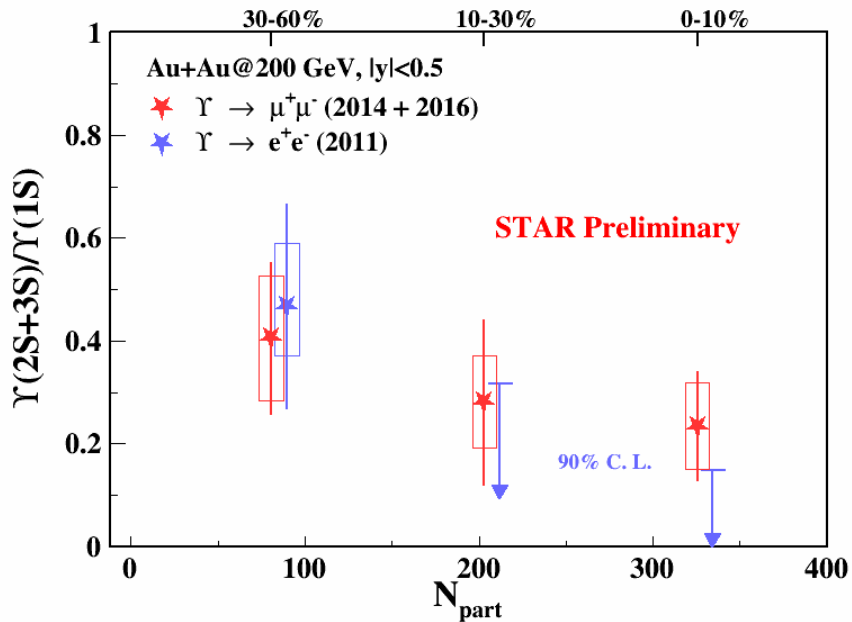


Back-up slides

Models at the LHC



Excited-to-ground ratio



world-wide measured ratio in p+p:

$$\Upsilon(2S+3S)/\Upsilon(1S) = 0.423 \pm 0.009$$

PRC 88 (2013) 067901

